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Biophysical modeling of in vitro and in vivo processes underlying regulated photoprotective mechanism in cyanobacteria

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Abstract Non-photochemical quenching (NPQ) is a mechanism responsible for high light tolerance in photosynthetic organisms. In cyanobacteria, NPQ is realized by the interplay between light-harvesting complexes, phycobilisomes (PBs), a light sensor and effector of NPQ, the photoactive orange carotenoid protein (OCP), and the fluorescence recovery protein (FRP). Here, we introduced a biophysical model, which takes into account the whole spectrum of interactions between PBs, OCP, and FRP and describes the experimental PBs fluorescence kinetics, unraveling interaction rate constants between the components involved and their relative concentrations in the cell. We took benefit from the possibility to reconstruct the photoprotection mechanism and its parts in vitro, where most of the parameters could be varied, to develop the model and then applied it to describe the NPQ kinetics in the *Synechocystis* sp. PCC 6803 mutant lacking photosystems.

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Our analyses revealed that while an excess of the OCP over PBs is required to obtain substantial PBs fluorescence quenching in vitro, in vivo the OCP/PBs ratio is less than unity, due to higher local concentration of PBs, which was estimated as $\sim 10^{-5}$ M, compared to in vitro experiments. The analysis of PBs fluorescence recovery on the basis of the generalized model of enzymatic catalysis resulted in determination of the FRP concentration in vivo close to 10% of the OCP concentration. Finally, the possible role of the FRP oligomeric state alteration in the kinetics of PBs fluorescence was shown. This paper provides the most comprehensive model of the OCP-induced PBs fluorescence quenching to date and the results are important for better understanding of the regulatory molecular mechanisms underlying NPQ in cyanobacteria.

Keywords Cyanobacteria · Non-photochemical quenching · Orange carotenoid protein · Fluorescence recovery protein · Biophysical model · Fluorescence

Introduction

Oxygenic photosynthetic organisms encounter a significant problem when they tame the energy of light and convert it into more usable forms such as the energy of chemical bonds. This problem deals with potentially hazardous excessive light conditions, in which the photosynthetic apparatus needs to be safely protected by dissipation of the excessive light energy. In cyanobacteria, the light energy is absorbed by water-soluble light-harvesting complexes, phycobilisomes (PBs), and normally reaches the reaction centers of the photosystem II on the initial steps of photosynthesis. Under supra-optimal light conditions, due to the strictly controlled process of non-photochemical quenching (NPQ), it becomes possible to reduce the amount of energy being delivered to the reaction centers (Kirilovsky [2015](#page-9-0)). This is thought to substantially prevent formation of damaging oxidation products, reactive oxygen species (Bartosz [1997](#page-9-1); Sedoud et al. [2014](#page-10-0)), but requires fine-tuned controlling of the components of NPQ. In vitro experiments (Gwizdala et al. [2011](#page-9-2)), aimed at recapitulation of processes known to occur in vivo, revealed that regulated photoprotective mechanism in cyanobacteria is realized by the set of three basic components: (1) light-harvesting antennae (PBs), (2) a sensor of elevated light levels and an effector which interacts with PBs and quenches their excessive fluorescence (the orange carotenoid protein, OCP), (3) the fluorescence recovery protein, FRP, which is required for abrogation of the OCP action and decouples the quenched OCP/ PBs system. Although different aspects of structural and functional organization of PBs, OCP, and FRP, have been successfully studied (Kerfeld et al. [2003](#page-9-3); Sutter et al. [2013](#page-10-1); Liu et al. [2013;](#page-10-2) Leverenz et al. [2015](#page-10-3); Gupta et al. [2015](#page-9-4); Sluchanko et al. [2017\)](#page-10-4), the complete picture of NPQ regulation is far from being fully understood, and some details of how this ternary OCP–PBs–FRP system operates require further intensive research.

It was proposed that as the major NPQ mediator the OCP exhibits a modular structure, with the two domains, N- and C-terminal (NTD and CTD), that hold a carotenoid molecule in the interdomain cavity (Kerfeld et al. [2003](#page-9-3)). The NTD and CTD, connected by a flexible linker, are believed to have a distinct, albeit still debatable, specialization: the NTD is suggested to be responsible for PBs binding, while the CTD is responsible for the interaction with FRP (Lev-erenz et al. [2015\)](#page-10-3). Despite the structures of individual components are relatively well studied (Kirilovsky and Kerfeld [2016](#page-10-5)), the field desperately requires structural knowledge about OCP/FRP and OCP/PBs complexes and currently relies on often indirect biochemical and biophysical information.

Figure [1](#page-1-0) summarizes our current understanding of the OCP, PBs, and FRP involvement in the PBs fluorescence quenching process from the kinetics point of view. During the first stage (see Fig. [1\)](#page-1-0), upon absorption of a blue–green light quantum by the OCP's carotenoid chromophore, inactive orange OCP form (OCP^O) is converted to an active red state (OCPR) (Wilson et al. 2008). During the second stage, the OCP^R forms a stable complex with the PB, leading to an efficient quenching of the excess excitation energy (Tian et al. [2011,](#page-10-7) [2012](#page-10-8); Maksimov et al. [2014\)](#page-10-9) and ultimately reducing the amount of excitation delivered to the reaction centers of photosystem II and protecting from formation of hazardous reactive oxygen species. Alternatively, OCP^R can spontaneously reconvert back to its inactive OCPO, or form a complex with the FRP (stage 7 in Fig. [1](#page-1-0)), which accelerates the conversion of the OCPR to OCP^O (stage 6) in Fig. [1](#page-1-0)) significantly (Boulay et al. [2010;](#page-9-5) Sluchanko et al. [2017](#page-10-4)). Finally, under dark/low light conditions, the OCP dissociates from PBs with a rate k_d . The FRP may also bind to the OCP–PB complex (stage 3 in Fig. [1](#page-1-0)), thus accelerating the OCP detachment from PBs (stage 5 in Fig. [1](#page-1-0)) and restoring an efficient energy flow from PBs to photosystem II reaction centers.

