**REVIEW**



# **New Development in Decarburization Research and Its Application to Spring Steels**

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# **Abstract**

Recently there was a new wave of research activities studying the decarburization behavior of spring steels with the main focus on the formation mechanism of a columnar ferrite layer within a certain temperature range which could not be explained by conventional decarburization theories. A new theory successfully developed recently in interpreting the oxide scale reduction mechanism on steel was then developed further and applied to interpret the observed columnar ferrite formation on spring steels. The essence of the new theory is that steel decarburization in the presence of a FeO scale on the steel surface is caused and governed by the reaction between the FeO scale and dissolved carbon in the steel, and therefore, the carbon concentration on the steel surface is determined by the FeO-steel interface equilibrium and cannot be treated as negligible within the temperature range where ferrite is able to form, because the equilibrium interface carbon concentration is in the same magnitude as the carbon solubility in ferrite. The new theory and available solutions for diferent decarburization scenarios using decarburization of 60Si2MnA as an example are summarized in this review. Explanations are given to interpret discrepancies between experimental observations and theoretical predictions. New areas for future research are also identifed.

**Keywords** Spring steel · 60Si2MnA · Decarburization · FeO · Columnar ferrite formation · Austenite · Carbon difusivity

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# **Introduction**

There has been a long history of studies of steel decarburization [\[1](#page-33-0)[–3](#page-33-1)]. In 1946, Pennington conducted a comprehensive study of carbon steel decarburization within the range of 691–927 °C in an atmosphere containing  $20\%H_2O-H_2$  under which steel oxidation could not occur [\[3](#page-33-1)]. In that study, a ferrite 'band' was found to develop on the steel at 732–893 °C but not at 691 and 927 °C with a maximum ferrite thickness observed at 790–815 °C. Pennington attributed the decreased ferrite thickness toward 905 °C to the decreasing solubility of carbon in ferrite to zero at 905 °C. One limitation of Pennington's analysis was that it was viewed that the carbon concentration in the ferrite formed was constant across the ferrite layer, and therefore, carbon difusion in ferrite did not contribute to its growth.

The simultaneous oxidation and decarburization were studied by Birks and coworkers in 1970 [\[4](#page-33-2)[–6\]](#page-33-3). Based on these studies, it was claimed that "the mechanism by which the decarburization of steels occur is well understood, particularly in the case of plain carbon and low alloy steels" [\[7](#page-33-4)]. However, the scope of Birks et al.'s studies was limited. First, the studies were focused on decarburization taking place in austenite only. No studies were directed to the situation where a ferrite layer could form. It was even claimed that as a result of the very low solubility of carbon in ferrite, the surface ferrite layer acted as a difusion barrier, leading to much lower decarburization rate than that at temperatures above  $910^{\circ}$  [\[5](#page-33-5), [8](#page-33-6)]. Secondly, while it was acknowledged that under an oxidizing atmosphere, decarburization took place via the reaction between the scale and dissolved carbon in the steel, Birks and co-workers continued to treat carbon concentration at the scale-steel interface as zero because of its low value. As will be demonstrated later, this assumption is acceptable only when decarburization takes place in the higher temperature range where the ferrite phase cannot form. Birks et al., citing Baud et al. [\[9\]](#page-33-7), also claimed that a carefully produced FeO scale could protect steel from decarburization because  $CO$  or  $CO<sub>2</sub>$  could not escape through the FeO layer. As will be revealed in this study, this again is not true. Finally, alloying efects on steel decarburization had not been addressed by Birks et al.

Despite the claim that the mechanism of decarburization was already well understood [[7\]](#page-33-4), there has been a new wave of spring steel decarburization studies recently [\[10](#page-33-8)[–26](#page-33-9)]. One reason for this was that when a spring steel was decarburized, a thick ferrite layer was always observed, same as that observed by Pennington [[3\]](#page-33-1) and there had been no consensus on the understanding of the mechanism governing its formation, and even on the temperature range where this ferrite layer could form. With the maximum ferrite thickness, some studies found that it occurred at 750  $^{\circ}$ C [[11–](#page-33-10)[16,](#page-33-11) [19](#page-33-12), [21](#page-33-13)], some found at 800 °C, [[22\]](#page-33-14) some said 850 °C [[17,](#page-33-15) [18](#page-33-16), [20](#page-33-17)] and another one claimed that it was within 900–950 °C [[25\]](#page-33-18).

On another front, in our recent study of the reduction mechanism of oxide scale in pure nitrogen, it was found that dissolved carbon in a low carbon steel could react with the FeO scale, leading to both scale reduction and decarburization of the steel [\[27](#page-33-19)]. The discovery led to the development of not only a new scale reduction theory, but also a new decarburization theory as a result of direct reaction between the FeO scale and dissolved carbon in the steel.

Subsequently, the newly developed theory was applied to the study of decarburization of a commonly used spring steel 60Si2MnA [[28–](#page-33-20)[30\]](#page-33-21). The studies were assisted by a thermodynamic computation package, the Pandat<sup>™</sup> program [[31\]](#page-33-22) and the PanFe data base [[32\]](#page-33-23) developed by CommuTherm. The current study provides an overview of the fndings of these studies, summarized the key steps of the new decarburization theory, and demonstrate how to apply the theory and available solutions to predict and interpret decarburization behavior under diferent conditions.

# **Experimental Studies for the Steel Grade 60Si2MnA**

The nominal compositions of two typical spring steel grades, 60Si2MnA and 55SiCr, are shown in Table [1.](#page-2-0) The 60Si2MnA steel examined in our recent studies contained Fe–0.602%C–1.787%Si–0.789%Mn–0.198%Cr. The 55SiCr grade contains a higher level of Cr and will be referred in this review for comparison.

Most of the previous experimental studies were conducted in ambient air with unknown moisture levels under isothermal conditions with the main objective to identify the ferrite formation temperature range and its formation mechanism. Moisture contents in the oxidizing gas can have a signifcant efect on silicon-containing steel oxidation [\[28](#page-33-20), [33\]](#page-33-24) and decarburization, which will be discussed when the issue arises.

#### **Isothermal Oxidation Studies**

Zhang et al.  $[13, 14]$  $[13, 14]$  $[13, 14]$  $[13, 14]$  $[13, 14]$  found that after oxidation in ambient air for 1.5 h, a columnar ferrite layer formed between 675 and 875 °C with a maximum thickness observed at 750 °C, as plotted in Fig. [1](#page-3-0). Another study conducted in a  $N_{2}$ –2%O<sub>2</sub> gas by Liu et al. also found that the ferrite layer formed after holding at 750 °C for 30 min was greater than that at 800 °C, whereas a ferrite layer did not form at 850 °C [\[20](#page-33-17), [34\]](#page-33-27). The absence of a ferrite layer at 850  $\degree$ C was unusual. There must be other reasons, such as gas moisture levels, contributing to this.

Under similar oxidation conditions, Yu et al. however found that the maximum ferrite thickness was observed at 800 °C after 60 min of oxidation in ambient air, [\[22](#page-33-14)] as compared in Fig. [1](#page-3-0). Zhao et al. also found a maximum ferrite thickness at 800 °C after oxidizing a 60Si2MnA steel in fowing air for 30 and 60 min [\[24](#page-33-28)]. On the other hand, in the study conducted by Y. Liu and X. Liu, very thick columnar ferrite layers were observed at 900–950  $\degree$ C after oxidation in ambient air in a muffle furnace for 90 min [\[25](#page-33-18)]. This observation was unusual because formation of a pure ferrite layer at 950 °C is thermodynamically impossible, as will be discussed later.

<span id="page-2-0"></span>



<span id="page-3-0"></span>**Fig. 1** Columnar ferrite formation on 60Si2MnA after oxidation and decarburization observed in diferent studies. Data sources are from Zhang et al. [[13,](#page-33-25) [14\]](#page-33-26), Yu et al. [\[22](#page-33-14)], and Chen et al. [\[28](#page-33-20)]

Recently we examined the decarburization behavior of 60Si2MnA in an atmos-phere containing 20%H<sub>2</sub>O-N<sub>2</sub> at 700–900 °C [\[28](#page-33-20)]. A scale layer of less than 6  $\mu$ m formed on the surface because the gas was not severely oxidizing. A columnar ferrite layer was observed within the range of 750–900 °C, but not at 700 °C, as plotted in Fig. [1](#page-3-0). Some ferrite morphologies are shown in Fig. [2](#page-3-1). It can be seen that the maximum ferrite thickness developed within 800–820 °C, consistent with the results obtained by Yu et al. [[34\]](#page-33-27) and Zhao et al. [[24\]](#page-33-28).

#### **Mechanism of Forming the Columnar Ferrite Structure**

Apart from the similar patterns of forming a maximum ferrite thickness at about 800 °C, the microstructures of the ferrite layers observed on 60Si2MnA were also similar to those observed by Pennington on a high carbon steel [[3\]](#page-33-1), with the ferrite



<span id="page-3-1"></span>**Fig. 2** Typical columnar ferrite layers formed after isothermal holding in a  $24.8H<sub>2</sub>O-N<sub>2</sub>$  gas for 20 min at **a** 750 °C, **b** 830 °C, and **c** 877 °C [[28\]](#page-33-20)

grains growing in a columnar structure spanning across the entire ferrite layer, as seen in Fig. [2](#page-3-1) as compared to the images presented by Pennington reproduced in Fig. [3](#page-4-0).

