

On some optical soliton structures to the Lakshmanan-Porsezian-Daniel model with a set of nonlinearities

Kalim U. Tariq¹ · Abdul-Majid Wazwaz² · Arslan Ahmed¹

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Abstract

In this work, the Lakshmanan-Porsezian-Daniel model is investigated which is the generalization of the non-linear Schrödinger model, to describes the dynamical behavior of optical solitons. The extended modified auxiliary equation mapping method is employed to develop some new exact solitary wave solutions to the complex model with the ker law, the parabolic law and the anti-cubic law nonlinearities. As a result, dark solitons, light solitons, singular solitons, solitary wave, periodic solitary wave, rational function, and elliptic function solutions are established. In the current era of communications network technology and nonlinear optics, the applied strategy appears to be a more powerful and efficient approach for achieving exact optical solutions to a number of diversified contemporary models.

Keywords Lakshmanan-Porsezian-Daniel model \cdot Ker law \cdot Parabolic law \cdot Anti-cubic law \cdot The extended modified auxiliary equation mapping method

1 Introduction

Nonlinear evolution equations (NLEES) are of key importance due to their significant role in diverse disciplines of science and technology. The nonlinear wave structures have fascinated many researchers in recent decades due to their diverse properties observed in various disciplines of contemporary sciences. In the presence of solitary waves, the nonlinear evolution models are utilized to simulate the effect of surface for deep water and weakly nonlinear dispersive long waves. Therefore, the exact solutions of such models play a vital role of study of dynamical structures and further properties of physical phenomenon occurring several fields to name a few, electromagnetism, physical chemistry, geophysics, ionised physics, elastic medium, fluid motion, fluid mechanics,

Kalim U. Tariq kalimulhaq@must.edu.pk

¹ Department of Mathematics, Mirpur University of Science and Technology (MUST), Mirpur, AJK 10250, Pakistan

² Department of Mathematics, Saint Xavier University, Chicago, IL 60655, USA

elastic medium, nuclear physics, electrochemistry, optical fibres, energy physics, chemical mechanics, gravity, biostatistics, statistical and natural physics (Chen et al. 2021; Ali and Gómez-Aguilar 2021; Yépez-Martínez et al. 2022; Islam et al. 2021a, b; Lü et al. 2021a, b; Yin et al. 2020; Li et al. 2010; Hosseini et al. 2018; Wazwaz 2014; Ivezić et al. 2019; Chen et al. 2021).

With the recent developments in various contemporary analytical methodologies, solitons plays a key role to understand the nonlinear phenomenon of many crucial structures in an exceptional way. The major feature of solitons is that they have nearly the same forms and speeds after colliding; also, the production of optical solitons is linked to optical frequency. Kink solutions are asymptotic waves that ascend or descend from one asymptotic state to the next, and they also approach a constant at infinity. Kink solutions, like classical particles, have a constant shape; nevertheless, their widths shrink, which can change. Solitons are transmitted as dark ones in the normal dispersion domain, but as bright ones in the anomalous dispersion domain. With the rapid advancement of information technology and telecommunications, the optical solitons play an important rule in understanding the dynamics of nonlinear wave propagation through a variety of wave-guides. The polarization of pulse propagation over trans-oceanic and trans-continental distances is an inherent problem with the dynamics of pulse propagation. This is due to a number of factors, including fiber diameter randomness, rough handling of optical fibers, and so on (Seadawy and Cheemaa 2019; Naher and Abdullah 2012; Wazwaz 2004; Ahmed et al. 2019; Kudryashov 2019; Sun et al. 2021; Zafar et al. 2022; Ahmad et al. 2021; Khater et al. 2021; Khodadad et al. 2021).

To investigate the behaviour of nonlinear models, researchers from all over the world discovered a variety of the numerical and analytical methods, such as the modified auxiliary equation method (Mahak and Akram 2020; Khater et al. 2019a, b), the optional decoupling condition approach (Lü and Chen 2021), the sub-equation method (Akinyemi et al. 2021), the extended trial function method (Ekici and Sonmezoglu 2019; Biswas et al. 2019; Nawaz et al. 2018), the modified kudryashov method (Kumar et al. 2018), the generalized riccati equation expansion method (Yong et al. 2003), the extended simple equation method (Lu et al. 2017; Zayed and Shohib 2019), the hirota bilinear method (Lü and Chen 2021; Jin-Ming and Yao-Ming 2011; El-Labany et al. 2018), the extended rational sinh-cosh method (Rezazadeh et al. 2019; Mahak and Akram 2019), the modified khater method (Khater et al. 2021) and so on.

The Lakshmanan-Porsezian-Daniel (LPD) model studies the optical solitons shine bright lights on the telecommunications industry (Inegbedion and Obadiaru 2019; Asima-kopoulos and Whalley 2017) to govern the dynamics of the pulse transmission (Li et al. 2018; Bandelow et al. 2020; Weckbrodt et al. 2018) through optical fibers (Buck 2004; Liu et al. 2019), the Photonic-crystal fiber (Wang et al. 2016; Bulbul et al. 2021) and the meta materials (Cui et al. 2010; Ma and Cui 2020) for the transcontinental (Wang et al. 2020) and transoceanic distance (Lindo 2020). Because of the important roles that models play in our day-to-day operations or activities, it is also crucial to discuss the characteristics of the models that occurs in ocean dynamics Olbers et al. (2012); Kamenkovich (2011).

A. Biswas et al. have applied two integration strategies namely the extended Jacobi's elliptic function approach and $exp(-\phi(\eta))$ -expansion method to construct dark and singular optical solitons solutions to the LPD model for Kerr nonlinearity (Biswas et al. 2018). In literature, a variety of interesting approaches have been implemented to analyse this model analytical such as the (Yépez-Martínez et al. 2022), the modified simple equation method (Biswas et al. 2018), the sine-Gordon equation method (Yildirim et al. 2021), the semi-inverse variational principle (Alzahrani and Belic 2021), the generalized projective Riccati

equations method (Akram et al. 2021), the modified extended direct algebraic method (Hubert et al. 2018) and several others.

In our work, the extended modified auxiliary equation mapping (AEM) method (Akram and Sarfraz 2021; Al-Munawarah and Arabia 2021) is employed to study the LPD model (Rezazadeh et al. 2018) reads

$$iq_{t} + \mu_{1}q_{xx} + \mu_{2}q_{xt} + \mu_{3}F(|q|^{2}q) = \alpha q^{*}(q_{x})^{2} + \sigma q_{xxxx} + \beta |q_{x}|^{2}q + \gamma |q|^{2}q_{xx} + \lambda q^{2}q_{xx}^{*} + \delta |q|^{4}q,$$
(1)

where q(x, t) represents the complex valued wave function. The first factor on the left side reflects the temporal evolution of the optical pulse, whereas the coefficients μ_1 and μ_2 represent the group-velocity dispersion and spatiotemporal dispersion respectively. Also σ is the fourth-order dispersion and δ is the two-photon absorption. The functional *F* is the source of nonlinearity which is a real-valued algebraic function.

2 The extended modified AEM method

Firstly consider the NPDE of the form

$$P(u, u_t, u_x, u_y, u_{xx}, ...) = 0, (2)$$

where *P* is the polynomial of u(x, y, t).

To attain the exact traveling wave solution, consider the traveling wave transformation of the form

$$u(x, y, t) = u(\xi), \quad \xi = \sum_{i=0}^{m} \kappa_i x_i,$$
 (3)

where κ_i , i = 0, 1, 2, ..., m are the constants. By using this transformation Eq. (2) reduce to nonlinear ODE of the form

$$Q(u, u', u'', u''', ...) = 0.$$
⁽⁴⁾

Where Q is the polynomial in $u(\xi)$ along with its derivatives w.r.t ξ .

Then we suppose the solution of Eq. (4) that will be expressed as

$$u(\xi) = \sum_{k=0}^{n} a_{k} \psi^{k}(\xi) + \sum_{k=-1}^{-n} b_{-k} \psi^{k}(\xi) + \sum_{k=2}^{n} c_{k} \psi^{k-2}(\xi) \psi'(\xi) + \sum_{k=1}^{n} d_{k} \left(\frac{\psi'(\xi)}{\psi(\xi)}\right)^{k}.$$
(5)

where a_k, b_k, c_k and d_k are the arbitrary constants to be determined.

Here $\psi(\xi)$ satisfies the following generalized solution,

$$\psi'(\xi)^{2} = \delta_{1}\psi(\xi)^{2} + \delta_{2}\psi(\xi)^{3} + \delta_{3}\psi(\xi)^{4},$$
(6)

where δ_1 , δ_2 and δ_3 are the constants to be determined. We have the following steps to attain the solution.

- 1. First we find the positive integer "*n*" by balancing the highest derivative and the highest nonlinear term of the Eq. (4).
- 2. Inserting Eq. (5) with its desired derivatives along with Eq. (6) into Eq. (4) and by collecting the same power terms $\psi'(\xi)\psi(\xi)$ and by equalizing them to zero, we have a system of algebraic equations. By solving the obtained system we attain a set of values of the constants a_k , b_k , c_k and d_k .
- 3. By substituting all the obtained values in Eq. (5) we find the solution of the Eq. (2).

