

# **Photonic crystal fber for robust orbital angular momentum transmission: design and investigation**

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### **Abstract**

In this paper, a novel design, photonic crystal fber (PCF) has been reported to support more orbital angular momentum (OAM) modes for the frst time. The cladding region formed by employing the three-layer chain shaped air cavities. The OAM properties and parameters are calculated by optimizing the design of the PCF. The simulation results can support 26 OAM modes with a wider bandwidth (750 nm). The confnement loss of the most OAM modes is below average  $10^{-7}$  dB/m and the nonlinear coefficient is less than 4 W<sup>-1</sup>/km. Some OAM modes with comparatively flat dispersion variation (such as 3.8684 ps/km-nm) for  $HE_{2,1}$  are noticed from the design. Moreover, fundamental optical characteristics have been rigorously computed by utilizing the fnite element method. Now, it can be anticipated that, these excellent optical characteristics confrm the introduced design as a prominent candidate for the OAM transmission and other relevant areas of optical communications.

**Keywords** Orbital angular momentum (OAM) · Photonic crystal fber (PCF) · Mode division multiplexing (MDM) · Space division multiplexing (SDM)

# **1 Introduction**

The orbital angular momentum (OAM) is the recently used consolidated multiplexing process concerning light for securing information. It is also a space division multiplexing (SDM) technology that is used in free space for optical communication (Huang et al. [2014\)](#page-12-0). For the optical transmissions, the OAM introduced with vacuum communication (Brunet et al. [2014\)](#page-12-1). Nowadays, optical communication confronts diferent types of

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challenges such as the transmission of big data, telecommunication, broadcast, community antenna television coaxial cable (CATV), high-defnition television (HDTV) for medium blocking and insufficient capacity of data transmission. Recently, for multimode fiber, the SDM technology provides the OAM method that enhances the capacity and efficiency of optical communication (Zhang et al. [2018;](#page-13-0) Richardson et al. [2013](#page-13-1); Zhu et al. [2016](#page-13-2); Li et al. [2018\)](#page-12-2).

The beams of OAM individualized by  $e^{(i\omega)}$  (Chen et al. [2016](#page-12-3)), where the topological charge is expressed by *l* and it is an integer number  $(..., -3, -2, -1, 0, +1, +2, +3,...)$ and azimuthal angle is expressed by  $\varphi$ . The azimuthal angle  $(\varphi)$  of a point in cylindrical or spherical coordinates is the positive angle between the projection of the vector to the XY-plane and positive X-axis (Zhu et al. [2016;](#page-13-2) Allen et al. [1992](#page-12-4)). Every positive (+) and negative (−) topological charges are spatially orthogonal to each and another (Jiao et al. [2019\)](#page-12-5). There are several methods to achieve an OAM beam like twisted fber (Wong et al. [2012\)](#page-13-3), helical lattice (Fang and Wang [2015](#page-12-6)), etc. There are various applications such as optical communication (Wang et al. [2012](#page-13-4)), particle trapping (Padgett and Bowman [2011](#page-12-7)), high-resolution imaging (Tamburini et al. [2006](#page-13-5)) and many more (Heng et al. [2017](#page-12-8)) of the OAM beams.

For improving the number of modes, Liu et al. [\(2019a\)](#page-12-9) designed a dual-cladding photonic quasi-crystal fber with 6 OAM modes for the OAM optical communication system. The quasi-periodic structure used in the photonic quasi-crystal fiber (Liu et al. [2019b](#page-12-10), [c](#page-12-11); Li et al. [2019](#page-12-12)) such as sensing (Shi et al. [2019;](#page-13-6) Liu et al. [2019d](#page-12-13)), flters (Liu et al. [2018](#page-12-14); Yan et al. [2018](#page-13-7)), nanocavity (Ren et al. [2018](#page-13-8)), metasurface (Tang et al. [2019](#page-13-9)), lens (Liu and Fan [2018\)](#page-12-15), etc. Zhang et al. ([2015\)](#page-13-10) proposed a C-PCF with silica as the background material supporting 10 OAM modes to transmit in one fber. A hollow-core fber was proposed by Li et al. ([2016\)](#page-12-16) that supports 16 OAM modes with high confnement loss. A C-PCF pro-posed by Hu et al. ([2016\)](#page-12-17) that supports 26 OAM modes with  $As_2S_3$  as background material. Wang et al. ([2018\)](#page-13-11) is also designed a PCF with 26 OAM modes. Xu et al. [\(2017](#page-13-12)) proposed a C-PCF supporting 30 OAM modes with negative dispersion. The proposed PCF of Lei et al. ([2018\)](#page-12-18) supports more OAM modes (76 OAM modes). But that was not enough for stable transmission. So, there is still room to design a novel PCF with larger OAM modes based on PCF with better output.

