

# **Structural and optical investigation of pure and Al doped ZnO annealed at diferent temperatures**

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## **Abstract**

The structural and optical properties of  $ZnO$  and  $Zn<sub>0.9</sub>Al<sub>0.1</sub>O$  (AlZnO) samples annealed in air at different temperatures of  $(400-900 \degree C)$  for 12 h are investigated. It is found that cell parameters, average crystalline diameter and average grain size are afected by both Al and annealing. The ultra violet (UV) absorption coefficient  $(\alpha)$  is increased by Al and can be represented by the relation ;( $\alpha$  (AlZnO) $\sim$ 2 $\alpha$  (ZnO), but it is generally decreased by annealing. Two diferent types of band gap are obtained for each sample in which their values are higher for AlZnO than that of ZnO. The first band gap is called fundamental energy gap  $(E_{oh})$  and its value above 3 eV, while the second is called optical energy gap  $(E_{gl})$  and its value below 2.75 eV. In addition, the values of  $E_{el}$  (ZnO),  $E_{eh}$  (AlZnO) and  $E_{el}$  (AlZnO) are increased by annealing up to 700  $\degree$ C, followed by a decrease at 900  $\degree$ C, while the values of  $E_{gh}$  (ZnO) are increased by annealing up to 900 °C. Furthermore, a unique energy gap of 3.33 eV is obtained at 312 °C for ZnO and AlZnO. On the other hand, the lattice dielectric constant  $(\varepsilon_L)$ , free carrier concentration  $(N/m^*)$ , dielectric loss (*tan*  $\delta$ ), optical conductivity ( $\sigma_{opt}$ ) and electrical conductivity ( $\sigma_{ele}$ ) are obtained for both series. It is found that  $\varepsilon_L$  is increased by annealing up to 900 °C for both series, while  $(N/m^*)$  is increased by annealing for ZnO, but it is decreased for AlZnO. But their values are higher for ZnO than that of AlZnO. Further, the value of  $\varepsilon_L$  (11.26) of the as-prepared ZnO is nearly similar to the value of  $\varepsilon_L$  (11.76) for AlZnO at 900 °C. Although *tan*  $\delta$  is similar for both series, it is decreased by annealing up to 900 °C for ZnO, and increased for AlZnO up to 700 °C followed by decrease at 900 °C. Similar behavior is obtained for  $\sigma_{\text{cont}}$ , while the vice is versa for  $\sigma_{ele}$ . The obtained results are interpreted in terms of the variance of carrier concentration, oxygen vacancies and interatomic spacing against annealing temperature and Al content. The present results confrm the semiconducting nature as well as assure that the high photo-response of the investigated samples.

**Keywords** ZnO · Doping · Thermal annealing · XRD · SEM · Optical properties

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## **1 Introduction**

ZnO has attracted much interest due to its versatile properties such as transparency in the visible range, direct band gap, large free-exciton binding energy and absence of toxicity (Gledhill et al. [2009;](#page-19-0) Shishiyanu et al. [2005;](#page-20-0) Anas et al. [2007\)](#page-18-0). Three values of energy gap  $(3.1, 3.2 \text{ and } 3.3 \text{ eV})$  and free exciton energy of  $(60-100 \text{ mV})$  have been reported for ZnO at room temperature (Hagemark [1976](#page-19-1); Azmi et al. [2005;](#page-18-1) Srikant and Clarke [1998\)](#page-20-1). The structure of ZnO is similar to GaN, but in contrast, ZnO is available as large bulk and single crystal. Generally, ZnO structure consists of alternating planes, containing tetrahedral coordinated  $O^{-2}$  and  $Zn^{2+}$  ions, which are stacked alternately along the *c*-axis. ZnO has a hexagonal structure (space group C6mc) with ideal lattice parameters a=0.3296 nm and c=0.52065 nm (Kim et al. [2008;](#page-19-2) Chaari et al. [2011](#page-18-2); Matsuoka et al. [1970](#page-19-3)).

ZnO is an *n*-type semiconductor due to excess of Zn metal through Zn twice ionization. The excess of Zn gives origin to intrinsic donors in ZnO, and can be assigned either to the interstitial Zn or to the oxygen vacancy (Houabes et al. [2005](#page-19-4); Ramanachalam et al. [1995;](#page-20-2) Secco and Moore [1957;](#page-20-3) Krèoger [1974;](#page-19-5) Mahan [1983\)](#page-19-6). However, the properties of ZnO are controlled by these intrinsic defects (oxygen vacancies) generated at high temperatures during sintering and annealing (Sedky et al. [2007\)](#page-20-4). During the quenching of the ZnO, these defects are usually annihilated at the grain boundaries. Since this migration process is normally slow as well as thermally activated, these defects are efectively eliminated (Han et al. [2002\)](#page-19-7). It is generally accepted that the localized states of band gap deplete the charge carriers from bulk ZnO on the either side of the grain boundaries, resulting in the formation of Schottky double barriers that control ZnO properties (Madou and Morrison [1989;](#page-19-8) Bonasewicz et al. [1986;](#page-18-3) Castro and Aldao [1998](#page-18-4)).