An important factor in the context of the PBs fluorescence kinetics is the oligomeric state of the involved components, the OCP and FRP in particular. For example, it was shown that the FRP forms stable dimeric

Fig. 1 Current understanding of the OCP-induced PBs fluorescence quenching process in cyanobacteria mediated by the interplay of the OCP, FRP, and PBs. The following stages are indicated: (*1*) the OCP photoactivation and spontaneous reconversion, (*2*) the OCPR binding to PB, (*3*) binding of the FRP dimer to the $OCP^R - PB$ complex, (*4*) monomerization of the FRP and stabilization of the OCP–PB–FRP complex, (*5*) the FRP-assisted detachment of the OCP from PB, (*6*) the FRP-assisted conversion of the OCP to its stable orange state, (*7*) binding of the FRP dimer to the OCPR, (*8*) monomerization of the FRP and stabilization of the FRP–OCP complex. See text for details

structures (Sutter et al. [2013](#page-10-1); Sluchanko et al. [2017](#page-10-4)) but forms complexes with the OCP at a 1:1 stoichiometry, implying dissociation of the FRP dimers upon interaction with the OCP (Sluchanko et al. [2017\)](#page-10-4). Most likely, the interaction between the OCP and FRP involves (i) initial binding of the FRP dimer to the OCP^R with a consequent detachment of the FRP monomer (stage 7 in Fig. [1](#page-1-0)), which may have a different binding affinity to the OCP^R , and (ii) further stabilization of the FRP/OCP complex (stage 8 in Fig. [1](#page-1-0)). The same stages of the FRP monomerization could be assumed for its interaction with the $OCP^{R}-PB$ complexes (stages 3 and 4 in Fig. [1](#page-1-0)). There are also observations in favor of the OCP self-association, especially at higher protein concentrations. It was suggested that the OCP may also exist in the dimeric form (Zhang et al. [2013\)](#page-10-10), and in this case the OCP monomerization caused by its photoactivation must be taken into account when describing kinetics of NPQ. However, it seems likely that at low (micromolar) concentrations the OCP^O and OCP^R are both monomeric (Maksimov et al. [2016;](#page-10-11) Sluchanko et al. [2017](#page-10-4)).

As two steps are required for PBs fluorescence quenching (Fig. [1\)](#page-1-0), under exposure to strong light, two stages can be detected in PBs fluorescence quenching upon exposure to strong light—light-dependent and light-independent (or "dark") (Gorbunov et al. [2011\)](#page-9-6). The lightdependent phase is detected on the onset of illumination and reflects binding of the activated OCP^R to PBs. Its rate is determined by the slower of the two processes: the OCP^O to OCP^R conversion (high light intensities) or binding of the OCP^R to PB (low light intensities). The "dark" phase of fluorescence quenching is observed after termination of strong illumination and reflects the binding of a portion of the activated OCP^R to PBs, which leads to further quenching of PBs fluorescence. The typical rate of the dark phase PBs fluorescence quenching (k_{dark}) is ~0.1 s⁻¹ (Gorbunov et al. [2011;](#page-9-6) Kuzminov et al. [2012\)](#page-10-12). The fluorescence recovery kinetics has been studied in less detail; however, it was shown that this process is strongly non-linear with its rate depending on the initial quenching level (Maksimov et al. [2015b\)](#page-10-13). However, a detailed biophysical model of processes, underlying PBs fluorescence kinetics during NPQ, is lacking.

This study is aimed at comparative kinetic analysis of the PBs fluorescence quenching processes in cyanobacteria in vivo and in vitro, as well as at determination of relative concentrations of the components involved in the regulated photoprotective mechanism in vivo. Here we demonstrate that the ratios of the OCP, FRP, and PBS in vivo, as well as the origin of processes responsible for quenching, are significantly different than that in vitro, suggesting the predominant role of the intracellular microenvironment on the rates of protein diffusion and interaction.

Materials and methods

In vivo experiments

The wild-type cells of the cyanobacterium *Synechocystis* sp. PCC 6803 were grown in modified BG-11 medium at 30° C constant white light illumination of 40 μ mol photons m−2 s−1. A ΔPSI/ΔPSII (Δ*psaAB*/Δ*psbDIC*/Δ*psbDII*) mutant of *Synechocystis* sp. PCC 6803, lacking photosystems I and II (Ermakova-Gerdes et al. [2015](#page-9-7)), was cultivated in BG-11 medium, containing glucose (15 mM), spectinomycin (25 µg/ml), erythromycin (20 µg/ml), and chloramphenicol (20 μ g/ml) at light exposure of 5 μ mol photons m^{-2} s⁻¹ as described in (Kuzminov et al. [2014;](#page-10-14) Maksimov et al. [2015b\)](#page-10-13).

