



# On a generalized Broer-Kaup-Kupershmidt system for the long waves in shallow water

Xin-Yi Gao · Yong-Jiang Guo · Wen-Rui Shan

Received: 19 November 2022 / Accepted: 25 January 2023 / Published online: 27 February 2023  
© The Author(s), under exclusive licence to Springer Nature B.V. 2023

**Abstract** Describing the long waves in shallow water, a generalized Broer-Kaup-Kupershmidt system is investigated in this paper. With respect to the horizontal velocity of the water wave and the height of the water surface, we use symbolic computation to build up (A) a scaling transformation, (B) a set of the hetero-Bäcklund transformations, from that generalized system to a known linear partial differential equation, as well as (C) two sets of the similarity reductions, each of which from that generalized system to a known ordinary differential equation. Our results depend on all the shallow-water coefficients for that generalized system.

**Keywords** Shallow water · Long waves · Generalized Broer-Kaup-Kupershmidt system · Scaling transformation · Hetero-Bäcklund transformations · Similarity reductions · Symbolic computation

## 1 Introduction

Being attractive, many researchers are now paying attention to the shallow water waves [1–15]. For example, Ref. [1], a recent paper in *Nonlinear Dyn.*, has presented some multi-soliton solutions of a generalized Broer-Kaup system for the shallow water waves.

In this paper, we plan to investigate a generalized Broer-Kaup-Kupershmidt system, which describes, e.g., the long waves in shallow water, i.e.,

$$u_t = \left( -\alpha u_x + \frac{\beta}{2} u^2 + \beta v \right)_x + \gamma u_x, \quad (1a)$$

$$v_t = (\alpha v_x + \beta uv)_x + \gamma v_x, \quad (1b)$$

with the real differentiable functions  $v(x, t)$  and  $u(x, t)$  denoting, e.g., the horizontal velocity of the water wave and the height of the water surface, respectively,  $x$  and  $t$  implying, e.g., the scaled space and time variables,  $\alpha, \beta \neq 0$  and  $\gamma$  as the real constants, while the subscripts being the partial derivatives. Shallow-water special cases of System (1) have been seen:

- for the long waves in the shallow water, a Broer-Kaup-Kupershmidt system,

$$u_t - u_{xx} + 2v_x + 2uu_x = 0, \quad (2a)$$

$$v_t + v_{xx} + 2(uv)_x = 0, \quad (2b)$$

when  $\gamma = 0$ ,  $\beta = -2$  and  $\alpha = -1$ , with  $t$  and  $x$  being the scaled time and space variables,  $v(x, t)$

---

X.-Y. Gao (✉) · Y.-J. Guo (✉) · W.-R. Shan (✉)  
State Key Laboratory of Information Photonics and Optical Communications, and School of Science, Beijing University of Posts and Telecommunications, Beijing 100876, China  
e-mail: xin\_yi\_gao@163.com

Y.-J. Guo  
e-mail: yongerguo@bupt.edu.cn

W.-R. Shan  
e-mail: shwr@bupt.edu.cn

representing the horizontal velocity of the water wave while  $u(x, t)$  meaning the height of the water surface [15];

- for the diffusion-involved shallow water waves, a generalized Broer-Kaup system,

$$u_t = \left(-\alpha u_x + \frac{1}{2}u^2 + v\right)_x, \tag{3a}$$

$$v_t = (\alpha v_x + uv)_x, \tag{3b}$$

when  $\gamma = 0$  and  $\beta = 1$ , with  $t$  meaning the time variable,  $x$  indicating the propagation direction,  $v(x, t)$  relevant to both the wave profile and the tangential fluid velocity at the surface while  $u(x, t)$  representing the tangential fluid velocity at the surface [16–18];

- for the long waves in the shallow water, a classical dispersiveless long-wave system, i.e.,

$$u_t + uu_x + v_x = 0, \tag{4a}$$

$$v_t + (uv)_x = 0, \tag{4b}$$

when  $\alpha = \gamma = 0$  and  $\beta = -1$ , with  $u(x, t)$  meaning the tangential fluid velocity at the surface,  $v(x, t)$  representing the wave profile,  $x$  denoting the propagation direction, while  $t$  being the time variable [17, 18] (and references therein);