Extensive research efforts have been directed towards the fabrication of ZnO and its properties investigation. It has been found that the properties of ZnO are closely related to their oxygen vacancies through either flling or ordering. These oxygen vacancies could be also controlled by diferent dopants such as Ni, Bi, Mn, Ga, Co, Fe and Al (Sedky et al. [2019;](#page-20-5) Sedky [2018;](#page-20-6) Greuter and Blatter [1990\)](#page-19-9).

Structural and electrical properties of  $\text{Zn}_{1-x} \text{Al}_x\text{O}$  have investigated by Sedky and El-Suheel [\(2012](#page-20-7)). It has been reported that Al has no effect on the XRD peaks corresponding to the wurtzite structure of ZnO, while the average crystalline diameter is decreased. Furthermore, the nonlinear coefficient and breakdown field are shifted to lower values, and consequently the conductivity of ZnO grains is improved. Moreover, it is found that Al, close to 10%, reacted with ZnO and caused depressing of the ZnO grain growth. But above 10%, the Al atoms move to the interstitial positions of ZnO lattice.

The optical and electronic properties of ZnO are very sensitive for any variation in the chemical structure of ZnO, which make these compounds having a wide range of applications (Zeyada et al. [2012](#page-20-8); Kumari et al. [2018](#page-19-10); Nursaadah Ahmad et al. [2018](#page-20-9); Shahrokh Abadi et al. [2015\)](#page-20-10). Although ZnO has a wide range in many applications, there is a lack in the data based on thermal stability of ZnO annealed with dopants, especially at high temperature above 500 °C (Ramanachalam et al. [1995](#page-20-2); Zhu et al. [2008;](#page-20-11) Kang et al. [2008;](#page-19-11) Shan et al. [2004;](#page-20-12) Kang et al. [2003](#page-19-12); Shinde et al. [2006\)](#page-20-13). However, evaluation of optical constants for pure and doped ZnO has a signifcant importance in optoelectronic and optical applications such as flters, modulators, solar cells, photodiode, switches, etc.

The stability of electrical properties in ZnO has been investigated by ion migration and defects models. They proposed that Zn interstitials can be formed and frozen in the depletion layer during cooling from sintering temperature down to room temperature (Eda et al. [1980;](#page-18-5) Sato et al. [1980\)](#page-20-14). It is also assumed that some of defects, such as donor-like positive charge and acceptor like negative charge, could be formed in the grain boundary and afect ZnO stability (Gupta et al. [1981](#page-19-13); Einzinger [1982\)](#page-18-6). Moreover, some of positive charges  $\text{Zn}^+$ and  $Zn^{++}$  could be formed in the depletion layer, but these ions could be removed by heat treatments such as thermal annealing (Gupta and Miller [1988](#page-19-14)). On the other hand, oxygen vacancies formed below the conduction band have been also identifed as donor defects and afect the degree of ZnO stability (Klug and Alexander [1962](#page-19-15)). In addition, the dopant of 3d transition metals such as Al has indirectly confrmed the grain boundary defects beside Zn in the depletion layer.

Based on the above, we tried to eliminate the possible efects of ion migration to the grain boundaries or the reaction of grain boundary defects with ambient oxygen on the ZnO and AlZnO optical properties. Consequently, this study focuses on the infuence of annealing at a temperature up to 900  $^{\circ}$ C on the structural and morphological of the two series through X-ray difraction (XRD), scanning electron microscope (SEM) characterization. Moreover, a detailed investigation of the optical constants as a result of Al doping and annealing is introduced. The present study reveals signifcant variations for ZnO and AlZnO series with respect to annealing temperature. Anyhow, by annealing up to 900 °C, we could increase the lattice dielectric constant, fundamental energy gap, and electrical conductivity for pure ZnO, but the optical conductivity and dielectric loss factor are decreased. While annealing of AlZnO generally improve the lattice dielectric constant, dielectric loss factor and optical conductivity, but the fundamental energy gap and electrical conductivity are decreased. However, these results gave us more hooping for ZnO stability optimization through both thermal annealing and doping by transition metals such as Al.