The samples were excited with 630 nm light, and fluorescence emission was measured in a 90° scheme using a Maya2000 Pro spectrometer (Ocean Optics), allowing fluorescence kinetic measurements with 0.1 s resolution. Nonphotochemical quenching was induced simultaneously with a 3000 µmol photons m⁻² s⁻¹ 455 nm LED.

Fluorescence recovery kinetics measured after prolonged illumination, which exceeded the times required to reach the maximal quenching $(>200 \text{ s})$, all had similar shape. This indicates the absence of irreversible photodegradation in the system during the experiments.

In vitro experiments

The OCP, FRP, and PBs preparations for the in vitro experiments were obtained as described previously (Sluchanko et al. [2017](#page-10-4)). N-terminally 6xHis-tagged wild-type orange carotenoid protein (OCP) from *Synechocystis* sp. PCC 6803 was expressed in *E. coli* NEB DH5α Turbo cells carrying the whole machinery for carotenoid synthesis and OCP/FRP were purified by a combination of immobilized metal-affinity and size-exclusion chromatography (SEC) as described in (Maksimov et al. [2016\)](#page-10-11). Protein concentrations were determined on a Lambda-25 spectrophotometer (Perkin-Elmer, USA) using molar extinction coefficients $\varepsilon^{\lambda = 280 \text{ nm}} = 15,220 \text{ M}^{-1} \text{ cm}^{-1}$ for FRP and $\varepsilon^{\lambda = 495 \text{ nm}} = 63,000 \text{ M}^{-1} \text{ cm}^{-1}$ for the OCP. PBs purification and concentration measurement was conducted according to the procedures described in (Stadnichuk et al. [2013](#page-10-15)).

Experiments on PB fluorescence quenching were conducted in a 0.75 M phosphate buffer at pH 7.4, the temperature was fixed to 25 °C. A M455L3 (Thorlabs, USA) 900 mW light-emitting diode with maximum emission at 455 nm was used for blue–green illumination of the samples (i.e., OCP photoswitching and PBs quenching induction). The steady-state fluorescence measurements were performed using FluoroMax-4 spectrofluorimeter (Horiba Jobin Yvon, Japan-France).

Results and discussion

OCP photoconversion and interaction with FRP in vitro

Although it includes three basic components, the model of OCP-induced PBs quenching in cyanobacteria (Fig. [1\)](#page-1-0) considers several transient complexes: OCPR, OCPR–FRP, $OCP^{R}-PB$, and $FRP-OCP^{R}-PB$, whose populations are governed by more than 10 parameters. In vitro experiments allow us to simplify this model by dividing it into several blocks of processes, which can be studied independently: (i) OCP photoconversion, (ii) OCP–FRP interactions, (iii) OCP–PB interaction, and (iv) FRP–OCP–PB interaction. The number of parameters in each sub-model is greatly reduced, thus permitting us to improve the precision of retrieving characteristics of each process.

First, we examined the in vitro kinetics of light-induced alterations of the OCP within the two-component system of OCP and FRP. Under illumination with strong blue–green light, the OCP exhibits a transition to its active state and the rate of this process is proportional to light intensity *F* and photoconversion cross-section *σ*:

$$
OCP^{O} \xrightarrow{F\sigma} OCP^{R}.
$$
 (1)

After the offset of illumination, reconversion to the OCP^O state occurs with the rate k_r :

$$
OCP^{R} \xrightarrow{k_r} OCP^{O}.
$$
 (2)

The combination of these two processes leads to the following equation describing the reversible photoconversion kinetics of the OCP (Maksimov et al. [2015a\)](#page-10-16):

$$
\frac{\partial R}{\partial t} = F\sigma(1 - R) - k_{\rm r}R,\tag{3}
$$

where R is the concentration of the OCP^R normalized to $[OCP]_T$. Hereinafter, $[OCP]_T$, $[FRP]_T$, and $[PBS]_T$ are the total concentrations of the OCP, FRP, and PBs, respectively.

Using this equation, kinetics of light-induced OCP transition was analyzed at different ionic strength, namely, 0.15 and 0.75 M phosphate buffer. The latter buffer was used further in the in vitro experiments to maintain PBs stability. In accordance with (Gwizdala et al. [2011](#page-9-2)), we observed that an increased ionic strength had a significant influence on both conversion level and rates, namely, only 10% of the OCP could be photoactivated in 0.75 M phosphate buffer (Fig. S1). The OCP^R relaxation rate in the 0.75 M phosphate buffer was k_r =0.04 s⁻¹. The photoconversion cross section was similar for both ionic strengths and equal to $\sigma = (2.0 \pm 0.2) \times 10^{-19}$ cm2 , that is in accordance with literature data (Maksimov et al. [2015a\)](#page-10-16).