The development of this ferrite layer at 760, 788, 816, 871  $\degree$ C was explained by Pennington with the assistance of an early binary Fe–C diagram and that at 760 °C as an example. Within the  $\alpha + Y$  zone at this temperature, the equilibrium carbon concentrations in the ferrite and austenite were 0.026% and 0.50%, respectively (note that the more correct compositions should be  $0.016\%$  and  $0.54\%$  at 760 °C as calculated using Pandat™). While the steel was decarburized, there was a local equilibrium at the steel-gas interface to take away carbon from the steel surface. The surface carbon content was then decreasing continuously from initially 0.85% toward 0.50%. Once the surface concentration reached 0.50%, it could no longer decrease gradually, but dropped all the way to 0.026% to form a thin ferrite layer.

Pennington did not explain why the ferrite grains developed into a columnar shape, but judging from the morphology observed, after the formation of ferrite nuclei on the austenite surface, the continuous growth must have been taking place via a directional ferrite growth mechanism, opposite to the carbon difusion direction through the ferrite layer, thus developing into a columnar structure, similar to that developed under directional annealing of very low carbon steels, [[35\]](#page-33-29) or the columnar zone developed in an ingot where directional solidifcation took place [\[36](#page-33-30)].

## **Decarburization in the Absence of a Columnar Ferrite Layer (at T>900 °C)**

When the decarburization temperature is sufficiently high, the decarburization zone shows a diferent structure, as shown in Fig. [4](#page-5-0) [\[37](#page-33-31)]. Generally, there is a complete decarburization zone near the steel surface (Fig. [4](#page-5-0)b) where only light gray ferrite grains are visible. The morphology of this zone is diferent from the columnar ferrite structure formed within the range of 700–900 °C. At higher temperatures, this zone is in an austenitic state, containing very little carbon. It transforms to equiaxed



<span id="page-4-0"></span>**Fig. 3** Columnar ferrite structure developed on a high carbon steel (0.85wt%C–0.17%Mn–0.26%Si) in  $20\%$ H<sub>2</sub>O-H<sub>2</sub> at 816 °C for 1, 2, and 6 h (B, C, and D, respectively) [[3\]](#page-33-1)

ferrite grains during cooling. Under this zone, there is a partial decarburization zone comprising a mixture of light gray ferrite grains and dark gray pearlite colonies. The depth of the partial decarburization is estimated by locating a zone where the light gray ferrite grains started to disappear. The estimated depths of the complete decarburization and partial decarburization zones developed at diferent temperatures are shown in Fig. [5.](#page-6-0) The corresponding consumptions of steel by oxidation at diferent temperatures are also shown in Fig. [6](#page-7-0) [\[37](#page-33-31)].



(a) 1000 °C (b) 1120 °C



(c) 1150 °C (d) 1160 °C



(e) 1170 °C (f) 1200 °C

<span id="page-5-0"></span>**Fig. 4** Decarburization zones formed after isothermal holding for 20 min in air-24.8%H<sub>2</sub>O at 1000–1200 °C [\[30](#page-33-21), [37](#page-33-31)]: **a** 1000 °C, **b** 1120 °C, **c** 1150 °C, **d** 1160 °C, **e** 1170 °C and **f** 1200 °C

As compared to steel consumptions, several observations of the decarburization tendency can be made. First, there is a steady increase of the total decarburi-zation depth from 100 μm to 165 μm within the range of 1000–1120 °C (Fig. [5\)](#page-6-0), corresponding to a steady increase in steel consumption from 70  $\mu$ m at 1000 °C to 150 μm at 1120 °C but then the steel consumption rate stays about the same within 1120 °C–1150 °C (Fig. [6](#page-7-0)) while the decarburization depth continues to increase to about 180 μm at 1150 °C. Second, corresponding to the rapid increase in the amount of steel consumption beyond 1150  $\degree$ C (Fig. [6](#page-7-0)), the decarburization zone depth on the contrary reduced significantly at 1160  $^{\circ}$ C (Fig. [4](#page-5-0)d) and then to zero at 1170  $^{\circ}$ C (Figs. [4](#page-5-0)e and [5](#page-6-0)). Finally, beyond 1170 °C, the decarburization zone disappeared altogether, while the steel consumption rate increased rapidly (Figs. [4f](#page-5-0), [5](#page-6-0) and [6\)](#page-7-0).

Judging from the morphologies, the steel surface is covered by a molten oxide phase at temperatures above 1160 °C [\[30](#page-33-21)].

The disappearance of the decarburization zone at temperatures above 1170  $^{\circ}$ C is consistent with the observations by Y. Liu et al. [[20\]](#page-33-17) and J. Liu et al. [[26\]](#page-33-9)

## **Development of the New Decarburization Theory**

#### **Direct Reaction Between Dissolved Carbon and FeO Scale**

A study was conducted recently to understand the reason why there was always a weight loss when a hot-rolled low carbon steel was heated in high purity nitrogen [\[27](#page-33-19)]. At the reaction temperatures, the scale had a wustite structure. After holding at a temperature within the range of  $600-900$  °C, different degrees of scale reduction



<span id="page-6-0"></span>**Fig. 5** Measured total decarburization depths as a function of oxidation temperature after isothermal holding for 20 min in air-24.8% $H<sub>2</sub>O$  [\[37](#page-33-31)]



<span id="page-7-0"></span>**Fig. 6** Amount of steel consumed as a function of oxidation temperature after isothermal holding for 20 min in air-24.8%H<sub>2</sub>O [[30,](#page-33-21) [37\]](#page-33-31)

were observed, as shown Fig. [7](#page-8-0)a. The reduction kinetics initially followed the parabolic law but then deviated at longer times.

Examination of the scale structure and the steel substrate revealed that while the scale was reduced (Fig. [7b](#page-8-0)), the steel substrate underneath was decarburized (Fig. [7c](#page-8-0), d). Reduction of the scale generated porosities at the interface, but the outer region remained compact and continuous, as seen in Fig. [7](#page-8-0)b. After 3 h of reaction at 800 °C, the entire specimen was decarburized  $[27]$  $[27]$ 

Because the steel surface was covered by a compact scale, which was reduced gradually over time, there was only one mechanism that could be responsible for the observed decarburization phenomenon, which was the direct reaction between the FeO scale and dissolved carbon in the steel. Generation of numerous porosities did not stop the decarburization process.

Our observations are consistent with the fndings of another study by Pennington [\[38](#page-33-32)]. In that study, a number of pickled and unpickled steel specimens were stacked alternately with spacings between them and then isothermally heated in an enclosed chamber at 788 °C for 28 h. The gas volume generated by the reactions taking place between the specimens and the atmosphere in the chamber was then carefully measured and the remaining carbon contents in both sets of specimens were analyzed.

The theory behind the experiments was that if the remaining  $O_2$  in the chamber reacted with dissolved carbon in the steel to form  $CO<sub>2</sub>$ , one  $O<sub>2</sub>$  molecule would generate one  $CO_2$  molecule without causing gas expansion. However, if  $CO_2$  reacted with dissolved carbon again to form CO, then one  $CO<sub>2</sub>$  molecule would generate two CO molecules, thus causing gas expansion. The increase of the gas volume then could be measured and recorded over time to monitor the kinetics of reaction. The specimens after the treatment could then be analyzed for carbon losses.



<span id="page-8-0"></span>**Fig. 7** FeO-steel reaction during isothermal holding in high purity  $N_2$  causing decarburization of a low carbon steel (0.055%C–0.23%Mn–<0.005%Si): **a** scale reduction kinetics at 650–900 °C, **b** scale-steel interface after holding for 60 min, **c** decarburization zone in the steel after holding for 30 min; **d** decarburization zone in the steel after holding for 90 min at 800  $^{\circ}$ C [\[27](#page-33-19)]

The experiments convincingly demonstrated that decarburization reactions took place over a period of 16 h and the remaining carbon contents in both the scaled and pickled samples were reduced from the original 0.062wt% to 0.010wt% [\[38](#page-33-32)].

Therefore, the oxide reduction kinetics shown in Fig. [7](#page-8-0)a directly represented the kinetics of reaction between the FeO scale and dissolved carbon in the steel. Several observations can be made from Fig. [7a](#page-8-0). First, the reaction between dissolved carbon with the scale was possible at 700–850 °C. Second, there is a maximum rate observed at 750 and 800  $^{\circ}$ C, which was followed by that at 700  $^{\circ}$ C and then at 850 °C. Third, no reaction took place at 650 °C and finally, there was in fact a small weight gain at 900 °C, which was recognized as the result of oxidation of the steel by residual oxygen and water vapour in the nitrogen gas used.