### **2 Experimental details**

ZnO and  $Zn_{0.9}Al_{0.1}O$  samples are prepared by using solid-state method through different steps. First, the  $Al_2O_3$  and ZnO powders (purity: 99.999, Aldrich) are mixed together according to the atomic percentages. In the next step, the mixed powders are calcined in air at 1000 °C for 12 h. After that, the pellets of calcined powders are sintered at 1000 °C for 10 h in air. In the fnal step of the synthesis, the samples are quenched in air at room temperature. However, a range of two similar sets of  $\text{Zn}_{1-x}\text{Fe}_{x}O$  samples with various x values are synthesized by Sedky et al. [\(2007](#page-20-4)). The frst set is quenched from the sintering temperature to room temperature, while the second set left in the furnace and slowly cooled to room temperature. Although the wurtzite structure and microstructure of pure ZnO are found to be similar in both cases, the breakdown field and nonlinear coefficient are three times higher in quenching than that of slow cooling. This could be attributed to the oxygen vacancies which are produced by quenching and supporting the potential barrier between ZnO grains. Therefore, in the present case, we prefer quenching than slow cooling. The pellets are divided into four pieces (one piece for as prepared and three pieces for annealing). The thermal annealing process for the studied samples is carried out individually in the furnace at different temperatures (400  $^{\circ}$ C, 700  $^{\circ}$ C and 900  $^{\circ}$ C) for 12 h in air. The structure of the samples is investigated by X-ray diffraction technique (Philips PW-1700) using  $Cu-K\alpha$ radiation of wave length  $(\lambda) = 1.5418$  Å at 40 kV and 30 mA settings, and at a diffraction angles range  $(20^{\circ}-70^{\circ})$ . The morphology of the samples is characterized by using scanning electron microscope (SEM) JSM 5400 LV. The optical absorbance and optical refectance are collected for the powder suspension of the samples using a thermo-evolution 300 UV

spectrophotometer with a scan speed of 1000 nm/min. The measurements are carried out at room temperature between (200–1000 nm) of photon wavelength.

## **3 Results and discussions**

#### **3.1 Structural analysis**

The crystal structure of the samples, shown by XRD pattern in Fig. [1a](#page-3-0), b, is wurtzite, and no other phases are observed. The density  $(\rho_{th})$  of the samples, calculated using XRD anal-ysis, is given in Table [1.](#page-4-0) It is clear that the annealing temperature does not affect  $\rho_{th}$  for all samples, but their values are generally higher for AlZnO than that of ZnO. The lattice

<span id="page-3-0"></span>

<span id="page-4-0"></span>

parameters, summarized in Table [1](#page-4-0), are plotted against annealing temperature and shown in Fig. [2a](#page-5-0). Although the lattice parameters are nearly unchanged by annealing temperature, their values are lower for AlZnO than that of ZnO. This could be attributed the ionic radius of  $Al^{3+}$  (0.51 Å), which is smaller than that of  $Zn^{2+}$  (0.74 Å).

The average crystallite size  $(D<sub>hkl</sub>)$  is calculated using the following expression (Klug and Alexander [1962](#page-19-15); Cullity and Stock [2001\)](#page-18-7);

$$
D_{hkl} = \frac{k\lambda}{\beta \cos \theta} \tag{1}
$$

where *k* is constant and equal 0.93 for ZnO,  $\lambda = 1.5418$  Å is the X-ray radiation wave length,  $\beta$  is half maximum line width and  $\theta$  is diffraction angle. The values of  $D_{h\nu}$ , summarized in Table [1,](#page-4-0) are plotted against annealing temperature in Fig. [2](#page-5-0)b. It is clear that  $D<sub>hkl</sub>$  values decrease by annealing for ZnO and increase for AlZnO. The values of  $D<sub>hkl</sub>$  are ranged between (42.11–58.76 nm) for ZnO, and from (44.78–63.16 nm) for AlZnO.

The morphology of the samples surface is shown in Fig. [3.](#page-6-0) The surface of the samples is rough and the grains of crystallites having very varied shapes of diferent sizes which are afected by annealing (see Fig. [3](#page-6-0)). Furthermore, the crystallites are randomly distributed and irregularly disoriented, and there is no additional phases formed at boundaries of the grain. The AlZnO sample revealed a decrease of the average grain size as compared to ZnO sample at the same annealing temperature. Moreover, the ZnO samples exhibit a uniform granular precipitation on the mother grains without black regions formation. In addition, the size of grains is increased by annealing up to 900  $^{\circ}$ C. On the other hand, the AlZnO show black regions with random distribution in the matrix structure and the size of grains are relatively small as compared to ZnO. Also, by increasing annealing at 900 °C, the shape of grains became not clear.