The high number of OAM mode is good for the encoding of data transmission. The same number of the OAM mode is useful for detecting receiving information with respect to sending information. When the sender sends data to the receiver, then it has to fnd out the similarities between sending and receiving data. By checking the mode properties of two transmission sides, the receiver ensures require data which transmit by the sender. In the OAM modulation process, the sender sends a carrier signal with the message signal; where the carrier signal carries the OAM mode properties. So, the OAM mode plays an important act in optical communication. The traditional optical fber fails to transmit the OAM mode thus we have to design a special PCF for supporting multimode and also to achieve good parameters of a fber. The diference between the two adjacent efective refractive indexes of the OAM *HE* and *EH* modes must be greater than  $10^{-4}$  ( $\Delta n_{eff}$ > $10^{-4}$ ) (Zhang et al. [2017a](#page-13-13)). The OAM transmission is supported by many unique PCF like spiral PCF (Nandam and Shin [2018](#page-12-19)), C-PCF with a square hole (Bai et al. [2018](#page-12-20)), dual guided ring-core PCF (Xu et al. [2018](#page-13-14)), C-PCF (Zhang et al. [2018\)](#page-13-0). These are separately good for confnement loss, dispersion, nonlinearity, and numerical aperture, but not better for the robust transmission.

This paper proposed a chain shaped photonic crystal fber that is composed of the three well-structured chain type air holes ring with a large air hole and the background material

is composed of fused silica. This structure of PCF is novel for OAM transmission. This novel design of PCF supports 26 modes with 750 nm bandwidth. This PCF supports a low dispersion variation of 3.8684 ps/km-nm for the  $HE_{2,1}$  mode, and 5.1421 ps/km-nm for the  $TE_{0.1}$  mode with low confinement loss that is  $3.19 \times 10^{-10}$  dB/m at wavelength 1.80 µm.

### **2 Model and algorithm**

As mention before, the proposed three-layer chain shaped air holes with a hollow-core C-PCF is designed and simulated by COMSOL Multiphysics 4.3b software. The hollow-core PCF is the part of PCFs that guide light through the core (Monfared [2018\)](#page-12-21). The three wellpatterned circular ring type chain shaped air holes would decrease the light outpouring of optical fber. The radius of the core of the proposed model is 3.7 μm. The large core increases the number of OAM modes (Zhang et al. [2017a](#page-13-13)). The air holes are made by the Bezier Polygons. Seven Cartesian points are needed to build this type of curve. The width of each polygon is 1.5  $\mu$ m ( $w_2 = w_3 = w_4 = 1.5 \mu$ m) and each co-ordinate has two levels that are totally eight levels. The area between the outer circle and the inner circle is called PML (Perfectly Matched Layer) and the width between them  $(w_1)$  is also 1.5  $\mu$ m. It is added here to truncate unwanted nonphysical scattering. All the numerical analysis are implemented by using the PML and the FEM. And the background material is fused silica with a refractive index of 1.446 (Monfared et al. [2013](#page-12-22)). The refractive index (RI) is calculated by the Sellmeier Eq. [\(1\)](#page-2-0) (Biswas et al. [2019\)](#page-12-23).

<span id="page-2-1"></span><span id="page-2-0"></span>
$$
n(\lambda) = \sqrt{1 + \sum_{i=1}^{m} \frac{A_i \lambda^2}{\lambda^2 - \lambda_i^2}}
$$
 (1)

where  $m, \lambda_i^2$  and  $A_i$  are the Sellmeier coefficients and  $n(\lambda)$  represents the RI that varies with wavelength. The cladding of the proposed PCF is constructed by fused silica. For fused silica,  $m=3$ . The complete design of the PCF is shown in Fig. [1](#page-3-0).