Based on above observations and analyses, a theory was developed to interpret the scale reduction and decarburization phenomena. The theory was based on the propositions that (a) the reaction between the scale and the dissolved carbon in the steel could only proceed if the solubility of carbon in the matrix phase, being ferrite at 659–859 °C for the steel examined (Fe–0.055%C–0.23%Mn) as will be shown later, was greater than the carbon concentration at the FeO-steel interface, (b) the carbon concentration at the FeO-steel interface is determined by the local equilibrium achieved at the scale-steel interface, and (c) the reaction products  $CO<sub>2</sub>$  and  $CO$ could escape through the FeO scale layer even when the FeO layer was compact and continuous, as discussed elsewhere [[27\]](#page-33-19). Based on these propositions, the proposed reaction takes place via the following steps.

First, it is assumed that when the FeO scale reacts with dissolved carbon [C] in the steel,

both CO and  $CO<sub>2</sub>$  are produced:

$$
[C] + 2FeO = 2Fe + CO2
$$
 (1)

<span id="page-9-2"></span><span id="page-9-1"></span>
$$
[C] + FeO = Fe + CO \tag{2}
$$

The CO and  $CO<sub>2</sub>$  gases then maintain a local equilibrium with FeO and the steel:

<span id="page-9-0"></span>
$$
CO + FeO = Fe + CO2
$$
 (3)

Using available thermodynamic data in the literature [\[39](#page-33-33), [40\]](#page-33-34), the standard Gibbs free energy of formation for Reaction [\(3](#page-9-0)) is given by:

$$
\Delta G_{(3)}^o = -22800 + 24.267 \ T \ (J/mole \ of \ CO)
$$
 (4)

When the reaction  $(3)$  $(3)$  reaches equilibrium:

$$
\frac{P_{\text{CO}}}{P_{\text{CO}_2}} = \exp\left(\frac{-22800 + 24.267 \text{ T}}{\text{RT}}\right) \tag{5}
$$

The  $\frac{P_{\text{CO}}}{P_{\text{CO}_2}}$  thus obtained can be used as the second step to determine the equilibrium carbon activity at the interface assuming  $P_{\text{CO}} + P_{\text{CO}_2} = 1$  atm if the experiments are conducted under ambient pressure from the following reaction:

<span id="page-9-3"></span>
$$
[C] + CO2(g) = 2CO(g)
$$
 (6)

The standard Gibbs free energy of formation for this reaction is given by: [\[39](#page-33-33), [40](#page-33-34)]

$$
\Delta G_{(6)}^o = -RT \ln \left[ \frac{P_{\text{CO}}^2}{a_c P_{\text{CO}_2}} \right] = 170700 - 174.5T \quad (J/\text{mole of C}) \tag{7}
$$

where  $a_c$  is the equilibrium carbon activity at the FeO-steel interface with graphite being the standard state. From Eq. ([7\)](#page-10-0), we obtain:

<span id="page-10-1"></span><span id="page-10-0"></span>
$$
a_c = \frac{P_{\text{CO}}^2}{P_{\text{CO}_2}} \exp\left[\frac{(170700 - 174.5 \ T) \left(\frac{J}{\text{mole of } C}\right)}{RT}\right]
$$
(8)

After the equilibrium carbon activity is determined, the corresponding carbon concentration in the steel at the scale-steel interface can be calculated using known relationships between carbon activities and carbon concentrations as the third step.

For dissolved carbon in  $\alpha$ -Fe, the relationship to express the activity coefficient of carbon in ferrite for carbon steel, Υ*C*(fer rite), is given by Lobo and Gaiger: [[41\]](#page-33-35)

<span id="page-10-2"></span>
$$
\log \Upsilon_{\text{C}}(\text{ferrite}) = \log \left(\frac{a_{\text{C}}}{X_{\text{C}}}\right)(\text{ferrite}) = \frac{5846}{T(\text{K})} - 2.687\tag{9}
$$

where  $X_{\Gamma}$  is the equilibrium molar fraction of carbon in the steel matrix at the scalesteel interface, which can be converted to carbon concentration in weight percent.

The carbon solubilities in the ferrite phase of the experimental steel Fe–0.055%C–0.23%Mn, calculated using Pandat<sup>™</sup> [[31\]](#page-33-22) and the PanFe [\[32](#page-33-23)] database were then plotted in Fig. [8](#page-11-0) as the fourth step for the determination of reaction possibility. For comparison and later use, the carbon solubilities in ferrite for 60Si2MnA and 55SiCr, also obtained from Pandat<sup>™</sup> and PanFe are plotted in Fig. [8.](#page-11-0)

It can be seen that while the equilibrium carbon concentration at the steel surface is very low  $(<0.01\%)$ , it is not negligible as compared to the solubilities of carbon in ferrite, and more importantly, it can be used to assess within what temperature range where the dissolved carbon in the ferrite phase can react with the adjacent FeO scale. It can also be used to determine whether a ferrite layer can develop on a steel when the steel is being decarburized, such as that shown in Fig. [7](#page-8-0)c, d.

For the Fe–0.055%C–0.23%Mn steel, it can be seen from Fig. [8](#page-11-0) that the carbon solubilities within the range of 659  $\degree$ C to 859  $\degree$ C are greater than the calculated interface equilibrium carbon concentration, but smaller outside this range. This explains the reason why the FeO scale could react with dissolved carbon in the steel at [7](#page-8-0)00–850 °C, but not at 650 °C and 900 °C, as seen in Fig. 7a, if the steel is a ferritic state.

Figure [8](#page-11-0) also shows that there is a maximum diference between the carbon solubility in the ferrite phase and the equilibrium carbon concentration at the interface at about 720 °C. However, the maximum scale reduction or decarburization rate was not observed at 720  $\degree$ C in Fig. [7a](#page-8-0). The reason for this was that the rate of decarburization was also afected by the carbon difusivity in the ferrite phase.



<span id="page-11-0"></span>**Fig. 8** Calculated carbon concentration at the FeO-steel interface as a function of temperature (solid line), as compared with the carbon solubilities (dotted and dashed lines) in ferrite calculated using Pandat™ [[31\]](#page-33-22) and PanFe [\[32](#page-33-23)] to determine the possibility of reaction between the FeO scale and dissolved carbon in ferrite adjacent to the FeO scale

To represent the combined efect of the concentration diference and difusivity, the concept of carbon permeability  $(P_C^{\alpha-\text{Fe}})$  through the ferrite layer was introduced as the ffth step [[42\]](#page-33-36):

$$
P_{\rm C}^{\alpha - \rm Fe} = D_{\rm C}^{\alpha - Fe} \cdot (C_{\alpha}^{\alpha - \rm Fe/Bulk} - C_{\rm C}^{\alpha - \rm Fe/FeO}) = D_{\rm C}^{\alpha - Fe} \cdot \Delta C_{\rm C}^{\alpha - \rm Fe} \tag{10}
$$

where  $D_{\text{C}}^{\alpha-\text{Fe}}$  is the diffusion coefficient of carbon through the ferrite layer,  $C_{\text{C}}^{\alpha-\text{Fe/Bulk}}$ the equilibrium carbon concentration on the ferrite side at the interface between the ferrite layer and the bulk of steel in wt%,  $C_{\text{C}}^{\alpha-\text{Fe}/\text{FeO}}$  is the equilibrium carbon concentration in ferrite at the steel-scale interface in wt%, and  $\Delta C_C^{\alpha-\text{Fe}}$  is the carbon concentration diference between the two interfaces of the surface ferrite layer.

When the ferrite layer is thin, the carbon concentration gradient in the ferrite layer is approximately linear and the carbon flux  $J_{\rm C}^{\alpha - {\rm Fe}}$  flowing through the ferrite layer can be expressed as:

$$
J_C^{\alpha-\text{Fe}} = A \cdot D_C^{\alpha-\text{Fe}} \cdot \frac{\Delta C_C^{\alpha-\text{Fe}}}{X} = A \cdot \frac{P_C^{\alpha-\text{Fe}}}{X} \quad \text{moles cm}^{-2} \text{ s}^{-1} \tag{11}
$$

where *A* is a constant used to to convert the concentration of carbon from wt% to mole/cm<sup>3</sup> and  $X$  is the thickness of the ferrite layer in cm. The calculated carbon permeability as a function of temperature, using the carbon concentrations shown in Fig. [8](#page-11-0) and the carbon difusivities shown in Fig. [9](#page-12-0)a, derived from the work of Smith [\[42](#page-33-36)], for the Fe–0.055%C–0.23%Mn steel at 650–850 °C are shown in Fig. [9](#page-12-0)b. It is seen that the calculated permeabilities for the several experimental temperatures match the observed scale reduction and hence steel decarburization rate very well, with the maximum rate observed within the range of 750–800 °C, followed by that at 700 °C, then by a very low rate at 850 °C, and finally a zero rate at 650 °C.