The average size of the grain (*D*) is determined by the linear intercept method with the help of expression,

$$
D = \frac{1.56L}{MN},
$$

where *L* is the length of random line on the micrograph, *M* is the micrograph magnification, and *N* is the number of grain boundaries intercepted by the lines (Wurst and Nel-son [1972](#page-20-15)). The values of the average grain size  $(D)$  given in Table [1,](#page-4-0) are plotted against annealing temperature in Fig. [2](#page-5-0)b. It is found that the values of *D* are about 13 times



<span id="page-5-0"></span>**Fig. 2 a** Lattice parameters versus annealing temperature for ZnO and AlZnO samples. **b** Grain size versus annealing temperature for ZnO and AlZnO samples

higher than  $D<sub>hkl</sub>$  determined from XRD. It is also apparent that average grain sizes are increased by annealing temperature for all studied samples. However, the reported TEM analysis of diluted magnetic semiconductors formed by partial substitution of ZnO with Ni showed that the size of most particles is about 60 nm (Pei et al. [2006\)](#page-20-16). However,  $D_{hkl}$ should be smaller than *D*. This is because XRD analysis determines the average diameter



ZnO (As-prepared)



Al ZnO (As-prepared)



ZnO  $(T=400 °C)$ 



Al ZnO  $(T=400 °C)$ 



 $ZnO(T=700 °C)$ 



 $ZnO(T=900 °C)$ 

<span id="page-6-0"></span>**Fig. 3** SEM images of the ZnO and Al ZnO samples



Al ZnO  $(T = 700 °C)$ 



Al ZnO (T=  $900 °C$ )

of crystallites which have the same orientation inside the grain, while SEM determines the average size of grain itself. It is well known that each grain contains lot of crystallites with the same orientations.

#### **3.2 Optical measurements and energy gap**

The optical absorbance (*A*) and optical refectance (*R*) versus the wavelength are shown in Fig. [4a](#page-7-0), b for all samples. A gradual decrease in the absorbance of ZnO with decreasing wavelength occurs. Also, there is an increase of absorbance by increasing photon energy near-infrared region (exciton peaks  $\sim$  400 nm), which can be attributed to an increase of free carrier concentration as the photon energy increases. Furthermore, the absorbance is gradually decreased by increasing annealing temperature up to 900 °C. On the other hand, a little bit increase in the absorbance of AlZnO samples with decreasing wavelength is observed. In addition, the values of absorbance of AlZnO are higher than that of ZnO (*A* of AlZnO  $\simeq$  2 *A* of ZnO). As compared to ZnO, it is observed that the absorbance of AlZnO is almost linear at the longer wavelengths, while it is nonlinear at the lower range of wavelength (300–400 nm). At such wave length range, a number of absorption peaks are recorded, and the intensity and number of these peaks are decreased by increasing annealing temperature. This can be attributed to the free carriers or some defects. Furthermore, the absorbance of AlZnO is increased by increasing annealing temperature up to 700  $^{\circ}$ C, followed by a decrease at 900 °C. The vice is versa for optical refectance *R*. This behavior may be related to oxygen vacancies and Al interstitials, which makes ZnO very suitable



(b) Absorption and reflectance versus wave length for annealed AlZnO samples

<span id="page-7-0"></span>**Fig. 4** Absorption and refectance versus wave length for annealed **a** ZnO samples, **b** AlZnO samples

for the use of oxygen sensor (Fergus [2003](#page-19-16)). The sensor output depends on the reactions between the defects in the ZnO oxide and the oxygen in the atmosphere (Ali and Abdel Hakeem [2010\)](#page-18-8).

However, ZnO acts as an n-type semiconductor and its conduction band is composed of 4 s orbital of Zn and wide enough for the efective transfer of the charges. Consequently, electrons are excited from the valence state due to absorption of light and generating electron-hole pairs necessary for conduction mechanism. Above the limit of the band gap, the signal of photon energy is due to non-radiative process, and below that limit, it is due to the Urbach tail (Shaaban et al. [2017\)](#page-20-17).

For optical band gap  $E_{\varrho}$  determination, the absorption coefficient  $\alpha$  is calculated using the formula,  $\alpha = 2.303$  ln ( $\dot{A}/t$ ), where *t* is the distance traveled by the photons into the solution of thickness,  $t = 12$  mm. After that,  $E<sub>g</sub>$  is calculated using the following Tauc's equation (Shen and Toyoda [2000](#page-20-18)):

$$
(\alpha h v)^{\frac{1}{m}} = A(hv - E_g) \tag{2}
$$

*A* is constant,  $\alpha$  is absorption coefficient,  $\nu$  is photon frequency and  $h$  is Planck's constant, respectively. *m* is a parameter which characterizes the type of electronic transition. It takes 1/2, 2, 1/3, 3 for direct and indirect allowed transition, and 1/3, 3 for direct and indirect forbidden transition, respectively. The present results showed that ZnO possess an allowed direct transition mechanism  $(m=1/2)$ . This is clearly observed through the exponential like behavior and the linear part charactering the optical band edge. Therefore,  $E<sub>o</sub>$  is determined by extrapolating the linear part of the plot of  $(ahv)^2$  versus photon energy *hv* to  $\alpha$ =0 (Abdalla et al. [2003](#page-18-9); Othman et al. [2017\)](#page-20-19). From Fig. [5a](#page-9-0) for ZnO, one can observe that  $\alpha$  values increase as the photon energy increases and it is significantly affected by the increase of annealing temperature up to 900 °C where the absorption edge also shifted to higher energy. While for AlZnO samples, shown in Fig. [5](#page-9-0)b, an increase in the values of *α* with increasing annealing temperature up to 700  $^{\circ}$ C is observed. But with further increase of annealing temperature up to  $700 \degree C$ , a decrease of the absorption coefficient is occurred. Thus the relation between the values of *α* between the two series can be represented by; *α*  $(AIZnO) \simeq 2\alpha$  (ZnO), indicating that AlZnO is convenient as absorber better than ZnO.