3 Applications to the LPD model

Consider the transformation

$$q(x,t) = u(\xi)e^{i\phi(x,t)} \text{ and } q^*(x,t) = U(\xi)e^{-i\phi(x,t)},$$
(7)

where

$$\xi = x - vt \text{ and } \phi = \theta + \omega t - \kappa x. \tag{8}$$

In Eq. (8) ϕ is the phase component of the soliton, κ is the frequency of soliton, while ω is the wave number, θ is the phase constant and v is the velocity of the soliton.

Substituting Eq. (8) into Eq. (1) and then splitting into real and imaginary parts yields a pair of relations. The real part gives

$$-U'' (6\kappa^2 \sigma - \mu_2 \nu + \mu_1) - U (-\kappa^4 \sigma - \kappa^2 \mu_1 + \kappa \mu_2 \omega - \omega) -\mu_3 F(u^2) u - \kappa^2 u^3 (\alpha - \beta + \gamma + \lambda) + \sigma U''' + \delta U^5 + u^2 (\gamma + \lambda) u'' + u (\alpha + \beta) (u')^2 = 0,$$
(9)

while the imaginary part gives

$$4\kappa\sigma u''' + u'\left(\kappa\mu_2\nu - 4\kappa^3\sigma - 2\kappa\mu_1 + \mu_2\omega - \nu\right) + 2\kappa u^2 u'(\alpha + \gamma - \lambda) = 0.$$
(10)

Setting the coefficients of the linearly independent functions to zero in Eqs. (9) and (10), we get

$$\begin{aligned} \alpha + \beta &= 0, \\ \gamma + \lambda &= 0, \\ \sigma &= 0, \\ \alpha + \gamma - \lambda &= 0, \end{aligned}$$

and consequently the soliton speed falls out to be

$$\nu = \frac{\mu_2 \omega - 2\kappa \mu_1}{1 - \kappa \mu_2}, \quad bk \neq 1.$$

Hence the Eq. (9) reduces to

$$\mu_3 u F(u^2) - \delta u^5 - 4\gamma \kappa^2 u^3 + (\mu_1 - \mu_2 \nu) u'' + u (\kappa \mu_2 \omega - \kappa^2 \mu_1 - \omega) = 0.$$
(11)

3.1 Ker law

Assume F(u) = u. The Eq. (1) reduces to

$$iq_{t} + \mu_{1}q_{xx} + \mu_{2}q_{xt} + \mu_{3}|q|^{2}q = -2\gamma q^{*}q_{x}^{2} + 2\gamma |q_{x}|^{2}q + \gamma |q|^{2}q_{xx} - \gamma q^{2}q_{xx}^{*} + \delta |q|^{4}q,$$
(12)

and Eq. (11) reduces to

$$(\mu_1 - \mu_2 \nu)u'' + u(-\kappa^2 \mu_1 + \kappa \mu_2 \omega - \omega) + u^3(\mu_3 - 4\gamma \kappa^2) - \delta u^5 = 0,$$
(13)

by setting $u = V^{\frac{1}{2}}$, the Eq. (13) becomes

$$(\mu_1 - \mu_2 \nu) (2VV'' - (V')^2) + 4V^2 (-\kappa^2 \mu_1 + \kappa \mu_2 \omega - \omega) + 4V^3 (\mu_3 - 4\gamma \kappa^2) - 4\delta V^4 = 0.$$
 (14)

By homogenous balance principle we get the positive integer n = 1, so we suppose the solution of the form

$$V = a_0 + a_1 \psi(\xi) + \frac{b_1}{\psi(\xi)} + d_1 \left(\frac{\psi'(\xi)}{\psi(\xi)}\right),$$
(15)

where a_0 , a_1 , b_1 and d_1 are the constants to be determined. Substituting Eq. (15) along with its desired derivatives into Eq. (14) and by collecting the coefficients of $\psi(\xi)\psi'(\xi)$ we get an algebraic system of equations. By solving that system, the following cases arise

Family I

ω

$$\begin{split} \psi(\xi) &= -\frac{4\delta_1 e^{\sqrt{\delta_1\xi}}}{\delta_2^2 \left(-e^{2\sqrt{\delta_1\xi}}\right) + 2\delta_2 e^{\sqrt{\delta_1\xi}} + 4\delta_1 \delta_3 e^{2\sqrt{\delta_1\xi}} - 1} \\ Case \ I \ a_0 &= \frac{3(\mu_2^2 \mu_3 - 4\gamma)}{8\delta\mu_2^2}, \ a_1 = 0, \ b_1 = 0, \ d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \ \kappa = -\frac{1}{\mu_2}, \\ \omega &= \frac{48\gamma^2 - 24\gamma\mu_3\mu_2^2 - 16\delta\mu_1\mu_2^2 + 3\mu_3^2\mu_2^4}{32\delta\mu_2^4}, \ \delta_1 &= \frac{3(\mu_2^2\mu_3 - 4\gamma)^2}{16\delta\mu_2^4 (\mu_1 - \mu_2 \nu)}, \ \delta_2 = 0, \ \delta_3 = 0, \end{split}$$

substituting values in Eq. (15) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_1(x,t) = e^{i\phi} \sqrt{\frac{3(\mu_2^2 \mu_3 - 4\gamma)}{8\delta\mu_2^2}} + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1},$$
(16)

$$q_1^*(x,t) = e^{-i\phi} \sqrt{\frac{3(\mu_2^2 \mu_3 - 4\gamma)}{8\delta \mu_2^2}} + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}.$$
 (17)

Case II $a_0 = \frac{3(\mu_2^2 \mu_3 - 4\gamma)}{8\delta \mu_2^2}, \ a_1 = 0, \ b_1 = 0, \ d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \ \kappa = -\frac{1}{\mu_2},$ $\omega = \frac{48\gamma^2 - 24\gamma\mu_3\mu_2^2 - 16\delta\mu_1\mu_2^2 + 3\mu_3^2\mu_2^4}{32\delta\mu_4^4}, \ \delta_1 = \frac{3(\mu_2^2\mu_3 - 4\gamma)^2}{16\delta\mu_4^4(\mu_1 - \mu_2\gamma)}, \ \delta_2 = 0, \ \delta_3 = 0,$

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$$q_2(x,t) = e^{i\phi} \sqrt{\frac{3(\mu_2^2 \mu_3 - 4\gamma)}{8\delta\mu_2^2}} - \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1},$$
(18)

$$q_{2}^{*}(x,t) = e^{-i\phi} \sqrt{\frac{3(\mu_{2}^{2}\mu_{3} - 4\gamma)}{8\delta\mu_{2}^{2}}} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}.$$
 (19)

Case III $a_1 = 0, \ b_1 = 0, \ d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \ \kappa = \frac{\sqrt{3\mu_3 - 8a_0\delta}}{2\sqrt{3}\sqrt{\gamma}}, \ \delta_1 = \frac{4a_0^2\delta}{3(\mu_1 - \mu_2 \nu)},$ $c_1 = \frac{(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3)(\mu_2\sqrt{9\mu_3 - 24a_0\delta} + 6\sqrt{\gamma})}{2\sqrt{\delta}}, \ \delta_2 = 0, \ \delta_3 = 0,$

 $\omega = \frac{(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3)(\mu_2\sqrt{9\mu_3 - 24a_0\delta} + 6\sqrt{\gamma})}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta - 3\mu_3) + 12\gamma)}, \ \delta_2 = 0, \ \delta_3 = 0,$ substituting values in Eq. (15) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{3}(x,t) = e^{i\phi} \sqrt{a_{0} + \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}},$$
(20)

$$q_{3}^{*}(x,t) = e^{-i\phi} \sqrt{a_{0} + \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}}.$$
(21)

 $\begin{aligned} Case \ IV \ a_1 &= 0, \ b_1 = 0, \ d_1 = \frac{\kappa \left(\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}\right)}{2\sqrt{\delta}}, \ \kappa = -\frac{\sqrt{3\mu_3 - 8a_0\delta}}{2\sqrt{3}\sqrt{\gamma}}, \ \delta_2 = 0, \ \delta_3 = 0, \\ \omega &= \frac{\left(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3\right)\left(6\sqrt{\gamma} - \mu_2\sqrt{9\mu_3 - 24a_0\delta}\right)}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta - 3\mu_3) + 12\gamma)}, \ \delta_1 &= \frac{4a_0^2\delta}{1^3(\mu_1 - \mu_2 \nu)}, \\ \text{substituting values in Eq. (15) also using } u = V^{\frac{1}{2}} \text{ and by applying the reverse transformation} \end{aligned}$

substituting values in Eq. (15) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_4(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}},$$
(22)

$$q_{4}^{*}(x,t) = e^{-i\phi} \sqrt{a_{0} + \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}}.$$
(23)

Case $Va_1 = 0$, $b_1 = 0$, $d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}$, $\kappa = \frac{\sqrt{3\mu_3 - 8a_0\delta}}{2\sqrt{3}\sqrt{\gamma}}$, $\delta_2 = 0$, $\delta_3 = 0$,

$$\omega = \frac{(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3)(\mu_2\sqrt{9}\mu_3 - 24a_0\delta + 6\sqrt{\gamma})}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta - 3\mu_3) + 12\gamma)}, \ \delta_1 = \frac{4a_0^2\delta}{3(\mu_1 - \mu_2\nu)},$$

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$$q_{5}(x,t) = e^{i\phi} \sqrt{a_{0} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}},$$
(24)

$$q_{5}^{*}(x,t) = e^{-i\phi} \sqrt{a_{0} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}\xi}}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}}.$$
(25)