<span id="page-12-0"></span>**Fig. 9 a** Difusivity of carbon in ferrite [[42\]](#page-33-36) and **b** carbon permeability through the ferrite layer within the range where FeO can reaction with dissolved carbon in the steel, as shown in Fig. [8](#page-11-0)

Thus, the combined effect of the carbon concentration difference  $\Delta C_{\text{C}}^{\alpha-\text{Fe}}$ which has a maximum value at about  $720\text{ °C}$  and the carbon diffusivity in ferrite  $D_{\rm C}^{\alpha-{\rm Fe}}$  which increases monotonically with temperature shifted the most rapid decarburization temperature to one within the range of  $750-800$  °C for the Fe–0.055%C–0.23%Mn steel.

It will be demonstrated later that the same mechanism is responsible for the development of a maximum columnar ferrite layer thickness at a certain temperature for the spring steel 60Si2MnA examined. Because the temperature range where the carbon solubility in ferrite is greater than the equilibrium carbon concentration at the FeO-ferrite interface is shifted higher to 679—912 °C for 60Si2MnA, as shown in Fig. [8](#page-11-0), this maximum ferrite thickness temperature is expected to be greater than the most rapid decarburization temperature observed for the Fe–0.055%C–0.23%Mn steel.

Finally, the A<sub>3</sub> line for the Fe–C–0.23%Mn system was computed using Pandat<sup>™</sup> and PanFe, as shown in Fig. [10](#page-13-0). It is seen that the Fe–0.055%C–0.23%Mn steel was in fact in an austenitic state at 900 °C.

In order to determine whether the dissolved carbon in the austenitic steel could be reacting with the scale, the interface carbon activity calculated from Eq. ([8\)](#page-10-1) was converted to the equilibrium carbon concentration in austenite using the equation given by Ellis et al. [[43\]](#page-34-0)

<span id="page-12-1"></span>
$$
\log a_C^{\gamma} = \log \left[ \frac{X_C^{\gamma}}{1 - 5X_C^{\gamma}} \right] + \frac{2080}{T} - 0.639 \tag{12}
$$

where  $X_C^{\gamma}$  is the molar fraction of carbon in austenite. The results thus obtained were converted to wt% and shown in Fig. [10](#page-13-0). It is seen that at 900  $^{\circ}C$ , the equilibrium carbon concentration at the interface is essentially identical to the carbon concentration in the steel. Therefore, Reactions  $(1)-(3)$  $(1)-(3)$  $(1)-(3)$  $(1)-(3)$  and  $(6)$  $(6)$  could not proceed at this temperature. This provides a satisfactory explanation for the observation that the FeO scale could not react with dissolved carbon in the steel at 900  $^{\circ}$ C as shown in Fig. [7a](#page-8-0).

The theory detailed above thus provides perfect interpretation of the observed reactions between the FeO scale on the steel surface and the dissolved carbon in



<span id="page-13-0"></span>**Fig. 10** Equilibrium carbon concentration at the scale-steel interface assuming that the steel phase in equilibrium is austenite, as compared to the high temperature ends of the  $A_3$  lines obtained from Fe–C– 1.787%Si–0.789%Mn–0.198%Cr and Fe–C–0.23%Mn isopleths

the Fe–0.055%C–0.23%Mn steel, leading to simultaneous scale reduction and steel decarburization. The principles then can be applied to the study of decarburization of other steels.

#### **Application of the New Theory to Decarburization of 60Si2MnA**

As shown in Table [1](#page-2-0), spring steels generally contain a much higher level of carbon (0.4–0.6%), a signifcant amount of silicon (1.5–2.0%) and various levels of Cr and Mn. As shown in Figs. [8](#page-11-0) and [10](#page-13-0), the alloying elements affect the temperature range of ferrite formation, carbon solubilities in ferrite, and the position of the  $A_3$  lines. For example, for the 60Si2MnA steel, formation of a columnar ferrite layer is thermodynamically impossible at temperature above 912 °C because the carbon solubilities at temperatures above 912 °C are below the equilibrium carbon concentrations at the FeO-ferrite interface, whereas for the 55SiCr steel, the critical temperature was at 890 °C.

However, ferrite formation is still possible in the 60Si2MnA steel examined (Fe–0.602%C–1.787%Si–0.789%Mn–0.198%Cr) at temperatures above 912 °C until the temperature exceeds 1025  $\degree$ C, but only as a component of a mixture with austenite if the carbon concentration in the surface layer is lowered to that under the  $A_3$ line shown in Fig. [10,](#page-13-0) as also shown in the isopleth computed using Pandat<sup>™</sup> and PanFe in Fig. [11a](#page-14-0).

As shown in Figs. [8](#page-11-0) and [10,](#page-13-0) when the 60Si2MnA steel is decarburized by the FeO scale, the equilibrium interface carbon concentration does not reduce to zero, whether the steel is in a ferritic or austenitic state within the range of 600–1050  $^{\circ}$ C. When a single-phase columnar ferrite forms at  $T=679-912$  °C, the interface carbon concentration is purely determined by Eq. ([9\)](#page-10-2) and Fig. [8](#page-11-0). When a single-phase



<span id="page-14-0"></span>**Fig. 11** Temperature zones where three diferent decarburization scenarios are operating: **a** full isopleth of the Fe–C–1.787%Si–0.789%Mn–0.198%Cr, **b** enlarged section for 700—800 °C of the Fe–C– 1.787%Si–0.789%Mn–0.198%Cr isopleth

austenite forms at temperatures above 990 °C, the interface composition is purely determined by Eq. [\(12](#page-12-1)) and Fig. [10.](#page-13-0) Between 912 and 990 °C, neither a single ferrite phase nor austenite phase can form on the surface of the steel. Therefore, within the temperature range of 912 °C and 990 °C, when the carbon concentration in the surface layer of the steel is decreased to that lower than the  $A_3$  value shown in Fig. [10,](#page-13-0)

there will be a two-phase ferrite–austenite structure developed on the steel surface, in equilibrium with the carbon activity determined by Eq.  $(8)$  $(8)$  with the concentrations of the ferrite and austenite determined by Eqs.  $(9)$  $(9)$  and  $(12)$  $(12)$  or Figs. [8](#page-11-0) and [10,](#page-13-0) respectively. This is a situation that had not been dealt with in previous studies. In this review, this will be treated as a situation where carbon difusion takes place in austenite only due to two reasons. First, the austenite phase is continuously present from the bulk to the interface. Second, the rate of carbon difusion is much more rapid in ferrite. Therefore, it is assumed that the presence of a ferrite phase in the surface layer does not impede on carbon transport from the bulk to the interface.

Based on above discussions, the three scenarios of carbon difusion in the steel during decarburization of a 60Si2MnA steel can be marked in Fig. [11a](#page-14-0) and schematically shown in Fig. [12,](#page-16-0) assuming no redistribution of the alloying elements during ferrite growth, namely under the para-equilibrium conditions [\[44](#page-34-1)].

Scenario 1—at T = 679 °C—774 °C:

The bulk of steel is a mixture of two-phase  $\alpha$ -Fe+γ-Fe at 742.4–774 °C,  $\alpha$ -Fe+ Fe<sub>3</sub>C at T < 731.6 °C, and three-phase  $\alpha$ -Fe+  $\gamma$ -Fe+ Fe<sub>3</sub>C state at 731.6–742.4 °C, as shown more clearly in Fig. [11](#page-14-0)b. The carbon concentration in the bulk equals to the original carbon concentration  $C_{\text{bulk}} = 0.602 \text{ wt\%}$ . The carbon concentration in ferrite at the ferrite-bulk interface is equal to the carbon concentration of the a-Fe phase in the bulk,  $C^*_{\alpha} = C_{\alpha \text{in bulk}}$ . The carbon concentration in ferrite at the FeO-ferrite phase is determined by the simultaneous equilibria of the interface Reactions  $(1)-(3)$  $(1)-(3)$  $(1)-(3)$  $(1)-(3)$  and  $(6)$  $(6)$ , as also indicated in Fig. [12a](#page-16-0).

Scenario II—at T=774–912 °C:

The carbon concentration in the bulk equals to the original carbon concentration,  $C_{\text{bulk}} = 0.602$  wt%. Carbon concentrations in austenite and ferrite at the α-γ interface,  $C^*_{\gamma}$  and  $C^*_{\alpha}$ , respectively, are determined by the equilibrium between the two phases, assuming no re-distribution of the alloying elements. The carbon concentration in ferrite at the FeO-ferrite interphase is determined by the simultaneous equilibria of the interface reactions as indicated, namely via Reactions ([1\)](#page-9-1)-([3](#page-9-2)) and ([6\)](#page-9-0).

Scenario III—at T>912 °C:

The carbon concentration in the bulk equals to the original carbon concentration,  $C_{\text{bulk}} = 0.602 \text{ wt\%}$ . The equilibrium carbon concentration in austenite at the FeOsteel interface,  $C_{\text{FeO-S}}$ , is determined by the simultaneous equilibria of the listed reactions at the interface, namely, Reactions  $(1)-(3)$  $(1)-(3)$  $(1)-(3)$  $(1)-(3)$  and  $(6)$  $(6)$ .

