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Fixed-time sliding mode attitude tracking control for a submarine-launched missile with multiple disturbances

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Abstract This paper studies a novel adaptive fixedtime sliding mode attitude tracking control for a submarine-launched missile, which is affected by sea winds, sea waves, ocean currents and other disturbances during the water-exit process. Firstly, the nonlinear water-exit dynamic model of the submarinelaunched missile is established, and then it is transformed into a simple second-order attitude tracking system. Subsequently, a novel non-singular fixedtime fast terminal sliding mode surface (NFFTSMS) with fixed-time convergence is presented, and the preestablished settling time is also developed. Moreover, a novel adaptive non-singular fixed-time fast terminal sliding mode control (ANFFTSMC) is presented by employing a fixed-time disturbance observer, a fixed-time differentiator and the proposed NFFTSMS. Closed-loop stability of the proposed controller is proved by utilizing the Lyapunov methodology. Finally, numerical simulations including two typical launch trajectories of the missile are carried out to demonstrate the strong robustness of the proposed control scheme.

Keywords Submarine-launched missile · Attitude tracking control · Non-singular fixed-time fast terminal sliding mode control · Multiple disturbances · Fixedtime disturbance observer · Fixed-time differentiator

1 Introduction

The submarine-launched missile, one of the most critical parts of the triad nuclear powers for every country, plays a crucial role in military strategies. The launching point of the submarine-launched missile is random in the presence of high-speed maneuver of the submarine underwater [1]. Meanwhile, the seawater is a strong barrier for the propagation of electromagnetic waves, such that the radar cannot detect the submarine effectively. Therefore, the submarine-launched missile has the characteristics of good concealment, strong ability to attack suddenly, high survivability, and formidable deterrence [2].

In most circumstances, the whole flight stage of the submarine-launched missile consists of three parts: the underwater stage, the water-exit stage and the aerial stage. The missile rises through the thrust of the underwater engine until its body is completely out of water in the water-exit stage [3]. However, the attitude and trajectory of the missile are seriously affected by the ocean currents, sea waves and sea winds during the water-exit process. This can make the missile deviate from the designed motion state and result in unstable flight or even an unsuccessful launch. Furthermore, the submarine-launched missile causes a series of complex changes in the flow field, such as the development and collapse of vacuoles, which produces complicated dynamic characteristics. Additionally, the flow filed evolution process adversely acts on the missile due to the pressure or even the impact load gener-

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ated by cavity collapses, which constantly changes the strained condition of the missile, thus forming a complex fluid solid coupling process [4]. Therefore, the water-exit dynamic model of the submarine-launched missile is complicated and difficult to describe accurately. Moreover, the attitude tracking control design for the submarine-launched missile is also one of the most challenging problems for engineers in the presence of highly complex nonlinear dynamics, external disturbances, couplings, and uncertainties.

The initial conditions for the aerial stage of the submarine-launched missile are essential. Thus, the attitude information needs to be accurately obtained at the end of the water-exit process. Although the classical proportional-integral-derivative (PID) controller has already been developed based on the dynamical models linearized around the fixed flight points, most of the existing approaches cannot satisfy the precision demands of attitude tracking [5-8]. To further improve the performance and reliability of the control system, an adaptive longitudinal trajectory tracking control method based on backstepping technique is presented in [9]. However, the steady-state error is comparatively large. A sliding mode attitude control law based on the feedback linearization is also presented for a submarine-launched missile in the presence of strong nonlinearity of submerged buoy space motion and interference uncertainty [10]. Nevertheless, system chattering is inevitable in this case. With the development of modern control theories, multiple new control laws are proposed to ensure the robust, accurate, and rapid response to the attitude tracking controller in the presence of large disturbances and uncertainties. Unfortunately, the research of attitude tracking control system design for submarine-launched missiles is rare in the existing literature. In other fields, finite-time attitude tracking control [11], smooth backstepping sliding mode control [12], adaptive output feedback integral sliding mode attitude tracking control [13], and disturbance observer-based attitude tracking control [14– 18] are widely utilized, especially for spacecrafts that have better performance than traditional ones. Many scholars, such as Utkin [19], Filippov [20], Roxin [21], Drakunov [22], Perruquetti [23], Defoort [24], and Baccioti [25], have also studied the finite-time stability theories to prove the stability of the corresponding controllers. In recent years, the fixed-time convergence sliding mode control has become a research focus [26– 28]. In fact, the fixed-time strategy has been applied to many areas, such as mobile robots [29], multiagent systems [30,31], fractional order chaotic systems [32] and power systems [33]. Different from finite-time stability, fixed-time stability can guarantee convergence within finite time, whose upper bound is a constant independent of initial conditions, which facilitates convergence time estimation and controller design to satisfy the severe settling time constraint [32]. This phenomenon was initially discovered in the context of systems that are homogeneous in the bi-limit [34]. Fixed-time stability was widely studied by Cruz-Zavala [35], Polyakov [36] and Fridman [37]. Moreover, fixed-time stability is an extension of finite-time stability and demands boundedness of the settling time function for a globally finite-time stable system [38]. The main tool for analysis of fixed-time stability is the Lyapunov function method [39,40]. Compared with other sliding mode controls, the fixed-time sliding mode control obtains faster convergence in both the reaching phase and the sliding phase [41]. Furthermore, it can guarantee that the steady-state error is within a smaller convergence region for attitude tracking especially in the field of submarine-launched missiles.