 $\begin{array}{l} Case \ VI \ a_1 = 0, \ b_1 = 0, \ d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \ \kappa = -\frac{\sqrt{3}\mu_3 - 8a_0\delta}{2\sqrt{3}\sqrt{\gamma}}, \ \delta_2 = 0, \ \delta_3 = 0, \\ \omega = \frac{\left(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3\right)\left(6\sqrt{\gamma} - \mu_2\sqrt{9\mu_3 - 24a_0\delta}\right)}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta - 3\mu_3) + 12\gamma)}, \ \delta_1 = \frac{4a_0^2\delta}{3(\mu_1 - \mu_2 \nu)}, \\ \text{substituting values in Eq. (15) also using } u = V^{\frac{1}{2}} \text{ and by applying the reverse trans-} \end{array}$

formation (7), we get our solutions

$$q_{6}(x,t) = e^{i\phi} \sqrt{a_{0} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}},$$
(26)

$$q_{6}^{*}(x,t) = e^{-i\phi} \sqrt{a_{0} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}}.$$
(27)

Family II

$$\psi(\xi) = \frac{4\delta_1 e^{\sqrt{\delta_1}\xi}}{-2\delta_2 e^{\sqrt{\delta_1}\xi} + e^{2\sqrt{\delta_1}\xi} + \delta_2^2 - 4\delta_1\delta_3}.$$
Case $Ia_0 = \frac{3(\mu_2^2\mu_3 - 4\gamma)}{8\delta\mu_2^2}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2\nu}}{2\sqrt{\delta}}, \kappa = -\frac{1}{\mu_2},$

$$\omega = \frac{48\gamma^2 - 24\gamma\mu_3\mu_2^2 - 16\delta\mu_1\mu_2^2 + 3\mu_3^2\mu_2^4}{32\delta\mu_2^4}, \delta_1 = \frac{3(\mu_2^2\mu_3 - 4\gamma)^2}{16\lambda_2^4(\mu_1 - \mu_2\nu)}, \delta_2 = 0, \delta_3 = 0,$$
substituting values in Eq. (15) also using $\mu = V_1^2$ and by applying the

$$q_{7}(x,t) = e^{i\phi} \sqrt{\frac{3(\mu_{2}^{2}\mu_{3} - 4\gamma)}{8\delta\mu_{2}^{2}}} + \frac{\sqrt{3}e^{\sqrt{\delta_{1}}\xi}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}},$$
(28)

$$q_{7}^{*}(x,t) = e^{-i\phi} \sqrt{\frac{3(\mu_{2}^{2}\mu_{3} - 4\gamma)}{8\delta\mu_{2}^{2}}} + \frac{\sqrt{3}e^{\sqrt{\delta_{1}}\xi}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}.$$
 (29)

$$\begin{aligned} Case \ II \ a_0 &= \frac{3(\mu_2^2 \mu_3 - 4\gamma)}{8\delta\mu_2^2}, \ a_1 = 0, \ b_1 = 0, \ d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \ \kappa = -\frac{1}{\mu_2}, \\ \omega &= \frac{48\gamma^2 - 24\gamma\mu_3\mu_2^2 - 16\delta\mu_1\mu_2^2 + 3\mu_3^2\mu_2^4}{32\delta\mu_2^4}, \ \delta_1 = \frac{3(\mu_2^2 \mu_3 - 4\gamma)^2}{16\delta\mu_2^4(\mu_1 - \mu_2 \nu)}, \ \delta_2 = 0, \ \delta_3 = 0, \end{aligned}$$

$$q_8(x,t) = e^{i\phi} \sqrt{\frac{3(\mu_2^2 \mu_3 - 4\gamma)}{8\delta\mu_2^2}} - \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1},$$
(30)

$$q_8^*(x,t) = e^{-i\phi} \sqrt{\frac{3(\mu_2^2\mu_3 - 4\gamma)}{8\delta\mu_2^2} - \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}}.$$
 (31)

Case III $a_1 = 0, \ b_1 = 0, \ d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \ \kappa = \frac{\sqrt{3\mu_3 - 8a_0\delta}}{2\sqrt{3}\sqrt{\gamma}}, \ \delta_1 = \frac{4a_0^2\delta}{3(\mu_1 - \mu_2 \nu)},$

$$\omega = \frac{(8a_0\delta(2a_0\gamma+\mu_1)-3\mu_1\mu_3)(\mu_2\sqrt{9}\mu_3-24a_0\delta+6\sqrt{\gamma})}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta-3\mu_3)+12\gamma)}, \ \delta_2 = 0, \ \delta_3 = 0,$$

substituting values in Eq. (15) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_9(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}},$$
(32)

$$q_{9}^{*}(x,t) = e^{-i\phi} \sqrt{a_{0} + \frac{\sqrt{3}e^{\sqrt{\delta_{1}}\xi}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1} - \mu_{2}\nu}}{2\sqrt{\delta}\delta_{1}}}.$$
(33)

$$\begin{aligned} Case \ IV \ a_1 &= 0, \ b_1 = 0, \ d_1 = \frac{\kappa \left(\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}\right)}{2\sqrt{\delta}}, \ \kappa = -\frac{\sqrt{3\mu_3 - 8a_0\delta}}{2\sqrt{3}\sqrt{\gamma}}, \ \delta_2 = 0, \ \delta_3 = 0, \\ \omega &= \frac{\left(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3\right)\left(6\sqrt{\gamma} - \mu_2\sqrt{9\mu_3 - 24a_0\delta}\right)}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta - 3\mu_3) + 12\gamma)}, \ \delta_1 &= \frac{4a_0^2\delta}{3(\mu_1 - \mu_2 \nu)}, \end{aligned}$$

substituting values in Eq. (15) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{10}(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}},$$
(34)

$$q_{10}^{*}(x,t) = e^{-i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}}.$$
(35)

 $Case \ Va_1 = 0, \ b_1 = 0, \ d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \ \kappa = \frac{\sqrt{3\mu_3 - 8a_0\delta}}{2\sqrt{3}\sqrt{\gamma}}, \ \delta_2 = 0, \ \delta_3 = 0,$

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$$\omega = \frac{(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3)(\mu_2\sqrt{9\mu_3 - 24a_0\delta} + 6\sqrt{\gamma})}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta - 3\mu_3) + 12\gamma)}, \ \delta_1 = \frac{4a_0^2\delta}{3(\mu_1 - \mu_2\nu)},$$

$$q_{11}(x,t) = e^{i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}},$$
(36)

$$q_{11}^*(x,t) = e^{-i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}}.$$
(37)

$$\begin{split} & Case \; VI \; a_1 = 0, \; b_1 = 0, \; d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}}, \; \kappa = -\frac{\sqrt{3\mu_3 - 8a_0\delta}}{2\sqrt{3}\sqrt{\gamma}}, \; \delta_2 = 0, \; \delta_3 = 0, \\ & \omega = \frac{\left(8a_0\delta(2a_0\gamma + \mu_1) - 3\mu_1\mu_3\right)\left(6\sqrt{\gamma} - \mu_2\sqrt{9\mu_3 - 24a_0\delta}\right)}{6\sqrt{\gamma}(\mu_2^2(8a_0\delta - 3\mu_3) + 12\gamma)}, \; \delta_1 = \frac{4a_0^2\delta}{3(\mu_1 - \mu_2 \nu)}, \end{split}$$

substituting values in Eq. (15) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{12}(x,t) = e^{i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}},$$
(38)

$$q_{12}^{*}(x,t) = e^{-i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta}\delta_1}}.$$
(39)

3.2 Parabolic law

Assume $F(u) = c_1 u + c_2 u^2$. The Eq. (1) reduces to

$$iq_t + \mu_1 q_{xx} + \mu_2 q_{xt} + (c_1 |q|^2 + c_2 |q|^4) * q =$$
(40)

$$-2\gamma q^{*}q_{x}^{2}+2\gamma |q_{x}|^{2}q+\gamma |q|^{2}q_{xx}-\gamma q^{2}q_{xx}^{*}+\delta |q|^{4}q,$$
(41)

and Eq. (11) reduces to

$$(\mu_1 - \mu_2 \nu) u'' + u(\kappa \mu_2 \omega - \kappa^2 \mu_1 - \omega) + u^3 (c_1 - 4\gamma \kappa^2) + u^5 (c_2 - \delta) = 0.$$
 (42)

By setting $u = V^{\frac{1}{2}}$, the Eq. (42) becomes

$$(\mu_1 - \mu_2 \nu) \left(2VV'' - (V')^2 \right) + 4V^2 \left(-\kappa^2 \mu_1 + \kappa \mu_2 \omega - \omega \right) + 4V^3 (c_1 - 4\gamma \kappa^2) + 4V^4 (c_2 - \delta) = 0.$$
(43)

By homogenous balance principle we get the positive integer n = 1, so we suppose the solution of the form

$$V = a_0 + a_1 \psi(\xi) + \frac{b_1}{\psi(\xi)} + d_1 \left(\frac{\psi'(\xi)}{\psi(\xi)}\right),$$
(44)

where a_0 , a_1 , b_1 and d_1 are the constants to be determined. Substituting Eq. (44) along with its desired derivatives into Eq. (43) and by collecting the coefficients of $\psi(\xi)\psi'(\xi)$ we get an algebraic system of equations. By solving that system, the following cases arise