- for the long waves in the shallow water, a Broer-Kaup system, i.e.,

$$u_t + \frac{1}{2}u_{xx} - v_x - uu_x = 0, \tag{5a}$$

$$v_t - \frac{1}{2}v_{xx} - (uv)_x = 0, \tag{5b}$$

when  $\alpha = \frac{1}{2}$ ,  $\beta = 1$  and  $\gamma = 0$ , with  $u(x, t)$  standing for the scaled wave horizontal velocity, while  $v(x, t)$  related to the wave height and wave horizontal velocity [19, 20] (and references therein);

- for the dispersion water waves in the shallow water, a generalized Broer-Kaup system [21] (and references therein), i.e.,

$$u_t = -\alpha \left(u_x - u^2 - 2v\right)_x + \gamma u_x, \tag{6a}$$

$$v_t = \alpha (v_x + 2uv)_x + \gamma v_x, \tag{6b}$$

when  $\beta = 2\alpha$ ;

- for the dispersion water waves in the shallow water, a Broer-Kaup system [21] (and references therein), i.e.,

$$u_t = \left(-u_x + u^2 + 2v\right)_x, \tag{7a}$$

$$v_t = (v_x + 2uv)_x, \tag{7b}$$

when  $\alpha = 1$ ,  $\beta = 2$  and  $\gamma = 0$ .

However, to our knowledge, for System (1), there have been no scaling-transformation work, Bäcklund-transformation work and similarity-reduction work published as yet. Hereby, for System (1), we employ symbolic computation [22–28], to construct a scaling transformation, a set of the hetero-Bäcklund transformations and two sets of the similarity reductions.

### 2 Scaling and hetero-Bäcklund transformations for System (1)

Similar to those in Refs. [29, 30], we work out a scaling transformation:

$$\begin{aligned} \alpha &\rightarrow \rho^0 \alpha, & \beta &\rightarrow \rho^0 \beta, & \gamma &\rightarrow \rho^{-1} \gamma, & x &\rightarrow \rho^1 x, \\ t &\rightarrow \rho^2 t, & u &\rightarrow \rho^{-1} u, & v &\rightarrow \rho^{-2} v, \end{aligned} \tag{8}$$

in which  $\rho$  stands for a positive constant.

Next, on the score of Scaling Transformation (8), assuming that<sup>1</sup>

$$\begin{aligned} u(x, t) &= \zeta_1 w_x(x, t) + \zeta_2, \\ v(x, t) &= \zeta_3 w_{xx}(x, t), \end{aligned} \tag{9}$$

making use of symbolic computation, integrating Eq. (1a) once with respect to  $x$  with the integration function equal to zero and making a choice of

$$\beta \zeta_3 - \alpha \zeta_1 = \frac{\beta}{2} \zeta_1^2, \tag{10}$$

we obtain

$$Y_t(w) - \frac{\beta}{2} \zeta_1 Y_{2x}(w) - (\beta \zeta_2 + \gamma) Y_x(w) = 0, \tag{11}$$

with  $\zeta_1 \neq 0$ ,  $\zeta_2$  and  $\zeta_3 \neq 0$  as three real constants, the Bell polynomials reported by Refs. [31, 32] as

<sup>1</sup> similar to those in Refs. [29, 30]

$$\begin{aligned}
 Y_{m,n}(w) &\equiv Y_{m,n}(w_{0,0}, \dots, w_{0,n}, \dots, w_{m,0}, \dots, w_{m,n}) \\
 &= e^{-w} \partial_x^m \partial_t^n e^w,
 \end{aligned}$$

$w(x, t)$  meaning a  $C^\infty$  function of  $x$  and  $t$ ,  $w_{k,l} = \partial_x^k \partial_t^l w$ ,  $k = 0, \dots, m$ ,  $l = 0, \dots, n$ , while  $m$  and  $n$  representing two non-negative integers.

Similarly, we use symbolic computation and Assumptions (9), integrate Eq. (1b) twice in relation to  $x$  with the integration functions equal to zero and choose

$$\alpha = \frac{\beta}{2} \zeta_1, \tag{12}$$

so as to find

$$Y_t(w) - \alpha Y_{2x}(w) - (\beta \zeta_2 + \gamma) Y_x(w) = 0, \tag{13}$$

which is the same as Expression (11).

Then, the assumption,

$$w(x, t) = \ln [h(x, t)], \tag{14}$$

helps us simplify System (1) into a linear partial differential equation, i.e.,

$$h_t(x, t) - \alpha h_{xx}(x, t) - (\beta \zeta_2 + \gamma) h_x(x, t) = 0, \tag{15}$$

with  $h(x, t)$  representing a positive differentiable function.

