

Phase transitions for a model with uncountable spin space on the Cayley tree: the general case

Golibjon Botirov2 · Benedikt Jahnel1

Received: 24 March 2018 / Accepted: 9 August 2018 / Published online: 17 August 2018 © Springer Nature Switzerland AG 2018

Abstract

In this paper we complete the analysis of a statistical mechanics model on Cayley trees of any degree, started in Botirov (Positivity 21(3):955–961, [2017\)](#page-10-0), Eshkabilov et al. (J Stat Phys 147(4):779–794, [2012\)](#page-10-1), Eshkabilov and Rozikov (Math Phys Anal Geom 13:275–286, [2010\)](#page-10-2), Botirov et al. (Lobachevskii J Math 34(3):256–263 [2013\)](#page-10-3) and Jahnel et al. (Math Phys Anal Geom 17:323–331 [2014\)](#page-10-4). The potential is of nearestneighbor type and the local state space is compact but uncountable. Based on the system parameters we prove existence of a critical value θ_c such that for $\theta \leq \theta_c$ there is a unique translation-invariant splitting Gibbs measure. For $\theta_c < \theta$ there is a phase transition with exactly three translation-invariant splitting Gibbs measures. The proof rests on an analysis of fixed points of an associated non-linear Hammerstein integral operator for the boundary laws.

Keywords Cayley trees · Hammerstein operators · Splitting Gibbs measures · Phase transitions

Mathematics Subject Classification 82B05 · 82B20 (primary); 60K35 (secondary)

1 Introduction

In the present note we complete a line of research about the phase-transition behavior of a nearest-neighbor model on Cayley trees with arbitrary degree $k \geq 2$. As first described in [\[4](#page-10-2)], for a given consistent family of finite-volume Gibbs measures, the

¹ Weierstrass Institute Berlin, Mohrenstr. 39, 10117 Berlin, Germany

B Golibjon Botirov botirovg@yandex.ru Benedikt Jahnel Benedikt.Jahnel@wias-berlin.de https://www.wias-berlin.de/people/jahnel/

² National University of Uzbekistan, University Street 4, Tashkent, Uzbekistan 100174

existence and multiplicity of a certain class of infinite-volume measures which are consistent with the prescribed finite-volume Gibbs measures, can be reduced to the analysis of fixed points of some non-linear integral equation of Hammerstein type. Every positive solution of the fixed point equation here corresponds to a measures which is called a splitting Gibbs measure. Every splitting Gibbs measure is also a Gibbs measure in the sense of the DLR formalism; see [\[1](#page-10-5)]. This approach has been successfully applied in the analysis of a variety of different models on Cayley trees with respect to their phase-transition properties; see [\[12\]](#page-10-6) for a comprehensive overview. In particular, starting with [\[3\]](#page-10-1), a phase-transition of multiple splitting Gibbs measures has been detected in a model with uncountable local state space [0, 1] and nearestneighbor interactions. This has motivated the subsequent analysis in [\[2](#page-10-0)[,5](#page-10-3)[,9](#page-10-4)], to further understand critical behavior of this model for all degrees of the underlying tree, where also new parameters are introduced. In $[6,7]$ $[6,7]$ $[6,7]$ the Potts model with a countable set of spin values is studied. It is the purpose of this note to complete the analysis of this model.

For nearest neighbors *x*, *y* on the *Cayley tree* Γ^k with degree $k \geq 2$ with local states $\sigma(x)$, $\sigma(y) \in [0, 1]$, we consider the potential

$$
\xi_{\sigma(x),\sigma(y)} = \log\left(1+\theta^{2n+1}\sqrt{4\left(\sigma(x)-\frac{1}{2}\right)\left(\sigma(y)-\frac{1}{2}\right)}\right) \tag{1.1}
$$

where $n \in \mathbb{N} \cup \{0\}$ and $0 \le \theta < 1$ are the system parameters. It can be interpreted as a certain symmetric pair-interaction with values in $[\log(1-\theta), \log(1+\theta)]$, admitting two distinct ground states given by the all-0 and the all-1 configuration. The main result is the existence of a sharp threshold

$$
\theta_{\rm c} = \frac{2n+3}{k(2n+1)}
$$

such that if $\theta_c < \theta < 1$, there are exactly three translation-invariant splitting Gibbs measures and otherwise there is only one.