In the presence of the FRP, the following processes must be taken into account to quantify the kinetics of OCP conversion: FRP binding to OCP^R , which occurs with a rate constant k_{FRP} ,

$$
OCP^{R} + FRP \xrightarrow{k_{FRP}} OCP^{R} - FRP,
$$
\n(4)

and the subsequent decay of the OCP^R –FRP complex (i.e., the accelerated FRP-assisted conversion of OCP^R to OCP^O) with a back conversion rate k_C :

OCP^R – FRP
$$
\xrightarrow{k_C}
$$
 OCP^O + FRP. (5)

This model is represented as a scheme in Fig. [2](#page-3-0)a and leads to the following system of equations describing the kinetics of light-induced OCP transitions in the presence of the FRP:

$$
\frac{\partial R}{\partial t} = F\sigma(1 - R - K) - k_{\text{FRP}}^* R(\beta - K) - k_r R
$$
\n
$$
\frac{\partial K}{\partial t} = k_{\text{FRP}}^* R(\beta - K) - k_{\text{C}} K,
$$
\n(6)

Fig. 2 a The scheme of OCP–FRP interaction, **b** the OCP conversion kinetics measured at different FRP concentrations. The conversion level in the absence of FRP was taken as 1. *Black lines* correspond to the approximation with Eq. ([6\)](#page-3-1). $[OCP]_T = 4 \times 10^{-6}$ M

where *R* and *K* are the concentrations of the OCP^R and OCP^R–FRP complexes normalized to $[OCP]_T$, $\beta = [FRP]_T$ / $[OCP]_T$, and $k_{\text{FRP}}^* = k_{\text{FRP}}[OCP]_T$.

Figure [2b](#page-3-0) shows the OCP photoconversion kinetics at different β values, i.e.. at different $[FRP]_T$. By fitting the measured kinetics to the model (6) (6) , we have obtained the following parameters: $k_{\text{FRP}}^* = 0.1 \text{ s}^{-1}$, $k_C = 0.15 \text{ s}^{-1}$. These results clearly suggest that the presence of the FRP accelerates the OCPR conversion to OCP^O in darkness by a factor of ~4, which is in agreement with the literature data (Boulay et al. [2010;](#page-9-5) Gwizdala et al. [2011,](#page-9-2) [2013;](#page-9-8) Sluchanko et al. [2017\)](#page-10-4). Interestingly, (Sluchanko et al. [2017](#page-10-4)) estimated the value of the dissociation constant for the OCP^R –FRP complex $(K_d \approx 2 \mu M)$ based on data obtained for the purple OCP-W288A mutant, which was shown to be an analogue of the red form of OCP. Given that by definition K_d is the inverse ratio of binding (k_{FRP}) and OCP^R–FRP back conversion (k_C) rate constants, the obtained values result in a 0.15 s⁻¹/(0.1 s⁻¹/4 × 10⁻⁶ M) = 6 μM value of K_d , which is in a qualitative agreement with the data of (Sluchanko et al. [2017](#page-10-4)). Therefore, modeling of the kinetic curves allowed us to obtain independent estimations of equilibrium constants.

Kinetics of PB fluorescence quenching in vitro

As shown in (Gwizdala et al. [2011\)](#page-9-2), reconstitution of the OCP-mediated photoprotective mechanism in cyanobacteria is possible in vitro using three components: OCP, PBs, and FRP. However, modeling of the initial phase of photoprotection, i.e., PBs fluorescence quenching, can also be done on the two-component system of OCP and PBs.

For the two-component OCP–PBs system, the lightinduced PBs quenching kinetics can be fully described by the model that takes into account the following processes: OCP photoconversion and relaxation, described by Eqs. ([1](#page-3-2)) and ([2](#page-3-3)), respectively; binding of the OCPR to PBs and formation of the OCP^R –PB complexes in the 1:1 stoichiometry (Wilson et al. [2006\)](#page-10-17) with a rate constant k_{PB} :

$$
OCP^{R} + PBs \xrightarrow{k_{PB}} OCP^{R} - PBs,
$$
 (7)

and decay of the OCP^{R} —PB complexes to the OCP^{O} and free PB with the rate constant k_d :

$$
OCP^{R} - PB \xrightarrow{k_d} OCP^{R} + PB.
$$
 (8)

Under the assumption of complete PBs fluorescence quenching upon their interaction with the OCP^R , the measured fluorescence signal is proportional to the number of free phycobilisomes, i.e., $([PBs]_T - [OCP_R - PBs])$. The processes are presented in Fig. [3a](#page-4-0) and can be described by the system of two differential equations:

$$
\frac{\partial R}{\partial t} = F\sigma \left(1 - R - \frac{C}{\alpha} \right) - k_{\text{PB}}^* R (1 - C) - k_{\text{r}} R
$$
\n
$$
\frac{\partial C}{\partial t} = k_{\text{PB}}^* R (1 - C) - k_{\text{d}} C. \tag{9}
$$

Here, R and C stand for the concentrations of OCP^R and OCP^R–PBs complexes normalized to $[PBS]_T$. Impor-tantly, Eqs. ([9\)](#page-4-1) contain the term $\alpha = [OCP]_T/[PBs]_T$, which can be used for estimation of the relative OCP concentration from a set of fluorescence quenching curves. Also, the rate constant $k_{\text{PB}}^* = k_{\text{PB}}[\text{PBs}]_T$. We note that the term that describes interaction between OCP^R and PB in the model ([9](#page-4-1)), is different from that described in (Gorbunov et al. [2011\)](#page-9-6) as it takes into account the depletion of free PBs with time caused by their occupation with the

Fig. 3 a The scheme of PBs fluorescence quenching by the OCP. **b** Light-induced PBs fluorescence quenching obtained at different $[OCP]_{T'}$ $[PBS]_T$ ratio at total PBs concentration $[PBS]_T = 5 \times 10^{-11}$ $[PBS]_T = 5 \times 10^{-11}$ $[PBS]_T = 5 \times 10^{-11}$ M and their fits with model (1). The curves were normalized to maximum

OCPR. Description of the dark phase of PB fluorescence quenching can be readily achieved with Eqs. [\(9](#page-4-1)) by setting *F* to zero.