Family I

$$\begin{split} \psi(\xi) &= -\frac{4\delta_1 e^{\sqrt{\delta_1}\xi}}{\delta_2^2 \left(-e^{2\sqrt{\delta_1}\xi}\right) + 2\delta_2 e^{\sqrt{\delta_1}\xi} + 4\delta_1\delta_3 e^{2\sqrt{\delta_1}\xi} - 1}.\\ Case \,I\,a_1 &= 0, b_1 = 0, d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_2}}, \delta_1 = \frac{4a_0^2 (\delta - c_2)}{3(\mu_1 - \mu_2 \nu)}, \delta_2 = 0, \delta_3 = 0,\\ \omega &= \frac{-16a_0^2 \gamma c_2 + 16a_0^2 \gamma \delta - 8a_0 c_2 \mu_1 + 8a_0 \delta \mu_1 - 3c_1 \mu_1}{2\sqrt{3}\sqrt{\gamma} \mu_2 \sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1} + 12\gamma}, \kappa = -\frac{\sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1}}{2\sqrt{3}\sqrt{\gamma}}, \end{split}$$

substituting values in Eq. (44) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{13}(x,t) = e^{i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}},$$
(45)

$$q_{13}^*(x,t) = e^{-i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 v}}{2\delta_1\sqrt{\delta - c_2}}}.$$
(46)

Case

$$II \qquad a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_2}}, \delta_1 = \frac{4a_0^2(\delta - c_2)}{3(\mu_1 - \mu_2 \nu)}, \delta_2 = 0, \delta_3 = 0, \delta_3 = 0, \delta_4 = \frac{\sqrt{3}}{2\sqrt{\delta - c_2}}, \delta_4 = \frac{\sqrt{3}}{2\sqrt{\delta - c_2}}, \delta_5 = 0, \delta_5 = 0,$$

$$\omega = \frac{-16a_0^2\gamma c_2 + 16a_0^2\gamma \delta - 8a_0c_2\mu_1 + 8a_0\delta\mu_1 - 3c_1\mu_1}{2\sqrt{3}\sqrt{\gamma}\mu_2\sqrt{8a_0c_2 - 8a_0\delta + 3c_1} + 12\gamma}, \quad \kappa = -\frac{\sqrt{8a_0c_2 - 8a_0\delta + 3c_1}}{2\sqrt{3}\sqrt{\gamma}}, \quad \text{substituting values in}$$

Eq. (44) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{14}(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}},$$
(47)

$$q_{14}^*(x,t) = e^{-i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}}.$$
(48)

$$\begin{aligned} Case \, III \, a_1 &= 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_2}}, \delta_1 = \frac{4a_0^2(\delta - c_2)}{3(\mu_1 - \mu_2 \nu)}, \delta_2 = 0, \delta_3 = 0, \\ \omega &= \frac{16a_0^2 \gamma c_2 - 16a_0^2 \gamma \delta + 8a_0 c_2 \mu_1 - 8a_0 \delta \mu_1 + 3c_1 \mu_1}{2\left(\sqrt{3}\sqrt{\gamma} \mu_2 \sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1} - 6\gamma\right)}, \kappa = \frac{\sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1}}{2\sqrt{3}\sqrt{\gamma}}, \end{aligned}$$

$$q_{15}(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}},$$
(49)

$$q_{15}^{*}(x,t) = e^{-i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}}.$$
(50)

Family II

$$\begin{split} \psi(\xi) &= \frac{4\delta_1 e^{\sqrt{\delta_1 \xi}}}{-2\delta_2 e^{\sqrt{\delta_1 \xi}} + e^{2\sqrt{\delta_1 \xi}} + \delta_2^2 - 4\delta_1 \delta_3}}.\\ Case \ I \ a_1 &= 0, \ b_1 = 0, \ d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_2}}, \ \delta_1 &= \frac{4a_0^2(\delta - c_2)}{3(\mu_1 - \mu_2 \nu)}, \ \delta_2 = 0, \ \delta_3 = 0,\\ \omega &= \frac{-16a_0^2 \gamma c_2 + 16a_0^2 \gamma \delta - 8a_0 c_2 \mu_1 + 8a_0 \delta \mu_1 - 3c_1 \mu_1}{2\sqrt{3}\sqrt{\gamma} \mu_2 \sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1} + 12\gamma}, \ \kappa &= -\frac{\sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1}}{2\sqrt{3}\sqrt{\gamma}}, \end{split}$$

substituting values in Eq. (44) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{16}(x,t) = e^{i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}},$$
(51)

$$q_{16}^{*}(x,t) = e^{-i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1 \sqrt{\delta - c_2}}}.$$
(52)

$$\begin{aligned} Case \ II \ a_1 &= 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_2}}, \delta_1 = \frac{4a_0^2(\delta - c_2)}{3(\mu_1 - \mu_2 \nu)}, \delta_2 = 0, \delta_3 = 0, \\ \omega &= \frac{-16a_0^2 \gamma c_2 + 16a_0^2 \gamma \delta - 8a_0 c_2 \mu_1 + 8a_0 \delta \mu_1 - 3c_1 \mu_1}{2\sqrt{3}\sqrt{\gamma} \mu_2 \sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1} + 12\gamma}, \\ \kappa &= -\frac{\sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1}}{2\sqrt{3}\sqrt{\gamma}}, \end{aligned}$$

$$q_{17}(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}},$$
(53)

$$q_{17}^*(x,t) = e^{-i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}}.$$
(54)

$$\begin{aligned} Case \, III \, a_1 &= 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_2}}, \delta_1 = \frac{4a_0^2(\delta - c_2)}{3(\mu_1 - \mu_2 \nu)}, \delta_2 = 0, \delta_3 = 0, \\ \omega &= \frac{16a_0^2 \gamma c_2 - 16a_0^2 \gamma \delta + 8a_0 c_2 \mu_1 - 8a_0 \delta \mu_1 + 3c_1 \mu_1}{2\left(\sqrt{3}\sqrt{\gamma} \mu_2 \sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1} - 6\gamma\right)}, \\ \kappa &= \frac{\sqrt{8a_0 c_2 - 8a_0 \delta + 3c_1}}{2\sqrt{3}\sqrt{\gamma}}, \end{aligned}$$

$$q_{18}(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_2}}},$$
(55)

$$q_{18}^*(x,t) = e^{-i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1 \sqrt{\delta - c_2}}}.$$
(56)

3.3 Anti-cubic law

Assume $F(u) = \frac{c_1}{u^2} + c_2 u + c_3 u^2$, the Eq. (1) reduces to $iq_t + \mu_1 q_{xx} + \mu_2 q_{xt} + (c_1 |q|^{-4} + c_2 |q|^2 + c_3 |q|^4) * q$ $= -2\gamma q^* q_x^2 + 2\gamma |q_x|^2 q + \gamma |q|^2 q_{xx} - \gamma q^2 q_{xx}^* + \delta |q|^4 q,$ (57)

and Eq. (11) reduces to

$$(\mu_1 - \mu_2 \nu) u'' + u (-\kappa^2 \mu_1 + \kappa \mu_2 \omega - \omega) + u^3 (c_2 - 4\gamma \kappa^2) + c_1 u^{-3} + u^5 (c_3 - \delta) = 0.$$
 (58)

By setting $u = V^{\frac{1}{2}}$, the Eq. (58) becomes

$$(\mu_1 - \mu_2 \nu) \left(2VV'' - (V')^2 \right) + 4V^2 \left(-\kappa^2 \mu_1 + \kappa \mu_2 \omega - \omega \right) + 4V^3 (c_2 - 4\gamma \kappa^2) + 4c_1 + 4V^4 (c_3 - \delta) = 0.$$
(59)

By homogenous balance principle we get the positive integer n = 1, so we suppose the solution of the form

$$V = a_0 + a_1 \psi(\xi) + \frac{b_1}{\psi(\xi)} + d_1 \left(\frac{\psi'(\xi)}{\psi(\xi)}\right),$$
(60)

where a_0 , a_1 , b_1 and d_1 are the constants to be determined. Substituting Eq. (60) along with its desired derivatives into Eq. (59) and by collecting the coefficients of $\psi(\xi)\psi'(\xi)$ we get an algebraic system of equations. By solving that system, the following cases arise