2 Setup

2.1 Gibbs measures on Cayley trees

The *Cayley tree* Γ^k of order $k \geq 1$ is an infinite tree, i.e., a graph without cycles, such that exactly $k + 1$ edges originate from each vertex. Let $\Gamma^k = (V, L)$ where *V* is the set of *vertices* and *L* is a symmetric subset of $V \times V$, called the *edge set*. The word "symmetric" means that $(x, y) \in L$ iff $(y, x) \in L$. Here, x and y are called the *endpoints* of the edge $\langle x, y \rangle$. Two vertices x and y are called *nearest neighbors* if there exists an edge $l \in L$ connecting them and we denote $l = \langle x, y \rangle$. For a fixed $x^0 \in V$, called the *root*, we defines *n*-spheres and *n*-disks in the graph distance $d(x, y)$ by

$$
W_n = \{x \in V | d(x, x^0) = n\}, \quad V_n = \bigcup_{i=0}^n W_i
$$

and denote for any $x \in W_n$ the set of *direct successors* of x by

$$
S(x) = \{ y \in W_{n+1} : d(x, y) = 1 \}.
$$

For $A \subset V$ let $\Omega_A = [0, 1]^A$ denote the set of all configurations σ_A on A. In particular, a configuration σ on *V* is then defined as a function $V \ni x \mapsto \sigma(x) \in [0, 1]$. According to the usual setup for Gibbs measure, we consider a (formal) Hamiltonian of the form

$$
H(\sigma) = -\sum_{\langle x, y \rangle \in L} \xi_{\sigma(x), \sigma(y)},\tag{2.1}
$$

where $\xi : (u, v) \in [0, 1]^2 \mapsto \xi_{u, v} \in \mathbb{R}$ is the interaction [\(1.1\)](#page-1-0) which assigns energy only to neighboring sites. Since ξ does not depend on the locations x and y, H is invariant under tree translations. Let λ be the Lebesgue measure on [0, 1] then, on the set of all configurations on *A* the a priori measure λ_A is introduced as the $|A|$ -fold product of the measure λ . Here and in the sequel, $|A|$ denotes the cardinality of A. We equip $\Omega = \Omega_V$ with the standard sigma-algebra *B* generated by the cylindrical subsets. A probability measure μ on (Ω, \mathcal{B}) is called a *Gibbs measure* (with Hamiltonian *H*) if it satisfies the DLR equation. That is, for any $n = 1, 2, \dots$ and bounded measurable test function f , we have that

$$
\int \mu(d\sigma) f(\sigma) = \int \mu(d\sigma) \int \gamma_{V_n}(d\tilde{\sigma}_{V_n} | \sigma_{W_{n+1}}) f(\tilde{\sigma}_{V_n} \sigma_{\Gamma^k \setminus V_n}), \tag{2.2}
$$

where $\gamma_{V_n}(d\sigma_{V_n}|\sigma_{\Gamma^k\setminus V_n})$ is the Gibbsian specification

$$
\gamma_{V_n}(d\tilde{\sigma}_{V_n}|\sigma_{\Gamma^k\setminus V_n})=\frac{1}{Z_{V_n}(\sigma_{W_{n+1}})}e^{-\beta H(\tilde{\sigma}_{V_n}\sigma_{W_{n+1}})}\lambda_{V_n}(d\tilde{\sigma}_{V_n}),
$$

with normalization Z_{V_n} and temperature parameter $\beta \geq 0$. Such a specification is also sometimes referred to as a Markov specification; see [\[8\]](#page-10-9).

2.2 Representation via Hammerstein operators

A subset of the infinite-volume Gibbs measures defined via the DLR equation [\(2.2\)](#page-2-0), called the *splitting Gibbs measures* or *Markov chains*, can be represented in terms of the fixed points of some nonlinear integral operator of Hammerstein type; see [\[4\]](#page-10-2) for details. More precisely, for every $k \in \mathbb{N}$ consider the integral operator H_k acting on the cone $C^+[0, 1] = \{f \in C[0, 1] : f(x) \ge 0\}$ given by

$$
(H_k f)(t) = \int_0^1 K(t, u) f^k(u) du.
$$
 (2.3)

 \mathcal{D} Springer

Then, the translation-invariant splitting Gibbs measures for the Hamiltonian [\(2.1\)](#page-2-1) correspond to positive fixed points of H_k with $K(t, u) = \exp(\beta \xi_{t, u})$, often called *boundary laws.* Note that H_k in general might generate ill-posed problems; see [\[10](#page-10-10)[,11](#page-10-11)].