The OAM modes in the PCF can be expressed by the combinations of even and odd modes of *EH* and *HE*. It can be defned by Eqs. [\(2\)](#page-2-1) and ([3](#page-2-2)) (Pakarzadeh and Sharif [2019;](#page-12-24) Zhang et al. [2017b\)](#page-13-15).

$$
\begin{cases}\nOAM_{\pm l,m}^{\pm} = HE_{l+1,m}^{even} \pm jHE_{l+1,m}^{odd} \\
OAM_{\pm l,m}^{\mp} = EH_{l-1,m}^{even} \pm jEH_{l-1,m}^{odd}\n\end{cases}; (l > 1)
$$
\n(2)

<span id="page-2-2"></span>
$$
\begin{cases}\nOAM_{\pm 1,m}^{\pm} = HE_{2,m}^{even} \pm jHE_{2,m}^{odd} \\
OAM_{\pm 1,m}^{\mp} = TM_{0,m} \pm jTE_{0,m}\n\end{cases}; (l = 1)
$$
\n(3)

where HE and EH are the composite mode of the OAM beam. The two main types of each mode are even and odd that obtains after  $\pi/2$  phase difference of the modes. The sign of " $\pm$ l" defnes the right or left wave front orbit direction. Where l is a topological charge, j denotes a  $\pi/2$  phase shift, and m denotes the radial order for the OAM mode, and equal to 1. The OAM mode does not support m > 1 because it causes "accidental degeneracies" (Pakarzadeh and Sharif [2019](#page-12-24)). Because it is a problem for the encoding, multiplexing and demultiplexing of the OAM modes. So, we use a lower-order radial mode  $(m=1)$  (Zhang et al. [2018](#page-13-0)).



<span id="page-3-0"></span>**Fig. 1** The cross-section view of the proposed PCF

By applying the Eqs. [\(2](#page-2-1)) and ([3](#page-2-2)), the proposed structure is supported 26 modes, these are  $OAM_{\pm1,1}^{\pm}$  { $HE_{2,1}$ },  $OAM_{\pm2,1}^{\pm}$  { $HE_{3,1}$ ,  $EH_{1,1}$ },  $OAM_{\pm3,1}^{\pm}$  { $HE_{4,1}$ ,  $EH_{2,1}$ },  $OAM_{\pm4,1}^{\pm}$  { $HE_{5,1}$ ,  $EH_{3,1}$ }  $A \rightarrow \text{L}$ <br>  $DAM^{\pm}_{\pm 5,1} \{ HE_{6,1}, EH_{4,1} \}$ ,  $OAM^{\pm}_{\pm 6,1} \{ HE_{7,1}, EH_{5,1} \}$ ,  $OAM^{\pm}_{\pm 7,1} \{ HE_{8,1}, EH_{6,1} \}$ .

The proposed PCF is supported by all the selected modes from the wavelength the ranges from 1.2 to 1.95 μm and simulations are done by those modes. It also supports some eigenmodes such as  $TE_{0,1}$  and  $TM_{0,1}$ , but they can't combine. Because the propagation constants of  $TE_{0,1}$  and  $TM_{0,1}$  modes are different.

The quality of the OAM based PCF depends on more numbers of the OAM modes. Because, it increases the capacity of transmission and also ensures the robust transmission. The encoding and multiplexing of lights are developed by the good quality of OAM mode. Therefore, the light intensity imbricate factor to qualitatively calculate the mode quality through the Eq.  $(4)$  (Zhang et al.  $2017b$ ).

<span id="page-3-1"></span>
$$
\eta = \frac{I_r}{I_c} = \frac{\iint_{rings} \left| \vec{E} \right|^2 dxdy}{\iint_{cross-section} \left| \vec{E} \right|^2 dxdy}
$$
(4)

where *I<sub>r</sub>* defines the normal mode concentration of the ring area where all the OAM modes will be confined and *I<sub>c</sub>* denotes the average mode intensity in full exploded view area of the C-PCF. In Fig. [2](#page-4-0), diferent OAM mode intensity order of eigenmodes and the wave front phase are exhibited.