Family I

$$\psi(\xi) = -\frac{4\delta_1 e^{\sqrt{\delta_1}\xi}}{\delta_2^2 \left(-e^{2\sqrt{\delta_1}\xi}\right) + 2\delta_2 e^{\sqrt{\delta_1}\xi} + 4\delta_1 \delta_3 e^{2\sqrt{\delta_1}\xi} - 1}$$

$$\begin{aligned} Case \ I \ a_1 &= 0, b_1 = 0, d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \gamma = \frac{8a_0(c_3 - \delta) + 3c_2}{12\kappa^2}, \\ \delta_1 &= \frac{4\left(a_0^2(\delta - c_3)(\mu_1 - \mu_2 \nu) + \sqrt{3}\sqrt{c_1(\delta - c_3)(\mu_1 - \mu_2 \nu)^2}\right)}{3(\mu_1 - \mu_2 \nu)^2}, \delta_2 = 0, \delta_3 = 0, \\ \omega &= \frac{(\mu_1 - \mu_2 \nu)(4a_0^2(\delta - c_3) - 3\kappa^2 \mu_1) - 2\sqrt{3}\sqrt{c_1(\delta - c_3)(\mu_1 - \mu_2 \nu)^2}}{3(\kappa \mu_2 - 1)(\mu_2 \nu - \mu_1)}, \end{aligned}$$

substituting values in Eq. (60) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{19}(x,t) = e^{i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}},$$
(61)

$$q_{19}^*(x,t) = e^{-i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}}.$$
(62)

$$\begin{aligned} \text{Case II } a_1 &= 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \gamma = \frac{8a_0(c_3 - \delta) + 3c_2}{12\kappa^2}, \\ \delta_1 &= \frac{4\left(a_0^2(\delta - c_3)(\mu_1 - \mu_2 \nu) + \sqrt{3}\sqrt{c_1(\delta - c_3)(\mu_1 - \mu_2 \nu)^2}\right)}{3(\mu_1 - \mu_2 \nu)^2}, \delta_2 &= 0, \delta_3 = 0, \\ \omega &= \frac{(\mu_1 - \mu_2 \nu)(4a_0^2(\delta - c_3) - 3\kappa^2 \mu_1) - 2\sqrt{3}\sqrt{c_1(\delta - c_3)(\mu_1 - \mu_2 \nu)^2}}{3(\kappa \mu_2 - 1)(\mu_2 \nu - \mu_1)}, \end{aligned}$$

$$q_{20}(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}},$$
(63)

$$q_{20}^{*}(x,t) = e^{-i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}}.$$
(64)

$$\begin{aligned} \text{Case III } a_0 &= -\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \left(4\left(\delta - c_3\right) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \right), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3) (\mu_2 \nu - \mu_1)}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{21}(x,t) = e^{i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}} -\frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) - 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(65)

$$q_{21}^{*}(x,t) = e^{-i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}} -\frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})-2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} \right)^{\frac{1}{2}}.$$

$$Wa_{0} = -\frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})-2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}}}, a_{1} = 0, b_{1} = 0, d_{2} = \frac{\sqrt{3}\sqrt{\mu_{1}-\mu_{2}\nu}}{\sqrt{2}}.$$
(66)

$$\begin{aligned} \text{Case IV } a_0 &= -\frac{1}{2} \sqrt{\frac{\gamma \mu_2(-3)}{\mu_2^4(\delta-c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{\gamma} \sqrt{\mu_1 - \mu_2}}{2\sqrt{\delta-c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \Biggl(4\Bigl(\delta - c_3) \sqrt{\frac{3\mu_1 \mu_2^2(\delta - c_3) - 2\sqrt{3}\sqrt{c_1 \mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta-c_3)^2}} + 3c_2 \Biggr), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2(c_3 - \delta) + 2\sqrt{3}\sqrt{c_1 \mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta-c_3)(\mu_2 \nu - \mu_1)}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{22}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 v}}{2\delta_1\sqrt{\delta - c_3}} - \frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) - 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(67)

$$q_{22}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}(-\xi)}}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}\xi}}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}} \right)^{\frac{1}{2}}$$
(68)

$$-\frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta-c_3)-2\sqrt{3}\sqrt{c_1\mu_2^8(\delta-c_3)^3}}{\mu_2^4(\delta-c_3)^2}}\right)^2.$$

$$\begin{aligned} Case \ V \ a_0 &= \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \Biggl(4 \Bigl(c_3 - \delta \Bigr) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \Biggr), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3) (\mu_2 \nu - \mu_1)}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{23}(x,t) = e^{i\phi} \times \left(\frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta-c_3) - 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta-c_3)^3}}{\mu_2^4(\delta-c_3)^2}} - \frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1-\mu_2\nu}}{2\delta_1\sqrt{\delta-c_3}}\right)^{\frac{1}{2}},$$
(69)

$$q_{23}^{*}(x,t) = e^{-i\phi} \times \left(\frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3}) - 2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}}\right)^{\frac{1}{2}}.$$
(70)

$$\begin{aligned} \text{Case VI } a_0 &= \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}} \\ \gamma &= \frac{1}{12} \mu_2^2 \left(4 \left(c_3 - \delta \right) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \right), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^3}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{24}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}} + \frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) - 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}}\right)^{\frac{1}{2}},$$
(71)

$$q_{24}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}v}}{2\delta_{1}\sqrt{\delta-c_{3}}} + \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})-2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(72)

$$\begin{aligned} \text{Case VII } a_0 &= -\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \Biggl(4\Bigl(\delta - c_3) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \Biggr), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^3}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{25}(x,t) = e^{i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}} - \frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) + 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(73)

$$q_{25}^{*}(x,t) = e^{-i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}} - \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})+2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(74)

Case VIII
$$a_0 = -\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}},$$

 $\gamma = \frac{1}{12} \mu_2^2 \left(4 \left(\delta - c_3 \right) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \right), \kappa = \frac{1}{\mu_2},$
 $\delta_{\nu} = \frac{\mu_1 \mu_2^2 (c_3 - \delta) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}, \delta_{\nu} = 0, \delta_{\nu} = 0.$

 $\delta_1 = \frac{\mu_{1/2}(s_3(s_2) - V + V + \mu_2(s_3(s_3))}{\mu_2^4(\delta - c_3)(\mu_2 v - \mu_1)}, \delta_2 = 0, \delta_3 = 0,$ substituting values in Eq. (60) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{26}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}} - \frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) + 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(75)

$$q_{26}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)} \sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}} \sqrt{\mu_{1} - \mu_{2}\nu}}{2\delta_{1}\sqrt{\delta - c_{3}}} - \frac{1}{2} \sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta - c_{3}) + 2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta - c_{3})^{3}}}{\mu_{2}^{4}(\delta - c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(76)

$$\begin{aligned} \text{Case IX } a_0 &= \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \left(4 \left(c_3 - \delta \right) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \right), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^4}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{27}(x,t) = e^{i\phi} \times \left(\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} - \frac{\sqrt{3} e^{\sqrt{\delta_1} (-\xi)} \sqrt{\delta_1^3 e^{2\sqrt{\delta_1} \xi}} \sqrt{\mu_1 - \mu_2 v}}{2\delta_1 \sqrt{\delta - c_3}} \right)^{\frac{1}{2}},$$
(77)

$$q_{27}^{*}(x,t) = e^{-i\phi} \times \left(\frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})+2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}}\right)^{\frac{1}{2}}.$$
(78)

$$\begin{aligned} Case & X \qquad a_0 = \frac{1}{2} \sqrt{\frac{3\mu_1\mu_2^2(\delta-c_3)+2\sqrt{3}\sqrt{c_1\mu_2^8(\delta-c_3)^3}}{\mu_2^4(\delta-c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1-\mu_2\nu}}{2\sqrt{\delta-c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \bigg(4 \big(c_3 - \delta \big) \sqrt{\frac{3\mu_1\mu_2^2(\delta-c_3)+2\sqrt{3}\sqrt{c_1\mu_2^8(\delta-c_3)^3}}{\mu_2^4(\delta-c_3)^2}} + 3c_2 \bigg), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1\mu_2^2(c_3-\delta)-2\sqrt{3}\sqrt{c_1\mu_2^8(\delta-c_3)^3}}{\mu_2^4(\delta-c_3)(\mu_2\nu-\mu_1)}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{28}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1}(-\xi)}\sqrt{\delta_1^3 e^{2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}} + \frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) + 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(79)

$$q_{28}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}}(-\xi)}\sqrt{\delta_{1}^{3}e^{2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}v}}{2\delta_{1}\sqrt{\delta-c_{3}}} + \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})+2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(80)

Family II

$$\begin{split} \psi(\xi) &= \frac{4\delta_1 e^{\sqrt{\delta_1}\xi}}{-2\delta_2 e^{\sqrt{\delta_1}\xi} + e^{2\sqrt{\delta_1}\xi} + \delta_2^2 - 4\delta_1\delta_3}.\\ Case \ I \ a_1 &= 0, b_1 = 0, d_1 = -\frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \gamma = \frac{8a_0(c_3 - \delta) + 3c_2}{12\kappa^2},\\ \delta_1 &= \frac{4\left(a_0^2(\delta - c_3)(\mu_1 - \mu_2 \nu) + \sqrt{3}\sqrt{c_1(\delta - c_3)(\mu_1 - \mu_2 \nu)^2}\right)}{3(\mu_1 - \mu_2 \nu)^2}, \delta_2 = 0, \delta_3 = 0,\\ \omega &= \frac{(\mu_1 - \mu_2 \nu)(4a_0^2(\delta - c_3) - 3\kappa^2 \mu_1) - 2\sqrt{3}\sqrt{c_1(\delta - c_3)(\mu_1 - \mu_2 \nu)^2}}{3(\kappa \mu_2 - 1)(\mu_2 \nu - \mu_1)}, \end{split}$$

$$q_{29}(x,t) = e^{i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}},$$
(81)

$$q_{29}^{*}(x,t) = e^{-i\phi} \sqrt{a_0 - \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}}.$$
(82)