# **3. Analytical Solutions**

Analytical solutions are obtained by solving Fick's second law. The general solution under the assumption that the carbon diffusivity  $D$  is not concentration dependent:

<span id="page-16-0"></span>**Fig. 12** Three diferent decarburization scenarios as marked in Fig. [11a](#page-14-0), showing diferent difusion processes involved: **a** Scenario I, **b** Scenario II, and **c** Scenario III



<span id="page-17-4"></span>
$$
C_{x,t} = A + B \cdot \text{erf}\left[\frac{x}{2\sqrt{D \cdot t}}\right] \tag{13}
$$

where the parameters A and B are determined by specific initial and boundary conditions.

Decarburization of spring steels generally takes place together with steel oxidation. The only scenario where an analytical solution is available under simultaneous oxidation and decarburization conditions is Scenario III where carbon difusion takes place within austenite only with only one moving scale-steel interface. For the other two scenarios with two moving interfaces involved, analytical solutions could only be obtained under the condition that the movement of the scale-steel interface is negligible, namely  $x = \sqrt{k_n t} \approx 0$ , as indicated in Fig. [12a](#page-16-0), b, where  $k_p$  is the rate constant of steel oxidation [[45\]](#page-34-2).

#### **Analytical Solution for Scenario I**

In Scenario I, the bulk of steel is always in a two-phase or three-phase state. We only need to consider difusion in the surface ferrite layer, as shown in Fig. [12](#page-16-0)a. Following the procedure Wagner used, [[46\]](#page-34-3) the carbon concentration as a function of the distance through the ferrite layer and time is expressed as:

$$
C_{x,t} = C_{\text{FeO}-\alpha} + \frac{C_{\alpha}^* - C_{\text{FeO}-\alpha}}{\text{erf}(\Phi)} \text{erf}\left(\frac{x}{2\sqrt{D_C^{\alpha} \cdot t}}\right)
$$
(14)

where  $C^*_{\alpha}$  is the carbon concentration in ferrite at the interface between the ferrite layer and the bulk of steel and is equal to the ferrite carbon concentration in the bulk  $C^*_{\alpha} = C_{\alpha \text{ in bulk}}$ ,  $D^{\alpha}_{\text{C}}$  is the diffusivity of carbon in ferrite,  $\Phi$  is a dimensionless parameter at a given temperature, determined by:

$$
F(\Phi) = \sqrt{\pi} \times \left( C_{\text{bulk}} - C_{\alpha}^{*} \right) - \frac{C_{\alpha}^{*} - C_{\text{FeO}-\alpha}}{\Phi e^{\Phi^{2}} \text{erf}(\Phi)} = 0 \tag{15}
$$

Once the Φ meeting Eq. ([15\)](#page-17-0) is found, the ferrite layer thickness, *M*, can be calculated using the following equation:

<span id="page-17-3"></span><span id="page-17-2"></span><span id="page-17-1"></span><span id="page-17-0"></span>
$$
M = 2\Phi \sqrt{D_{\rm C}^{\alpha} \cdot t} \tag{16}
$$

Equation  $(15)$  $(15)$  had been then simplified by Smith to become the permeability equation:

$$
P_C^{\alpha - \text{Fe}} = D_C^{\alpha} C_{\alpha}^* = \frac{\left(C_{\text{bulk}} - \frac{2}{3} C_{\alpha}^*\right)}{2} \times \frac{M^2}{t}
$$
 (17)

or

 $\circled{2}$  Springer

<span id="page-18-0"></span>
$$
M = \sqrt{\frac{6 \cdot C_{\alpha}^* \cdot D_C^{\alpha} \cdot t}{3C_{\text{bulk}} - 2 \cdot C_{\alpha}^*}}
$$
(18)

Equation ([18\)](#page-18-0) will be called the Smith's equation, or Smith's permeability equation. In deriving this equation, Smith had assumed that the surface carbon concentration was negligible, *i.e.*  $C_s \cong 0$ . It had been the equation used by many researchers to calculate the ferrite layer thickness caused by decarburization in their studies [[10,](#page-33-8) [19](#page-33-12), [21](#page-33-13), [47,](#page-34-4) [48\]](#page-34-5). Because it had been simplifed from Eq. [\(15](#page-17-0)), it can only be used to calculate ferrite thickness under Scenario I conditions, where only carbon difusion through the ferrite layer contributes to the growth of the ferrite layer.

If however, the carbon concentration on the surface is not negligible, namely  $C_s \neq 0$ , Eq. [\(17](#page-17-1)) then takes the form of:

$$
P_C^{\alpha - \text{Fe}} = D_C^{\alpha} (C_{\alpha}^* - C_s) = \frac{\left( C_{\text{bulk}} - \frac{2}{3} C_{\alpha}^* \right)}{2} \times \frac{M^2}{t}
$$
(19)

and Eq. [\(18](#page-18-0)) becomes:

<span id="page-18-2"></span>
$$
M = \sqrt{\frac{6 \cdot (C_{\alpha}^* - C_{\text{FeO}-\alpha}) \cdot D_{\text{C}}^{\alpha} \cdot t}{3C_{\text{bulk}} - 2 \cdot C_{\alpha}^*}}
$$
(20)

This equation will be called the modifed Smith's equation to be used under the conditions where the interface carbon concentration is not negligible.

#### **Analytical Solution for Scenario II**

In this scenario, two difusion processes are contributing to the formation of the ferrite layer, as shown in Fig. [12b](#page-16-0). Carbon difusion through the ferrite layer leads to its own thickening, whereas carbon difusion from the bulk of the steel to the ferrite–austenite interface tends to slow down the thickening of the ferrite layer.

As the boundary conditions for the ferrite layer growth are the same at all temperatures, except that now  $C^*_{\alpha} = C^{\alpha-\gamma}_{\alpha}$ , regardless of whether carbon diffusion in the austenite contributes to the ferrite growth or not, the carbon concentration in the ferrite layer can still be described by Eq. ([14\)](#page-17-2) but takes the form of,

<span id="page-18-1"></span>
$$
C_{x,t}^{\alpha} = C_{\text{FeO}-\alpha} + \frac{C_{\alpha}^{\alpha-\gamma} - C_{\text{FeO}-\alpha}}{\text{erf}(\Phi)} \text{erf}\left(\frac{x}{2\sqrt{D_{C}^{\alpha} \cdot t}}\right)
$$
(21)

after replacing  $C^*_{\alpha}$  with  $C^{\alpha-\gamma}_{\alpha}$ . Following the procedure described by Wagner [\[46](#page-34-3)] the carbon concentration in austenite as a function of distance and time can be expressed as:

$$
C_{x,t}^{\gamma} = C_{\text{bulk}} - \frac{C_{\text{bulk}} - C_{\gamma}^{\alpha - \gamma}}{\text{erfc}\left(\cdot \frac{\Phi}{\beta}\right)} \text{erfc}\left(\frac{x}{2\sqrt{D_{\text{C}}^{\gamma} \cdot t}}\right) \tag{22}
$$

where  $D_{\text{C}}^{\gamma}$  is the carbon diffusivity in austenite, in cm<sup>2</sup>/sec, and the constant  $\beta$  is defned as,

<span id="page-19-0"></span>
$$
\beta = \frac{\sqrt{D_{\rm C}^{\gamma}}}{\sqrt{D_{\rm C}^{\alpha}}} \tag{23}
$$

The  $\Phi$  value in Eqs. [\(21](#page-18-1)) and [\(22](#page-19-0)) can be obtained by solving the following equation, which was derived by considering mass balance at the  $\alpha$ -γ interface,

$$
F(\Phi) = \sqrt{\pi} \cdot (C_{\gamma}^{\alpha-\gamma} - C_{\alpha}^{\alpha-\gamma}) - \left[ \frac{C_{\gamma}^{\alpha-\gamma} - C_{\text{bulk}}}{\frac{\Phi}{\beta} \cdot e^{\left(\frac{\Phi}{\beta}\right)^{2}} \cdot \text{erfc}\left(\frac{\Phi}{\beta}\right)} + \frac{C_{\alpha}^{\alpha-\gamma} - C_{\text{FeO}-\alpha}}{\Phi \cdot e^{\Phi^{2}} \cdot \text{erf}(\Phi)} \right] = 0
$$
\n(24)

which can then be used to calculate the ferrite layer thickness using Eq. [\(16](#page-17-3)).

Using the Wagner's equations with the consideration of interface equilibrium, namely Eqs. [\(14](#page-17-2))-[\(15](#page-17-0)) for the temperature range of 700–774 °C and Eqs. [\(21](#page-18-1))-[\(24](#page-19-0)) for the range of 774–907 °C to represent the difusion processes involved at diferent temperatures, the ferrite layer thickness for the range of 700–907 °C can be calculated using Eq.  $(16)$  $(16)$ , as shown by the purple solid line in Fig. [13.](#page-20-0) It is seen that it has a maximum value at 800  $^{\circ}$ C, a much smaller value at 700  $^{\circ}$ C and a very small value at 900 °C, having a similar trend to that of the experimental results shown in Fig. [1](#page-3-0).