Thinking about all the above together, under the constraint

$$\alpha \neq 0, \tag{16}$$

for System (1), we construct one set of the hetero-Bäcklund transformations, i.e.,

$$u(x, t) = \frac{2\alpha}{\beta} \frac{h_x(x, t)}{h(x, t)} + \zeta_2, \tag{17a}$$

$$v(x, t) = \frac{4\alpha^2}{\beta^2} \left[ \frac{h_{xx}(x, t)}{h(x, t)} - \frac{h_x(x, t)^2}{h(x, t)^2} \right], \tag{17b}$$

$$h_t(x, t) - \alpha h_{xx}(x, t) - (\beta \zeta_2 + \gamma) h_x(x, t) = 0. \tag{17c}$$

We note that

- Eqs. (17) are one set of the hetero-Bäcklund transformations, which could couple the solutions  $h(x, t)$  of Eq. (17c) and the solutions  $u(x, t)$  and  $v(x, t)$  of System (1);
- Eq. (17c) is an already-investigated linear partial differential equation, as seen in Refs. [33, 34];
- Hetero-Bäcklund Transformations (17) are related to  $\gamma$ ,  $\beta$  and  $\alpha$ , the shallow-water coefficients for System (1), under Constraint (16);

- the above work is on the long waves in the shallow water, concerning the wave profile and tangential fluid velocity at the surface.

### 3 Similarity reductions for System (1)

Making use of symbolic computation and substituting the assumptions,<sup>2</sup>

$$u(x, t) = \theta(x, t) + \omega(x, t)p[z(x, t)], \tag{18a}$$

$$v(x, t) = \delta(x, t) + \kappa(x, t)q[z(x, t)], \tag{18b}$$

into System (1) result in

$$\begin{aligned}
 &p''\alpha\omega z_x^2 - pp'\beta\omega^2 z_x - p^2\beta\omega\omega_x - q'\beta\kappa z_x - q\beta\kappa_x \\
 &+ p'[\omega z_t + \alpha(2\omega_x z_x + \omega z_{xx}) - \gamma\omega z_x - \beta\omega\theta z_x] \\
 &+ p[\omega_t + \alpha\omega_{xx} - \beta(\omega\theta_x + \omega_x\theta) - \gamma\omega_x] \\
 &+ \theta_t + \alpha\theta_{xx} - \beta\delta_x - \beta\theta\theta_x - \gamma\theta_x = 0, \tag{19a}
 \end{aligned}$$

$$\begin{aligned}
 &-q''\alpha\kappa z_x^2 - (p'q + pq')\beta\omega\kappa z_x - pq\beta(\omega\kappa_x + \omega_x\kappa) \\
 &+ q'[\kappa z_t - \alpha(2\kappa_x z_x + \kappa z_{xx}) - \gamma\kappa z_x - \beta\theta\kappa z_x] \\
 &+ q[\kappa_t - \alpha\kappa_{xx} - \beta(\theta\kappa_x + \theta_x\kappa) - \gamma\kappa_x] \\
 &- p'\beta\delta\omega z_x - p\beta(\delta_x\omega + \delta\omega_x) + \delta_t - \alpha\delta_{xx} \\
 &- \beta(\theta\delta_x + \theta_x\delta) - \gamma\delta_x = 0, \tag{19b}
 \end{aligned}$$

with  $z(x, t) \neq 0$ ,  $\kappa(x, t) \neq 0$ ,  $\delta(x, t)$ ,  $\omega(x, t) \neq 0$  and  $\theta(x, t)$  meaning the to-be-determined differentiable real functions with  $z_x z_t \neq 0$ , the prime sign standing the differentiation on  $z$ , while  $q(z)$  and  $p(z)$  denoting for two differentiable non-zero real functions of  $z$ .

Considering Eqs. (19) as a couple of the ordinary differential equations (ODEs) concerning  $q(z)$  and  $p(z)$ , one might require the ratios of the coefficients of different derivatives/powers of  $q(z)$  and  $p(z)$  to merely represent certain functions of  $z$ .