Case II $a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3}\sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \gamma = \frac{8a_0(c_3 - \delta) + 3c_2}{12\kappa^2},$

$$\delta_{1} = \frac{4\left(a_{0}^{2}(\delta-c_{3})(\mu_{1}-\mu_{2}\nu)+\sqrt{3}\sqrt{c_{1}(\delta-c_{3})(\mu_{1}-\mu_{2}\nu)^{2}}\right)}{3(\mu_{1}-\mu_{2}\nu)^{2}}, \delta_{2} = 0, \delta_{3} = 0,$$

$$\omega = \frac{(\mu_{1}-\mu_{2}\nu)(4a_{0}^{2}(\delta-c_{3})-3\kappa^{2}\mu_{1})-2\sqrt{3}\sqrt{c_{1}(\delta-c_{3})(\mu_{1}-\mu_{2}\nu)^{2}}}{3(\kappa\mu_{2}-1)(\mu_{2}\nu-\mu_{1})},$$

substituting values in Eq. (60) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{30}(x,t) = e^{i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}},$$
(83)

$$q_{30}^{*}(x,t) = e^{-i\phi} \sqrt{a_0 + \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}}}.$$
(84)

$$\begin{aligned} \text{Case III } a_0 &= -\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \Biggl(4\Bigl(\delta - c_3) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \Biggr), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^3}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{31}(x,t) = e^{i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1\sqrt{\delta - c_3}} -\frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) - 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(85)

$$q_{31}^{*}(x,t) = e^{-i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_{1}\xi}}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}\xi}}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}} - \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})-2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(86)

$$Case IV a_{0} = -\frac{1}{2} \sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})-2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}}, a_{1} = 0, b_{1} = 0, d_{1} = \frac{\sqrt{3}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\sqrt{\delta-c_{3}}},$$
$$\gamma = \frac{1}{12}\mu_{2}^{2} \left(4\left(\delta-c_{3}\right)\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})-2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} + 3c_{2} \right), \kappa = \frac{1}{\mu_{2}},$$
$$\delta = -\frac{\mu_{1}\mu_{2}^{2}(c_{3}-\delta)+2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{2}}} + 3c_{2} \right), \kappa = \frac{1}{\mu_{2}},$$

 $\delta_1 = \frac{\mu_1 \mu_2^2 (c_3 - 0) + 2\sqrt{3} \sqrt{c_1 \mu_2^2 (\delta - c_3)^2}}{\mu_2^4 (\delta - c_3) (\mu_2 v - \mu_1)}, \delta_2 = 0, \delta_3 = 0,$ substituting values in Eq. (60) also using $u = V^{\frac{1}{2}}$ and by applying the reverse transformation (7), we get our solutions

$$q_{32}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1\xi}} \sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}} \sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1 \sqrt{\delta - c_3}} - \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(87)

$$q_{32}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}\xi}}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}\xi}}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}} - \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3}) - 2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}}} \right)^{\frac{1}{2}}.$$
(88)

$$\begin{aligned} Case \ V a_0 &= \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \Biggl(4\Bigl(c_3 - \delta) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \Biggr), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^3}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{33}(x,t) = e^{i\phi} \times \left(\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} - \frac{\sqrt{3} e^{\sqrt{\delta_1} \xi} \sqrt{\delta_1^3 e^{-2\sqrt{\delta_1} \xi}} \sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1 \sqrt{\delta - c_3}} \right)^{\frac{1}{2}},$$
(89)

$$q_{33}^{*}(x,t) = e^{-i\phi} \times \left(\frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3}) - 2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}\xi}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}}\right)^{\frac{1}{2}}.$$
(90)

$$\begin{aligned} Case \ VI \ a_0 &= \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}} \\ \gamma &= \frac{1}{12} \mu_2^2 \Biggl(4 \Bigl(c_3 - \delta \Bigr) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \Biggr), \kappa = \frac{1}{\mu_2}, \end{aligned}$$

$$\delta_1 = \frac{\mu_1 \mu_2^2(c_3 - \delta) + 2\sqrt{3}\sqrt{c_1 \mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)(\mu_2 \nu - \mu_1)}, \delta_2 = 0, \delta_3 = 0,$$

$$q_{34}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1}\xi} \sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}} \sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1 \sqrt{\delta - c_3}} + \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(91)

$$q_{34}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}\xi}} \sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}\xi}}} \sqrt{\mu_{1} - \mu_{2}v}}{2\delta_{1}\sqrt{\delta - c_{3}}} + \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta - c_{3}) - 2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta - c_{3})^{3}}}{\mu_{2}^{4}(\delta - c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(92)

$$\begin{aligned} \text{Case VII } a_0 &= -\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \left(4 \left(\delta - c_3 \right) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \right), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3) (\mu_2 \nu - \mu_1)}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{35}(x,t) = e^{i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_1}\xi}\sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}}\sqrt{\mu_1 - \mu_2 v}}{2\delta_1\sqrt{\delta - c_3}} - \frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta - c_3) + 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta - c_3)^3}}{\mu_2^4(\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(93)

$$q_{35}^{*}(x,t) = e^{-i\phi} \times \left(-\frac{\sqrt{3}e^{\sqrt{\delta_{1}\xi}}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}\xi}}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}} - \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})+2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(94)



Fig. 1 Absolute graphical solution of $q_1(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 4$, $\mu_2 = -2$, $\mu_3 = 1$, $\delta = 4$



Fig. 2 Real graphical solution of $q_2(x, y, t)$ for the values of $\gamma = -3$, $\theta = 2$, $\mu_1 = 4$, $\mu_2 = -2$, $\mu_3 = 1$, $\delta = 4$



Fig. 3 Imaginary graphical solution of $q_2^*(x, y, t)$ for the values of $\gamma = -3$, $\theta = 2$, $\mu_1 = 4$, $\mu_2 = -2$, $\mu_3 = 1$, $\delta = 4$



Fig. 4 Absolute graphical solution of $q_3(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 1$, $\mu_2 = 1$, $\mu_3 = 2$, $\delta = 4$, $a_0 = -0.1$



Fig.5 Absolute graphical solution of $q_4(x, y, t)$ for the values of $\gamma = 3, \theta = 2, \mu_1 = 2, \mu_2 = 1, \mu_3 = 3, \delta = 0.1, a_0 = -1$



Fig.6 Absolute graphical solution of $q_4^*(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 2$, $\mu_2 = 1$, $\mu_3 = 3$, $\delta = 0.1$, $a_0 = -1$

$$\begin{aligned} \text{Case VIII } a_0 &= -\frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 v}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \left(4 \left(\delta - c_3 \right) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \right), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^3}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

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Fig. 7 Absolute graphical solution of $q_5(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 1$, $\mu_2 = 1$, $\mu_3 = 1$, $\delta = 4$, $a_0 = 0.1$



Fig. 8 Absolute graphical solution of $q_5^*(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 1$, $\mu_2 = 1$, $\mu_3 = 1$, $\delta = 4$, $a_0 = 0.1$



Fig. 9 Absolute graphical solution of $q_6(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 2$, $\mu_2 = 1$, $\mu_3 = 3$, $\delta = 0.1$, $a_0 = 1$



Fig. 10 Absolute graphical solution of $q_6^*(x, y, t)$ for the values of $\gamma = 3, \theta = 2, \mu_1 = 2, \mu_2 = 1, \mu_3 = 3, \delta = 0.1, a_0 = 1$



Fig. 11 Real graphical solution of $q_7(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 1$, $\mu_2 = 3$, $\mu_3 = 1$, $\delta = 0.5$



Fig. 12 Imaginary graphical solution of $q_7^*(x, y, t)$ for the values of $\gamma = 3, \theta = 2, \mu_1 = 1, \mu_2 = 3, \mu_3 = 1, \delta = 0.5$

$$q_{36}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1\xi}} \sqrt{\delta_1^3 e^{-2\sqrt{\delta_1\xi}}} \sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1 \sqrt{\delta - c_3}} - \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(95)

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Fig. 13 Real graphical solution of $q_8(x, y, t)$ for the values of $\gamma = -3$, $\theta = 2$, $\mu_1 = 4$, $\mu_2 = -2$, $\mu_3 = 1$, $\delta = 4$



Fig. 14 Imaginary graphical solution of $q_8^*(x, y, t)$ for the values of $\gamma = -3$, $\theta = 2$, $\mu_1 = 4$, $\mu_2 = -2$, $\mu_3 = 1$, $\delta = 4$



Fig. 15 Real graphical solution of $q_{10}(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 2$, $\mu_2 = 1$, $\mu_3 = 3$, $\delta = 0.1$, $a_0 = -1$

$$q_{36}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}\xi}}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}\xi}}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}} - \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})+2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(96)



Fig. 16 Imaginary graphical solution of $q_{10}^*(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 2$, $\mu_2 = 1$, $\mu_3 = 3$, $\delta = 0.1$, $a_0 = -1$



Fig. 17 Real graphical solution of $q_{12}(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 2$, $\mu_2 = 1$, $\mu_3 = 3$, $\delta = 0.1$, $a_0 = 1$



Fig. 18 Imaginary graphical solution of $q_{12}^*(x, y, t)$ for the values of $\gamma = 3$, $\theta = 2$, $\mu_1 = 2$, $\mu_2 = 1$, $\mu_3 = 3$, $\delta = 0.1$, $a_0 = 1$

$$\begin{aligned} Case \ IX \ a_0 &= \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = -\frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \left(4 \left(c_3 - \delta \right) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \right), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3) (\mu_2 \nu - \mu_1)}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