3 Main results

The main result of this note is the following characterization of phase-transition regimes of the model [\(1.1\)](#page-1-0) with $\beta = 1$.

Theorem 3.1 *For all n* ∈ \mathbb{N} ∪ {0} *and k* ≥ 2 *let* $\theta_c = (2n + 3)/(k(2n + 1))$ *, then the model [\(1.1\)](#page-1-0) has*

- *(1) a unique translation-invariant splitting Gibbs measure if* $0 \le \theta \le \theta_c$ *and*
- *(2) exactly three translation-invariant splitting Gibbs measures if* $\theta_c < \theta < 1$ *.*

The proof is based on a characterization of solutions to the fixed point equation for the associated Hammerstein integral operator [\(2.3\)](#page-2-2) as given in Proposition [3.2](#page-3-0) below. In case of the model at hand, then the analysis can be reduced to finding the fixed points of the following 2-dimensional operator $V_k : \mathbb{R}^2 \to \mathbb{R}^2$

$$
V_{k,n}(x, y) = \begin{cases} \sum_{i=0,2,\dots, \lfloor k \rfloor_{\text{even}}} \binom{k}{i} \frac{2n+1}{2n+1+i} 2^{\frac{i}{2n+1}} x^{k-i} (\theta y)^i\\ \sum_{i=1,3,\dots, \lfloor k \rfloor_{\text{odd}}} \binom{k}{i} \frac{2n+1}{2n+2+i} 2^{\frac{i-1}{2n+1}} x^{k-i} (\theta y)^i \end{cases} (3.1)
$$

with $k \ge 2$, which is then the content of Proposition [3.3.](#page-3-1) Here we use the notation

$$
\lfloor k \rfloor_{\text{even}} = \begin{cases} k, & \text{if } k \text{ is even} \\ k - 1, & \text{if } k \text{ is odd} \end{cases} \qquad \text{and} \qquad \lfloor k \rfloor_{\text{odd}} = \begin{cases} k, & \text{if } k \text{ is odd} \\ k - 1, & \text{if } k \text{ is even.} \end{cases}
$$

Proposition 3.2 *A function* $\varphi \in C[0, 1]$ *is a solution of the Hammerstein equation*

$$
H_k f = f \tag{3.2}
$$

with Hk defined in [\(2.3\)](#page-2-2) *for our model* [\(1.1\)](#page-1-0)*, iff* ϕ *has the following form*

$$
\varphi(t) = c_1 + c_2 \theta \sqrt[2n+1]{4(t-\frac{1}{2})},
$$

where $(c_1, c_2) \in \mathbb{R}^2$ *is a fixed point of the operator* $V_{k,n}$ *as defined in* [\(3.1\)](#page-3-2)*.*

In the following proposition we characterize the fixed points of $V_{k,n}$ which readily implies Theorem [3.1](#page-3-3) using Proposition [3.2.](#page-3-0)

Proposition 3.3 *Let* $\theta_c = (2n+3)/(k(2n+1))$ *, then there exist uniquely defined points* $x_0, y_0 \in (0, \infty)$ *such that the number and form of the fixed points of the operator* $V_{k,n}$ *are as presented in the following Table* [1](#page-4-0)*.*

k even	Fixed points if $0 \le \theta \le 1$			Additional fixed points if $\theta_c < \theta < 1$			
	(0, 0)	(1,0)		(x_0, y_0)		$(x_0, -y_0)$	
k odd	(0, 0)	(1, 0)	$(-1, 0)$	(x ₀ , y ₀)	$(-x_0, -y_0)$	$(xo, -yo)$	$(-x_o, y_o)$

Table 1 Set of 2-dimensional fixed points of $V_{k,n}$

Further, only the fixed points $(1, 0)$, (x_o, y_o) and $(x_o, -y_o)$ give rise to positive solutions for the Hammerstein equation [\(3.2\)](#page-3-4).

