

Optical electroosmotic magnetic density with antiferromagnetic model

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Abstract

In our manuscript, we obtain optical angular momentum of $\phi(\alpha)$, $\phi(\mathbf{t})$, $\phi(\mathbf{s})$ by using optical spherical frame. Also, we construct magnitude of optical angular momentum and spherical magnetic moment of $\phi(\alpha)$, $\phi(\mathbf{t})$, $\phi(\mathbf{s})$. Thus, we illustrate optical $\tau_{\phi(\alpha)}$, $\tau_{\phi(\mathbf{t})}$, $\tau_{\phi(\mathbf{s})}$ magnetic torque phase microscale. Moreover, we have electroosmotic microfluidic $\tau_{\phi(\alpha)}$, $\tau_{\phi(\mathbf{t})}$, $\tau_{\phi(\mathbf{s})}$, $\tau_{\phi(\mathbf{s})}$, magnetic torque density. Finally, we design electroosmotic magnetic torque density with antiferromagnetic model.

Keywords Optical angular momentum \cdot Spherical magnetic moment \cdot Magnetic torque density \cdot Torque phase

1 Introduction

The most generally optical models for electromagnetic energy introduce electromagnetic flux, thermal drift energy, extrusion modeling. Optical thermal energy engages numerical design and macroscopic applications with thermal energy photonics. Optical thermal energy applications have conducted electroosmotic illustrations with different multi optical optimistic energy. Thermal energy applications associate photonic geometric phase, optoelectronic geometric phase, piezoelectric phase, photodiode phase. Semirecursional optical flux is constructed in optical microscale sensing methods (Yamada et al. 2011; Amjadi et al. 2016; Yan et al. 2018; Tao et al. 2012; Abouraddy et al. 2007; Fink et al. 1998; Korpinar and Körpinar 2021; Körpinar 2021; Körpinar and Korpinar 2021a, b).

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Optical energy improvement of quasi electronics and spherical designs with energy modeling machinery is inflation important optical applications. Optical hybrid flux is decisive significance in mathematical energy physics, geometric applications, optical geometry in optical energy (Garcia de Andrade 2006c, a, b; Vieira and Horley 2012; Guo and Ding 2008; Hasimoto 1972; Jones 1941; Biener et al. 2002; Berry and Klein 1996; Zygelman 1987; Son and Yamamoto 2012; Smit 1955; Körpınar and Demirkol 2019; Körpınar et al. 2019, 2019a, b; Cao et al. 2018; Körpınar and Demirkol 2020; Cao et al. 2017; Erb et al. 2016; Körpınar et al. 2020).

Microfluidic optical energy with recursional sensing phases are regularly experienced in optical spherical applications, optical design, PDT, optical signal theory, PTT. Modeling of spherical flux systems are associated by diverse optical and physical systems with antiferromagnetic Heisenberg optical model, Landau–Lifshitz optical model, Schrödinger optical model, binormal optical model, Da Rios optical model, localized induction optical model, etc. (Furlani and Ng 2006; Tomita and Chiao 1986; Körpinar et al. 2020; Balakrishnan et al. 1990; Wassmann and Ankiewicz 1998; Balakrishnan et al. 1993; Balakrishnan and Dandoloff 1999; Körpinar and Demirkol 2018a, b; Körpinar 2020; Dandoloff and Zakrzewski 1989; Kugler and Shtrikman 1988; Yamashita 2012; Satija and Balakrishna 2009; Lamb 1977; Yamashita 2012; Murugesh and Balakrishnan 2001; Gürbüz 2005; Körpinar et al. 2021a; Körpinar et al. 2021b; Körpinar and Körpinar 2021a; Körpinar et al. 2021b; Körpinar et al. 2022).

The establishment of our paper is as follows. First, we obtain optical angular momentum of $\phi(\alpha)$, $\phi(\mathbf{t})$, $\phi(\mathbf{s})$ by using optical spherical frame. Also, we construct magnitude of optical angular momentum and spherical magnetic moment of $\phi(\alpha)$, $\phi(\mathbf{t})$, $\phi(\mathbf{s})$. Thus, we illustrate optical $\tau_{\phi(\alpha)}$, $\tau_{\phi(\mathbf{t})}$, $\tau_{\phi(\mathbf{s})}$ magnetic torque phase microscale. Moreover, we have electroosmotic microfluidic $\tau_{\phi(\alpha)}$, $\tau_{\phi(\mathbf{t})}$, $\tau_{\phi(\mathbf{s})}$ magnetic torque density. Finally, we design electroosmotic magnetic torque density with antiferromagnetic model.

2 Spherical optical magnetic torque

Spherical frame of $\alpha : \mathbb{I} \to \mathbb{S}^2$ is given by

$$\nabla_{\sigma} \alpha = \mathbf{t}$$
$$\nabla_{\sigma} \mathbf{t} = -\alpha + \varepsilon \mathbf{s}$$
$$\nabla_{\sigma} \mathbf{s} = -\varepsilon \mathbf{t},$$

and

$$\alpha = \mathbf{t} \times \mathbf{s}, \ \mathbf{t} = \mathbf{s} \times \boldsymbol{\alpha}, \ \mathbf{s} = \boldsymbol{\alpha} \times \mathbf{t}.$$

Lorentz forces and magnetic field are given

$$\phi(\alpha) = \mathbf{t} + \rho \mathbf{s},$$

$$\phi(\mathbf{t}) = -\alpha + \varepsilon \mathbf{s},$$

$$\phi(\mathbf{s}) = -\rho\alpha - \varepsilon \mathbf{t},$$

$$\mathcal{B} = \varepsilon \alpha - \rho \mathbf{t} + \mathbf{s}$$

Optical angular momentum is obtained by

$$\otimes = f \times m^f \Psi(f).$$

Optical magnitude of angular momentum is

$$\|\otimes\| = \left\| f \times m^f \Psi(f) \right\|.$$

Optical magnetic moment is given

$$\theta_{\phi(f)} = \frac{\omega^{f}}{2m^{f}} \otimes .$$

Optical magnetic torque given by

$$\tau_{\phi(f)} = \theta \times \mathcal{B}.$$

Also, we have

$$\begin{split} \theta_{\phi(\alpha)} &= \frac{1}{2} \rho \omega^{\alpha} \alpha, \\ \theta_{\phi(t)} &= \frac{1}{2} \varepsilon \omega^{t} \alpha + \frac{1}{2} \omega^{t} \mathbf{s}, \\ \theta_{\phi(s)} &= \frac{1}{2} \rho \omega^{s} \mathbf{s}. \end{split}$$

Putting

$$\frac{\partial \alpha}{\partial t} = \mathcal{L}\mathbf{t} + \mathcal{F}\mathbf{s}.$$

• Flows of forces are obtained by

$$\nabla_{t}\phi(\boldsymbol{\alpha}) = -(\mathcal{L}+\rho\mathcal{F})\boldsymbol{\alpha} - \rho\left(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial\sigma}\right)\mathbf{t} + \left(\frac{\partial\rho}{\partial t} + \left(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial\sigma}\right)\right)\mathbf{s},$$

$$\nabla_{t}\phi(\mathbf{t}) = -\mathcal{F}\varepsilon\boldsymbol{\alpha} - \left(\mathcal{L} + \left(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial\sigma}\right)\varepsilon + \right)\mathbf{t} + \left(-\mathcal{F} + \frac{\partial\varepsilon}{\partial t}\right)\mathbf{s},$$

$$\nabla_{t}\phi(\mathbf{s}) = \left(-\frac{\partial\rho}{\partial t} + \mathcal{L}\right)\boldsymbol{\alpha} - \left(\rho\mathcal{L} + \frac{\partial\varepsilon}{\partial t}\right)\mathbf{t} - \left(\varepsilon\left(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial\sigma}\right) + \rho\mathcal{F}\right)\mathbf{s}.$$