Figure [3](#page-4-0)b shows the kinetics of PBs fluorescence quenching under exposure to strong light (3000 μ M s⁻¹ cm⁻² 455 nm) obtained in the samples with different $[OCP]_{T}$ / $[PBS]_T$ ratios. On the offset of the illumination, a "dark" phase of PBs fluorescence quenching is observed (Fig. S2), which is due to interaction of the OCP molecules, activated during the light phase, with PBs (Gorbunov et al. [2011](#page-9-6)). Fluorescence quenching was measured for varying $[OCP]_{T}$ / $[PBS]_T$ for three fixed $[PBS]_T$ values until fluorescence signal reached a steady-state, that corresponded to a stationary concentration of $OCP^{R}-PBs$ complexes. As shown in Fig. [3b](#page-4-0), an efficient PBs fluorescence quenching was observed only at a significant excess of the OCP over PBs, $[OCP]_T > [PBs]_T$, that is in agreement with the results of (Gwizdala et al. [2011](#page-9-2); Sluchanko et al. [2017](#page-10-4)).

The maximum possible number of OCP^R –PBs complexes is $[PBS]_T$ if the OCP^R decoupling is neglected. Under such conditions, zero fluorescence would be observed for 100% quenching efficiency, however, only a 50–60% quenching occurred in our experiments. This relatively small amplitude of fluorescence quenching could be due to the presence of a fraction of free phycocyanine rods, which are a part of PBs, in the solution, which cannot be quenched by the OCP (Gwizdala et al. [2011;](#page-9-2) Tian et al. [2012\)](#page-10-8). This suggestion is supported by the fact that the spectral maximum of PBs fluorescence emission without the OCP was at 659 nm—in comparison to the previously published steady-state fluorescence emission spectra (Gwizdala et al. [2011](#page-9-2)), the observed spectrum was blue-shifted by more than 10 nm. It suggests that the PBs used in the in vitro experiments could be partly dissociated, and a mixture of PBs species was investigated. Moreover, incomplete quenching could be caused by less than 100% excitation energy transfer from PBs to the Histagged OCP due to the geometry of OCP^R –PBs complex. To take this fact into account while modeling the quenching kinetics, we used the following expression for fluorescence intensity estimation:

$$
I \sim 1 - (1 - \varphi)\mathcal{C},\tag{10}
$$

where ϕ is the fluorescence quantum yield of the OCP^R–PBs complexes. To assess the value of ϕ , the dependence of stationary fluorescence levels on $[PBS]_{T}$ / $[OCP]_T$ was obtained for each series of measurements and then extrapolated to zero, where fluorescence corresponded to an infinite OCP concentration, i.e., to the impact of the non-fully quenched complexes.

We performed the global approximation of PBs fluorescence quenching kinetics, including both light and dark phases, recorded at 10 varying values of $\alpha = [OCP]_{T}/[PBs]_{T}$ with the biophysical model (Eqs. [9,](#page-4-1) [10\)](#page-5-0) for three different $[PBS]_T$ concentrations. As the OCP^R reconversion rate k_r =0.04 s⁻¹ and the photoconversion cross-section σ were determined in independent experiments (Fig. [2](#page-3-0)b), the fitting procedure resulted in determination of two remaining parameters, k_{PB}^* and k_d (Table [1](#page-5-1)).

Our results (Table [1](#page-5-1)) revealed the following trends. First, the rate k_{PB}^* , which determines the OCP interaction with free PB, is proportional to $[PB]_T$. This dependence is indicative of a predominant role of diffusion in this process. Second, the k_d values are virtually independent on $[PBS]_{T}$. This rate of k_d (~5 × 10⁻⁵ s⁻¹) is extremely slow and is in agreement with the previous observations of slow PBs fluorescence recovery $({\sim}10^{-3} \text{ s}^{-1})$ in the absence of the FRP (Boulay et al. [2010](#page-9-5); Gwizdala et al. [2011\)](#page-9-2). Finally, approximately the same maximum quenching levels $(-55%)$ were observed in different series of experiments, possibly indicating heterogeneity in PBs preparation.

We also note that the apparent rate of the dark phase of fluorescence quenching (k_{dark}) , obtained by monoexponential approximation, was independent of $[PBS]_T$ (Fig. S2). As this rate is determined by two processes, namely, mutual diffusion of the OCP and PBs and OCP binding to PBs, one could expect the concentration dependence of k_{dark} on $[PBS]_T$ similar to that for k_{PB}^* . However, it can be shown that the rate of dark phase is governed by the $k_{\text{PB}}^{*}/k_{\text{r}}$ ratio, and as $k_{\text{PB}}^* \ll k_r$ in case of in vitro experiments, quenching kinetics is governed by the OCP^R relaxation, and diffusion processes are not pronounced. It was reported previously (Gorbunov et al. [2011](#page-9-6); Kuzminov et al. [2012\)](#page-10-12) that k_{dark} for *Synechocystis* sp. PCC 6803 is ~0.1 s⁻¹, that is similar to our data obtained in vitro (Table [1\)](#page-5-1); however, as will be shown below, the rate of the dark phase in vivo is dependent on the quenching level and is not limited by k_r .