If the interface carbon concentration is assumed to be zero when the Wagner's equations are used, the calculated ferrite thickness as a function of temperature is shown as the shorter green dash line in Fig. [13.](#page-20-0) It is seen that the maximum ferrite thickness is now at 829 °C, and the predicted ferrite thicknesses are signifcant greater at all temperatures. Particularly at  $900\text{ °C}$ , the calculated ferrite thickness is about 80  $\mu$ m, much greater than the observed thickness at 900 °C of 30  $\mu$ m (Fig. [1](#page-3-0)).

If the Smith's equation without considering the interface equilibrium, namely Eq. ([18\)](#page-18-0), is used, the calculated ferrite thicknesses as shown by the longer blue dash line in Fig. [13](#page-20-0) became much greater beyond 774 °C, the maximum ferrite thickness temperature is moved to 859  $\degree$ C, and the calculated ferrite thickness at 900  $\degree$ C is even greater at about 100 µm.

If the modifed Smith's equation, Eq. [\(20](#page-18-2)), is used to take interface equilibrium into consideration, the calculated thicknesses are also greater than those calculated using the Wagner's equation beyond 770  $\degree$ C, as shown by the red dot line in Fig. [13,](#page-20-0) although the predicted ferrite thickness at 900  $^{\circ}$ C is very much reduced at about 40  $\mu$ m, closer to the observed thickness of 30  $\mu$ m (Fig. [1](#page-3-0)).

The calculated results shown in Fig. [13](#page-20-0) have several indications. First, the interface carbon content signifcantly afects the calculated results. If this content is



<span id="page-20-0"></span>**Fig. 13** Calculated ferrite thickness as a function of decarburization temperature using diferent equations

assumed to be zero, the calculated thicknesses are much greater at all temperatures, the calculated maximum ferrite thickness temperatures are shifted to higher temperatures, and most importantly, the trends of thickness variation with temperature deviate significantly from the observed trend shown in Fig. [1](#page-3-0).

Second, carbon diffusion in austenite at temperatures above  $774\text{ °C}$  affects the calculated ferrite layer thickness. Without considering this efect, namely, when the Smith equations are used, the calculated ferrite layer growth within the range of 774–900  $\degree$ C can deviate from the calculated ferrite thicknesses as much as more than 15  $\mu$ m. Because the formation of austenite at temperatures above the A<sub>3</sub> temperature for the 60Si2MnA steel examined is a reality, as seen in Fig. [11a](#page-14-0), the use of the Smith equations is therefore inappropriate for the prediction of ferrite growth for temperatures greater than the  $A_3$  temperature.

Thus, the most reliable calculated results are those shown by the sold purple line in Fig. [13](#page-20-0) as reproduced as the red solid line in Fig. [14](#page-21-0) to compare with experimental results (blue dash line). It is seen that while the predicted trend of ferrite thickness variation with temperature matches with the experimental results well, with a maximum thickness observed at around 800 °C and very small thicknesses at 700 and 900  $\degree$ C, the observed ferrite thicknesses are 40  $\mu$ m greater than the calculated ferrite thicknesses at around 800 °C. To obtain a perfect match, both in the trends and in the absolute values, we need to identify the reasons responsible for the thickness discrepancy near 800 °C. One reason that could be responsible is that the diffusivity data used for calculating the ferrite layer thickness could be incorrect. There are two possible causes for this.

First, in calculating the ferrite thickness using Eq. [\(16](#page-17-3)), the Smith's difusivity data were used [[42\]](#page-33-36). A closer examination of the technique used by Smith revealed



<span id="page-21-0"></span>**Fig. 14** Comparison of measured thicknesses of the columnar ferrite layers at diferent temperatures and the thicknesses calculated using the Wagner's equations assuming that the carbon concentrations are determined by interface equilibrium between FeO and dissolved carbon in the steel. The steel examined contains Fe–0.602%C–1.787%Si–0.789%Mn–0.198%Cr

that the method employed may not be reliable because Smith used Eq.  $(18)$  $(18)$  to compute the diffusion coefficient. As discussed earlier, in deriving Eq.  $(18)$  $(18)$ , Smith assumed that the carbon concentration on the steel surface was zero. Such an assumption could lead to error because Smith used a  $H_2O-H_2$  mixture produced by 'burning metered amounts of hydrogen and oxygen using a gas burner' as the decarburizing gas [[42\]](#page-33-36). While the exact gas composition was not given, it was likely that the gas mixture would have a relatively high oxygen potential that corresponded to an equilibrium surface carbon concentration that was not negligible. If this was the case, then Eq.  $(20)$  $(20)$ , rather than Eq.  $(18)$  $(18)$  should be used to derive the carbon diffusivity. When  $C_s \neq 0$ ,  $(C^*_\alpha - C_s)$  is always smaller than  $C^*_\alpha$ , and therefore, a greater  $D_C^{\alpha}$  value must be obtained from Eq. [\(20](#page-18-2)) for a given ferrite layer thickness. The actual deviation would depend on the exact gas composition applied by Smith in his experiment [[42\]](#page-33-36)

The second factor is the possible alloying efect on carbon difusivity in ferrite. The steel examined in the experimental study contained 1.787%Si–0.789%Mn–0.198%Cr [[28\]](#page-33-20). Based on the data presented by Krishtal, the presence of Si increases carbon difusivity in ferrite, and the increase is more significant at higher temperatures, as shown in Fig.  $15$  [\[49](#page-34-6)]. The effect of Cr is even greater at higher temperatures (Fig. [15](#page-22-0)) [\[49](#page-34-6)], and there could be the combined efect of Cr and Si, which could be even more profound, similar to that observed in austenite  $[55]$  $[55]$ . By calculation using Eq.  $(16)$  $(16)$ , if the true carbon diffusivity at around 800 °C is twice as large, then the predicted maximum ferrite thickness can match the observed values well. Future studies will be required to understand the true reasons of the observed discrepancy.



<span id="page-22-0"></span>**Fig. 15** Efect of Si and Cr additions on carbon difusivity in ferrite, expressed as the ratio between the carbon difusivity in an Fe-Si or Fe–Cr alloy and carbon difusivity in pure iron [\[49](#page-34-6)]

#### **Analytical Solution for Scenario III**

Applying the boundary condition specifed in Fig. [12](#page-16-0)c and the initial condition of  $C_{x,t} = C_C^{\gamma - bulk}$  at  $x > 0$  at  $t = 0$  to Eq. [\(13](#page-17-4)), one obtains: [\[37](#page-33-31)]

<span id="page-22-1"></span>
$$
C_{x,t} = C_{\text{bulk}} - \frac{(C_{\text{bulk}} - C_{\text{FeO}-\gamma})}{\text{erfc}\left[\frac{\sqrt{k_P}}{2\sqrt{D'_C t}}\right]} \text{erfc}\left[\frac{x}{2\sqrt{D'_C t}}\right]
$$
(25)

where  $C_{x,t}$  is the carbon concentration at the location *x* at time *t* and  $D_C^{\gamma}$  is the carbon difusivity in austenite, assumed to be a constant at a certain temperature. This equation is equivalent to that derived by Birks et al. [[4–](#page-33-2)[7\]](#page-33-4) where the scale-steel interface position is expressed as  $x^2 = 2k<sub>c</sub>t$ , rather than  $x^2 = k<sub>p</sub>t$  as defined in Fig. [12](#page-16-0), namely  $k_P = 2k_C$ . The assumption of parabolic scaling is convenient in arriving at a simple solution because it allows one to eliminate the time variable *t* from the term *X*  $\sqrt{D_{\epsilon}^{\gamma}}$  $\frac{v}{c} \cdot t} = \frac{\sqrt{k_p t}}{2\sqrt{D_c^{\gamma}}}$  $rac{\sqrt{k_p t}}{2\sqrt{D'_C \cdot t}} = \frac{\sqrt{k_p}}{2\sqrt{L}}$  $\frac{\sqrt{k_p}}{2\sqrt{D'_C}}$ , when the boundary condition  $C_{x,t} = C_C^{\text{FeO}-\gamma}$  at the scalesteel interface  $x = X$  is applied.