For all modes, proposed PCF gains fat dispersion, less nonlinearity, better numerical aperture, less confnement loss and other OAM properties.

Dispersion is evaluated by Eq. ([5](#page-4-1)) (Zhang et al. [2018\)](#page-13-0).



<span id="page-4-0"></span>**Fig. 2 a** The intensity ordinations of  $EH_{1,1}$ ,  $EH_{2,1}$ ,  $HE_{3,1}$ ,  $HE_{4,1}$ ,  $EH_{5,1}$ ,  $EH_{6,1}$ ,  $HE_{7,1}$ ,  $HE_{8,1}$ ,  $TE_{0,1}$  and  $TM_{0,1}$ modes, **b** the phase ordinations of  $EH_{2,1}HE_{3,1}$ ,  $EH_{4,1}HE_{5,1}$  and  $HE_{8,1}$  modes

<span id="page-4-2"></span><span id="page-4-1"></span>
$$
D = -\frac{\lambda}{c} \frac{d^2 Re[n_{\text{eff}}]}{d\lambda^2} \tag{5}
$$

where  $Re[n_{\text{eff}}]$  denotes the real values of the effective refractive index and c defines the velocity of light in the vacuum.

The diferences of the refractive index can be evaluated through Eq. [\(6\)](#page-4-2) (Pakarzadeh and Sharif [2019](#page-12-24)). This diference is performed on two adjacent vector mode.

$$
\Delta n_{\text{eff}} = \left| n_{\text{eff}_{\text{HE}_{l+1,m}}} - n_{\text{eff}_{\text{EH}_{l-1,m}}} \right| > 10^{-4} \tag{6}
$$

So,  $\Delta n_{\text{eff}}$  measures, the absolute value within the *HE* and *EH* OAM modes. It must be larger than  $10^{-4}$ .

The effective OAM mode area is symbolized by  $A_{\text{eff}}$  and can be evaluated through Eq. [\(7\)](#page-5-0) (Jia et al. [2018](#page-12-25)).

$$
A_{\text{eff}} = \frac{\left(\iint |E(x, y)|^2 dxdy\right)^2}{\iint |E(x, y)|^4 dxdy} \tag{7}
$$

where  $E(x, y)$  is an electric field allocation of the electromagnetic (EM) transverse field.

Numerical aperture (NA) can be calculated by Eq. ([8](#page-5-1)) (Biswas et al. [2019](#page-12-23)).

<span id="page-5-0"></span>
$$
NA = \left[1 + \frac{\pi A_{\text{eff}}}{\lambda^2}\right]^{-\frac{1}{2}}
$$
 (8)

The nonlinear property of PCF is expressed through Eq. ([9\)](#page-5-2) and it is denoted by the  $\gamma$ (Jia et al. [2018](#page-12-25)).

<span id="page-5-3"></span><span id="page-5-2"></span><span id="page-5-1"></span>
$$
\gamma = \frac{2\pi n_2}{\lambda A_{\text{eff}}} \tag{9}
$$

where  $n_2$  defines the nonlinear coefficient of fused silica and  $n_2 = 2.6 \times 10^{-20}$  m<sup>2</sup>/W.

Confinement loss  $(C_{loss})$  can be evaluated by Eq. ([10](#page-5-3)) (Xu et al. [2017](#page-13-12)).

$$
C_{loss} = \frac{40\pi}{\ln(10)\lambda} Im(n_{\text{eff}})(dB/m)
$$
 (10)

where  $Im(n_{\text{eff}})$  is the unreal part of the refractive index of the OAM transmission mode and *λ* is the wavelength of the OAM mode.

## **3 Results and discussion**

The proposed PCF supports 26 OAM modes with low confnement loss, fat dispersion, better nonlinearity, low numerical aperture and the refractive index diference greater than 10<sup>−</sup><sup>4</sup> between two adjacent OAM modes. Now, here is described the OAM properties such as an efective refractive index, dispersion, efective refractive index diference, efective mode area, numerical aperture, nonlinearity, and confnement loss with appropriate discussion, equation, and the necessary fgure.