• Flows of magnetic moments are

$$\begin{split} \nabla_{t}\theta_{\phi(\alpha)} &= \frac{1}{2}\frac{\partial\rho}{\partial t}\omega^{\alpha}\alpha + \frac{1}{2}\rho\mathcal{L}\omega^{\alpha}\mathbf{t} + \frac{\rho}{2}\omega^{\alpha}\mathcal{F}\mathbf{s}, \\ \nabla_{t}\theta_{\phi(\mathbf{t})} &= \left(\frac{1}{2}\omega^{\mathbf{t}}\frac{\partial\varepsilon}{\partial t} - \frac{\mathcal{F}}{2}\omega^{\mathbf{t}}\right)\alpha + \left(\frac{1}{2}\varepsilon\omega^{\mathbf{t}}\mathcal{L} - \frac{1}{2}\left(\frac{\partial\mathcal{F}}{\partial\sigma} + \mathcal{L}\varepsilon\right)\omega^{\mathbf{t}}\right)\mathbf{t} + \frac{\varepsilon}{2}\omega^{\mathbf{t}}\mathcal{F}\mathbf{s}, \\ \nabla_{t}\theta_{\phi(\mathbf{s})} &= -\frac{1}{2}\mathcal{F}\rho\omega^{\mathbf{s}}\alpha - \frac{1}{2}\omega^{\mathbf{s}}\rho\left(\frac{\partial\mathcal{F}}{\partial\sigma} + \mathcal{L}\varepsilon\right)\mathbf{t} + \frac{1}{2}\frac{\partial}{\partial t}\rho\omega^{\mathbf{s}}\mathbf{s}. \end{split}$$

3 Spherical magnetic torque phase microscale

3.1 Microfluidic antiferromagnetic $au_{\phi(lpha)}$ magnetic torque phase

• $\tau_{\phi(\alpha)}$ magnetic torque phase microscale is

$${}^{\mathcal{B}}\mathcal{M}_{\phi(\mathbf{t}_{\boldsymbol{\omega}})} = \tau_0^{\boldsymbol{\alpha}} \int \int_{\mathcal{W}} \nabla_t \phi(\boldsymbol{\alpha}) \cdot \nabla_t \tau_{\phi(\boldsymbol{\alpha})} d\mathcal{W}.$$

Optical angular momentum of $\phi(\alpha)$ is given

$$\otimes_{\phi(\alpha)} = \rho m^{\alpha} \alpha.$$

The magnitude of optical angular momentum is

$$\left\|\bigotimes_{\phi(\alpha)}\right\| = \rho m^{\alpha}.$$

Spherical magnetic moment of $\phi(\alpha)$ is

$$\theta_{\phi(\alpha)} = \frac{\omega^{\alpha}}{2} \rho \alpha$$

 $\tau_{\phi(\alpha)}$ magnetic torque of $\phi(\alpha)$ is constructed by

$$\tau_{\phi(\alpha)} = -\frac{1}{2}\rho\omega^{\alpha}\mathbf{t} - \frac{1}{2}\rho^{2}\omega^{\alpha}\mathbf{s}.$$

Thus, fluid of $\tau_{\phi(\alpha)}$ magnetic torque is

$$\nabla_{t}\tau_{\phi(\alpha)} = \left(\frac{1}{2}\rho\mathcal{L}\omega^{\alpha} + \frac{1}{2}\mathcal{F}\omega^{\alpha}\rho^{2}\right)\alpha + \left(\frac{1}{2}\rho^{2}\omega^{\alpha}\left(\frac{\partial\mathcal{F}}{\partial\sigma} + \mathcal{L}\varepsilon\right) - \frac{1}{2}\frac{\partial\rho}{\partial t}\omega^{\alpha}\right)\mathbf{t} \\ - \left(\frac{1}{2}\rho\omega^{\alpha}\left(\frac{\partial\mathcal{F}}{\partial\sigma} + \varepsilon\mathcal{L}\right) + \rho\omega^{\alpha}\frac{\partial\rho}{\partial t}\right)\mathbf{s}.$$

• Electroosmotic microfluidic $\tau_{\phi(\alpha)}$ magnetic torque density is

$$\begin{split} \mathcal{V}_{\tau_{\phi(\alpha)}} &= -(\mathcal{L} + \rho \mathcal{F}) \Big(\frac{1}{2} \rho \omega^{\alpha} \mathcal{L} + \frac{1}{2} \mathcal{F} \rho^{2} \omega^{\alpha} \Big) \\ &\quad - \rho \Big(\mathcal{L} \varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \Big) \Big(\frac{1}{2} \rho^{2} \omega^{\alpha} \Big(\frac{\partial \mathcal{F}}{\partial \sigma} + \mathcal{L} \varepsilon \Big) - \frac{1}{2} \omega^{\alpha} \frac{\partial \rho}{\partial t} \Big) \\ &\quad - \Big(\frac{1}{2} \rho \omega^{\alpha} \Big(\frac{\partial \mathcal{F}}{\partial \sigma} + \varepsilon \mathcal{L} \Big) + \rho \frac{\partial \rho}{\partial t} \omega^{\alpha} \Big) \Big(\Big(\mathcal{L} \varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \Big) + \frac{\partial \rho}{\partial t} \Big). \end{split}$$

• Microfluidic $\tau_{\phi(\alpha)}$ magnetic torque phase microscale is given

$$\begin{split} \mathcal{T}_{\tau_{\phi(\alpha)}} &= \tau_0^{\alpha} \int \int_{\mathcal{W}} \left(-\rho \Big(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \Big) \Big(\frac{1}{2} \rho^2 \omega^{\alpha} \Big(\frac{\partial \mathcal{F}}{\partial \sigma} + \mathcal{L}\varepsilon \Big) - \frac{1}{2} \frac{\partial \rho}{\partial t} \omega^{\alpha} \right) \\ &- (\mathcal{L} + \rho \mathcal{F}) \Big(\frac{1}{2} \rho \omega^{\alpha} \mathcal{L} + \frac{1}{2} \mathcal{F} \rho^2 \omega^{\alpha} \Big) \\ &- \Big(\frac{1}{2} \rho \omega^{\alpha} \Big(\frac{\partial \mathcal{F}}{\partial \sigma} + \varepsilon \mathcal{L} \Big) + \rho \frac{\partial \rho}{\partial t} \omega^{\alpha} \Big) \Big(\Big(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \Big) + \frac{\partial \rho}{\partial t} \Big) \Big) d\mathcal{W}, \end{split}$$

where τ_0^{α} is magnetic torque constant for $\phi(\alpha)$. Microfluidic antiferromagnetic $\tau_{\phi(\alpha)}$ magnetic torque phase microscale is given ٠

$${}^{\mathcal{B}}\mathcal{M}_{\tau_{\phi(\alpha)}} = \tau_0^{\alpha} \int \int_{\mathcal{W}} \nabla_t \phi(\alpha) \cdot \tau_{\phi(\alpha)} \times \nabla_{\sigma} \tau_{\phi(\alpha)} d\mathcal{W}.$$

From spherical frame, we have

$$\tau_{\phi(\alpha)} \times \nabla_{\sigma} \tau_{\phi(\alpha)} = \left(\frac{1}{2}\rho\omega^{\alpha} \left(\rho\omega^{\alpha}\frac{\partial\rho}{\partial\sigma} + \frac{\omega^{\alpha}}{2}\varepsilon\right) + \left(\varepsilon\rho^{2} - \frac{\partial\rho}{\partial\sigma}\right) \times \left(\frac{1}{2}\omega^{\alpha}\right)^{2}\rho^{2}\right) \alpha - \left(\frac{1}{2}\omega^{\alpha}\right)^{2}\rho^{2}\mathbf{t} + \left(\frac{1}{2}\omega^{\alpha}\right)^{2}\rho\mathbf{s}.$$