Let us finally give the references to the special cases considered prior to this work. [\[5](#page-10-3), Theorem 4.2 and Theorem 5.2] proves the cases $k = 2, 3$ with $n = 1$ of [\(1.1\)](#page-1-0) whereas in [\[9,](#page-10-4) Theorem 3.2.] the cases $k \ge 2$ with $n = 1$ are given. Finally, in [\[2,](#page-10-0) Theorem 2.3] the cases $k = 2$ with general $n \ge 1$ is provided.

4 Proofs

In order to ease notation, let us write $m = 2n + 1$ in this section. Note that for the model [\(1.1\)](#page-1-0) with $\beta = 1$, the kernel $K(t, u)$ of the Hammerstein operator H_k is given by

$$
K(t, u) = 1 + \theta \sqrt[m]{4 \left(t - \frac{1}{2} \right) \left(u - \frac{1}{2} \right)}.
$$

Proof of Proposition [3.2](#page-3-0) Let us start with necessity. Assume $\varphi \in C[0, 1]$ to be a solution of the equation (3.2) . Then we have

$$
\varphi(t) = c_1 + c_2 \theta \sqrt[m]{4 \left(t - \frac{1}{2} \right)},
$$
\n(4.1)

where

$$
c_1 = \int_0^1 \varphi^k(u) du \quad \text{and} \quad c_2 = \int_0^1 \sqrt[m]{u - \frac{1}{2}} \varphi^k(u) du. \tag{4.2}
$$

Substituting $\varphi(t)$ into the first equation of [\(4.2\)](#page-4-1) we get

$$
c_1 = \int_0^1 \left(c_1 + c_2 \theta \sqrt[m]{4 \left(u - \frac{1}{2} \right)} \right)^k du
$$

=
$$
\int_0^1 \sum_{i=0}^k {k \choose i} c_1^{k-i} \left(c_2 \theta \sqrt[m]{4} \sqrt[m]{u - \frac{1}{2}} \right)^i du
$$

$$
= \sum_{i=0}^k {k \choose i} c_1^{k-i} (\theta c_2)^i 2^{\frac{2i}{m}} \int_0^1 {u - \frac{1}{2}}^{\frac{i}{m}} du.
$$

Now, we use the following equality

$$
\int_{0}^{1} \left(u - \frac{1}{2} \right)^{\frac{i}{m}} du = \begin{cases} 0, & \text{if } i \text{ is odd and} \\ \frac{m}{m+i} 2^{-\frac{i}{m}}, & \text{if } i \text{ is even.} \end{cases}
$$
(4.3)

Then we get

$$
c_1 = \sum_{i=0,2,\dots, \lfloor k \rfloor_{\text{even}}} {k \choose i} \frac{m}{m+i} 2^{\frac{i}{m}} c_1^{k-i} (\theta c_2)^i
$$

and substituting the function φ into the second equation of [\(4.2\)](#page-4-1) we have

$$
c_2 = \int_0^1 \left(u - \frac{1}{2} \right)^{\frac{1}{m}} \left(c_1 + \theta c_2 \sqrt[m]{4 \left(u - \frac{1}{2} \right)} \right)^k du
$$

=
$$
\int_0^1 \left(u - \frac{1}{2} \right)^{\frac{1}{m}} \sum_{i=0}^k {k \choose i} c_1^{k-i} \left(\theta c_2 \sqrt[m]{4} \sqrt[m]{u - \frac{1}{2}} \right)^i du
$$

=
$$
\sum_{i=0}^k {k \choose i} c_1^{k-i} (c_2 \theta)^i 2^{\frac{2i}{m}} \int_0^1 \left(u - \frac{1}{2} \right)^{\frac{i+1}{m}} du.
$$

Now, using the following equality

$$
\int_{0}^{1} \left(u - \frac{1}{2} \right)^{\frac{i+1}{m}} du = \begin{cases} 0, & \text{if } i \text{ is even and} \\ \frac{m}{m+1+i} 2^{-\frac{i+1}{m}}, & \text{if } i \text{ is odd} \end{cases}
$$
(4.4)

we arrive at the equation

$$
c_2 = \sum_{i=1,3,\dots,\lfloor k \rfloor_{\text{odd}}} {k \choose i} \frac{m}{m+1+i} 2^{\frac{i-1}{m}} c_1^{k-i} (\theta c_2)^i.
$$

In particular, the point $(c_1, c_2) \in \mathbb{R}^2$ must be a fixed point of the operator $V_{k,n}$ from [\(3.1\)](#page-3-2).