#### **3.1 The efective refractive index of eigenmodes**

The refractive index of fused silica is calculated by Eq. [\(1\)](#page-2-0). Higher modes are obtained by decreasing the search around of the refractive index at a fxed wavelength.

Figure [3](#page-6-0) shows the relationship between all selected OAM modes efective refractive index  $(n_{\text{eff}})$  and wavelength within the range from 1.2 to 1.95  $\mu$ m in the proposed PCF. The changing rate of the refractive index is downward with respect to the increment of wavelength. The efective refractive index is related to the dispersion. The fat dispersion is obtained for the small diferences of the refractive index at any specifc mode.



<span id="page-6-0"></span>Fig. 3 The effective refractive index of different eigenmodes with respect to the wavelength

#### **3.2 Dispersion**

The dispersion parameter is one of the most important parameters of PCF. Dispersion is the expansion of the light pulse at transmitting time. It decreases the performance of an optical fber. It contains material and waveguide dispersion (Zhang et al. [2018\)](#page-13-0). Material dispersion has little efect on the total dispersion. So, the calculation is only for the waveguide dispersion.

The dispersion is inversely proportional to the wavelength. The  $TM_{0,1}$ ,  $TE_{0,1}$ ,  $EH_{1,1}$ ,  $HE_{1,1}$ ,  $EH_{2,1}$ ,  $HE_{2,1}$  and others OAM modes are exhibited in Fig. [4.](#page-6-1)



<span id="page-6-1"></span>**Fig. 4** Dispersion of diferent OAM modes with respect to the wavelength

The dispersion of  $TE_{0,1}$ ,  $EH_{1,1}$ ,  $HE_{1,1}$ ,  $EH_{2,1}$ ,  $HE_{2,1}$ , and  $TM_{0,1}$  modes are nearly horizontal while other OAM modes show flat dispersion. The high dispersion is shown in  $HE_{61}$ and  $EH_{61}$  modes of the OAM. The proposed PCF shows the dispersion variation of different OAM modes. For example, the dispersion variation for the  $TE_{0.1}$  mode is 5.1421 ps/ km-nm, the  $HE_{2,1}$  mode is 3.8684 ps/km-nm, the  $HE_{3,1}$  mode is 18.0606 ps/km-nm, the  $HE_{4,1}$  mode is 39.5561 ps/km-nm, the  $HE_{6,1}$  mode is 64.8221 ps/km-nm, the  $HE_{7,1}$  mode is 47.2286 ps/km-nm, the  $EH_{2,1}$  mode is 3.8566 ps/km-nm, the  $EH_{3,1}$  mode is 11.6646 ps/ km-nm, the  $EH_{4,1}$  mode is 26.7569 ps/km-nm, the  $EH_{5,1}$  mode is 43.4371 ps/km-nm, the  $EH_{61}$  mode is 81.6529 ps/km-nm, the  $EH_{71}$  mode is 56.7001 ps/km-nm within the wavelength from 1.2 to 1.85  $\mu$ m.

Especially, for the  $TE_{0,1}$  and  $HE_{2,1}$  mode, the variation of dispersion is much better than indicated values (Nandam and Shin [2018;](#page-12-19) Zhang et al. [2016](#page-13-16); Yue et al. [2012\)](#page-13-17).

The dispersion of the *HE* mode is larger than the *EH* mode. Because, the efective index of the *HE* mode is smaller than the *EH* mode in the same order of the OAM modes (Bai et al. [2018\)](#page-12-20). In Fig. [4,](#page-6-1) the dispersion of  $HE_{4,1}$  and the  $EH_{4,1}$  mode is 142.6587 ps/km-nm and 121.1491 ps/km-nm respectively. So, the higher orders of the OAM modes have a higher dispersion.

The dispersion of the proposed PCF is improved by modifying the basic element of the PCF. The improvement of dispersion is also dependent on the small refractive index diferences.