Electroosmotic antiferromagnetic microfluidic $au_{\phi(\alpha)}$ *magnetic torque density is* .

$$\begin{split} \mathcal{V}^{\mathcal{A}}_{\tau_{\phi(\alpha)}} &= -(\mathcal{L} + \rho \mathcal{F}) \bigg(\frac{1}{2} \rho \bigg(\rho \omega^{\alpha} \frac{\partial \rho}{\partial \sigma} + \frac{\omega^{\alpha}}{2} \varepsilon \bigg) \omega^{\alpha} \\ &+ \bigg(\varepsilon \rho^{2} - \frac{\partial \rho}{\partial \sigma} \bigg) \bigg(\frac{1}{2} \omega^{\alpha} \bigg)^{2} \rho^{2} \bigg) + \bigg(\frac{1}{2} \omega^{\alpha} \bigg)^{2} \rho^{3} \bigg(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \bigg) \\ &+ \bigg(\frac{1}{2} \omega^{\alpha} \bigg)^{2} \rho \bigg(\bigg(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \bigg) + \frac{\partial \rho}{\partial t} \bigg). \end{split}$$

Microfluidic antiferromagnetic $au_{\phi(\alpha)}$ magnetic torque phase microscale is

$$\begin{split} \mathcal{T}^{\mathcal{A}}_{\tau_{\phi(\alpha)}} &= \tau_{0}^{\alpha} \int \int_{\mathcal{W}} \left(\left(\frac{1}{2}\omega^{\alpha}\right)^{2} \rho^{3} \left(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial\sigma}\right) - \left(\frac{1}{2}\omega^{\alpha}\rho \left(\rho\omega^{\alpha}\frac{\partial\rho}{\partial\sigma} + \frac{1}{2}\varepsilon\omega^{\alpha}\right)\right) \\ &+ \left(\varepsilon\rho^{2} - \frac{\partial\rho}{\partial\sigma}\right) \left(\frac{1}{2}\omega^{\alpha}\right)^{2}\rho^{2} \right) (\mathcal{L} + \rho\mathcal{F}) \\ &+ \left(\frac{1}{2}\omega^{\alpha}\right)^{2}\rho \left(\left(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial\sigma}\right) + \frac{\partial\rho}{\partial t}\right) \right) d\mathcal{W}, \end{split}$$

where τ_0^{α} is magnetic torque constant for $\phi(\alpha)$. Also, we easily get

$$\phi(\boldsymbol{\alpha}) \times \nabla_{\sigma} \phi(\boldsymbol{\alpha}) = \left(\frac{\partial \rho}{\partial \sigma} + \rho^2 \varepsilon + \varepsilon\right) \boldsymbol{\alpha} - \rho \mathbf{t} + \mathbf{s}.$$

Electroosmotic $\tau_{\phi(\alpha)}$ magnetic torque density with antiferromagnetic $\phi(\alpha)$ is •

Antiferromagnetic $\tau_{\phi(\alpha)}$ magnetic torque phase microscale with antiferromagnetic $\phi(\alpha)$ • is

$$\begin{split} {}^{\phi(\alpha)}\mathcal{T}^{\mathcal{A}}_{\tau_{\phi(\alpha)}} &= \tau_{0}^{\alpha} \int \int_{\mathcal{W}} \left(\left(\frac{1}{2}\omega^{\alpha}\right)^{2} \rho^{3} + \left(\frac{\partial\rho}{\partial\sigma} + \rho^{2}\varepsilon + \varepsilon\right) \left(\frac{1}{2}\rho\omega^{\alpha} \left(\rho\omega^{\alpha}\frac{\partial\rho}{\partial\sigma} + \frac{1}{2}\varepsilon\omega^{\alpha}\right) + \left(\varepsilon\rho^{2} - \frac{\partial\rho}{\partial\sigma}\right) \left(\frac{1}{2}\omega^{\alpha}\right)^{2}\rho^{2} \right) + \left(\frac{1}{2}\omega^{\alpha}\right)^{2}\rho \right) d\mathcal{W}, \end{split}$$

where τ_0^{α} is magnetic torque constant for $\phi(\alpha)$. Optical effect of microfluidic antiferromagnetic $\tau_{\phi(\alpha)}$ magnetic torque phase microscale is constructed vortex reactors in Fig. 1.

3.2 Microfluidic antiferromagnetic $au_{\phi(t)}$ magnetic torque phase

 $\tau_{\phi(\mathbf{t})}$ magnetic torque phase microscale is

$${}^{\mathcal{B}}\mathcal{M}_{\phi(\mathbf{t})} = \tau_0^{\mathbf{t}} \int \int_{\mathcal{W}} \nabla_t \phi(\mathbf{t}) \cdot \nabla_t \tau_{\phi(\mathbf{t})} d\mathcal{W}.$$

Optical angular momentum of $\phi(\mathbf{t})$ is given

 $\otimes_{\phi(\mathbf{t})} = m^{\mathbf{t}}\mathbf{s} + \varepsilon m^{\mathbf{t}}\boldsymbol{\alpha}.$

The magnitude of optical angular momentum is



Fig. 1 Microfluidic antiferromagnetic $\tau_{\phi(\alpha)}$

$$\left\|\bigotimes_{\phi(\mathbf{t})}\right\| = m^{\mathbf{t}}(1+\varepsilon^2)^{\frac{1}{2}}.$$

Spherical magnetic moment of $\phi(\mathbf{t})$ is

$$\theta_{\phi(\mathbf{t})} = \frac{1}{2} \varepsilon \omega^{\mathbf{t}} \boldsymbol{\alpha} + \frac{1}{2} \omega^{\mathbf{t}} \mathbf{s}.$$

 $\tau_{\phi(\mathbf{t})}$ magnetic torque of $\phi(\mathbf{t})$ is constructed by

$$\tau_{\phi(\mathbf{t})} = \frac{1}{2}\rho\omega^{\mathbf{t}}\boldsymbol{\alpha} - \frac{1}{2}\varepsilon\omega^{\mathbf{t}}\rho\mathbf{s}.$$

Thus, fluid of $\tau_{\phi(t)}$ magnetic torque is

$$\begin{split} \nabla_t \tau_{\phi(\mathbf{t})} &= \left(\frac{1}{2}\frac{\partial\rho}{\partial t}\omega^{\mathbf{t}} + \frac{1}{2}\mathcal{F}(\varepsilon\rho)\right)\omega^{\mathbf{t}}\boldsymbol{\alpha} + \left(\frac{\omega^{\mathbf{t}}}{2}(\varepsilon\rho)\left(\frac{\partial\mathcal{F}}{\partial\sigma} + \mathcal{L}\varepsilon\right)\right. \\ &\left. + \frac{1}{2}\rho\omega^{\mathbf{t}}\mathcal{L}\right)\mathbf{t} + \left(\frac{1}{2}\rho\omega^{\mathbf{t}}\mathcal{F} - \frac{1}{2}\frac{\partial}{\partial t}(\varepsilon\rho)\omega^{\mathbf{t}}\right)\mathbf{s}. \end{split}$$

Electroosmotic microfluidic $\tau_{\phi(t)}$ magnetic torque density is •

$$\begin{split} \mathcal{V}_{\tau_{\phi(\mathbf{t})}} &= -\varepsilon \mathcal{F} \bigg(\frac{1}{2} \frac{\partial \rho}{\partial t} \omega^{\phi(\mathbf{t})} + \frac{1}{2} \mathcal{F}(\varepsilon \rho) \omega^{\phi(\mathbf{t})} \bigg) \\ &- \Big(\varepsilon \Big(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \Big) + \mathcal{L} \Big) \Big(\frac{1}{2} (\varepsilon \rho) \omega^{\phi(\mathbf{t})} \Big(\frac{\partial \mathcal{F}}{\partial \sigma} + \mathcal{L}\varepsilon \Big) \\ &+ \frac{1}{2} \rho \omega^{\phi(\mathbf{t})} \mathcal{L} \Big) + \Big(\frac{1}{2} \rho \mathcal{F} \omega^{\phi(\mathbf{t})} - \frac{1}{2} \frac{\partial}{\partial t} (\varepsilon \rho) \omega^{\phi(\mathbf{t})} \Big) \Big(\frac{\partial \varepsilon}{\partial t} - \mathcal{F} \Big). \end{split}$$