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$$q_{37}(x,t) = e^{i\phi} \times \left(\frac{1}{2}\sqrt{\frac{3\mu_1\mu_2^2(\delta-c_3) + 2\sqrt{3}\sqrt{c_1\mu_2^8(\delta-c_3)^3}}{\mu_2^4(\delta-c_3)^2}} - \frac{\sqrt{3}e^{\sqrt{\delta_1\xi}}\sqrt{\delta_1^3e^{-2\sqrt{\delta_1\xi}}}\sqrt{\mu_1-\mu_2\nu}}{2\delta_1\sqrt{\delta-c_3}}\right)^{\frac{1}{2}},$$
(97)

$$q_{37}^{*}(x,t) = e^{-i\phi} \times \left(\frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta-c_{3})+2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta-c_{3})^{3}}}{\mu_{2}^{4}(\delta-c_{3})^{2}}} - \frac{\sqrt{3}e^{\sqrt{\delta_{1}}\xi}\sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}}\xi}}\sqrt{\mu_{1}-\mu_{2}\nu}}{2\delta_{1}\sqrt{\delta-c_{3}}}\right)^{\frac{1}{2}}.$$
(98)

$$\begin{aligned} Case \ X \ a_0 &= \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}}, a_1 = 0, b_1 = 0, d_1 = \frac{\sqrt{3} \sqrt{\mu_1 - \mu_2 \nu}}{2\sqrt{\delta - c_3}}, \\ \gamma &= \frac{1}{12} \mu_2^2 \bigg(4 \big(c_3 - \delta \big) \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} + 3c_2 \bigg), \kappa = \frac{1}{\mu_2}, \\ \delta_1 &= \frac{\mu_1 \mu_2^2 (c_3 - \delta) - 2\sqrt{3} \sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^3}, \delta_2 = 0, \delta_3 = 0, \end{aligned}$$

$$q_{38}(x,t) = e^{i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_1}\xi} \sqrt{\delta_1^3 e^{-2\sqrt{\delta_1}\xi}} \sqrt{\mu_1 - \mu_2 \nu}}{2\delta_1 \sqrt{\delta - c_3}} + \frac{1}{2} \sqrt{\frac{3\mu_1 \mu_2^2 (\delta - c_3) + 2\sqrt{3}\sqrt{c_1 \mu_2^8 (\delta - c_3)^3}}{\mu_2^4 (\delta - c_3)^2}} \right)^{\frac{1}{2}},$$
(99)

$$q_{38}^{*}(x,t) = e^{-i\phi} \times \left(\frac{\sqrt{3}e^{\sqrt{\delta_{1}\xi}} \sqrt{\delta_{1}^{3}e^{-2\sqrt{\delta_{1}\xi}}} \sqrt{\mu_{1} - \mu_{2}v}}{2\delta_{1}\sqrt{\delta - c_{3}}} + \frac{1}{2}\sqrt{\frac{3\mu_{1}\mu_{2}^{2}(\delta - c_{3}) + 2\sqrt{3}\sqrt{c_{1}\mu_{2}^{8}(\delta - c_{3})^{3}}}{\mu_{2}^{4}(\delta - c_{3})^{2}}} \right)^{\frac{1}{2}}.$$
(100)

4 Discussion and results

In this section, the graphical representations of LPD model have been illustrated. The extended modified AEM method is applied to get the exact solitary wave solutions for the set of values. The results attained here are the soliton wave solutions, bright solitons, dark solitons, multi solitons, periodic solitary wave, rational function and elliptic function solutions for some appropriated values of parameters. The 3D, contour and 2D graphs visualizes the nature of nonlinear waves constructed from Eq. (1).

5 Conclusion

In this work, the extended modified auxiliary equation mapping approach is utilized to study the Lakshmanan-Porsezian-Daniel model with the Ker law, parabolic law and anticubic law of nonlinearity. As a result, a variety of exact solutions are established in the form of dark solitons, light solitons, solitary wave, periodic solitary wave, rational function, and complex function solutions, for details see Figs. 1, 2, 3, 4, 5, 6, 7, 8, 9, 10, 11, 12, 13, 14, 15, 16, 17 and 18. The obtained solutions conforms that the applied technique is more powerful and efficient approach for obtaining the exact optical solutions to a variety of nonlinear problems emerging in the contemporary era of communications network technology and nonlinear optics. In future, other methodologies and nonlinearity laws may be used to explore such model, thus there is still a lot of fresh work to be done on it.

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Declarations

Conflict of interest The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

References

- Ahmad, I., Hijaz, A., Inc, M., Hadi, R., Ali Akbar, M., Mostafa, M.A.K., Lanre, A., Adil, J.: Solution of fractional-order Korteweg-de Vries and Burgers' equations utilizing local meshless method. J. Ocean Eng. Sci. (2021)
- Ahmed, I., Seadawy, A.R., Lu, D.: M-shaped rational solitons and their interaction with kink waves in the fokas-lenells equation. Phys. Scr. 94(5), 055205 (2019)
- Akinyemi, L., Senol, M., Iyiola, O.S.: Exact solutions of the generalized multidimensional mathematical physics models via sub-equation method. Math. Comput. Simul. 182, 211–233 (2021)
- Akram, G., Sadaf, M., Arshed, S., Sameen, F.: Bright, dark, kink, singular and periodic soliton solutions of Lakshmanan-Porsezian-Daniel model by generalized projective Riccati equations method. Optik 241, 167051 (2021)
- Akram, G., Sarfraz, M.: Multiple optical soliton solutions for cgl equation with kerr law nonlinearity via extended modified auxiliary equation mapping method. Optik 242, 167258 (2021)
- Al-Munawarah, A.-M., Arabia, S.: Optical soliton perturbation with fractional temporal evolution by extended modified auxiliary equation mapping. Rev. Mex. Fis. 67(3), 403–414 (2021)
- Ali, K.K., Gómez-Aguilar, J.F.: New solitary wave solutions of the space-time fractional coupled equal width wave equation (CEWE) and coupled modified equal width wave equation (CMEWE). Int. J. Appl. Computat. Math. 7(4), 1–19 (2021)
- Alzahrani, A.K., Belic, M.R.: Cubic-quartic optical soliton perturbation with Lakshmanan-Porsezian-Daniel model by semi-inverse variational principle. Ukr. J. Phys. Opt. 22(123) (2021)
- Asimakopoulos, G., Whalley, J.: Market leadership, technological progress and relative performance in the mobile telecommunications industry. Technol. Forecast. Soc. Change 123, 57–67 (2017)
- Bandelow, U., Amiranashvili, S., Pickartz, S.: Stabilization of optical pulse transmission by exploiting fiber nonlinearities. J. Light. Technol. 38(20), 5743–5747 (2020)
- Biswas, A., Ekici, M., Sonmezoglu, A., Belic, M.R.: Solitons in optical fiber bragg gratings with dispersive reflectivity by extended trial function method. Optik 182, 88–94 (2019)
- Biswas, A., Ekici, M., Sonmezoglu, A., Triki, H., Majid, F.B., Zhou, Q., Moshokoa, S.P., Mirzazadeh, M., Belic, M.: Optical solitons with Lakshmanan-Porsezian-Daniel model using a couple of integration schemes. Optik 158, 705–711 (2018)
- Biswas, A., Yildirim, Y., Yasar, E., Zhou, Q., Moshokoa, S.P., Belic, M.: Optical solitons for Lakshmanan-Porsezian-Daniel model by modified simple equation method. Optik 160, 24–32 (2018)
- Buck, J.A.: Fundamentals of Optical Fibers. Wiley, New Jersey (2004)
- Bulbul, A.A., Jibon, R.H., Das, S.K., Roy, T., Saha, A., Hossain, M.B.: Pcf based formalin detection by exploring the optical properties in thz regime. Nanosci. Nanotechnol.-Asia 11(3), 314–321 (2021)
- Chen, S., Lü, X., Li, M.-G., Wang, F.: Derivation and simulation of the M-lump solutions to two (2+1)-dimensional nonlinear equations. Phys. Scr. 96(9), 095201 (2021)
- Chen, S.-J., Lü, X., Tang, X.F.: Novel evolutionary behaviors of the mixed solutions to a generalized Burgers equation with variable coefficients. Commun. Nonlinear Sci. Numer. Simul. 95, 105628 (2021)
- Cui, T.J., Smith, D.R., Liu, R.: Metamaterials. Springer, New York (2010)
- Ekici, M., Sonmezoglu, A.: Optical solitons with biswas-arshed equation by extended trial function method. Optik 177, 13–20 (2019)
- El-Labany, S., El-Taibany, W., Behery, E., Fouda, S.: Collision of dust ion acoustic multisolitons in a nonextensive plasma using hirota bilinear method. Phys. Plasmas 25(1), 013701 (2018)
- Hosseini, K., Mayeli, P., Kumar, D.: New exact solutions of the coupled sine-gordon equations in nonlinear optics using the modified kudryashov method. J. Mod. Opt. 65(3), 361–364 (2018)
- Hubert, M.B., Gambo, B., Mibaile, J., Serge, Y.D., Kofane, T.C., Anjan, B., Qin, Z., et al.: Optical solitons with Lakshmanan-Porsezian-Daniel model by modified extended direct algebraic method. Optik 162, 228–236 (2018)
- Inegbedion, H., Obadiaru, E.: Modelling brand loyalty in the nigerian telecommunications industry. J. Strateg. Mark. 27(7), 583–598 (2019)
- Islam, T., Gómez-Aguilar, J.F., Ali, A., Fernández-Anaya, G.: Diverse soliton structures for fractional nonlinear Schrodinger equation, KdV equation and WBBM equation adopting a new technique. Opt. Quantum Electron. 53(12), 1–27 (2021)
- Islam, M.T., Md Ali, A., Gómez-Aguilar, J.F., Bonyah, E., Fernandez-Anaya, G.: Assorted soliton structures of solutions for fractional nonlinear Schrodinger types evolution equations. J. Ocean Eng. Sci. (2021)
- Ivezić, Ż, Kahn, S.M., Tyson, J.A., Abel, B., Acosta, E., Allsman, R., Alonso, D., AlSayyad, Y., Anderson, S.F., Andrew, J., et al.: Lsst: from science drivers to reference design and anticipated data products. Astrophys. J. 873(2), 111 (2019)