Anyhow, two diferent values of energy are evident from the plot shown in Fig. [5.](#page-9-0) The first band gap  $(E_{gh})$  is obtained at higher photon energy (UV region), while the second one  $(E_{gl})$  is obtained at lower photon energy (visible region). The behavior of  $E_g$  against annealing temperature for ZnO and AlZnO samples is shown in Fig. [5](#page-9-0)c, and similar values are summarized in Table [2.](#page-10-0) It is clear that the values of  $E_{gh}$  and  $E_{gl}$  are, respectively 3.1, 3.4, 3.5, 3.8 eV, and 2.05, 1.75, 1.53, 1.35 eV for ZnO. This indicates that *Egh* is increased, whereas  $E_{gl}$  is decreased as the annealing temperature increases up to 900 °C. On the other hand, *Egh* and *Egl* values for AlZnO are 3.48, 3.3, 3.19, 3.57 eV and 2.75, 2.65, 2.17, 2.25 eV, respectively. This of course means that both  $E_{gh}$  and  $E_{gl}$  are decreased by increasing annealing temperature up to 700 °C, followed by increase at 900 °C. Although the values of *Eg* agree with reported for ZnO compound, two optical energy gaps are also reported for *n*-type semiconductors (Ali et al. [2005;](#page-18-10) Darzi et al. [2009;](#page-18-11) Ottaviano et al. [2004;](#page-20-20) El-Nahass et al. [2016\)](#page-19-17). This can be in principle interpreted in terms of the shape of conduction and valence bands as follows. At higher photon energy and by assuming the geometrical symmetry, the electronic transition could take place from the valence band (VB) top or bottom to conduction band (CB) top or bottom, which is consistent with the higher absorptions values (homo-lumo). But when the phonon energy is decreased, the energy might be not enough for the above transition to occur. Therefore the transition of the electron is only



(a)  $(\text{ahv})^2$  versus photon energy at annealed ZnO samples



(c) Energy gap versus annealing temperature for ZnO and AlZnO samples

<span id="page-9-0"></span>**Fig. 5 a**  $(\alpha h\nu)^2$  versus photon energy at annealed ZnO samples. **b**  $(\alpha h\nu)^2$  versus photon energy of annealed AlZnO samples. **c** Energy gap versus annealing temperature for ZnO and AlZnO samples

Annealing temp.		$E_{gh} (eV)$ $E_{gl} (eV)$		$\Delta E$ (eV) N/m <sup>*</sup> (g <sup>-1</sup> cm <sup>-3</sup> ) N (cm <sup>-3</sup> )		$R$ (cm)	$\varepsilon_{\rm L}$
ZnO							
As-prepared	3.1	2.05	1.05	$3.69 \times 10^{55}$	$14.09 \times 10^{30}$	$4.24 \times 10^{-11}$	11.26
400	3.4	1.75	1.65	$3.81 \times 10^{55}$	$14.55 \times 10^{30}$	$4.19 \times 10^{-11}$	15.53
700	3.5	1.53	2.15	$3.87 \times 10^{55}$	$14.78 \times 10^{30}$	$4.17 \times 10^{-11}$	16.26
900	3.8	1.35	2.45	$4.02 \times 10^{55}$	$15.36 \times 10^{30}$	$4.12 \times 10^{-11}$	18.55
AlZnO							
As-prepared	3.48	2.75	0.73	$1.23 \times 10^{55}$	$4.70 \times 10^{30}$	$6.11 \times 10^{-11}$	7.44
400	3.30	2.65	0.65	$1.11 \times 10^{55}$	$4.24 \times 10^{30}$	$6.33 \times 10^{-11}$	8.74
700	3.19	2.16	1.03	$0.49 \times 10^{55}$	$1.87 \times 10^{30}$	$8.31 \times 10^{-11}$	9.36
900	3.57	2.25	1.32	$0.01 \times 10^{55}$	$0.04 \times 10^{30}$	$29.89 \times 10^{-11}$	11.74

<span id="page-10-0"></span>**Table 2**  $E_{gh}$ ,  $E_{gl}$ ,  $\Delta E$ , slope, N/m<sup>\*</sup>, N, R and  $\varepsilon_L$  of studied samples

take place from the top of VB to bottom of CB, which agrees with the lower absorptions values. Generally, this behavior is in a good agreement with the reported data based on n-type semiconductors (El-Denglawey [2018](#page-19-18); Yu et al. [2001](#page-20-21)).