Microfluidic $\tau_{\phi(t)}$ magnetic torque phase microscale is given

$$\begin{split} \mathcal{T}_{\tau_{\phi(t)}} &= \tau_0^{\mathbf{t}} \int \int_{\mathcal{W}} \left(\left(\frac{1}{2} \rho \omega^{\mathbf{t}} \mathcal{F} - \frac{1}{2} \omega^{\mathbf{t}} \frac{\partial}{\partial t} (\varepsilon \rho) \right) \left(\frac{\partial \varepsilon}{\partial t} - \mathcal{F} \right) \\ &- \varepsilon \mathcal{F} \bigg(\frac{1}{2} \omega^{\mathbf{t}} \frac{\partial \rho}{\partial t} + \frac{1}{2} \mathcal{F} (\varepsilon \rho) \omega^{\mathbf{t}} \bigg) - \bigg(\varepsilon \bigg(\mathcal{L} \varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \bigg) + \mathcal{L} \bigg) \\ &\times \bigg(\frac{1}{2} (\varepsilon \rho) \bigg(\frac{\partial \mathcal{F}}{\partial \sigma} + \mathcal{L} \varepsilon \bigg) \omega^{\mathbf{t}} + \frac{1}{2} \rho \omega^{\mathbf{t}} \mathcal{L} \bigg) \bigg) d\mathcal{W}, \end{split}$$

where $\tau_0^{\mathbf{t}}$ is magnetic torque constant for $\phi(\mathbf{t})$. Microfluidic antiferromagnetic $\tau_{\phi(\mathbf{t})}$ magnetic torque phase microscale is given •

$${}^{\mathcal{B}}\mathcal{M}_{\tau_{\phi(\mathbf{t})}} = \tau_0^{\mathbf{t}} \int \int_{\mathcal{W}} \nabla_t \phi(\mathbf{t}) \cdot \tau_{\phi(\mathbf{t})} \times \nabla_\sigma \tau_{\phi(\mathbf{t})} d\mathcal{W}.$$

From spherical frame, we have

$$\begin{split} \tau_{\phi(\mathbf{t})} \times \nabla_{\sigma} \tau_{\phi(\mathbf{t})} &= \left(\frac{1}{2} (\varepsilon^{2} \rho) \omega^{\mathbf{t}} + \frac{1}{2} \omega^{\mathbf{t}} \rho\right) \frac{1}{2} \varepsilon \omega^{\mathbf{t}} \rho \alpha + \left(\rho \left(\frac{1}{2} \omega^{\mathbf{t}}\right)^{2} \frac{\partial}{\partial \sigma} (\varepsilon \rho) \right. \\ &\left. - \varepsilon \rho \left(\frac{1}{2} \omega^{\mathbf{t}}\right)^{2} \frac{\partial \rho}{\partial \sigma}\right) \mathbf{t} + \frac{1}{2} \rho \omega^{\mathbf{t}} \left(\frac{1}{2} (\varepsilon^{2} \rho) \omega^{\mathbf{t}} + \frac{1}{2} \rho \omega^{\mathbf{t}}\right) \mathbf{s} \end{split}$$

Electroosmotic antiferromagnetic microfluidic $\tau_{\phi(t)}$ magnetic torque density is

$$\begin{split} \mathcal{V}_{\tau_{\phi(t)}}^{\mathcal{A}} &= -\left(\rho \Big(\frac{1}{2}\omega^{\mathbf{t}}\Big)^2 \frac{\partial}{\partial \sigma}(\varepsilon\rho) - \varepsilon\rho \Big(\frac{1}{2}\omega^{\mathbf{t}}\Big)^2 \frac{\partial\rho}{\partial \sigma}\Big) \Big(\varepsilon \Big(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial \sigma}\Big) + \mathcal{L}\Big) \\ &- \Big(\frac{1}{2}(\varepsilon^2\rho)\omega^{\mathbf{t}} + \frac{1}{2}\rho\omega^{\mathbf{t}}\Big) \frac{1}{2}\varepsilon^2\rho\omega^{\mathbf{t}}\mathcal{F} \\ &+ \Big(\frac{\partial\varepsilon}{\partial t} - \mathcal{F}\Big) \frac{1}{2}\rho\omega^{\mathbf{t}} \Big(\frac{1}{2}(\varepsilon^2\rho)\omega^{\mathbf{t}} + \frac{1}{2}\omega^{\mathbf{t}}\rho\Big). \end{split}$$

Microfluidic antiferromagnetic $au_{\phi(\mathbf{t})}$ magnetic torque phase microscale is •

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$$\begin{split} \mathcal{T}^{\mathcal{A}}_{\tau_{\phi(t)}} &= \tau_{0}^{\mathsf{t}} \int \int_{\mathcal{W}} \left(\left(\frac{\partial \varepsilon}{\partial t} - \mathcal{F} \right) \frac{1}{2} \rho \omega^{\mathsf{t}} \left(\frac{1}{2} (\varepsilon^{2} \omega^{\mathsf{t}} \rho) + \frac{1}{2} \omega^{\mathsf{t}} \rho \right) \right. \\ & - \left(\rho \left(\frac{1}{2} \omega^{\mathsf{t}} \right)^{2} \frac{\partial}{\partial \sigma} (\varepsilon \rho) - \varepsilon \rho \left(\frac{1}{2} \omega^{\mathsf{t}} \right)^{2} \frac{\partial \rho}{\partial \sigma} \right) \left(\varepsilon \left(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \right) + \mathcal{L} \right) \\ & - \left(\frac{1}{2} (\varepsilon^{2} \rho) \omega^{\mathsf{t}} + \frac{1}{2} \rho \omega^{\mathsf{t}} \right) \frac{1}{2} \varepsilon^{2} \omega^{\mathsf{t}} \rho \mathcal{F} \right) d\mathcal{W}, \end{split}$$

where $\tau_0^{\mathbf{t}}$ is magnetic torque constant for $\phi(\mathbf{t})$. Also, we easily get

$$\phi(\mathbf{t}) \times \nabla_{\sigma} \phi(\mathbf{t}) = -\varepsilon (\varepsilon^2 + 1) \boldsymbol{\alpha} + \frac{\partial \varepsilon}{\partial \sigma} \mathbf{t} - (\varepsilon^2 + 1) \mathbf{s}.$$

Electroosmotic $\tau_{\phi(\mathbf{t})}$ magnetic torque density with antiferromagnetic $\phi(\mathbf{t})$ is

$$\begin{split} {}^{\phi(\mathbf{t})}\mathcal{V}^{\mathcal{A}}_{\tau_{\phi(\mathbf{t})}} &= \frac{\partial \varepsilon}{\partial \sigma} \left(\rho \Big(\frac{1}{2}\omega^{\mathbf{t}}\Big)^2 \frac{\partial}{\partial \sigma} (\varepsilon\rho) - \varepsilon\rho \Big(\frac{1}{2}\omega^{\mathbf{t}}\Big)^2 \frac{\partial \rho}{\partial \sigma} \Big) \\ &- \varepsilon (\varepsilon^2 + 1) \Big(\frac{1}{2} (\varepsilon^2 \rho) \omega^{\mathbf{t}} + \frac{1}{2} \rho \omega^{\mathbf{t}} \Big) \frac{1}{2} \varepsilon \omega^{\mathbf{t}} \rho \\ &- (\varepsilon^2 + 1) \frac{1}{2} \rho \omega^{\mathbf{t}} \Big(\frac{1}{2} \omega^{\mathbf{t}} (\varepsilon^2 \rho) + \frac{1}{2} \rho \omega^{\mathbf{t}} \Big). \end{split}$$