The coefficients of  $p''$  in Eq. (19a) and  $q''$  in Eq. (19b), as the normalizing coefficients in Eqs. (19), respectively, come to

$$\Omega_1(z)\alpha\omega z_x^2 = -\beta\omega^2 z_x, \tag{20a}$$

$$\Omega_2(z)\alpha\omega z_x^2 = -\beta\omega\omega_x, \tag{20b}$$

$$\Omega_3(z)\alpha\omega z_x^2 = -\beta\kappa z_x, \tag{20c}$$

$$\Omega_4(z)\alpha\omega z_x^2 = -\beta\kappa_x, \tag{20d}$$

<sup>2</sup> similar to those in Refs. [35–41]

$$\Omega_5(z)\alpha\omega z_x^2 = \omega z_t + \alpha(2\omega_x z_x + \omega z_{xx}) - \gamma\omega z_x - \beta\omega\theta z_x, \tag{20e}$$

$$\Omega_6(z)\alpha\omega z_x^2 = \omega_t + \alpha\omega_{xx} - \beta(\omega\theta_x + \omega_x\theta) - \gamma\omega_x, \tag{20f}$$

$$\Omega_7(z)\alpha\omega z_x^2 = \theta_t + \alpha\theta_{xx} - \beta\delta_x - \beta\theta\theta_x - \gamma\theta_x, \tag{20g}$$

and

$$-\Gamma_1(z)\alpha\kappa z_x^2 = -\beta\omega\kappa z_x, \tag{21a}$$

$$-\Gamma_2(z)\alpha\kappa z_x^2 = -\beta(\omega\kappa_x + \omega_x\kappa), \tag{21b}$$

$$-\Gamma_3(z)\alpha\kappa z_x^2 = -\beta\delta\omega z_x, \tag{21c}$$

$$-\Gamma_4(z)\alpha\kappa z_x^2 = \kappa z_t - \alpha(2\kappa_x z_x + \kappa z_{xx}) - \gamma\kappa z_x - \beta\theta\kappa z_x, \tag{21d}$$

$$-\Gamma_5(z)\alpha\kappa z_x^2 = -\beta(\delta_x\omega + \delta\omega_x), \tag{21e}$$

$$-\Gamma_6(z)\alpha\kappa z_x^2 = \kappa_t - \alpha\kappa_{xx} - \beta(\theta\kappa_x + \theta_x\kappa) - \gamma\kappa_x, \tag{21f}$$

$$-\Gamma_7(z)\alpha\kappa z_x^2 = \delta_t - \alpha\delta_{xx} - \beta(\theta\delta_x + \theta_x\delta) - \gamma\delta_x, \tag{21g}$$

with  $\Omega_i(z)$ 's ( $i=1, \dots, 7$ ) and  $\Gamma_j(z)$ 's ( $j=1, \dots, 7$ ) meaning certain real to-be-determined functions of  $z$  only.

On account of the remarks in Ref. [35], each set of  $\theta(x, t)$ ,  $\omega(x, t)$ ,  $\delta(x, t)$ ,  $\kappa(x, t)$  and  $z(x, t)$  can turn to, at least, a similarity reduction.

Ground on the second freedom in Remark 3 in Ref. [35],<sup>3</sup> Eqs. (20a), (20c) and (21a) give rise to

$$\omega(x, t) = \pm \frac{\alpha}{\beta} z_x, \quad \kappa(x, t) = \mp \frac{\alpha^2}{\beta^2} z_x^2, \tag{22}$$

$$\Omega_1(z) = \mp 1, \quad \Omega_3(z) = 1, \quad \Gamma_1(z) = \pm 1,$$

according to the first freedom in Remark 3 in Ref. [35], Eq. (20b) indicates

$$z(x, t) = \lambda_1 x + \lambda_2 t + \lambda_3, \quad \Omega_2(z) = 0, \tag{23}$$

and then Eqs. (20d) and (21b) bring about

$$\Omega_4(z) = \Gamma_2(z) = 0, \tag{24}$$

in which  $\lambda_1 \neq 0$ ,  $\lambda_2 \neq 0$  and  $\lambda_3$  mean the real constants.