In the in vitro experiments with the three-component system (PBs, OCP, and FRP), the FRP addition resulted

in a dramatic reduction of PBs fluorescence quenching efficiency (Fig. S3A). However, addition of FRP to the OCP–PBs system significantly complicates the model [\(9](#page-4-1)), as FRP is involved into four processes: binding to the OCP^R [Eq. (4) (4)], reconversion of the OCP^R to OCP^O [Eq. (5) (5)], binding to the OCP^R –PBs complex with the rate constant $k_{\text{FRP-PB}}$

and a subsequent FRP-accelerated decay of the OCP^R -PB complex with a rate constant k_{d-PBS} : $OCP^R - PBs + FRP \xrightarrow{k_{FRP-PB}} OCP^R - PBs - FRP,$ (11)

OCP^R – PBS – FRP
$$
\xrightarrow{k_{d-PB}}
$$
 OCP^O + PBS + FRP. (12)

As shown above, the rate of the OCP^R binding to PB is $k_{PR} \approx 2 \times 10^5$ $k_{PR} \approx 2 \times 10^5$ $k_{PR} \approx 2 \times 10^5$ M⁻¹ s⁻¹ (Table 1), while the rate of the OCP^R binding to FRP is an order of magnitude lower, $k_{\text{FRP}} \approx 2 \times 10^4 \text{ M}^{-1} \text{ s}^{-1}$. Considering that $k_{\text{FRP-PB}}$ and k_{FRP} are of the same order of magnitude, the role of FRP in PBs fluorescence quenching in vivo can be described as follows. During the in vitro experiments, $[OCP]_T>>[PBs]_T$, hence, FRP binding to the OCP^R is highly efficient, i.e., the corresponding rate exceeds the rate of the OCP binding to PBs, and an addition of low $\left(<\leq[OCP]_T\right)$ concentrations of the FRP significantly reduces the quenching level (Fig. S3B). Assuming that in vivo $[OCP]_T \sim [PBs]_T$, and taking into account that $k_{FRP} \ll k_{FRP-PB}$, it can be expected that the FRP binding to the OCP can be neglected when processing the quenching curves in vivo.

Kinetics of PB fluorescence quenching in vivo

PB fluorescence quenching kinetics obtained for intact cells of the *Synechocystis* sp. PCC 6803 ΔPSI/ΔPSII mutant at different illumination intensities are presented in Fig. [4.](#page-6-0)

Similarly to the in vitro experiments, the in vivo PBs fluorescence quenching kinetics measured in the ΔPSI/ ΔPSII mutant of *Synechocystis* sp. PCC 6803 were fitted to Eqs. [\(9](#page-4-1), [10\)](#page-5-0), but the fitting procedure itself was essentially different. While for the in vitro experiments the $[OCP]_{T}$ / $[PBS]$ _T ratio (*α*) was fixed, it was unknown for cells, and the reconversion rate for OCP^R could be determined in independent measurements, while it had to be varied in vivo. As the in vitro experiments were performed at specific conditions required to maintain the PBs structure (in a 0.75 M buffer, pH 7.4), properties of microenvironment could be significantly different for a cell, thus setting limitations on a direct extrapolation of the parameters obtained in vitro for the in vivo case. Hence, we performed global fitting of in vivo PBs quenching curves (Fig. [4\)](#page-6-0) which exhibit different dependencies on the fitting parameters. Totally, ~30 quenching curves were simultaneously fitted for each experiment. Table [2](#page-7-0) shows parameters used in the fitting procedure and summarize the retrieved values.

Our results revealed the following striking differences between PBs fluorescence quenching processes in vitro and in vivo. First, the comparison of OCP^R –PBs interaction rate constants obtained in vitro and in vivo allows estimation of the effective $[PBS]_T$ in cell as ~10⁻⁵ M, that is four orders higher than the maximum $[PBS]_T$ value used in the in vitro experiments, and local intracellular concentrations of PBs in cell could be even higher. While it was demonstrated above (Table [1\)](#page-5-1) that the OCP^R –PBs interaction in vitro is governed by free diffusion, as its rate is proportional to $[PBS]_T$, OCP^R–PBs interaction in the cell can be influenced by other processes such as slower diffusion or increased viscosity, thus making $[PBS]_T$ estimation from kinetic data speculative. However, the number of PBs per cell can be estimated as ~1000 from the cryo electron microscopy data (Wilson et al. 2006), that for a volume of 1 μ m³ gives a

Fig. 4 PBs fluorescence quenching curves measured **a** at different illumination levels and **b** different illumination durations at a fixed intensity in intact cells of the ΔPSI/ΔPSII mutant. *Solid lines* correspond to fits to Eqs. [\(9](#page-4-1), [10](#page-5-0))