Antiferromagnetic $\tau_{\phi(\mathbf{t})}$ magnetic torque phase microscale with antiferromagnetic $\phi(\mathbf{t})$ is

$$\begin{split} {}^{\phi(\mathbf{t})} T^{\mathcal{A}}_{\tau_{\phi(\mathbf{t})}} &= \tau^{\mathbf{t}}_{0} \int \int_{\mathcal{W}} \left(- \left(\frac{1}{2} (\varepsilon^{2} \omega^{\mathbf{t}} \rho) + \frac{1}{2} \rho \omega^{\mathbf{t}} \right) \frac{1}{2} \varepsilon^{2} (\varepsilon^{2} + 1) \omega^{\mathbf{t}} \rho \right. \\ &+ \frac{\partial \varepsilon}{\partial \sigma} \left(\rho \left(\frac{1}{2} \omega^{\mathbf{t}} \right)^{2} \frac{\partial}{\partial \sigma} (\varepsilon \rho) - \varepsilon \rho \left(\frac{1}{2} \omega^{\mathbf{t}} \right)^{2} \frac{\partial \rho}{\partial \sigma} \right) \\ &- (\varepsilon^{2} + 1) \frac{1}{2} \rho \omega^{\mathbf{t}} \left(\frac{1}{2} (\varepsilon^{2} \rho) \omega^{\mathbf{t}} + \frac{1}{2} \rho \omega^{\mathbf{t}} \right) \right) d\mathcal{W}, \end{split}$$

where $\tau_0^{\mathbf{t}}$ is magnetic torque constant for $\phi(\mathbf{t})$. Optical effect of microfluidic antiferromagnetic $\tau_{\phi(\mathbf{t})}$ magnetic torque phase microscale is constructed vortex reactors in Fig. 2.

3.3 Microfluidic antiferromagnetic $au_{\phi(s)}$ magnetic torque phase

 $\tau_{\phi(s)}$ magnetic torque phase microscale is

$${}^{\mathcal{B}}\mathcal{M}_{\phi(\mathbf{s})} = \tau_0^{\mathbf{s}} \int \int_{\mathcal{W}} \nabla_t \phi(\mathbf{s}) \cdot \nabla_t \tau_{\phi(\mathbf{s})} d\mathcal{W}.$$

Optical angular momentum of $\phi(\mathbf{s})$ is given

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Fig. 2 Microfluidic antiferromagnetic $\tau_{\phi(t)}$

$$\otimes_{\phi(\mathbf{s})} = \rho m^{\mathbf{s}} \mathbf{s}.$$

The magnitude of optical angular momentum is

$$\left\|\bigotimes_{\phi(\mathbf{s})}\right\| = \rho m^{\mathbf{s}}.$$

Spherical magnetic moment of $\phi(\mathbf{s})$ is

$$\theta_{\phi(\mathbf{s})} = \frac{1}{2}\rho\omega^{\mathbf{s}}\mathbf{s}.$$

 $\tau_{\phi(\mathbf{s})}$ magnetic torque of $\phi(\mathbf{s})$ is constructed by

$$\tau_{\phi(\mathbf{s})} = \frac{1}{2}\rho^2 \omega^{\mathbf{s}} \boldsymbol{\alpha} + \frac{1}{2}\rho \varepsilon \omega^{\mathbf{s}} \mathbf{t}.$$

Thus, fluid of $\tau_{\phi(s)}$ magnetic torque is

$$\nabla_{t}\tau_{\phi(\mathbf{s})} = \left(\rho\omega^{\mathbf{s}}\frac{\partial\rho}{\partial t} - \mathcal{L}\frac{1}{2}(\rho\omega^{\mathbf{s}}\varepsilon)\right)\boldsymbol{\alpha} + \left(\frac{1}{2}\omega^{\mathbf{s}}\frac{\partial}{\partial t}(\rho\varepsilon) + \frac{1}{2}\rho^{2}\omega^{\mathbf{s}}\mathcal{L}\right)\mathbf{t} + \left(\mathcal{F}\frac{1}{2}\rho^{2}\omega^{\mathbf{s}} + \left(\frac{\partial\mathcal{F}}{\partial\sigma} + \varepsilon\mathcal{L}\right)\frac{1}{2}(\varepsilon\omega^{\mathbf{s}}\rho)\right)\mathbf{s}.$$

• Electroosmotic microfluidic $\tau_{\phi(s)}$ magnetic torque density is

$$\begin{split} \mathcal{V}_{\tau_{\phi(s)}} &= \left(\mathcal{L} - \frac{\partial \rho}{\partial t}\right) \left(\rho \omega^{s} \frac{\partial \rho}{\partial t} - \mathcal{L} \frac{1}{2} \varepsilon \omega^{s}(\rho)\right) \\ &- \left(\frac{1}{2} \frac{\partial}{\partial t} (\rho \varepsilon) \omega^{s} + \frac{1}{2} \rho^{2} \omega^{s} \mathcal{L}\right) \left(\frac{\partial \varepsilon}{\partial t} + \rho \mathcal{L}\right) - \left(\varepsilon \left(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma}\right) + \rho \mathcal{F}\right) \\ &\times \left(\mathcal{F} \frac{1}{2} \rho^{2} \omega^{s} + \left(\frac{\partial \mathcal{F}}{\partial \sigma} + \varepsilon \mathcal{L}\right) \frac{1}{2} (\varepsilon \omega^{s} \rho)\right). \end{split}$$

• Microfluidic $\tau_{\phi(s)}$ magnetic torque phase microscale is given

$$\begin{split} \mathcal{T}_{\tau_{\phi(s)}} &= \tau_0^{\mathbf{s}} \int \int_{\mathcal{W}} \left(- \Big(\frac{1}{2} \frac{\partial}{\partial t} (\rho \varepsilon) \omega^{\mathbf{s}} + \frac{1}{2} \rho^2 \omega^{\mathbf{s}} \mathcal{L} \Big) \Big(\frac{\partial \varepsilon}{\partial t} + \rho \mathcal{L} \Big) \\ &+ \Big(\mathcal{L} - \frac{\partial \rho}{\partial t} \Big) \Big(\rho \omega^{\mathbf{s}} \frac{\partial \rho}{\partial t} - \mathcal{L} \frac{1}{2} (\rho \varepsilon) \omega^{\mathbf{s}} \Big) - \Big(\varepsilon \Big(\mathcal{L} \varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \Big) + \rho \mathcal{F} \Big) \\ &\times \Big(\mathcal{F} \frac{1}{2} \rho^2 \omega^{\mathbf{s}} + \omega^{\mathbf{s}} \Big(\frac{\partial \mathcal{F}}{\partial \sigma} + \varepsilon \mathcal{L} \Big) \frac{1}{2} (\varepsilon \rho) \Big) \Big) d\mathcal{W}, \end{split}$$

where $\tau_0^{\mathbf{s}}$ is magnetic torque constant for $\phi(\mathbf{s})$.