In this study, we present a novel adaptive nonsingular fixed-time fast terminal sliding mode control (ANFFTSMC) for a submarine-launched missile. The proposed ANFFTSMC consists of a non-singular fixedtime fast terminal sliding mode surface (NFFTSMS), a fixed-time disturbance observer and a fixed-time differentiator. To the best of our knowledge, this paper might be the first attempt to design a fixed-time attitude tracking control for submarine-launched missiles. The main contribution of this paper includes:

- The complicated nonlinear water-exit dynamic model of submarine-launched missiles is completely established in the presence of engine thrust and multiple disturbances for the first time. Then, the attitude tracking control system design for the complicated dynamic model is transformed into a simple second-order state equation, which simplifies the control system.
- The derivative of the attitude command, which can be constructed to be a novel NFFTSMS, is obtained by a fixed-time differentiator. Moreover, the proposed NFFTSMS can avoid singularity without adding additional parameters, and achieve faster convergence and a more compact structure than the existing approaches.



Fig. 1 Two typical trajectories of the submarine-launched missile

3. Unlike the existing attempts to design the sliding mode controller, the proposed ANFFTSMC is developed based on the total disturbances that are estimated by the fixed-time disturbance observer. The proposed controller does not require the bounded disturbances in advance. Also, the control parameters are not required to be equal to or greater than the disturbances.

This paper is organized as follows. In Sect. 2, the waterexit dynamic model of submarine-launched missiles is developed. After that, Sect. 3 presents a novel adaptive non-singular fixed-time fast terminal sliding mode control. Then, Sect. 4 provides the simulations and results. Lastly, Sect. 5 briefly concludes this work.

Notation In this paper, the bold italic variable X is denoted as a vector and $X = \begin{bmatrix} X_x & X_y & X_z \end{bmatrix}^T$. The italic variable X is the norm of the vector X.

2 Water-exit dynamic model

In most cases, a submarine-launched missile is launched at a depth of 30 m or even 100 m beneath the ocean.



Fig. 2 Definition of two coordinate systems and force analysis of the water-exit process

Meanwhile, the submarine-launched missile has two typical trajectories in the water-exit process, which correspond to the vertical launching mode and the oblique launching mode. The two typical trajectories are shown by the yellow and blue lines, respectively, in Fig. 1.

This section establishes the nonlinear attitude kinematics and dynamics model of the submarine-launched missile in the water-exit process. Figure 2 defines the body coordinate system $Ox_1y_1z_1$ and the ground coordinate system $O_ex_0y_0z_0$. It also presents the force analysis of the water-exit process.