 $\hat{2}$ Springer

For the sufficiency, assume that, a point $(c_1, c_2) \in \mathbb{R}^2$ is a fixed point of the operator $V_{k,n}$ and define the function $\varphi \in C[0, 1]$ by the equality

$$
\varphi(t) = c_1 + c_2 \theta \sqrt[m]{4 \left(t - \frac{1}{2} \right)}.
$$

Then, we can calculate

$$
(H_k \varphi)(t) = \int_0^1 \left(1 + \sqrt[m]{4} \theta \sqrt[m]{(t - \frac{1}{2})(u - \frac{1}{2})} \right) \varphi^k(u) du
$$

\n
$$
= \int_0^1 \varphi^k(u) du + \sqrt[m]{4} \theta \sqrt[m]{t - \frac{1}{2}} \int_0^1 \sqrt[m]{u - \frac{1}{2}} \varphi^k(u) du
$$

\n
$$
= \int_0^1 \left(c_1 + c_2 \theta \sqrt[m]{4 \left(u - \frac{1}{2} \right)} \right)^k du
$$

\n
$$
+ \sqrt[m]{4} \theta \sqrt[m]{t - \frac{1}{2}} \int_0^1 \sqrt[m]{u - \frac{1}{2}} \left(c_1 + c_2 \theta \sqrt[m]{4 \left(u - \frac{1}{2} \right)} \right)^k du
$$

\n
$$
= \sum_{i=0}^k {k \choose i} c_1^{k-i} (\theta c_2)^i 2^{\frac{2i}{m}} \int_0^1 \left(u - \frac{1}{2} \right)^{\frac{i}{m}} du + \sqrt[m]{4} \theta \sqrt[m]{t - \frac{1}{2}}
$$

\n
$$
\times \sum_{i=0}^k {k \choose i} c_1^{k-i} (c_2 \theta)^i 2^{\frac{2i}{m}} \int_0^1 \left(u - \frac{1}{2} \right)^{\frac{i+1}{m}} du.
$$

\n(4.5)

Now, we using (4.3) and (4.4) , from (4.5) we get

$$
(H_k \varphi)(t) = \sum_{i=0,2,\dots,\lfloor k \rfloor_{\text{even}}} \binom{k}{i} \frac{m}{m+i} 2^{\frac{i}{m}} c_1^{k-i} (\theta c_2)^i
$$

+ $\theta \sqrt[m]{4(t-\frac{1}{2})} \sum_{i=1,3,\dots,\lfloor k \rfloor_{\text{odd}}} \binom{k}{i} \frac{m}{m+1+i} 2^{\frac{i-1}{m}} c_1^{k-i} (\theta c_2)^i = \varphi(t).$

Thus, φ is a solution of the equation [\(3.2\)](#page-3-4).

Proof of Proposition [3.3](#page-3-1) We determine the number and form of solutions to $V_{k,n}$ in equation [\(3.1\)](#page-3-2). For $\theta = 0$, the fixed point equation for $V_{k,n}$ reduces to $x = x^k$ and $y = 0$ and hence for *k* even, the solutions are given by $(1, 0)$ and $(0, 0)$ and for *k* odd, the solutions are given by $(1, 0)$, $(-1, 0)$ and $(0, 0)$. Further note that in both cases, if $x = 0$, then also $y = 0$. So from now on we assume that $\theta > 0$ and $x \neq 0$.

$$
\Box
$$

Let us start by considering the case where *k* is even. Inspecting the equation for *x* in $V_{k,n}$, we see that $x \geq 0$. We introduce a reduction of the 2-dimensional fixed point equation to a 1-dimensional one. Writing $z = \theta y/x$, then for any fixed point (x, y) of $V_{k,n}$, *z* necessarily is a solution to the fixed point equation

$$
z = \theta \frac{\sum_{i=1,3,\dots,k-1} {k \choose i} \frac{m}{m+1+i} 2^{\frac{i-1}{m}} z^i}{\sum_{i=0,2,\dots,k} {k \choose i} \frac{m}{m+i} 2^{\frac{i}{m}} z^i} = \theta \frac{F_1^{\text{even}}(z)}{F_2^{\text{even}}(z)} = f_{\text{even}}(z).
$$