Parameter	In vitro		In vivo	
	Type	Value	Type	Value
OCP ^R –PBs interaction rate constant, k_{PR} (M ⁻¹ s ⁻¹)	Fitted	2×10^5	Fitted	$(2.25 \pm 0.13)/[PBs]_{T}$
OCPR reconversion rate, k_r	Fixed (obtained independently on OCP solution)	0.04	Fitted	0.06 ± 0.02
OCP ^R –PBs dissociation rate, k_d	Fitted	$\rm < 10^{-4} \, s^{-1}$	Fitted	$0.005 + 0.002$ s ⁻¹
PBs quenching efficiency by OCP ^R , $(1 - \phi)$	Fitted	$0.55 + 0.05$	Fitted	$0.98 + 0.01$
$\alpha = [OCP]_{T} / [PBs]_{T}$	Fixed (known for each curve)	>>1	Fitted	$0.647 + 0.004$

Table 2 Summary of fixed and varied parameters used for fitting of PBs fluorescence quenching in vitro and in vivo with Eqs. ([2,](#page-3-3) [3](#page-3-6))

 $\sim 10^{-5}$ M concentration. As a result, a similar quenching level is reached in both cases, but in vitro an excess of the OCP over PBs is required due to low local concentrations, and the quenching efficiency for a single OCPR–PBs complex is relatively low. In contrast, in cells almost complete quenching in OCP^R –PBs complexes takes place, and high local concentrations make the overall quenching efficient even at $[OCP]_T/[PBs]_T<1$. Finally, almost 10 times higher value of k_d was obtained from fluorescence quenching curves in vivo, that is explained by the FRP-assisted acceleration of PBs fluorescence recovery.

When processing the PBs fluorescence quenching phase, we made an assumption about the absence of interactions with the FRP, which resulted in a satisfactory approximation of the data. During the PBs recovery phase, such an assumption is invalid, and FRP-related processes must be taken into account, thus increasing the number of parameters and making the modeling more complex. However, one of the main goals of this work was to verify whether assessment of PBs, OCP, and FRP concentrations is possible on the basis of mathematical description of PBs fluorescence kinetics in cyanobacteria cells; hence, it was necessary to suggest a simplified model of fluorescence recovery, as the FRP, in opposite to the OCP, cannot be quantified solely from the quenching kinetics.

Kinetics PBs fluorescence recovery in vivo

Figure [5](#page-8-0)a demonstrates the time courses of PBs fluorescence recovery in darkness observed in vivo in *Synechocystis* sp. PCC 6803 ΔPSI/ΔPSII mutant cells on the offset of illumination. These curves were obtained for different durations pre-exposure to strong light and thus different initial levels of PB fluorescence quenching prior to recovery.

The dependence of the recovery rate on the initial quenching level was strongly non-linear (Fig. [5\)](#page-8-0). The higher was the initial amplitude of fluorescence quenching, the slower was the subsequent fluorescence recovery. This non-linearity cannot be explained by a commonly used first-order kinetics approach, which assumes that the recovery rate is independent on the initial concentration of quenchers. Moreover, our analysis of the time course of the rate of fluorescence recovery d[PBs]/dt (Fig. [5b](#page-8-0)) revealed that this dependence is non-monotonous, suggesting the presence of several processes that influence the fluorescence recovery rate.

To quantify the non-linear pattern of the PBs fluorescence recovery, we applied an approach that was previously used for description of the kinetics of catalytic reactions. The process of recovery can be presented as a sequence of the following two reactions: (i) $OCP^R - PB$ complex (substrate, S) slowly decays to OCP^O and free PBs (product, P), (ii) this conversion is accelerated by the FRP (F) by forming a ternary OCP^R–PB–FRP complex (SF). Collectively, these processes can be represented by the standard scheme of enzymatic catalysis:

$$
S + F \leftrightarrow SF \rightarrow P + F,\tag{13}
$$

In order to describe the system behavior and elucidate the origin of fluorescence recovery non-linearity, we further applied the generalized approach for enzyme kinetics investigation (Cao and Enrique [2013\)](#page-9-9) that allows one to fit the time course of product formation with the following equation:

$$
[P] = \frac{v_0}{\eta} (1 - e^{-\eta t}),\tag{14}
$$

where ν_0 represents the initial product formation rate, whereas η indicates non-linearity of the curve caused by the reduction in the enzyme cycling velocity.

Figure [5](#page-8-0)c, d shows the dependence of the retrieved ν_0 and *η* values obtained by fitting PBs fluorescence recovery kinetics to Eq. (14) (14) on the initial amplitude of quenching, which is proportional to concentration of the substrate (i.e., OCPR—PB complexes).

The observed saturation of the initial velocity (Fig. [5](#page-8-0)c) corresponds to the "enzyme" (FRP) saturation by the substrate, which is consistent with the suggested excess of OCP molecules over the FRP in vivo. Above the initial quenching levels of 30%, the maximum rate of recovery is constant and equal to 14 s^{-1} . This is due to the fact that all FRPs are attached to $OCP^R - PB$ complexes, thus

Fig. 5 a The time course of fluorescence recovery in intact cells of the ΔPSI/ΔPSII mutant at different initial levels of fluorescence quenching. **b** The dependence of fluorescence recovery rate on the OCPR–PBs complex concentration (PB fluorescence quenching level)

obtained at different temperatures. **c, d** The dependence of the initial enzyme cycling velocity ν_0 and enzyme cycling velocity reduction η on the initial level of PBs fluorescence quenching

providing the maximum possible rate of recovery. As the initial quenching level and thus the initial substrate concentration decrease, the maximum recovery rate goes down. For instance, at the initial quenching of 7%, ν_0 was equal to 7 s⁻¹, i.e., ~50% of the maximum recovery rate. Under such conditions, the substrate concentration becomes comparable to the FRP concentration, and $[FRP]_T$ can be estimated as the concentration of OCP–PBs complexes at the corresponding quenching level. As the maximum quenching level corresponds to the $[OCP]_T/[PBs]_T$ ratio, and 7% quenching corresponds to ~10% of the maximum quenching, $[FRP]_T$ can be estimated as ~0.1 $[OCP]_T$.