Because the first freedom in Remark 3 in Ref. [35] makes us transform Eqs. (21c) and (21d) into

$$\delta(x, t) = 0, \quad \theta(x, t) = \frac{\lambda_2 - \gamma\lambda_1}{\beta\lambda_1}, \tag{25}$$

$$\Gamma_3(z) = \Gamma_4(z) = 0,$$

Eqs. (20e), (20f), (20g), (21e), (21f) and (21g) come to

$$\Omega_6(z) = \Omega_5(z) = \Omega_7(z) = \Gamma_5(z) = \Gamma_6(z) = \Gamma_7(z) = 0. \tag{26}$$

Since System (1) can be simplified to the following ODEs:

$$p'' \mp pp' + q' = 0, \tag{27a}$$

$$q'' \pm (p'q + pq') = 0, \tag{27b}$$

integrating ODEs (27) once with respect to  $z$ , separately, we are capable of transforming ODEs (27) into a single ODE, i.e.,

$$p'' - \frac{1}{2}p^3 \mp \phi_1 p + \phi_2 = 0, \tag{28}$$

based on

$$q = -p' \pm \frac{1}{2}p^2 + \phi_1, \tag{29}$$

with  $\phi_1$  and  $\phi_2$  representing two real constants of integration.

Taking into consideration all the above in this section, under Constraint (16), we build up the following two sets of the similarity reductions for System (1):

$$u(x, t) = \frac{\lambda_2 - \gamma\lambda_1}{\beta\lambda_1} \pm \frac{\alpha}{\beta}\lambda_1 p[z(x, t)], \tag{30a}$$

$$v(x, t) = \pm \frac{\alpha^2}{\beta^2}\lambda_1^2 \left\{ p'[z(x, t)] \mp \frac{1}{2}p[z(x, t)]^2 - \phi_1 \right\}, \tag{30b}$$

$$z(x, t) = \lambda_1 x + \lambda_2 t + \lambda_3, \tag{30c}$$

$$p'' - \frac{1}{2}p^3 \mp \phi_1 p + \phi_2 = 0. \tag{30d}$$

Each of ODEs (30d), as a known ODE, has been presented in Refs. [42, 43]. The reason for the appearance of two sets of Similarity Reductions (30) is that the “ $\pm$ ” signs present.

As for the long waves in the shallow water, with respect to the wave profile and tangential fluid velocity at the surface, under Constraint (16), Similarity Reductions (30) are dependent on  $\alpha$ ,  $\beta$  and  $\gamma$ , the shallow-water coefficients for System (1).

### 4 Conclusions

People are studying hard the shallow water waves. For example, a recent paper in Nonlinear Dyn., Ref. [1],

<sup>3</sup> There exist three freedoms in Remark 3, as seen in Ref. [35].

has presented a generalized Broer-Kaup system for the shallow water waves.

In this paper, we have investigated System (1), a generalized Broer-Kaup-Kupershmidt system for the long waves in shallow water. Concerning the horizontal velocity of the water wave as well as the height of the water surface, we have symbolically computed<sup>4</sup> the following:

- Scaling Transformation (8) and Hetero-Bäcklund Transformations (17), from System (1) to an already-investigated linear partial differential equation, under Constraint (16);
- Similarity Reductions (30), each of which from System (1) to a known ODE, under Constraint (16).

We have found the reason for the appearance of two sets of Similarity Reductions (30), i.e., the presence of the “±” signs.

Both Hetero-Bäcklund Transformations (17) and Similarity Reductions (30) have been seen to be dependent on  $\alpha$ ,  $\beta$  and  $\gamma$ , the shallow-water coefficients for System (1).

**Acknowledgements** We express our sincere thanks to the Editors and Reviewers for their valuable comments.

**Funding Information** This work has been supported by the National Natural Science Foundation of China under Grant Nos. 11871116 and 11772017, and by the Fundamental Research Funds for the Central Universities of China under Grant No. 20S19XD-A11. X. Y. Gao also thanks the National Scholarship for Doctoral Students of China.

**Data availability** Data sharing is not applicable to this article as no datasets were generated or analyzed during the current study.

## Declarations

**Conflict of interest** The authors have no relevant financial or non-financial interests to disclose.

## References

1. Zhang, S., Zheng, X.W.: N-soliton solutions and nonlinear dynamics for two generalized Broer-Kaup systems. *Nonlinear Dyn.* **107**, 1179 (2022)
2. Wazwaz, A.M.: Painlevé integrability and lump solutions for two extended (3 + 1)- and (2 + 1)-dimensional Kadomtsev-Petviashvili equations. *Nonlinear Dyn.* **111**, 3623 (2023)
3. Wazwaz, A.M.: New integrable (2+1)- and (3+1)-dimensional shallow water wave equations: multiple soliton