where *O* is the barycenter, *O_e* is the launch point, *c* is the center of buoyancy, and *G* is gravity. *N* denotes the buoyancy, and *P* represents the thrust of the main engine for the submarine-launched missile. $\mathbf{R}_{\alpha} = \begin{bmatrix} R_{\alpha x} \ R_{\alpha y} \ R_{\alpha z} \end{bmatrix}^{\mathrm{T}}$ represents the fluid viscosity force. $\mathbf{R}_{\omega} = \begin{bmatrix} R_{\omega x} \ R_{\omega y} \ R_{\omega z} \end{bmatrix}^{\mathrm{T}}$ denotes the fluid viscosity damping force and $\mathbf{R} = \begin{bmatrix} R_{x} \ R_{y} \ R_{z} \end{bmatrix}^{\mathrm{T}}$ stands for the fluid inertia force. The coordinate transform matrix from $O_{e}x_{0}y_{0}z_{0}$ to $Ox_{1}y_{1}z_{1}$ is given by

$$\boldsymbol{B}_{G} = \begin{bmatrix} \cos\varphi\cos\psi & \sin\varphi\cos\psi & -\sin\psi\\ \cos\varphi\sin\psi\sin\gamma - \sin\varphi\cos\gamma\sin\psi\sin\gamma + \cos\varphi\cos\gamma\cos\psi\sin\gamma\\ \cos\varphi\sin\psi\cos\gamma + \sin\varphi\sin\gamma\sin\psi\sin\gamma - \cos\varphi\sin\gamma\cos\psi\sin\gamma \end{bmatrix}$$
(1)

where φ , ψ and γ represent the pitch angle, yaw angle and roll angle, respectively, for the submarine-launched missile [42]. Subsequently, these forces or torques will be analyzed in detail in the following parts.

2.1 Fluid inertia force and torque

The submarine-launched missile is considered to be axisymmetric, and the seawater is assumed as the ideal fluid. Therefore, the fluid momentum Q_f and fluid moment of momentum K_f are described as follows in [43].

$$\begin{bmatrix} Q_{fx} \\ Q_{fy} \\ Q_{fz} \\ K_{fx} \\ K_{fy} \\ K_{fz} \end{bmatrix} = \begin{bmatrix} \lambda_{11} & 0 & 0 & 0 & 0 & 0 \\ 0 & \lambda_{22} & 0 & 0 & 0 & \lambda_{26} \\ 0 & 0 & \lambda_{33} & 0 & \lambda_{35} & 0 \\ 0 & 0 & 0 & \lambda_{44} & 0 & 0 \\ 0 & 0 & \lambda_{53} & 0 & \lambda_{55} & 0 \\ 0 & \lambda_{62} & 0 & 0 & 0 & \lambda_{66} \end{bmatrix} \begin{bmatrix} V_x \\ V_y \\ V_z \\ \omega_x \\ \omega_y \\ \omega_z \end{bmatrix}$$
(2)

where V_i (i = x, y, z) is the velocity of the submarinelaunched missile in the body coordinate system $Ox_1y_1z_1$, ω_i (i = x, y, z) is the angular velocity and λ_{ij} (i = 1, 2, ..., 6, j = 1, 2, ..., 6) is the added mass. By employing the momentum theorem and the moment of momentum theorem, the inertia force and torque of the ideal fluid for the submarine-launched missile are described as

$$-R = \frac{\mathrm{d}\boldsymbol{Q}_{f}}{\mathrm{d}t} + \boldsymbol{\omega} \times \boldsymbol{Q}_{f}$$
$$-M = \frac{\mathrm{d}\boldsymbol{K}_{f}}{\mathrm{d}t} + \boldsymbol{\omega} \times \boldsymbol{K}_{f} + \boldsymbol{V} \times \boldsymbol{Q}_{f}$$
(3)

2.2 Fluid viscosity force and torque

The definition of fluid viscosity force R_{α} and torque M_{α} is similar to that of aerodynamic force and torque for the aircraft, which yields

$$\boldsymbol{R}_{\alpha} = \begin{bmatrix} -C_x \frac{1}{2} \rho V^2 S & \frac{1}{2} \rho V^2 S C_y^{\alpha} \alpha & \frac{1}{2} \rho V^2 S C_z^{\beta} \beta \end{bmatrix}^{\mathrm{T}} \\ \boldsymbol{M}_{\alpha} = \begin{bmatrix} \frac{1}{2} \rho V^2 S L m_x^{\beta} \beta & \frac{1}{2} \rho V^2 S L m_y^{\beta} \beta & \frac{1}{2} \rho V^2 S L m_z^{\alpha} \alpha \end{bmatrix}^{\mathrm{T}}$$
(4)

where ρ , *S*, *V* and *L* represent the density of seawater, the cross-sectional area of the missile, the velocity of the missile and the length of the missile, respectively. α and β are the angle of attack and angle of sideslip, respectively. C_x , C_y^{α} and C_z^{β} denote fluid viscosity force coefficients. m_x^{β} , m_y^{β} and m_z^{α} represent fluid viscosity torque coefficients.