The calculated carbon concentration profles in the partial decarburization zones in a 60Si2MnA steel 1000–1150 °C after 20 min of oxidation in wet air that contains  $24.8\%$ H<sub>2</sub>O are shown in Fig. [16a](#page-23-0) [\[37](#page-33-31)]. Consistent with the experimental observations shown in Figs. [4](#page-5-0) and [5](#page-6-0), the calculated decarburization depth increased with increasing temperature within this range, as summarized in Table [2.](#page-24-0)





<span id="page-23-0"></span>**Fig. 16** Calculated carbon concentration gradients in austentite at diferent temperatures within the range of 1000–1150 °C: **a** for 60Si2MnA using carbon difusivity data derived for Fe–0.2%C–1.2%Si–0.2%Cr, and **b** for 55SiCr using difusivity data derived for Fe–0.2%C–1.2%Si–0.7%Cr (showing the combined efect of Si with a higher Cr content)

# **Alloying Efects on Decarburization**

# **Efects on Carbon Difusivity**

The analytical solutions given in Eqs.  $(13)$  $(13)$ ,  $(14)$  $(14)$ ,  $(21)$  $(21)$ ,  $(22)$  $(22)$  and  $(25)$  $(25)$  were obtained under the assumption that carbon difusivity was not a function of carbon

<span id="page-24-0"></span>



concentration in the relevant phases. It is well known that the difusivity of carbon in austenite is a function of carbon concentration  $[50-54]$  $[50-54]$ . It is also affected significantly by the presence of alloying elements [\[49](#page-34-6), [55](#page-34-7)[–61](#page-34-10)]. The most comprehensive equation currently available was that derived by Lee et al., [\[57](#page-34-11)] who incorporated the binary alloying efects provided by Krishtal: [[55\]](#page-34-7)

<span id="page-24-2"></span>
$$
D_{\rm C}^{\gamma} = \left[ 0.146 - 0.036C(1 - 1.075C_{\rm F}) \right] + k_{1,\rm Si}M_{\rm Si} + k_{1,\rm Mn}M_{\rm Mn} + k_{1,\rm Cr}M_{\rm Cr} \right] \cdot \exp\left[ -\frac{144.3 - 15.0C + 0.37C^2 + k_{2,\rm Si}M_{\rm Si} + k_{2,\rm Mn}M_{\rm Mn} + k_{2,\rm Cr}M_{\rm Cr}}{R_{kj}T} \right]
$$
(26)

where  $k_{1,i}$  ( $i = \text{Si}$ , Mn and Cr) and  $k_{2,i}$  ( $i = \text{Si}$ , Mn and Cr) are given in Table [3,](#page-24-1) and  $M_i(i = Si, Mn)$  and Cr) is the concentration for Si, Mn and Cr, respectively, in  $wt\%$ .

However, in a closer examination, it was found that the authors had used experimental data for ternary Fe-M-C systems (M being a metallic element) only [\[55](#page-34-7), [61\]](#page-34-10). More valuable data for quaternary systems Fe–1.2wt%Si–0.7%C–xCr also provided by Krishtal [[55\]](#page-34-7) were not incorporated in the derivation of Eq. ([26\)](#page-24-2). One important fnding reported by Krishtal was that while Cr and Si alone had retarding efects to carbon difusion, the combined efect of Cr and Si was much more profound, as shown in Table [4.](#page-25-0) Taking the steels containing Fe–0.7C for example. An addition of 1.6%Si decreased the carbon difusivity by about 10%, whereas if 1.6%Si and 0.2%Cr were simultaneously added to Fe–0.7%C, the carbon difusivity was reduced by more than six times for the temperature range of 1000–1200 °C.

The much lower carbon difusivity can produce a much shallower decarburization depth, as shown in Table [2,](#page-24-0) which also listed the calculated results for 55SiCr shown in Fig. [16b](#page-23-0) and the calculated results for 60Si2MnA using difusivity data calculated using Eq. [\(26](#page-24-2)) for comparison. It can be seen that if the calculated results

<span id="page-24-1"></span>



and ternary systems Fe–0.2C–1.6Si and Fe–0.7C–1.6Si systems to factor in the effects of chromium and carbon concentrations, with reference to the binary Fe–0.2C and

<span id="page-25-0"></span>Fe–0.7C systems and other ternary systems

using difusivity data derived from Eq. [\(26](#page-24-2)) had represented the true decarburization depths when the combined efect of Cr and Si were absent, then the presence of the combined Cr-Si efect in 60Si2MnA decreases this depth by 60%. The predicted results for 55SiCr also show that by further increasing the Cr content to 0.7%, the depths can be deceased by a further 40%.

Very few experimental studies were conducted to address the efect of carbon concentration and alloying elements on the difusivity of carbon in ferrite [\[62](#page-34-13)[–67](#page-34-14)] Based on available experimental data, Silva and McLellan derived an equation to calculate  $D_C^{\alpha}$ : [\[68](#page-34-15)]

<span id="page-26-0"></span>
$$
Ln(D_C^{\alpha}) = -2.087 - 1.197 \cdot \chi + 0.037 \chi^2 \text{ cm}^2/\text{s}
$$
 (27)

where  $\chi = \frac{10000}{T}$  °K<sup>-1</sup>. Because Smith data were used as the most important source for the derivation, the predicted ferrite layer growth using data calculated from Eq. [\(27\)](#page-26-0) was essentially the same as those obtained when Smith's original data were used.

As discussed earlier, Smith's data were unlikely to be reliable because the surface carbon concentration was assumed to be zero, and even if they were reliable, they were only for carbon difusion in Fe–C system only. For spring steels, alloying effects must be also considered, as shown in Fig. [16](#page-23-0). The data for plotting Fig. [15](#page-22-0) are shown in Table [5](#page-26-1) together with data from some other alloy systems. It is seen that both Si and Cr increased the activation energy of carbon difusion, with Cr having a more profound efect. However, because the data were obtained by an indirect method using  $C^{14}$  through intermittent carburization over the temperature range 500–800  $\degree$ C, they can only be used as an indication of the alloying effects, rather than for quantitative predictions [\[49](#page-34-6)].

More importantly, as the combined efect of Cr and Si shown for the carbon diffusivity in austenite suggested, there could be the combined efect of Cr and Si on the carbon difusivity in ferrite that cannot be predicted by simply adding up the separate efects of Si and Cr. If this efect is understood, the discrepancy between the observed and calculated results shown in Fig. [14](#page-21-0) then may be resolved.

#### **Efects of Molten Oxide Formation**

carbon diffu

The most commonly used alloying element in spring steels is silicon. It is known that silicon in steel can provide good protection to high temperature oxidation below 1170  $\degree$ C, but at temperatures above 1173  $\degree$ C, formation of a molten oxide phase

<span id="page-26-1"></span>

becomes possible as the reaction product of  $FeSi<sub>2</sub>O<sub>4</sub>$  and FeO, which generally coexist at the scale-steel interface during oxidation of the steel [[30\]](#page-33-21), as indicated in the FeO-SiO<sub>2</sub> phase diagram shown in Fig. [17](#page-27-0)  $[69, 70]$  $[69, 70]$  $[69, 70]$ . Formation of the molten oxide drastically accelerated steel oxidation [[30,](#page-33-21) [37\]](#page-33-31), as shown in Fig. [6](#page-7-0). If Fe<sub>3</sub>O<sub>4</sub> can form at the scale-steel interface, the molten oxide formation temperature can be even lowered to 1140–1150 °C [[71–](#page-34-18)[73\]](#page-34-19).

When liquid oxide formed, it can erode the steel surface piece by piece, even before the steel was completely decarburized, as seen in Fig. [18.](#page-28-0) 'Erosion' of the steel in this way will consume a signifcant part of the decarburized layer. However, some degree of partial decarburization would still be expected if the oxide can react with dissolved carbon in the steel because the carbon difusion rate in the steel is also increased signifcantly if there is a loss of carbon on the steel surface. This view is supported by the observation of oxygen difusion into the steel substrate causing internal oxidation to accompany steel erosion by the molten oxide, as seen in Fig. [19](#page-29-0). If the more difficult oxygen diffusion could take place, so could carbon, if carbon had been lost at the scale-steel interface. Therefore, the complete absence of partial decarburization suggested that there must be some other mechanisms operating to prevent decarburization reactions at the scale-steel interface.

Thermodynamic calculation indicated that the equilibrium carbon activity at the  $Fe<sub>2</sub>SiO<sub>4</sub>/steel$  interface was always significantly lower than the carbon activity in the bulk of the steel (60Si2MnA), as shown in Fig. [20](#page-29-1) [[37\]](#page-33-31). The assessments required the use of silicon activity in the steel, which was computed using Pandat™. For the 60Si2MnA steel assessed. the silicon activity is a function of temperature, given by:



<span id="page-27-0"></span>**Fig.** 17  $\text{FeO-SiO}_2$  phase diagram in the presence of iron  $[70, 71]$  $[70, 71]$  $[70, 71]$  $[70, 71]$ 



<span id="page-28-0"></span>**Fig. 18** Steel pieces eroded from the steel surface and encapsulated by the molten oxide after oxidation of a 60Si2MnA steel at 1190 °C in 24.8%H<sub>2</sub>O-air for 20 min [[30,](#page-33-21) [37\]](#page-33-31)

$$
a_{\rm Si} = 0.255 \times e^{-\frac{116930}{RT}}
$$
 (28)

where T is temperature in K and R the gas constant in  $J \cdot \text{mol}^{-1} \cdot \text{T}^{-1}$ .

Diferent degrees of internal oxidation may consume silicon from the surface layer to diferent degrees, and therefore, the assessments were conducted by comparing the carbon activity in the bulk of steel and that at the interface with diferent degrees of remaining silicon activities, as shown in Fig. [20](#page-29-1) [[37\]](#page-33-31). The results indicated that the formation of a continuous and compact  $FeSi<sub>2</sub>O<sub>4</sub>$  oxide formation on the steel surface could provide full protection from decarburization below 1084 °C. This temperature decreases with decreased level of silicon activity on the surface layer caused by internal oxidation [[37\]](#page-33-31).