As the initial rate of fluorescence recovery reaches a plateau at moderate quenching amplitudes, the rate of this process is determined by the turnover rate of the enzyme. The values of the enzyme cycling velocity reduction exceed the reciprocal observation time *τ: η*>*τ*−1 (*τ*−1 is shown with a horizontal dashed line in Fig. [5b](#page-8-0)), hence the fluorescence recovery process is strongly non-linear. When *η* decreases with substrate concentration, the non-linearity mainly comes from the substrate depletion (Cao and Enrique [2013](#page-9-9)). However, fluorescence recovery is non-linear even in the substrate concentration region corresponding to ν_0 saturation. This fact indicates that substrate depletion is not the only source of non-linearity in the system.

To further investigate the FRP functioning in vivo, we examined the dependence of fluorescence recovery rate on the OCPR–PBs concentration (Fig. $5b$). The measured curves exhibit sigmoidal shape and two deflection points, that, according to the enzyme catalysis literature, corresponds to allosteric functioning of an enzyme. This pattern means that an enzyme is capable of changing its properties, such as its oligomeric state, during the time course of product formation. Another feature of the allosteric influence is that it becomes less pronounced at elevated temperatures (Teipel and Koshland [1969\)](#page-10-18), that is also shown in Fig. [5](#page-8-0)b.

The FRP normally exists in the dimeric form, while upon interaction with the OCP^R it forms monomers (Sluchanko et al. [2017\)](#page-10-4). We propose the following working model to explain the non-linear kinetics of PBs fluorescence recovery. Dimers and monomers of the FRP have different affinity to the OCP^R –PBs complex. In this case,

the FRP, mainly existing in dimeric form, forms complexes with the $OCP^{R}-PBs$ with lower affinity, and after that all dimeric FRP splits into bound and free monomeric fractions. We also note that the sigmoidal shape of PBs fluorescence recovery curve, which is observed at high quenching levels (see e.g., the curve corresponding to 56% quenching in Fig. [5a](#page-8-0)), can be explained as a consequence of dimer–monomer transition of the FRP, which accompanies its interaction with the OCP–PBs complexes. Following this, free FRP is mainly in the monomeric form, and the FRP binding to the OCP–PBs complexes occurs with the higher rates, which correspond to the FRP monomers. Finally, when PBs fluorescence approaches its complete recovery, dimers of the FRP start to reappear, thus slowing down the reaction rate again. This hypothesis explains the behavior of PBs fluorescence recovery (Fig. [5b](#page-8-0)) and strong non-linearity (Fig. [5c](#page-8-0)); however, further research is necessary to prove this model.

Conclusions

In this work, we presented a biophysical model, which allows for the description of experimental PBs fluorescence kinetics following cyanobacteria illumination with a strong blue–green light. This process is governed by the interplay between three components, PBs, OCP, and FRP, as well as their transient complexes. Using the possibility to reconstruct the OCP-mediated photoprotective mechanism and its parts in vitro, we performed an independent testing of the suggested model on the OCP–FRP and OCP–PBs systems. It was shown that that the dissociation constant for the $OCP^{R}-FRP$ complex estimated from the modeling data as the ratio of OCP and FRP interaction rate constants was equal to 6 μM, that is in a good agreement with the experimental data (Sluchanko et al. [2017\)](#page-10-4). The rate constant of OCP^R binding to PBs was shown to be proportional to $[PBS]_T$, as expected for the bimolecular diffusion-limited interaction. Application of the biophysical model to the analysis of PBs fluorescence quenching kinetics in vivo provided for the estimations of relative OCP concentration ($[OCP]_{T}/[PBs]_{T}\approx 0.6$) and the local PBs concentration ($[PBS]_{T}$ ~10⁻⁵ M) in cells, as well as the efficiency of a single PB fluorescence quenching by the OCP^R (98%). Next, description of PBs fluorescence recovery kinetics on the basis of the generalized model of enzymatic catalysis (Cao and Enrique [2013\)](#page-9-9) provided for the estimation of the FRP concentration in cells as $[FRP]_T/[OCP]_T \approx 0.1$. Our proposed approach allows one to determine relative OCP, FRP, and PBs concentrations in vivo, and can be used to quantify the NPQ characteristics in different strains of cyanobacteria, as well as in cyanobacteria grown/living in different environmental conditions. Indeed, the time course of non-photochemical quenching is determined by $[PBS]_{T}$, $[FRP]_T$, or $[OCP]_T$ concentrations. In vivo, the OCP to FRP ratio is regulated by the *slr1963* (OCP) and *slr1964* (FRP) gene expression levels (Boulay et al. [2010\)](#page-9-5), which may vary depending on the environmental conditions. Hence, the approach, suggested in this work, could be used for the examination of fine tuning of the molecular machinery responsible for photoprotection in cyanobacteria.

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