The influence of thermal annealing up to 500  $^{\circ}$ C on the absorption coefficient of asdeposited semiconducting thin flms as a function of photon energy has been investigated (Wang et al. [2003;](#page-20-22) Abdel-Khalek et al. [2019](#page-18-12)). Two peaks have been obtained in the UV region; the frst at 398.7 nm (3.11 eV) and was red-shifted with increasing annealing temperature, while the second at 279.9 nm (4.43 eV). There is also shoulder exists at 223 nm (5.56 eV) which is attributed to a transition from trap level near the LUMO band to LUMO band (Zeyada et al. [2013\)](#page-20-23). It is approved that the UV emission peaks (300–380 nm) are fundamental and normally resulted from the transfer of electron from conduction band to valence band (Su et al. [1984\)](#page-20-24). While the violet emission peaks close to 450 nm (2.75 eV) are attributed to the transition of electron from conduction band to zinc vacancies (VZn) (Kundu et al. [2011\)](#page-19-19). The blue emissions peaks close to 500 nm (2.48 eV) are attributed to the transition of electron from zinc interstitial  $(Zn_i)$  to zinc vacancies (VZn). The green emission close 550 nm  $(2.25 \text{ eV})$  is attributed to the transition of electron from conduction band to the oxygen anisite vacancy defect  $(OZn)$  (Kundu et al. [2011\)](#page-19-19). Therefore, the annealed crystalline flms emits UV, violet, blue, and green emissions, but the deposited amorphous flm emit UV emission only which indicated that the density of point defect in the annealed flm is higher than the deposited at RT. These results strengthens the presence of two energy gap in pure and Al doped ZnO of the present work.

On the other hand, by looking to the behavior of  $E_{gh}$  against the annealing temperature, we noted that  $E_{gh}$  increased for as-prepared ZnO from 3.1 eV to 3.48 eV for as-prepared AlZnO. The beauty of Fig. [5c](#page-9-0) that  $E_{gh}$  equal 3.33 eV at a temperature of 314 °C for both series. Furthermore, an increase of temperature above 700 °C results an increase of *Egh* for both samples. While, *Egl* is generally decreased by increasing annealing temperature up to 900 °C for the two series, but it is always higher for AlZnO, which is completely diferent than that of  $E_{\phi h}$  behavior.

The crystalline and oxygen vacancies defects have signifcant efects on the optical band gap of ZnO. The as prepared ZnO sample is usually characterized by structure disorder and normally contains some defects which produce localized states (Lin et al. [2013](#page-19-20); Ali and Hakeem [2015](#page-18-13)). These defects are higher for the as prepared AlZnO, and consequently  $E_{oh}$ is 3.48 eV, which is higher than *Egh* for ZnO (3.1 eV). Additionally, the increase of optical

band gap by Al may be due to decreasing free carrier concentration, thereby resulting in a widening of the band gap which is well-known as Burstein–Moss efect (Burstein [1954](#page-18-14)). This will be discussed in the next paragraph. Anyhow, the higher values of  $E_{gh}$  for  $ZnO$ by annealing can be ascribed by structural changes from disordered to ordered structure, which is supported by increasing their crystalline size (1711–2371 nm). The vice is versa for *Egl* obtained in the visible region. On the other hand, both Al and annealing temperature generally decreased the crystalline size of ZnO (693–652 nm), and consequently  $E_{oh}$ decreased by annealing up to 700 °C, followed by an increase at 900 °C. Similar behavior is obtained for  $E_{ol}$ .

#### **3.3 Dielectric constant and optical conductivity**

The dielectric constant real part  $(\varepsilon_1)$  is determined by using the equation (El-Denglawey [2011;](#page-18-15) Zaki and El-Amin [2017](#page-20-25));

$$
\varepsilon_1 = n^2 - k^2 = \varepsilon_L - \frac{e^2 N}{4\pi^2 \epsilon_0 c^2 m^*} \lambda^2
$$
 (3a)

where  $e$  is the electronic charge,  $N$  is concentration of the free carriers,  $m^*$  is effective mass,  $\epsilon_o$  is permittivity of free space, *c* is light speed and  $\epsilon_L$  is residual lattice dielectric constant which represents high frequency component of the relative permittivity at  $\lambda = 0$ . *n* is refractive index given from the values of refectance *R* as follows;

$$
n = \frac{1+R}{1-R} + \left(\frac{4R}{(1-R)^2} - K^2\right)^{\frac{1}{2}}
$$
(3b)