Without the formation of a molten oxide layer, the inner scale layer was found to comprise a mixture of  $Fe<sub>2</sub>SiO<sub>4</sub>$  and FeO [[30\]](#page-33-21) and therefore, the interface equilibrium was efectively one between FeO and the steel. When a molten oxide forms, however, the interface carbon activity is always signifcantly lower than that in the bulk, as seen in Fig. [20](#page-29-1), and therefore, decarburization would be still possible if the molten oxide can react with the dissolved carbon and the reaction products CO and or  $CO<sub>2</sub>$  gases can escape through the molten oxide. The complete absence of decarburization thus suggested that the molten oxide was impermeable to  $CO$  and  $CO<sub>2</sub>$ gases. The reason could be that a liquid phase does not have micro-cracks of grain boundaries for the gases to travel through. In addition, because the molten oxide could wet both the steel and the FeO phase well, as shown in Figs. [4](#page-5-0)d–f, [18](#page-28-0) and [19,](#page-29-0) it could then spread along the interface to cover the entire steel surface during steel oxidation, thus preventing the steel from decarburization completely.



<span id="page-29-0"></span>**Fig. 19** SEM image of the scale-steel interface in the specimen oxidized in wet air at 1200 °C. The different phases indicated were determined by EDS analyses [[30,](#page-33-21) [37\]](#page-33-31)



<span id="page-29-1"></span>Fig. 20 Effect of Fe<sub>2</sub>SiO<sub>4</sub> formation on decarburization at different levels of remaining Si activity in the surface layer of the steel caused by internal oxidation

#### **Effect of Possible Formation of SiO<sub>2</sub> on the Steel Surface**

One interesting fnding of an earlier study was that when the 60Si2MnA steel was exposed to a dry  $O<sub>2</sub>$ -containing atmosphere, decarburization could be prevented or eliminated with very little scale formed [[28\]](#page-33-20). The observed phenomenon was attributed to the possible formation of a highly protective  $SiO<sub>2</sub>$  scale covering the steel surface, thus providing protection from both oxidation and decarburization. The possible reactions dominating decarburization under this condition are:

$$
[C] + SiO2 = [Si] + CO2
$$
 (29)

<span id="page-30-0"></span>
$$
[C] + SiO2 = [Si] + CO \tag{30}
$$

The reaction products  $CO<sub>2</sub>$  and CO then maintain their equilibrium with [Si] and [C] via the following reaction:

$$
2CO_2 + [Si] = 2CO + SiO_2 \tag{31}
$$

and the reaction shown in Eq. ([6\)](#page-9-3). The free energy of formation for Reaction [\(31](#page-30-0)) is:

$$
\Delta G_T^0 = -306690 + 7.531T = -RTln \frac{P_{CO}^2}{P_{CO_2}^2 \cdot a_{Si}}
$$
(32)

or

$$
\frac{P_{\rm CO}}{P_{\rm CO_2}} = e^{-\left[\frac{(-306690 + 7.531T)}{2RT} - \frac{ln a_{\rm Si}}{2}\right]}
$$
(33)

The calculated  $\frac{P_{\text{CO}}}{P_{\text{CO}_2}}$  then can be used to calculate the equilibrium carbon activity using Eq. ([8\)](#page-10-1), which then can be converted equilibrium carbon concentration at the interface assuming that the steel is in an austenitic state using Eq.  $(12)$  $(12)$ . Table [6](#page-31-0) compares the calculated equilibrium carbon concentrations at the  $SiO<sub>2</sub>$ -steel interface and that in the steel, as well as those in equilibrium with FeO and  $Fe<sub>2</sub>SiO<sub>4</sub>$ , for the range of 800–1250 °C. It is seen that the calculated equilibrium carbon concentrations are extremely high, much higher than that in the bulk, indicating that when a compact  $SiO<sub>2</sub>$  formed on the surface preventing direct contact between the steel and the atmosphere, decarburization can be prevented, as already observed [[28\]](#page-33-20)

# **Conclusions**

Steel decarburization is a very old research topic and also an unresolved problem for medium and high carbon steel manufacturers. Despite the claim by Birks et al. in 1983 that the mechanism of decarburization for plain carbon and low alloy steels was already well understood, conventional theories could not explain many aspects of steel decarburization. One such aspect was the formation mechanism of a columnar ferrite layer on steel which became a hot topic recently in the study of spring

T. <sup>o</sup> C	FeO	Solid Fe <sub>2</sub> SiO <sub>4</sub>	Molten $Fe2SiO4$	SiO <sub>2</sub>	Bulk of steel—aus- tenite
800	$\sharp N/A$	4.788	$\sharp N/A$	5.086	0.602
900	0.0548	3.490	$\sharp N/A$	5.000	0.602
1000	0.0245	1.494	$\sharp N/A$	4.548	0.602
1100	0.0123	0.560	$\sharp N/A$	3.244	0.602
1150	0.00897	0.341	$\sharp N/A$	2.376	0.602
1200	0.00670	$\sharp N/A$	0.211	1.571	0.602
1250	0.00509	$\sharp N/A$	0.137	$\sharp N/A$	0.602

<span id="page-31-0"></span>**Table 6** Calculated equilibrium carbon concentrations at the oxide-steel interface covered by diferent types of oxides

steel decarburization. This review summarizes the fndings of our recent studies on this topic, particularly the development of a new decarburization theory to interpret the formation mechanism of the columnar ferrite layer, leading to the following conclusions:

- (1) A columnar ferrite layer is generally observed in the surface area of a spring steel if the steel is exposed to an oxidizing atmosphere within the temperature range of 700–900 °C. Diferent studies claimed diferent temperatures at which a maximum columnar ferrite layer thickness developed and there had been no consensus on the mechanism responsible for its formation. Our studies found that this temperature is within the range of 800–820 °C and the presence of such a maximum ferrite thickness could be explained by the new simultaneous scale reduction and steel decarburization theory we developed recently.
- (2) The new theory proposes that when a steel is covered by a FeO scale, the FeO scale, whether it is in contact or detached from the steel surface, could react with dissolved carbon in the steel, causing simultaneous reduction of the FeO scale itself and decarburization of the steel. Whether the reaction can take place or not depends on whether the carbon activity in the steel is greater than the equilibrium carbon activity at the FeO-steel interface when the surface oxides are solid which were found to be permeable to  $CO$  and  $CO<sub>2</sub>$ . However, at temperatures above 1170 °C, formation of a liquid oxide phase is possible, which is able to spread and cover the steel surface thus preventing steel decarburization altogether.
- (3) The new theory not only satisfactorily explains the observed reduction kinetics of the FeO scale, but more importantly for this study provides a good explanation of the columnar ferrite formation phenomenon with a maximum thickness observed at 800–820 °C. The columnar ferrite structure is believed to be caused by directional growth of the ferrite phase toward to steel substrate opposite to the carbon difusion direction as a result of gradual carbon loss from the surface of the steel, whereas the presence of a maximum ferrite thickness is the combined efect of three factors: the presence of a maximum carbon solubility in ferrite at

about 720 °C, gradual decreasing trend of the equilibrium carbon concentration at the ferrite-scale interface with temperature, and the rapid increase of carbon difusivity in ferrite with temperature, leading to the presence of a maximum carbon permeability through the ferrite layer, which is defned as the product of the carbon concentration diference between the two interfaces of the ferrite layer and the carbon difusivity in ferrite.

- (4) Historically in dealing with steel decarburization under oxidizing conditions, it was always assumed that the carbon concentration on the steel surface was negligible and treated as zero. Based on the new theory, this is no longer the case when the formation of a ferrite layer is possible, because within the temperature range where ferrite is stable, the magnitude of the equilibrium carbon concentration at the FeO-steel interface is in the same magnitude as that of carbon solubility in ferrite.
- (5) Incorporating the new interface conditions determined by the FeO-steel equilibrium, a set of equations for calculating carbon concentration gradient in the austenite are summarized in this study. However, it was found that there was signifcant discrepancy in the predicted and experimentally measured decarburization depths in austenite when the carbon difusivity data summarized by Lee et al*.* were used.
- (6) Further examination of available date showing the alloying efects on carbon difusivity revealed that in the presence of both Cr and Si in the steel, the Lee's equation should not be used as it did not consider the combined efect of these two elements.
- (7) It was also found that the difusivity data derived by Smith may not be reliable either because Smith had used an overly simplifed Wagner's equation to derive the difusivity data. In Smith's study, the carbon concentration at the steel surface was assumed to be zero. Based on the fnding of the current study, this could have led to error because the equilibrium carbon concentration on the steel surface could be in the same magnitude as the carbon solubility in ferrite.
- (8) The study also demonstrates how the theory can be applied to predict and explain decarburization tendency when the steel surface is covered by other oxides, such as  $Fe<sub>2</sub>SiO<sub>4</sub>$ , whether solid or molten, and SiO<sub>2</sub>.

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**Author contributions** The manuscript and most of the fgures and all tables were prepared by Dr YRC. Dr FZ computed thermodynamic data, assisted in the computation of the phase diagrams and assisted in the interpretation of thermodynamic principles related to the interpretation of the research results. Both authors reviewed the manuscript.

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# **Declarations**

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