In order to find solutions for $z = f_{\text{even}}(z)$, we have to find roots of the polynomial

$$
P_{\text{even}}(z) = \sum_{i=1,3,\dots,k+1} {k \choose i-1} \frac{m}{m+i-1} 2^{\frac{i-1}{m}} z^i - \theta \sum_{i=1,3,\dots,k-1} {k \choose i} \frac{m}{m+i+1} 2^{\frac{i-1}{m}} z^i
$$

= $r_{\theta}(k, k+1) z^{k+1} + \sum_{i=1,3,\dots,k-1} r_{\theta}(k, i) z^i$ (4.6)

where $r_{\theta}(k, k+1) = \frac{m}{m+k} 2^{\frac{k}{m}}$ and

$$
r_{\theta}(k,i) = {k \choose i} \frac{m}{m+i+1} 2^{\frac{i-1}{m}} \Big[\frac{i}{k-i+1} \frac{m+i+1}{m+i-1} - \theta \Big].
$$

Moreover,

$$
r_{\theta}(k, i) \begin{cases} < 0 & \text{if } \theta > \frac{i}{k - i + 1} \frac{m + i + 1}{m + i - 1} \\ = 0 & \text{if } \theta = \frac{i}{k - i + 1} \frac{m + i + 1}{m + i - 1} \\ > 0 & \text{if } \theta < \frac{i}{k - i + 1} \frac{m + i + 1}{m + i - 1} \end{cases}
$$

and we denote the critical θ by $\theta_{k,i}$. Further note that $i \mapsto \theta_{k,i}$ is increasing. Indeed, the derivative of the continuous version (where $i \in \mathbb{N}$ is for a moment replaced by $x \in \mathbb{R}$) is given by

$$
\frac{(1+k)(m-1)^2 + (3+k)i^2 + 2(1+k)(m-1)(1+i))}{(1+k-i)^2(m-1+i)^2}
$$

which is non-negative. Hence, for θ below the lowest critical value, $\theta_c = \theta_{k,1} = \frac{m+2}{km}$, all coefficients are positive and hence there is no positive real root of *P*_{even} by Descartes' rule of sign. In particular, for $0 \le \theta \le \theta_c$, there can not be any solutions to the fixed point equation associated to $V_{k,n}$ of the form (x, y) with $0 < x, y$. Since $P_{\text{even}}(-z) = -P_{\text{even}}(z)$, there can also not exist any fixed point (x, y) with $y < 0 < x$. This settles the subcritical and critical case for even *k*.

Further, for the supercritical case $\theta > \theta_c$ with even *k*, again by Descartes' rule of sign, there is exactly one sign change in *P*even and hence exactly one non-trivial positive real root of P_{even} exists which we denote $z_0 > 0$. By the point symmetry of the polynomial P_{even} , with z_0 also $-z_0$ is a root. In order to uniquely recover the solution (x_0, y_0) from the positive non-trivial solution z_0 , note that

$$
V_{k,n}(x, y) = \begin{cases} x^k \sum_{i=0,2,...,k} {k \choose i} \frac{m}{m+i} 2^{\frac{i}{m}} (\theta^{\frac{y}{x}})^i = x^k F_2^{\text{even}} (\theta^{\frac{y}{x}}) \\ x^k \sum_{i=1,3,...,k-1} {k \choose i} \frac{m}{m+i+1} 2^{\frac{i-1}{m}} (\theta^{\frac{y}{x}})^i = x^k F_1^{\text{even}} (\theta^{\frac{y}{x}}) \end{cases}
$$

and hence $x_o = F_2^{\text{even}}(z_o)^{1/(1-k)} > 0$ and $y_o = F_1^{\text{even}}(z_o)F_2^{\text{even}}(z_o)^{k/(1-k)} > 0$ solve the 2-dimensional equation. Note that (x_o, y_o) is the only solution with $\theta y/x = z_o$. Indeed, any other such solution would be $x_1 = cx_0$ and $y_1 = cy_0$ for some $c \in \mathbb{R} \setminus \{0\}$, but then $cx_o = x_1 = x_1^k F_2(z_o) = c^k x_o^k F_2(z_o)$ which implies that $c = c^k$. But this is true if and only if $c = 1$ for even k .