Figure [6](#page-12-0)a, b shows the behavior of  $(n^2 - k^2)$  against  $\lambda^2$  at different annealing temperature for the samples. It is observed that this behavior is almost linear at the lower photon energy, while it loses its linearity at the higher energy.  $N/m^*$  and  $\varepsilon_L$  are calculated from the slope of linear parts and extrapolation of the plot to  $\lambda^2 = 0$ , respectively. The behaviors of  $(N/m^*)$  and  $\varepsilon_L$  against annealing temperature for ZnO and AlZnO are shown in Fig. [7](#page-13-0)a, b, and similar values are listed in Table [2.](#page-10-0) It is clear that the values of  $\varepsilon_L$  are increased as annealing temperature increases up to 900 °C for both series. While (*N/m*\* ) are increased for ZnO, but it is decreased for AlZnO. In addition, one can see that the value of  $\varepsilon_L$  ( $\varepsilon_L$ =11.26) for the asprepared ZnO, which is nearly similar that  $(\varepsilon_L = 11.76)$  for AlZnO at 900 °C. Furthermore, there is a dramatic decrease in the values of (*N/m*\* ) above 400 °C for AlZnO as compared to the increase for ZnO. By looking to the presnt results, one can see that the optical energy gap,  $E_{oh}$ , and  $\varepsilon_L$  seemed to be changed in the same manner for both series. The increase of free carrier concentration up to 900 °C for ZnO can be achieved to the decrease of oxygen deficiency or electrons. While the decrease of free carrier concentration for AlZnO can be attributed to the increase of holes carriers as a result of substitution of  $\text{Zn}^{2+}$  by  $\text{Al}^{3+}(\text{Ali})$ and Hakeem [2015](#page-18-13)). For more clarifcation, free carrier concentration *N* is calculated using the average value of  $m^*$  which has been reported for ZnO ( $m^*$  = 3.82 × 10<sup>-25</sup> g) (El-Desoky et al.  $2014$ ), and therefore the inter-atomic distance  $(R)$  can be calculated by using the relation;  $R = (1/N^{(1/3)})$  (El-Desoky [2005;](#page-19-22) Wemple [1973\)](#page-20-26).

The obtained values of *N* and *R* are listed in Table [2.](#page-10-0) It is clear that *N* is increased by annealing up to 900 °C for ZnO, while it is decreased for AlZnO. The vice is versa for inter-atomic distance *R*.



<span id="page-12-0"></span>**Fig. 6 a** The relation between  $(n^2 - k^2)$  and  $\lambda^2$  of annealed ZnO samples. **b**  $(n^2 - k^2)$  versus  $\lambda^2$  of annealed AlZnO samples

The dielectric constant  $\varepsilon$  is very important optical parameter because it is usually provide information for the electronic structure of materials and helps for its design as optoelectronic devices. The dielectric constant is written as function of the imaginary parts  $\varepsilon_1$  and  $\varepsilon_2$  as follows (Swanepoel [1984;](#page-20-27) Li et al. [2010](#page-19-23));



(b)  $(N/m^*)$  against annealing temperature of ZnO and AlZnO samples

<span id="page-13-0"></span>**Fig. 7 a** Lattice dielectric versus annealing temperature of ZnO and AlZnO samples. **b** (N/m\*) against annealing temperature of ZnO and AlZnO samples

$$
\varepsilon = \varepsilon_1 - i\varepsilon_2
$$
  
\n
$$
\tan \delta = \frac{\varepsilon_2}{\varepsilon_1} = \frac{2nk}{n^2 - k^2}
$$
\n(4)

where  $tan \delta$  is the dielectric loss factor generated from the dielectric constant components  $\varepsilon_2$  and  $\varepsilon_1$ . These components are given in terms of extinction coefficient (*k*) and refractive index *n*. *k* is responsible for attenuation of light and  $\varepsilon_1$  is responsible for attenuation





(b) Dielectric loss versus hv of annealed AlZnO samples

<span id="page-14-0"></span>**Fig. 8 a** The relation between Dielectric loss and hν of annealed ZnO samples. **b** Dielectric loss versus hν of annealed AlZnO samples

of electrical feld causing dielectric loss. The relation between *tan δ* and *hν* is shown in Fig. [8](#page-14-0)a, b. It is found that  $tan \delta$  is almost constant and increases sharply as the photon energy increases. While, it is gradually decreased by annealing up to 900 °C for ZnO, but it is increased for AlZnO up to 700 °C, followed by decrease at 900 °C. The drop of *tanδ* at 900 °C for AlZnO can be attributed to the increase in polarization probability occurred by increasing temperature, in agreement with the reported data (Dongol et al. [2006](#page-18-16), [2007](#page-18-17)).