#### **3.3 The efective refractive index diference of OAM modes**

Figure [5](#page-7-0) shows the relationship between the  $\Delta n_{\text{eff}}$  and wavelength. For the disparity, in Fig. [5,](#page-7-0) a blue straight line is used to indicate for the 10<sup>−</sup>4. If diferences of OAM modes (*HE* and *EH*) are not greater than  $10^{-4}$ , then this mode will be coupled with one another, and also command to linearly polarized (LP). The  $\Delta n_{\text{eff}}$  *is* increased by increasing the wavelength. The higher order of the OAM modes will be highly effective refractive index differences ( $\Delta n_{eff}$ ). The highest  $\Delta n_{eff}$  is 4.4302 × 10<sup>-2</sup> at 1950 nm wavelength for  $|HE_{8,1}-EH_{6,1}|$ .



<span id="page-7-0"></span>**Fig. 5** Efective refractive index diferences of diferent OAM modes with respect to the wavelength

Therefore, the PCF has the benefts of decrement of the OAM mode coupling as well as ensuring the robust transmission of OAM modes.

#### **3.4 The efective mode area**

The  $A_{\text{eff}}$  of the PCF has a significant response to the nonlinearity parameter. The effective area is increased by increasing with the wavelength. The effective area of different OAM modes is shown in Fig. [6.](#page-8-0)

Here,  $TM_{0,1}$  and  $TE_{0,1}$  achieves the lowest and the highest effective mode area respectively. The indicated part of Fig. [6](#page-8-0) is the zoom view of the dashed shaped area that clarifes the efective mode area of diferent OAM modes. Low numerical aperture and nonlinearity are obtained for the fat efective area.

#### **3.5 Numerical aperture (NA)**

NA is acknowledged by the total amount of optical power. The higher values of NA have various applications in the optical feld. Besides, the low values of NA provide depth of feld (DOF).

From the Eq. ([8\)](#page-5-1), NA is dependent on the efective area and the wavelength. The NA is a unit less parameter. It is increased by increasing the wavelength as well as decreasing the efective area.

The indicated part of Fig. [7](#page-9-0) is the zoom view of the dashed shaped area that clarifes the numerical aperture of diferent OAM modes.

The low numerical aperture of the proposed PCF is obtained for the random variations of the refractive index.



<span id="page-8-0"></span>Fig. 6 The effective area of different OAM modes with respect to the wavelength



<span id="page-9-0"></span>**Fig. 7** The numerical aperture of diferent OAM modes with respect to the wavelength

#### **3.6 Nonlinearity**

Figure [8](#page-9-1) displays the relationship between the nonlinearity and wavelength of all selected OAM modes. Nonlinearity is inversely proportional to the effective mode area and wavelength. In the proposed PCF, all selected OAM modes are low nonlinearity within the range from 1.2 to 1.95 μm. Especially, the nonlinearity of most OAM modes are below 4  $\text{W}^{-1}/\text{km}$  and particularly, the lowest nonlinearity of  $TE_{0.1}$  OAM mode is 1.58 W<sup>-1</sup>/km at the wavelength 1.8  $\mu$ m.



<span id="page-9-1"></span>**Fig. 8** Nonlinearity of diferent OAM modes with respect to the wavelength

The indicated part of Fig. [8](#page-9-1) is the zoom view of the dashed shaped area that clarifes the nonlinearity of diferent OAM modes.

The proposed PCF is achieved lower nonlinearity. Because, lower nonlinearity is obtained for the large core (Hu et al. [2016\)](#page-12-17).

#### **3.7 Confnement loss (CL)**

The confnement loss is caused by structural imperfection, micro bending, and intrinsic material absorption characteristics when light is transmitted over the PCF (Shi et al. [2019](#page-13-6)). The CL is determined for the different modes of the OAM. The CL for the  $TM_{0,1}$ ,  $TE_{0,1}$ ,  $EH_{1,1}$ ,  $HE_{1,1}$ ,  $EH_{2,1}$ ,  $HE_{2,1}$ , and other modes are represented in the Fig. [9.](#page-10-0)

In the scenario of Fig. [9](#page-10-0), the inconstancy of CL of all diferent OAM modes with respect to the wavelength does not display signifcant regularity. It is mainly adequate to the high probability of leakage to cladding for the larger wavelength. The indicated part of Fig. [9](#page-10-0) is the zoom view of the dashed shaped area that clarifes the CL of diferent OAM modes from the wavelength  $1.5$  to  $1.6 \mu m$  of step size 1 nm.