• Microfluidic antiferromagnetic $\tau_{\phi(s)}$ magnetic torque phase microscale is given

$${}^{\mathcal{B}}\mathcal{M}_{\tau_{\phi(\mathbf{s})}} = \tau_0 \int \int_{\mathcal{W}} \nabla_t \phi(\mathbf{s}) \cdot \tau_{\phi(\mathbf{s})} \times \nabla_\sigma \tau_{\phi(\mathbf{s})} d\mathcal{W}.$$

From spherical frame, we have

$$\begin{split} \tau_{\phi(\mathbf{s})} \times \nabla_{\sigma} \tau_{\phi(\mathbf{s})} &= \left(\frac{1}{2}\omega^{\mathbf{s}}\right)^{2} (\rho^{2}\varepsilon^{3})\boldsymbol{\alpha} - \rho^{2} \left(\frac{1}{2}\omega^{\mathbf{s}}\right)^{2} (\rho\varepsilon^{2}) \mathbf{t} \\ &+ \left(\frac{1}{2}\omega^{\mathbf{s}}\rho^{2} \left(\frac{1}{2}\omega^{\mathbf{s}}\frac{\partial}{\partial\sigma}(\rho\varepsilon) + \frac{1}{2}\rho^{2}\omega^{\mathbf{s}}\right) \\ &- \frac{1}{2}\rho\omega^{\mathbf{s}}\varepsilon \left(\omega^{\mathbf{s}}\rho\frac{\partial\rho}{\partial\sigma} - \frac{1}{2}(\rho\omega^{\mathbf{s}}\varepsilon)\right) \right) \mathbf{s}. \end{split}$$

• Electroosmotic antiferromagnetic microfluidic $au_{\phi(s)}$ magnetic torque density is

$$\begin{split} \mathcal{V}^{\mathcal{A}}_{\tau_{\phi(s)}} &= \rho^2 \Big(\frac{1}{2}\omega^s\Big)^2 (\rho\varepsilon^2) \Big(\frac{\partial\varepsilon}{\partial t} + \rho\mathcal{L}\Big) + \Big(\frac{1}{2}\omega^s\Big)^2 (\rho^2\varepsilon^3) \bigg(\mathcal{L} - \frac{\partial\rho}{\partial t}\bigg) \\ &- \Big(\varepsilon\Big(\mathcal{L}\varepsilon + \frac{\partial\mathcal{F}}{\partial\sigma}\Big) + \rho\mathcal{F}\Big) \Big(\frac{1}{2}\rho^2\omega^s\Big(\frac{1}{2}\frac{\partial}{\partial\sigma}(\rho\varepsilon)\omega^s + \frac{1}{2}\rho^2\omega^s\Big) \\ &- \frac{1}{2}\omega^s\rho\varepsilon\bigg(\omega^s\rho\frac{\partial\rho}{\partial\sigma} - \frac{1}{2}(\rho\omega^s\varepsilon)\bigg)\bigg). \end{split}$$

• Microfluidic antiferromagnetic $\tau_{\phi(s)}$ magnetic torque phase microscale is

$$\begin{split} T^{\mathcal{A}}_{\tau_{\phi(s)}} &= \tau^{s}_{0} \int \int_{\mathcal{W}} \left(- \Big(\varepsilon \Big(\mathcal{L}\varepsilon + \frac{\partial \mathcal{F}}{\partial \sigma} \Big) + \rho \mathcal{F} \Big) \Big(\frac{1}{2} \omega^{s} \rho^{2} \Big(\frac{1}{2} \frac{\partial}{\partial \sigma} (\rho \varepsilon) \omega^{s} + \frac{1}{2} \rho^{2} \omega^{s} \Big) \right. \\ &\left. - \frac{1}{2} \rho \omega^{s} \varepsilon \Big(\omega^{s} \rho \frac{\partial \rho}{\partial \sigma} - \frac{1}{2} (\rho \varepsilon) \omega^{s} \Big) \Big) \right. \\ &\left. + \rho^{2} \Big(\frac{1}{2} \omega^{s} \Big)^{2} (\rho \varepsilon^{2}) \Big(\frac{\partial \varepsilon}{\partial t} + \rho \mathcal{L} \Big) + \Big(\frac{1}{2} \omega^{s} \Big)^{2} (\rho^{2} \varepsilon^{3}) \Big(\mathcal{L} - \frac{\partial \rho}{\partial t} \Big) \Big) d\mathcal{W}, \end{split}$$

where $\tau_0^{\mathbf{s}}$ is magnetic torque constant for $\phi(\mathbf{s})$.

Also, we easily get

$$\phi(\mathbf{s}) \times \nabla_{\sigma} \phi(\mathbf{s}) = \varepsilon^3 \boldsymbol{\alpha} + \left(\rho \left(\frac{\partial \varepsilon}{\partial \sigma} + \rho \right) - \rho \varepsilon^2 \right) \mathbf{t} + \left(\varepsilon - \frac{\partial \rho}{\partial \sigma} \right) \varepsilon \mathbf{s}.$$

 $\textit{Electroosmotic} au_{\phi(\mathbf{s})}$ magnetic torque density with antiferromagnetic $\phi(\mathbf{t})$ is

$$\begin{split} {}^{\phi(\mathbf{s})}\mathcal{V}^{\mathcal{A}}_{\tau_{\phi(\mathbf{s})}} &= \varepsilon^{6}\rho^{2} \Big(\frac{1}{2}\omega^{\mathbf{s}}\Big)^{2} - \rho^{2} \Big(\rho\Big(\frac{\partial\varepsilon}{\partial\sigma} + \rho\Big) - \rho\varepsilon^{2}\Big)\Big(\frac{\omega^{\mathbf{s}}}{2}\Big)^{2}(\rho\varepsilon^{2}) + \Big(\varepsilon - \frac{\partial\rho}{\partial\sigma}\Big)\varepsilon) \\ &\times \Big(\frac{\omega^{\mathbf{s}}}{2}\rho^{2}\Big(\frac{\omega^{\mathbf{s}}}{2}\frac{\partial}{\partial\sigma}(\rho\varepsilon) + \frac{\omega^{\mathbf{s}}}{2}\rho^{2}\Big) \\ &- \frac{\omega^{\mathbf{s}}}{2}\rho\varepsilon\Big(\omega^{\mathbf{s}}\rho\frac{\partial\rho}{\partial\sigma} - \frac{\omega^{\mathbf{s}}}{2}(\rho\varepsilon)\Big)\Big). \end{split}$$

Antiferromagnetic $\tau_{\phi(s)}$ magnetic torque phase microscale with antiferromagnetic $\phi(s)$ • is

$$\begin{split} {}^{\phi(\mathbf{s})}\mathcal{T}^{\mathcal{A}}_{\tau_{\phi(\mathbf{s})}} &= \tau^{\mathbf{s}}_{0} \int \int_{\mathcal{W}} \left(-\rho^{2} \Big(\rho \Big(\frac{\partial \varepsilon}{\partial \sigma} + \rho \Big) - \rho \varepsilon^{2} \Big) \Big(\frac{\omega^{\mathbf{s}}}{2} \Big)^{2} (\rho \varepsilon^{2}) \\ &+ \varepsilon^{6} \rho^{2} \Big(\frac{\omega^{\mathbf{s}}}{2} \Big)^{2} + \Big(\varepsilon - \frac{\partial \rho}{\partial \sigma} \Big) \varepsilon \Big) \Big(\frac{\omega^{\phi(\mathbf{s})}}{2} \rho^{2} \Big(\frac{\omega^{\mathbf{s}}}{2} \frac{\partial}{\partial \sigma} (\rho \varepsilon) + \frac{\omega^{\mathbf{s}}}{2} \rho^{2} \Big) \\ &- \frac{\omega^{\mathbf{s}}}{2} \rho \varepsilon \Big(\omega^{\phi(\mathbf{s})} \rho \frac{\partial \rho}{\partial \sigma} - \frac{\omega^{\mathbf{s}}}{2} (\rho \varepsilon) \Big) \Big) \Big) d\mathcal{W}, \end{split}$$

where τ_0^s is magnetic torque constant for $\phi(s)$. Optical effect of microfluidic antiferromagnetic $\tau_{\phi(s)}$ magnetic torque phase microscale is constructed vortex reactors in Fig. 3.