Further, note that $F_2^{\text{even}}(-z_o)^{1/(1-k)} = F_2^{\text{even}}(z_o)^{1/(1-k)} = x_o$ and $F_1^{\text{even}}(-z_o)$ $F_2^{\text{even}}(-z_o)^{k/(1-k)} = -F_1^{\text{even}}(z_o)F_2^{\text{even}}(z_o)^{k/(1-k)} = -y_o$ and hence also $(x_o, -y_o)$ is a solution to the 2-dimensional fixed point equation. Finally, as in the case of z_o , the tuple $(x_o, -y_o)$ is the only solution with $\theta y/x = -z_o$ since any other solution would be of the form $x_1 = cx_o$ and $y_1 = -cy_o$ for some $c \in \mathbb{R} \setminus \{0\}$. But then again $cx_o = x_1 = x_1^k F_2^{\text{even}}(z_o) = c^k x_o^k F_2^{\text{even}}(z_o)$ which implies that $c = c^k$ which is true if and only if $c = 1$ for even *k*. So, for even *k*, we have now completely charaterized solutions of the fixed point equation associated with also in the super critical regime.

For odd *k* and $\theta > 0$, the situations is slightly more complicated, since we can not exclude certain signs from the fixed points of

$$
V_{k,n}(x, y) = \begin{cases} x^k \sum_{i=0,2,...,k-1} {k \choose i} \frac{m}{m+i} 2^{\frac{i}{m}} (\theta^y_x)^i \\ x^k \sum_{i=1,3,...,k} {k \choose i} \frac{m}{m+i+1} 2^{\frac{i-1}{m}} (\theta^y_x)^i. \end{cases}
$$

Writing again $z = \theta y/x$, the fixed point equation for [\(3.1\)](#page-3-2) becomes

$$
z = \theta \frac{\sum_{i=1,3,\dots,k} {k \choose i} \frac{m}{m+1+i} 2^{\frac{i-1}{m}} z^i}{\sum_{i=0,2,\dots,k-1} {k \choose i} \frac{m}{m+i} 2^{\frac{i}{m}} z^i}
$$

and the corresponding point symmetric polynomial is given by

$$
P_{\text{odd}}(z) = \sum_{i=1,3,\dots,k} r_{\theta}(k,i)z^i.
$$

Hence, the exact same arguments as in the case of even *k* apply, yielding again no roots for $\theta \leq \theta_c$ and two roots z_o and $-z_o$ for $\theta > \theta_c$. In contrast to the case for even k, for odd *k*, both (x_0, y_0) and $(-x_0, -y_0)$ are 2-dimensional fixed points corresponding to z_o since $c = c^k$ can be solved by ± 1 . Finally, following the exact same arguments as in the case of even *k*, the fixed points $(x_o, -y_o)$ and $(-x_o, y_o)$ correspond to $-z_o$. The complete list of fixed points is recorded in Table [1.](#page-4-0)

For $(\pm x_0, \pm y_0)$ to give rise to a positive solution of the Hammerstein fixed point equation, by the form of solutions φ we must have that for all $t \in [0, 1]$

$$
\pm x_o \pm y_o \theta \sqrt[m]{4(t - 1/2)} > 0. \tag{4.7}
$$

Clearly, for $-x_o$, in $t = 1/2$, the inequality is violated and it suffices to consider the points $(x_0, \pm y_0)$. Note that for all $t \in [0, 1]$

$$
-2^{1/m} = \sqrt[m]{4(0-1/2)} \le \sqrt[m]{4(t-1/2)} \le \sqrt[m]{4(1-1/2)} = 2^{1/m}
$$

and hence, it suffices to show that

$$
2^{-1/m} > \theta y_o / x_o = z_o \tag{4.8}
$$

where z_o is the unique positive root of the polynomial $P \in \{P_{\text{even}}, P_{\text{odd}}\}$. Note that the sign change of *P* in z_0 must be from minus to plus, i.e., $P(z) < 0$ for $z < z_0$ and $P(z) > 0$ for $z > z_0$, it suffices to show that $P(2^{-1/m}) > 0$. Let us do this here only for P_{even} , the case of P_{odd} can be solved identically. Note that