The optical response is mainly obtained for any material in terms of optical conductivity  $(\sigma_{opt})$  which is determined using the following formula (Ali and Abdel-Rahim [2013;](#page-18-18) Girisun and Dhanuskodi [2009](#page-19-24));

$$
\sigma_{opt} = \frac{\alpha n c}{4\pi} \tag{5a}
$$

Also, the electrical conductivity  $(\sigma_{e|e})$  due optical response of the material is given by Ali and Abdel-Rahim [\(2013](#page-18-18)) and Girisun and Dhanuskodi ([2009\)](#page-19-24)

$$
\sigma_{ele} = \left(\frac{2\lambda}{\alpha}\right) \sigma_{opt} \tag{5b}
$$

The behavior of  $\sigma_{opt}$  against photon energy and at different values of annealing tempera-ture for ZnO and AlZnO is shown in Fig. [9a](#page-16-0), b. It is noticed that  $\sigma_{\text{out}}$  decreases by annealing up to 900 °C for ZnO, while it increases up to 700 °C, followed by a decrease at 900 °C for AlZnO. This behavior probably confrms semi-metallic nature up to 700 °C, followed by semiconducting nature at 900 °C. Similar behavior is obtained for the behaviors of *σele* against annealing temperature for ZnO and AlZnO as shown in Figs. [10a](#page-17-0), b. However, the higher magnitude of  $\sigma_{opt}$  as the photon energy increases assures the high photo-response of the annealed samples. This behavior is due to the increase of electron excitation by photon energy and annealing temperature (Fouad and El-Radaf [2018](#page-19-25); El-Bana and Fouad [2017](#page-18-19)). It can be also noticed that  $\sigma_{ele}$  increases as the photon energy increases, which confirms the semiconducting nature of the samples (Fouad et al. [2016\)](#page-19-26). Furthermore, the values of  $\sigma_{opt}$ and *σele* for ZnO are slightly higher than that of AlZnO, in agreement with the structural analysis.

Anyhow, It has been reported that Al reacted with  $ZnO$  and formed  $ZnAl_2O_4$  on the grain boundaries, and small amount of  $ZnAl_2O_4 > 1\%$  depressed the grain growth of ZnO, and can signifcantly decreased its dielectric constant (Li et al. [2016\)](#page-19-27). Furthermore, it is obtained from our previous work that breakdown feld decreased by 0.1 Al, while electrical conductivity is improved. However, for ZnO doped more than  $10\%$  Al,  $Al^{3+}$  ions tend to move to the interstitial positions of Zn sites. At such interstitial position, the following reaction can be applied;  $\frac{1}{2}A l_2 O_3 \rightarrow \frac{3}{4}O_2 + Al^* + P$ , where an electron is absorbed (Matsuoka et al.  $1970$ ; Houabes et al.  $2005$ ; Gupta [1992\)](#page-19-28). Al<sup>\*</sup> is the ionized Al atom in interstitial sites and P is positive hole charge. In this case, the Al acts as an acceptor which dominates the donor efect, resulting in a decrease of the electrical conductivity of ZnO. Furthermore, the structural distortion becomes stronger than relaxation time  $\tau$  and consequently the conductivity is decreased ( $\sigma_{300}$   $\alpha$  *nt*), n is carrier concentration. These could be the reasons for the variation of the optical properties of AlZnO samples as compared to ZnO ones.

### **4 Conclusion**

Efect of Annealing temperature on structural and optical investigation of ZnO and  $Zn_0A_0A_1O$  is investigated. The cell parameters and average crystalline diameters are affected by both Al and annealing temperature. The ultra violet  $(UV)$  absorption coefficient  $(a)$  is increased by Al, but it is decreased by annealing, and can be represented by



<span id="page-16-0"></span>**Fig. 9** Optical conductivity versus hν of annealed **a** ZnO samples. **b** AlZnO samples

the relation;(*α* (AlZnO)~2 *α* (ZnO) between the two series. Furthermore,  $E_{eL}$  (ZnO),  $E_{eh}$ (AlZnO) and  $E_{gL}$  (AlZnO) are increased by annealing up to 700 °C, followed by a decrease at 900 °C, while *Egh* (ZnO) is increased up to 900 °C. Moreover, a unique energy gap of 3.33 eV could be obtained at 312 °C for both series. On the other hand,  $\varepsilon_L$  is increased by annealing up to 900 °C for both series, while (*N/m\** ) is increased for ZnO and decreased for AlZnO. But their values are higher for ZnO than that of AlZnO. Interestingly, the value of  $ε_L$  (11.26) of the as-prepared ZnO is equal to the value of  $ε_L$  (11.76) for AlZnO at 900 °C.



<span id="page-17-0"></span>**Fig. 10** Electrical conductivity versus hν of annealed **a** ZnO samples, **b** AlZnO samples

Although *tan*  $\delta$  is similar for both series, it is decreased by annealing up to 900 °C for ZnO, and increased for AlZnO up to 700  $\degree$ C, followed by decrease at 900  $\degree$ C. Similar behavior is obtained for  $\sigma_{opt}$  and the vice is versa for  $\sigma_{ele}$ . The obtained results are interpreted in terms of the variance of carrier concentration, oxygen vacancies and interatomic spacing against annealing temperature and Al content. Moreover, they confrm the semiconducting nature as well as assure that the high photo-response of the investigated samples.

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