Jin-Ming, Z., Yao-Ming, Z.: The hirota bilinear method for the coupled burgers equation and the highorder boussinesq-burgers equation. Chin. Phys. B 20(1), 010205 (2011)

Kamenkovich, V.M.: Fundamental of Ocean Dynamics. Elsevier, Amsterdam (2011)

- Khater, M., Adil, J., Hadi, R., Lanre, A., Ali Akbar, M., Inc, M., Hijaz, A.: New kinds of analytical solitary wave solutions for ionic currents on microtubules equation via two different techniques. Opt. Quantum Electron. 53(11), 1–27 (2021)
- Khater, M.M., Anwar, S., Tariq, K.U., Mohamed, M.S.: Some optical soliton solutions to the perturbed nonlinear schrödinger equation by modified khater method. AIP Adv. **11**(2), 025130 (2021)
- Khater, M., Attia, R.A., Lu, D.: Modified auxiliary equation method versus three nonlinear fractional biological models in present explicit wave solutions. Math. Comput. Appl. 24(1), 1 (2019b)
- Khater, M.M., Lu, D., Attia, R.A.: Dispersive long wave of nonlinear fractional wu-zhang system via a modified auxiliary equation method. AIP Adv. 9(2), 025003 (2019a)
- Khodadad, F.S., Mirhosseini-Alizamini, S.M., Günay, B., Akinyemi, L., Inc, M., Rezazadeh, H.: Abundant optical solitons to the Sasa-Satsuma higher-order nonlinear Schrödinger equation. Opt. Quantum Electron. 53(12), 1–17 (2021)
- Kudryashov, N.A.: The painlevé approach for finding solitary wave solutions of nonlinear nonintegrable differential equations. Optik 183, 642–649 (2019)
- Kumar, D., Seadawy, A.R., Joardar, A.K.: Modified kudryashov method via new exact solutions for some conformable fractional differential equations arising in mathematical biology. Chin. J. Phys. 56(1), 75–85 (2018)
- Li, L.-X., Li, E.-Q., Wang, M.-L.: The (g'/g, 1/g)-expansion method and its application to travelling wave solutions of the zakharov equations. Appl. Math. J. Chin. Univ. 25(4), 454–462 (2010)
- Li, W., Schmitt, D.R., Zou, C., Chen, X.: A program to calculate pulse transmission responses through transversely isotropic media. Comput. Geosci. **114**, 59–72 (2018)
- Lindo, Z.: Transoceanic dispersal of terrestrial species by debris rafting. Ecography **43**(9), 1364–1372 (2020)
- Liu, W., Zhang, Y., Triki, H., Mirzazadeh, M., Ekici, M., Zhou, Q., Biswas, A., Belic, M.: Interaction properties of solitonics in inhomogeneous optical fibers. Nonlinear Dyn. 95(1), 557–563 (2019)
- Lu, D., Seadawy, A., Arshad, M.: Applications of extended simple equation method on unstable nonlinear schrödinger equations. Optik 140, 136–144 (2017)
- Lü, X., Chen, S.-J.: New general interaction solutions to the KPI equation via an optional decoupling condition approach. Commun. Nonlinear Sci. Numer. Simul. 103, 105939 (2021)
- Lü, X., Chen, S.-J.: Interaction solutions to nonlinear partial differential equations via Hirota bilinear forms: one-lump-multi-stripe and one-lump-multi-soliton types. Nonlinear Dyn. 103(1), 947–977 (2021)
- Lü, X., Hua, Y.-F., Chen, S.-J., Tang, X.-F.: Integrability characteristics of a novel (2+ 1)-dimensional nonlinear model: Painlevé analysis, soliton solutions, Bäcklund transformation, Lax pair and infinitely many conservation laws. Commun. Nonlinear Sci. Numer. Simul. 95, 105612 (2021a)
- Lü, X., Hui, H.-W., Liu, F.-F., Bai, Y.-L.: Stability and optimal control strategies for a novel epidemic model of COVID-19. Nonlinear Dyn. 106(2), 1491–1507 (2021b)
- Ma, Q., Cui, T.J.: Information metamaterials: bridging the physical world and digital world. PhotoniX 1(1), 1–32 (2020)
- Mahak, N., Akram, G.: Exact solitary wave solutions by extended rational sine-cosine and extended rational sinh-cosh techniques. Phys. Scr. 94(11), 115212 (2019)
- Mahak, N., Akram, G.: The modified auxiliary equation method to investigate solutions of the perturbed nonlinear schrödinger equation with kerr law nonlinearity. Optik **207**, 164467 (2020)
- Naher, H., Abdullah, F.A.: The modified benjamin-bona-mahony equation via the extended generalized riccati equation mapping method. Appl. Math. Sci. 6(111), 5495–5512 (2012)
- Nawaz, B., Rizvi, S.T.R., Ali, K., Younis, M.: Optical soliton for perturbed nonlinear fractional schrödinger equation by extended trial function method. Opt. Quantum Electron. 50(5), 1–14 (2018)
- Olbers, D., Willebrand, J., Eden, C.: Ocean Dynamics. Springer, New York (2012)
- Rezazadeh, H., Korkmaz, A., Khater, M.M., Eslami, M., Lu, D., Attia, R.A.: New exact traveling wave solutions of biological population model via the extended rational sinh-cosh method and the modified khater method. Mod. Phys. Lett. B 33(28), 1950338 (2019)
- Rezazadeh, H., Mirzazadeh, M., Mirhosseini-Alizamini, S.M., Neirameh, A., Eslami, M., Zhou, Q.: Optical solitons of lakshmanan-porsezian-daniel model with a couple of nonlinearities. Optik 164, 414–423 (2018)

- Seadawy, A.R., Cheemaa, N.: Applications of extended modified auxiliary equation mapping method for high-order dispersive extended nonlinear schrödinger equation in nonlinear optics. Mod. Phys. Lett. B 33(18), 1950203 (2019)
- Sun, Y.-L., Ma, W.-X., Yu, J.-P.: N-soliton solutions and dynamic property analysis of a generalized threecomponent hirota-satsuma coupled kdv equation. Applied . Lett. 120, 107224 (2021)
- Wang, Q., Kong, L., Dang, Y., Xia, F., Zhang, Y., Zhao, Y., Hu, H., Li, J.: High sensitivity refractive index sensor based on splicing points tapered smf-pcf-smf structure mach-zehnder mode interferometer. Sens. Actuators B Chem. 225, 213–220 (2016)
- Wang, Y., Rehren, T., Tan, Y., Cong, D., Jia, P.W., Henderson, J., Ma, H., Betts, A., Chen, K.: New evidence for the transcontinental spread of early faience. J. Archaeol. Sci. 116, 105093 (2020)
- Wazwaz, A.-M.: A sine-cosine method for handlingnonlinear wave equations. Math. Comput. Model. 40(5– 6), 499–508 (2004)
- Wazwaz, A.-M.: Gaussian solitary waves for the logarithmic-kdv and the logarithmic-kp equations. Phys. Scr. 89(9), 095206 (2014)
- Weckbrodt, J., Ginot,N., Batard, C., Azzopardi, S.: Short pulse transmission for sic communicating gate driver under high dv/dt. In: PCIM Europe 2018; International Exhibition and Conference for Power Electronics, Intelligent Motion, Renewable Energy and Energy Management, VDE, pp. 1–6. (2018)
- Yildirim, Y., Topkara, E., Biswas, A., Triki, H., Ekici, M., Guggilla, P., Khan, S., Belic, M.R.: Cubic-quartic optical soliton perturbation with Lakshmanan-Porsezian-Daniel model by sine-Gordon equation approach. J. Opt. 50(2), 322–329 (2021)
- Yin, Y.-H., Xing, L., Si-Jia, C.: Localized characteristics of lump and interaction solutions to two extended Jimbo-Miwa equations. Chin. Phys. B 29(12), 120502 (2020)
- Yong, C., Biao, L., Hong-Qing, Z.: Generalized riccati equation expansion method and its application to the bogoyavlenskii's generalized breaking soliton equation. Chin. Phys. 12(9), 940 (2003)
- Yépez-Martínez, H., Inc, M., Hadi, R., Mehmet, A.A., Gómez-Aguilar, J.F.: Analytical solutions to the fractional Lakshmanan-Porsezian-Daniel model. Opt. Quantum Electr. 54(1), 1–41 (2022)
- Yépez-Martínez, H., Rezazadeh, H., Inc, M., Gómez-Aguilar, J.F.: A new local fractional derivative applied to the analytical solutions for the nonlinear Schrödinger equation with third-order dispersion. J. Nonlinear Opt. Phys. Mater. 31, 2250011 (2022)
- Zafar, A., Shakeel, M., Ali, A., Akinyemi, L., Rezazadeh, H.: Optical solitons of nonlinear complex Ginzburg-Landau equation via two modified expansion schemes. Opt. Quantum Electron. 54(1), 1–15 (2022)
- Zayed, E.M., Shohib, R.M.: Optical solitons and other solutions to biswas-arshed equation using the extended simplest equation method. Optik 185, 626–635 (2019)

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