In this proposed PCF design, the average CL of the most OAM modes is varied between  $4 \times 10^{-10}$  and  $5 \times 10^{-6}$  dB/m with the wavelength from 1.2 to 1.95 µm. Such as, The *HE*<sub>41</sub> mode has the lowest confinement loss that is numerically  $3.19 \times 10^{-10}$  dB/m at wavelength 1.8  $\mu$ m. In Fig. [9](#page-10-0), the confinement loss of the  $H\overline{E}_{1,1}$  mode is  $1.01 \times 10^{-9}$  dB/m, the *EH*<sub>1,1</sub> mode is  $6.62 \times 10^{-9}$  dB/m, the *HE*<sub>2,1</sub> mode is  $5 \times 10^{-8}$  dB/m, the *EH*<sub>2,1</sub> mode is  $3.46\times10^{-7}$  dB/m, the  $EH_{3,1}$  mode is  $5.28\times10^{-7}$  dB/m, the  $HE_{4,1}$  mode is  $3.22\times10^{-7}$  dB/m, the  $EH_{4,1}$  mode is  $4.57 \times 10^{-8}$  dB/m, the  $EH_{6,1}$  mode is  $3.39 \times 10^{-7}$  dB/m, the  $HE_{7,1}$ mode is  $7.86 \times 10^{-8}$  dB/m, the  $EH_{71}$  mode is  $6.15 \times 10^{-9}$  dB/m and the  $TE_{01}$  mode is  $1.5 \times 10^{-7}$  dB/m at the wavelength of 1.45 µm.

The CL is destructive for the optical communication system. The CL is generated for structural imperfection and internal absorption. The CL of the proposed PCF is decreased



<span id="page-10-0"></span>**Fig. 9** Confnement loss of diferent OAM modes with respect to the wavelength

References (prior $PCFs$ )	Operating wave- length for CL (nm)	Dispersion variation of $TE_{0.1}$ (ps/km-nm)	Dispersion variation of $HE_{21}$ (ps/km-nm)	Confinement loss (dB/m)
Nandam and Shin (2018)	1550	10.35	12.1	$5.0951 \times 10^{-3}$
Zhang et al. $(2017b)$	1550			$9.52 \times 10^{-9}$
Zhang et al. $(2016)$	1550	<46.38		$3.434 \times 10^{-9}$
Yue et al. (2012)	>2000	<60		$3 \times 10^{-2}$
This work	1800	5.1421	3.8684	$3.19 \times 10^{-10}$

<span id="page-11-0"></span>**Table 1** The comparison chart of dispersion variations for TE, HE modes and confnement loss of the proposed article and the recent published articles

for periodic cladding. For all OAM modes confnement loss does not show the regularity. It is caused for a greater chance of leakage to the cladding (fused silica) for higher wavelength.

The comparison of confnement loss and dispersion variation of OAM modes with some previous article of OAM based PCF is shown in Table [1.](#page-11-0)

By comparing with some OAM modes PCF parameters that are previously reported and it can be investigated that our proposed PCF gives much better optical characteristics than other OAM based PCF articles.

# **4 Conclusion**

In this article, a novel chain-shape PCF design is proposed with low confnement loss, fat dispersion, better nonlinearity and high refractive index diferences between the wide bandwidth the range of 750 nm (from 1200 to 1950 nm). The proposed fber parameters are numerically calculated by applying FEM and PML as a boundary condition. The design can support up to 26 OAM modes. The numerical exploration shows that the efective refractive index diferences of all selected OAM modes are above than 10<sup>−</sup>4. The maximum  $\Delta n_{\text{eff}}$  is 4.4302×10<sup>-2</sup> at 1950 nm wavelength for HE<sub>8.1</sub> to EH<sub>6.1</sub> mode. The low confinement loss is 3.19×10<sup>-10</sup> dB/m at optical wavelength 1800 nm. The dispersion variation for the HE<sub>2,1</sub> mode is 3.8684 ps/km nm and 5.1421 ps/km nm for  $TE_{0.1}$  mode. So, the proposed PCF will be a strong candidate in stable and high capacity transmissions, encoding and MDM applications.

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# **Compliance with ethical standards**

**Confict of interest** All the authors have read the manuscript and approved this for submission as well as no competing interests.

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