Fig. 3 Microfluidic antiferromagnetic $\tau_{\phi(s)}$

4 Conclusions

Quasi optical electromagnetic energy and flux are illustrated by flexible elastic curves, optical waves and biharmonic sonics. The results of optical modelling of biharmonic magnetic curves with optical applications are characterized (Körpınar 2020; Körpınar et al. 2020, 2021c, d; Körpınar and Körpınar 2021c, d; Ashkin et al. 1986; Ashkin 1970; Dholakia and Zemánek 2010; Schief and Rogers 2005; Dong et al. 2019; Seung 2015; Körpinar and Körpinar 2023d, b; Körpinar et al. 2023; Körpinar and Körpinar 2023d, a; Korpinar et al. 2023; Körpinar and Körpinar 2023c; Körpinar et al. 2023; Körpinar and Körpinar 2023c; Körpinar et al. 2023; Körpinar and Körpinar 2023; Körpinar et al. 2023; Körpinar and Körpinar 2023; Körpinar et al. 2023; Körpinar and Körpinar 2023; Körpinar and Körpinar 2023; Körpinar and Körpinar 2023; Körpi

In this paper, we illustrate optical $\tau_{\phi(\alpha)}, \tau_{\phi(t)}, \tau_{\phi(s)}$ magnetic torque phase microscale. Moreover, we have electroosmotic microfluidic $\tau_{\phi(\alpha)}, \tau_{\phi(t)}, \tau_{\phi(s)}$ magnetic torque density. Finally, we design electroosmotic magnetic torque density with antiferromagnetic model.

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Declarations

Conflict of interest The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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References

- Abouraddy, A.F., Bayindir, M., Benoit, G., Hart, S.D., Kuriki, K., Orf, N., Shapira, O., Sorin, F., Temelkuran, B., Fink, Y.: Towards multimaterial multifunctional fibres that see, hear, sense and communicate. Nat. Mater. 6, 336–347 (2007)
- Amjadi, M., Kyung, K.-U., Park, I., Sitti, M.: Stretchable, skin-mountable, and wearable strain sensors and their potential applications: a review. Adv. Funct. Mater. 26, 1678 (2016)
- Ashkin, A.: Acceleration and trapping of particles by radiation pressure. Phys. Rev. Lett. 24, 156–159 (1970)

Ashkin, A., Dziedzic, J.M., Bjorkholm, J.E., Chu, S.: Observation of a single-beam gradient force optical trap for dielectric particles. Opt. Lett. 11, 288–290 (1986)

Balakrishnan, R., Dandoloff, R.: The Schrodinger equation as a moving curve. Phys. Lett. A 260, 62 (1999)

Balakrishnan, R., Bishop, R., Dandoloff, R.: Geometric phase in the classical continuous antiferromagnetic Heisenberg spin chain. Phys. Rev. Lett. 64, 2107 (1990)

Balakrishnan, R., Bishop, R., Dandoloff, R.: Anholonomy of a moving space curve and applications to classical magnetic chains. Phys. Rev. B 47, 3108 (1993)

Berry, M.V., Klein, S.: Geometric phases from stacks of crystal plates. J. Mod. Opt. 43, 165-180 (1996)

Biener, G., Niv, A., Kleiner, V., Hasman, E.: Formation of helical beams by use of Pancharatnam–Berry phase optical elements. Opt. Lett. 27, 1875–1877 (2002)

Cao, Q., Liu, M., Wang, Z., Han, X., Li, L.: Dynamic motion analysis of magnetic particles in microfluidic systems under an external gradient magnetic field. Microfluid. Nanofluid. 21(2), 24 (2017)

- Cao, Q., Li, Z., Wang, Z., Han, X.: Rotational motion and lateral migration of an elliptical magnetic particle in a microchannel under a uniform magnetic field. Microfluid. Nanofluid. 22, 3 (2018)
- Dandoloff, R., Zakrzewski, W.J.: Parallel transport along a space curve and related phases. J. Phys. A: Math. Gen. 22(11), L461 (1989)
- Dholakia, K., Zemánek, P.: Colloquium: gripped by light: optical binding. Rev. Mod. Phys. 82, 1767–1791 (2010)
- Dong, C., Page, A.G., Yan, W., Nguyen-Dang, T., Sorin, F.: Microstructured multimaterial fibers for microfluidic sensing. Adv. Mater. Technol. 1900417 (2019)
- Erb, R.M., Martin, J.J., Soheilian, R., Pan, C., Barber, J.R.: Actuating soft matter with magnetic torque. Adv. Funct. Mater. 26(22), 3859–3880 (2016)
- Fink, Y., Winn, J., Fan, S., Chen, C., Michel, J., Joannopoulos, J., Thomas, E.: A dielectric omnidirectional reflector. Science 282, 1679 (1998)
- Furlani, E.P., Ng, K.C.: Analytical model of magnetic nanoparticle capture in the microvasculature. Phys. Rev. E 73(6), 061919 (2006)
- Garcia de Andrade, L.C.: Non-Riemannian geometry of twisted flux tubes. Braz. J. Phys. 36(5A), 1290– 1295 (2006a)
- Garcia de Andrade, L.C.: Vortex filaments in MHD. Physica Scripta, 73(5), 484 (2006b)
- Garcia de Andrade, L.C.: Riemannian geometry of twisted magnetic flux tubes in almost helical plasma flows. Phys. Plasmas **13**(2), 022309–022309 (2006c)
- Guo, B., Ding, S.: Landau–Lifshitz Equations. World Scientific, Singapore (2008)
- Gürbüz, N.: The differential formula of Hasimoto transformation in Minkowski 3-space. Int. J. Math. Math. Sci. 2005, 542381 (2005)
- Hasimoto, H.: A soliton on a vortex filament. J. Fluid Mech. 51(3), 477–485 (1972)
- Jones, R.C.: A new calculus for the treatment of optical systems I. Description and discussion of the calculus. J. Opt. Soc. Am. 31, 488–493 (1941)
- Körpinar, T.: Optical directional binormal magnetic flows with geometric phase: Heisenberg ferromagnetic model. Optik 219, 165134 (2020)
- Körpinar, T.: Optical electromotive force with Heisenberg spherical ferromagnetic spin. Optik 245, 167521 (2021)
- Körpınar, T., Demirkol, R.C.: Electromagnetic curves of the linearly polarized light wave along an optical fiber in a 3D semi-Riemannian manifold. J. Mod. Opt. 66(8), 857–867 (2019)
- Körpinar, T., Demirkol, R.C.: Electromagnetic curves of the linearly polarized light wave along an optical fiber in a 3D Riemannian manifold with Bishop equations. Optik 200, 163334 (2020)
- Körpınar, Z., Körpınar, T.: Optical hybrid electric and magnetic B₁-phase with Landau Lifshitz approach. Optik 247, 1–14 (2021a)
- Körpinar, Z., Körpinar, T.: Optical tangent hybrid electromotives for tangent hybrid magnetic particle. Optik 247, 1–11 (2021b)
- Körpınar, T., Körpınar, Z.: Timelike spherical magnetic flux flows with Heisenberg spherical ferromagnetic spin with some solutions. Optik (2021c). https://doi.org/10.1016/j.ijleo.2021.166745
- Körpinar, T., Körpinar, Z.: Spherical electric and magnetic phase with Heisenberg spherical ferromagnetic spin by some fractional solutions. Optik 242, 1–11 (2021d)
- Körpinar, Z., Korpinar, T.: Optical hybrid electric and magnetic B₁-phase with Landau Lifshitz approach. Optik 247, 167917 (2021)
- Körpınar, T., Demirkol, R.C., Körpınar, Z.: Soliton propagation of electromagnetic field vectors of polarized light ray traveling in a coiled optical fiber in the ordinary space. Int. J. Geom. Methods Mod. Phys. 16(8), 1950117 (2019)
- Körpinar, T.: Optical directional binormal magnetic flows with geometric phase: Heisenberg ferromagnetic model. Optik Int. J. Light Electron Opt. 219, 1–19 (2020)
- Körpinar, T., Demirkol, R.C.: Gravitational magnetic curves on 3D Riemannian manifolds. Int. J. Geom. Methods Mod. Phys. 15, 1–19 (2018a)
- Körpinar, T., Demirkol, R.C.: Frictional magnetic curves in 3D Riemannian manifolds. Int. J. Geom. Methods Mod. Phys. 15, 1–13 (2018b)
- Körpinar, T., Körpinar, Z.: Antiferromagnetic viscosity model for electromotive microscale with second type nonlinear heat frame. Int. J. Geom. Mod. Phys. 2350163 (2023a) (in press)
- Körpinar, T., Körpinar, Z.: New optical geometric recursional electromagnetic ferromagnetic microscale. Int. J. Mod. Phys. B 2450092 (2023b) (in press)
- Körpinar, Z., Körpinar, T.: New optical recursional spherical ferromagnetic flux for optical sonic microscale. J. Nonlinear Opt. Phys. Mater. 2350051 (2023c) (in press)
- Körpinar, T., Körpinar, Z.: Optical visco microfluidic optimistic hybrid optical electromotive microscale. International Journal of Modern Physics B, 2450159 (2023d) (in press)