$$
P_{\text{even}}(2^{-1/m}) = 2^{-1/m} \left[\sum_{i=1,3,\dots,k+1} \binom{k}{i-1} \frac{m}{m+i-1} - \theta \sum_{i=1,3,\dots,k-1} \binom{k}{i} \frac{m}{m+1+i} \right] > 0
$$

is implied by

$$
\sum_{i=1,3,\dots,k+1} {k \choose i-1} \frac{1}{m+i-1} - \sum_{i=1,3,\dots,k-1} {k \choose i} \frac{1}{m+1+i} > 0 \qquad (4.9)
$$

since θ < 1. We can further bound the left hand side of [\(4.9\)](#page-9-0) from below by

$$
\sum_{i=0}^{k} (-1)^{i} {k \choose i} \frac{1}{m+1+i} = \frac{k!m!}{(m+1+k)!}
$$

which is positive for all m, k . By direct computation we also see that $(1, 0)$ satisfies inequality (4.7) whereas $(0, 0)$ and $(-1, 0)$ do not. This completes the proof.

Proof of Theorem [3.1](#page-3-3) According to [\[4\]](#page-10-2), the translation-invariant splitting Gibbs measures are in one-to-one correspondence with the positive solutions of equation (2.3) for $K(t, u) = \exp(\xi_{t, u})$. By Proposition [3.2](#page-3-0) solutions to [\(2.3\)](#page-2-2) are characterized by 2-dimensional fixed points of the operator *Vk*,*n*. Now, Proposition [3.3](#page-3-1) provides a complete list of these fixed points in the two parameter regimes for θ . Proposition [3.3](#page-3-1) further asserts that only exactly one respectively three of them give rise to positive solutions. This completes the proof. **Acknowledgements** Golibjon Botirov thanks the DAAD program for the financial support and the Weierstrass Institute Berlin for its hospitality. Benedikt Jahnel thanks the Leibniz program 'Probabilistic methods for mobile ad-hoc networks' for the support.

References

- 1. Baxter, R.J.: Exactly Solved Models in Statistical Mechanics. Academic, London (1982)
- 2. Botirov, G.I.: A model with uncountable set of spin values on a Cayley tree: phase transitions. Positivity **21**(3), 955–961 (2017)
- 3. Eshkabilov, Y.K., Haydarov, F.H., Rozikov, U.A.: Non-uniqueness of Gibbs measure for models with uncountable set of spin values on a Cayley tree. J. Stat. Phys. **147**(4), 779–794 (2012)
- 4. Eshkabilov, Y.K., Rozikov, U.A.: On models with uncountable set of spin values on a Cayley tree: integral equations. Math. Phys. Anal. Geom. **13**, 275–286 (2010)
- 5. Botirov, G.I., Eshkabilov, Y.K., Rozikov, U.A.: Phase transition for a model with uncountable set of spin values on Cayley tree. Lobachevskii J. Math. **34**(3), 256–263 (2013)
- 6. Ganikhodjaev, N.N.: Potts model on Z*^d* with countable set of spin values. J. Math. Phys. **45**, 1121–1127 (2004)
- 7. Ganikhodjaev, N.N., Rozikov, U.A.: The Potts model with countable set of spin values on a Cayley tree. Lett. Math. Phys. **74**, 99–109 (2006)
- 8. Georgii, H.-O.: Gibbs Measures and Phase Transitions. De Gruyter, New York (2011)
- 9. Jahnel, B., Külske, C., Botirov, G.I.: Phase transition and critical values of a nearest-neighbor system with uncountable local state space on Cayley trees. Math. Phys. Anal. Geom. **17**, 323–331 (2014)
- 10. Krasnosel'skii, M.A.: Topological Methods in the Theory of Nonlinear Integral Equations. Macmillan, Basingstoke (1964)
- 11. Krasnosel'skii, M.A., Zabreiko, P.P.: Geometrical Methods of Nonlinear Analysis. Springer, Berlin (1984)
- 12. Rozikov, U.A.: Gibbs Measures on Cayley Trees. World Scientific, Singapore (2013)