- solutions and lump solutions. *Int. J. Numer. Method. H.* **32**, 138 (2022)
4. Mandal, U.K., Malik, S., Kumar, S., Das, A.: A generalized (2+1)-dimensional Hirota bilinear equation: integrability, solitons and invariant solutions. *Nonlinear Dyn.* **111**, 4593 (2023)
5. Ismael, H.F., Akkilic, A.N., Murad, M.A., Bulut, H., Mahmoud, W., Osman, M.S.: Boiti-Leon-Manna-Pempinelli equation including time-dependent coefficient (vcBLMPE): a variety of nonautonomous geometrical structures of wave solutions. *Nonlinear Dyn.* **110**, 3699 (2022)
6. Shen, Y., Tian, B., Liu, S.H., Zhou, T.Y.: Studies on certain bilinear form, N-soliton, higher-order breather, periodic-wave and hybrid solutions to a (3+1)-dimensional shallow water wave equation with time-dependent coefficients. *Nonlinear Dyn.* **108**, 2447 (2022)
7. Liu, F.Y., Gao, Y.T., Yu, X., Ding, C.C.: Wronskian, Gramian, Pfaffian and periodic-wave solutions for a (3+1)-dimensional generalized nonlinear evolution equation arising in the shallow water waves. *Nonlinear Dyn.* **108**, 1599 (2022)
8. Zhou, T.Y., Tian, B., Chen, Y.Q., Shen, Y.: Painlevé analysis, auto-Bäcklund transformation and analytic solutions of a (2+1)-dimensional generalized Burgers system with the variable coefficients in a fluid. *Nonlinear Dyn.* **108**, 2417 (2022)
9. Cheng, C.D., Tian, B., Shen, Y., Zhou, T.Y.: Bilinear form and Pfaffian solutions for a (2+1)-dimensional generalized Konopelchenko-Dubrovsky-Kaup-Kupershmidt system in fluid mechanics and plasma physics. *Nonlinear Dyn.* (2023). <https://doi.org/10.1007/s11071-022-08189-6>
10. Liu, F.Y., Gao, Y.T., Yu, X.: Rogue-wave, rational and semi-rational solutions for a generalized (3 + 1)-dimensional Yu-Toda-Sasa-Fukuyama equation in a two-layer fluid. *Nonlinear Dyn.* **111**, 3713 (2023)
11. Gao, X.Y., Guo, Y.J., Shan, W.R.: Letter to the Editor on a (2 + 1)-dimensional variable-coefficient Sawada-Kotera system in plasma physics and fluid dynamics. *Results Phys.* **44**, 106099 (2023)
12. Shen, Y., Tian, B., Cheng, C.D., Zhou, T.Y.: Pfaffian solutions and nonlinear waves of a (3 + 1)-dimensional generalized Konopelchenko-Dubrovsky-Kaup-Kupershmidt system in fluid mechanics. *Phys. Fluids* **35**, 025103 (2023)
13. Cheng, C.D., Tian, B., Ma, Y.X., Zhou, T.Y., Shen, Y.: Pfaffian, breather and hybrid solutions for a (2 + 1)-dimensional generalized nonlinear system in fluid mechanics and plasma physics. *Phys. Fluids* **34**, 115132 (2022)
14. Shen, Y., Tian, B.: Bilinear auto-Bäcklund transformations and soliton solutions of a (3 + 1)-dimensional generalized nonlinear evolution equation for the shallow water waves. *Appl. Math. Lett.* **122**, 107301 (2021)
15. Li, L.Q., Gao, Y.T., Yu, X., Deng, G.F., Ding, C.C.: Gramian solutions and solitonic interactions of a (2+1)-dimensional Broer-Kaup-Kupershmidt system for the shallow water. *Int. J. Numer. Method. H.* **32**, 2282 (2022)
16. Whitham, G.B.: Variational methods and applications to water waves. *Proc. Roy. Soc. Lond. A* **299**, 6 (1967)
17. Broer, L.J.: Approximate equations for long water waves. *Appl. Sci. Res.* **31**, 377 (1975)
18. Kupershmidt, B.A.: Mathematics of dispersive water waves. *Commun. Math. Phys.* **99**, 51 (1985)

<sup>4</sup> More related symbolic-computation studies on other nonlinear evolution equations have been reported, i.e., in Refs. [44–51].