- Körpinar, T., Körpinar, Z.: Antiferromagnetic complex electromotive microscale with first type Schrödinger frame. Opt. Quantum Electron. 55(6), 2–15 (2023e)
- Körpınar, T., Körpınar, Z.: Optical phase of recursional hybrid visco ferromagnetic electromagnetic microscale. Phys. Lett. A 462, 1–12 (2023f)
- Körpinar, T., Körpinar, Z.: Antiferromagnetic Schr ödinger electromotive microscale in Minkowski space. Opt. Quantum Electron. 55(8), 681 (2023)
- Körpinar, Z., Korpinar, T.: Optical spherical electromotive density with some fractional applications with Laplace transform in spherical Heisenberg space S_{44}^2 . Optik **245**, 167596 (2021)
- Korpinar, T., Körpinar, Z.: Timelike spherical magnetic S_N flux flows with Heisenberg spherical ferromagnetic spin with some solutions. Optik 242, 166745 (2021)
- Körpınar, T., Demirkol, R.C., Körpınar, Z.: Soliton propagation of electromagnetic field vectors of polarized light ray traveling in a coiled optical fiber in Minkowski space with Bishop equations. Eur. Phys. J. D 73, 1–11 (2019a)
- Körpinar, T., Demirkol, R.C., Körpinar, Z.: Soliton propagation of electromagnetic field vectors of polarized light ray traveling along with coiled optical fiber on the unit 2-sphere S². Rev. Mex. Fis. 65, 626–633 (2019b)
- Körpınar, T., Demirkol, R.C., Körpınar, Z., Asil, V.: Maxwellian evolution equations along the uniform optical fiber in Minkowski space. Rev. Mex. Fis. 66(4), 431 (2020)
- Körpınar, T., Demirkol, R.C., Körpınar, Z., Asil, V.: Maxwellian evolution equations along the uniform optical fiber in Minkowski space. Optik 217, 164561 (2020)
- Körpinar, T., Ünlütürk, Y., Körpinar, Z.: A novel approach to the motion equations of null Cartan curves via the compatible Hasimoto map. Optik 1–9 (2023)
- Körpınar, T., Körpınar, Z., Demirkol, R.C.: Binormal schrodinger system of wave propagation field of light radiate in the normal direction with q-HATM approach. Optik 235, 166444 (2020)
- Körpinar, T., Körpinar, Z., Yeneroğlu, M.: Optical energy of spherical velocity with optical magnetic density in Heisenberg sphere space S²_{Heis}. Optik 247, 1–7 (2021a)
- Körpinar, T., Sazak, A., Körpinar, Z.: Optical effects of some motion equations on quasi-frame with compatible Hasimoto map. Optik 247, 1–8 (2021b)
- Körpınar, T., Demirkol, R.C., Körpınar, Z.: Polarization of propagated light with optical solitons along the fiber in de-sitter space. Optik 226, 1–9 (2021c)
- Körpinar, T., Demirkol, R.C., Körpinar, Z.: Approximate solutions for the inextensible Heisenberg antiferromagnetic flow and solitonic magnetic flux surfaces in the normal direction in Minkowski space. Optik 238, 1–12 (2021d)
- Körpinar, T., Körpinar, Z., Asil, V.: New approach for optical electroostimistic phase with optical quasi potential energy. Optik 251, 168291 (2022)
- Körpinar, T., Demirkol, R.C., Körpinar, Z.: On the new conformable optical ferromagnetic and antiferromagnetic magnetically driven waves. Opt. Quantum Electron. 55(6), 496 (2023)
- Korpinar, Z., Inc, M., Korpinar, T.: Ferromagnetic recursion for geometric phase timelike SN-magnetic fibers. Opt. Quantum Electron. 55(4), 382 (2023)
- Körpinar, T., Körpinar, Z., Asil, V.: Optical electromotive microscale with first type Schrödinger frame. Optik 276, 170629 (2023)
- Kugler, M., Shtrikman, S.: Berry's phase, locally inertial frames, and classical analogues. Phys. Rev. D 37(4), 934 (1988)
- Lamb, G.L.: Solitons on moving space curves. J. Math. Phys. 18, 1654 (1977)
- Murugesh, S., Balakrishnan, R.: New connections between moving curves and soliton equations. Phys. Lett. A 290, 81 (2001)
- Satija, I.I., Balakrishan, R.: Geometric phases in twisted strips. Phys. Lett. A 373(39), 3582 (2009)
- Schief, W.K., Rogers, C.: The Da Rios system under a geometric constraint: the Gilbarg problem. J. Geom. Phys. 54(3), 286–300 (2005)
- Seung, W., et al.: Nanopatterned textile-based wearable triboelectric nanogenerator. ACS Nano 9, 3501– 3509 (2015)
- Smit, J.: The spontaneous Hall effect in ferromagnetics I. Physica 21, 877 (1955)
- Son, D.T., Yamamoto, N.: Berry curvature, triangle anomalies, and the chiral magnetic effect in Fermi liquids. Phys. Rev. Lett. 109, 81602 (2012)
- Tao, G., Stolyarov, A.M., Abouraddy, A.F.: Multi-material fibers. Int. J. Appl. Glass Sci. 3, 349 (2012)
- Tomita, A., Chiao, Y.: Observation of Berry's topological phase by use of an optical fiber. Phys. Rev. Lett. 57, 937 (1986)
- Vieira, V.R., Horley, P.P.: The Frenet–Serret representation of the Landau–Lifshitz–Gilbert equation. J. Phys. A: Math. Theor. 45(6), 065208 (2012)

- Wassmann, F., Ankiewicz, A.: Berry's phase analysis of polarization rotation in helicoidal fibers. Appl. Opt. 37, 3902 (1998)
- Yamada, T., Hayamizu, Y., Yamamoto, Y., Yomogida, Y., Izadi-Najafabadi, A., Futaba, D.N., Hata, K.: A stretchable carbon nanotube strain sensor for human-motion detection. Nat. Nanotechnol. 6, 296 (2011)
- Yamashita, O.: Effect of the geometrical phase shift on the spin and orbital angular momenta of light traveling in a coiled optical fiber with optical activity. Opt. Commun. **285**, 3740 (2012)
- Yamashita, O.: Geometrical phase shift of the extrinsic orbital angular momentum density of light propagating in a helically wound optical fiber. Opt. Commun. 285, 3061 (2012)
- Yan, W., Page, A., Nguyen-Dang, T., Qu, Y., Sordo, F., Wei, L., Sorin, F.: Advanced multimaterial electronic and optoelectronic fibers and textiles. Adv. Mater. 31, 1802348 (2018)
- Zygelman, B.: Appearance of gauge potentials in atomic collision physics. Phys. Lett. A **125**, 476–481 (1987)

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