19. Zhao, Z.L., Han, B.: On optimal system, exact solutions and conservation laws of the Broer-Kaup system. *Eur. Phys. J. Plus* **130**, 223 (2015)
20. Cao, X.Q., Guo, Y.N., Hou, S.H., Zhang, C.Z., Peng, K.C.: Variational Principles for two kinds of coupled nonlinear equations in shallow water. *Symmetry-Basel* **12**, 850 (2020)
21. Malik, S., Kumar, S., Kumari, P., Nisar, K.S.: Some analytic and series solutions of integrable generalized Broer-Kaup system. *Alex. Eng. J.* **61**, 7067 (2022)
22. Zhou, T.Y., Tian, B.: Auto-Bäcklund transformations, Lax pair, bilinear forms and bright solitons for an extended (3+1)-dimensional nonlinear Schrödinger equation in an optical fiber. *Appl. Math. Lett.* **133**, 108280 (2022)
23. Wu, X.H., Gao, Y.T., Yu, X., Ding, C.C., Li, L.Q.: Modified generalized Darboux transformation, degenerate and bound-state solitons for a Lakshmanan-Porsezian-Daniel equation in a ferromagnetic spin chain. *Chaos Solitons Fract.* **162**, 112399 (2022)
24. Yang, D.Y., Tian, B., Tian, H.Y., Wei, C.C., Shan, W.R., Jiang, Y.: Darboux transformation, localized waves and conservation laws for an M-coupled variable-coefficient nonlinear Schrödinger system in an inhomogeneous optical fiber. *Chaos Solitons Fract.* **156**, 111719 (2022)
25. Wu, X.H., Gao, Y.T., Yu, X., Ding, C.C.: N-fold generalized Darboux transformation and soliton interactions for a three-wave resonant interaction system in a weakly nonlinear dispersive medium. *Chaos Solitons Fract.* **165**, 112786 (2022)
26. Zhou, T.Y., Tian, B., Zhang, C.R., Liu, S.H.: Auto-Bäcklund transformations, bilinear forms, multiple-soliton, quasi-soliton and hybrid solutions of a (3 + 1)-dimensional modified Korteweg-de Vries-Zakharov-Kuznetsov equation in an electron-positron plasma. *Eur. Phys. J. Plus* **137**, 912 (2022)
27. Yang, D.Y., Tian, B., Hu, C.C., Liu, S.H., Shan, W.R., Jiang, Y.: Conservation laws and breather-to-soliton transition for a variable-coefficient modified Hirota equation in an inhomogeneous optical fiber. *Wave. Random Complex.* (2023). <https://doi.org/10.1080/17455030.2021.1983237>
28. Shen, Y., Tian, B., Zhou, T.Y., Gao, X.T.: Nonlinear differential-difference hierarchy relevant to the Ablowitz-Ladik equation: Lax pair, conservation laws, N-fold Darboux transformation and explicit exact solutions. *Chaos Solitons Fract.* **164**, 112460 (2022)
29. Gao, X.Y., Guo, Y.J., Shan, W.R.: On a Whitham-Broer-Kaup-like system arising in the oceanic shallow water. *Chin. J. Phys.* (2023). <https://doi.org/10.1016/j.cjph.2022.11.005>
30. Gao, X.Y., Guo, Y.J., Shan, W.R.: Symbolically computing the shallow water via a (2+1)-dimensional generalized modified dispersive water-wave system: similarity reductions, scaling and hetero-Bäcklund transformations. *Qual. Theor. Dyn. Syst.* **22**, 17 (2023)
31. Bell, E.T.: Exponential polynomials. *Ann. Math.* **35**, 258 (1934)
32. Lambert, F., Loris, I., Springael, J., Willer, R.: On a direct bilinearization method: Kaup's higher-order water wave equation as a modified nonlocal Boussinesq equation. *J. Phys. A* **27**, 5325 (1994)
33. Rodrigo-Ilarri, J., Rodrigo-Clavero, M.E., Cassiraga, E., Ballesteros-Almonacid, L.: Assessment of groundwater contamination by terbuthylazine using vadose zone numerical models. Case study of Valencia province (Spain). *Int. J. Environ. Res. Public Health* **17**, 3280 (2020)
34. Gao, X.Y., Guo, Y.J., Shan, W.R.: Scaling transformation, hetero-Bäcklund transformation and similarity reduction on a (2+1)-dimensional generalized variable-coefficient Boiti-Leon-Pempinelli system for water waves. *Rom. Rep. Phys.* **73**, 111 (2021)
35. Clarkson, P., Kruskal, M.: New similarity reductions of the Boussinesq equation. *J. Math. Phys.* **30**, 2201 (1989)
36. Gao, X.T., Tian, B.: Water-wave studies on a (2+1)-dimensional generalized variable-coefficient Boiti-Leon-Pempinelli system. *Appl. Math. Lett.* **128**, 107858 (2022)
37. Gao, X.Y., Guo, Y.J., Shan, W.R.: Reflecting upon some electromagnetic waves in a ferromagnetic film via a variable-coefficient modified Kadomtsev-Petviashvili system. *Appl. Math. Lett.* **132**, 108189 (2022)
38. Gao, X.T., Tian, B., Shen, Y., Feng, C.H.: Considering the shallow water of a wide channel or an open sea through a generalized (2 + 1)-dimensional dispersive long-wave system. *Qual. Theory Dyn. Syst.* **21**, 104 (2022)
39. Gao, X.Y., Guo, Y.J., Shan, W.R.: Oceanic shallow-water symbolic computation on a (2 + 1)-dimensional generalized dispersive long-wave system. *Phys. Lett. A* **457**, 128552 (2023)
40. Gao, X.Y., Guo, Y.J., Shan, W.R.: Letter to the Editor on a shallow water wave equation in *Results Phys.* **43**, 106048 (2022) and its generalization. *Results Phys.* **44**, 106199 (2023)
41. Gao, X.T., Tian, B., Feng, C.H.: In oceanography, acoustics and hydrodynamics: investigations on an extended coupled (2+1)-dimensional Burgers system. *Chin. J. Phys.* **77**, 2818 (2022)
42. Ince, E.: *Ordinary Differential Equations*. Dover, New York (1956)
43. Zwillinger, D.: *Handbook of Differential Equations*, 3rd edn. Acad, San Diego (1997)
44. Wu, X.H., Gao, Y.T., Yu, X., Liu, L.Q., Ding, C.C.: Vector breathers, rogue and breather-rogue waves for a coupled mixed derivative nonlinear Schrödinger system in an optical fiber. *Nonlinear Dyn.* **111**, 5641 (2023)
45. Shen, Y., Tian, B., Zhou, T.Y., Gao, X.T.: N-fold Darboux transformation and solitonic interactions for the Kraenkel-Manna-Merle system in a saturated ferromagnetic material. *Nonlinear Dyn.* **111**, 2641 (2023)
46. Yang, D.Y., Tian, B., Wang, M., Zhao, X., Shan, W.R., Jiang, Y.: Lax pair, Darboux transformation, breathers and rogue waves of an N-coupled nonautonomous nonlinear Schrödinger system for an optical fiber or plasma. *Nonlinear Dyn.* **107**, 2657 (2022)
47. Liu, F.Y., Gao, Y.T.: Lie group analysis for a higher-order Boussinesq-Burgers system. *Appl. Math. Lett.* **132**, 108094 (2022)
48. Wu, X.H., Gao, Y.T., Yu, X., Ding, C.C., Hu, L., Li, L.Q.: Binary Darboux transformation, solitons, periodic waves and modulation instability for a nonlocal Lakshmanan-Porsezian-Daniel equation. *Wave Motion* **114**, 103036 (2022)
49. Gao, X.Y., Guo, Y.J., Shan, W.R., Du, Z., Chen, Y.Q.: Magneto-optic studies on a ferromagnetic material via an extended (3+1)-dimensional variable-coefficient modified

- Kadomtsev-Petviashvili system. Qual. Theory Dyn. Syst. **21**, 153 (2022)
50. Yang, D.Y., Tian, B., Hu, C.C., Zhou, T.Y.: The generalized Darboux transformation and higher-order rogue waves for a coupled nonlinear Schrödinger system with the four-wave mixing terms in a birefringent fiber. Eur. Phys. J. Plus **137**, 1213 (2022)
51. Wu, X.H., Gao, Y.T., Yu, X., Ding, C.C., Liu, F.Y., Jia, T.T.: Darboux transformation, bright and dark-bright solitons of an N-coupled high-order nonlinear Schrödinger system in an optical fiber. Mod. Phys. Lett. B **36**, 2150568 (2022)

**Publisher's Note** Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Springer Nature or its licensor (e.g. a society or other partner) holds exclusive rights to this article under a publishing agreement with the author(s) or other rightsholder(s); author self-archiving of the accepted manuscript version of this article is solely governed by the terms of such publishing agreement and applicable law.