2.3 Fluid viscosity damping force and torque

The fluid viscosity damping force R_{ω} and torque M_{ω} are defined as

$$\begin{aligned} \boldsymbol{R}_{\omega} &= \begin{bmatrix} 0 & C_{y}^{\bar{\omega}_{z}} \frac{1}{2} \rho V^{2} S \omega_{z} & C_{z}^{\bar{\omega}_{y}} \frac{1}{2} \rho V^{2} S \omega_{y} \end{bmatrix}^{T} \\ \boldsymbol{M}_{\omega} &= \begin{bmatrix} \frac{1}{2} \rho V^{2} S L \left(m_{x}^{\bar{\omega}_{x}} \omega_{x} + m_{x}^{\bar{\omega}_{y}} \omega_{y} \right) \\ & \frac{1}{2} \rho V^{2} S L \left(m_{y}^{\bar{\omega}_{x}} \omega_{x} + m_{y}^{\bar{\omega}_{y}} \omega_{y} \right) & m_{z}^{\bar{\omega}_{z}} \frac{1}{2} \rho V^{2} S L \omega_{z} \end{bmatrix}^{T} \end{aligned}$$

$$(5)$$

where $C_y^{\bar{\omega}_z}$ and $C_z^{\bar{\omega}_y}$ are damping force coefficients. $m_x^{\bar{\omega}_x}, m_x^{\bar{\omega}_y}, m_y^{\bar{\omega}_x}, m_y^{\bar{\omega}_y}$ and $m_z^{\bar{\omega}_z}$ are damping torque coefficients.

2.4 Buoyancy and its torque

The buoyancy of the submarine-launched missile is given by

$$N = \rho V_L g \tag{6}$$

where V_L is the volume of displacement for the submarine-launched missile and g is the acceleration of gravity. It is worth noting that N is pointing along the $O_e y_0$ direction; thus we have

$$N = B_G \begin{bmatrix} 0\\N\\0 \end{bmatrix} = \begin{bmatrix} N\sin\varphi\cos\psi\\N(\sin\varphi\sin\psi\sin\gamma + \cos\varphi\cos\gamma)\\N(\sin\varphi\sin\psi\cos\gamma - \cos\varphi\sin\gamma) \end{bmatrix}$$
(7)

Similarly, the torque of buoyancy is defined as

$$\boldsymbol{M}_{N} = \begin{bmatrix} N \left(\sin\varphi \sin\psi \sin\gamma + \cos\varphi \cos\gamma \right) z_{c} - N \left(\sin\varphi \sin\psi \cos\gamma - \cos\varphi \sin\gamma \right) y_{c} \\ -N \sin\varphi \cos\psi z_{c} + N \left(\sin\varphi \sin\psi \cos\gamma - \cos\varphi \sin\gamma \right) x_{c} \\ N \sin\varphi \cos\psi y_{c} - N \left(\sin\varphi \sin\psi \sin\gamma + \cos\varphi \cos\gamma \right) x_{c} \end{bmatrix}$$
(8)



Fig. 3 Main engine distribution of the submarine-launched missile

where x_c , y_c and z_c represent the position of buoyancy in the body coordinate system $Ox_1y_1z_1$. M_N is the torque vector of buoyancy.

2.5 Gravity

Gravity G is pointing along the negative direction of $O_e y_0$; thus we have

$$G = B_G \begin{bmatrix} 0\\ -G\\ 0 \end{bmatrix}$$
$$= \begin{bmatrix} -G\sin\varphi\cos\psi\\ -G(\sin\varphi\sin\psi\sin\gamma + \cos\varphi\cos\gamma)\\ -G(\sin\varphi\sin\psi\cos\gamma - \cos\varphi\sin\gamma) \end{bmatrix}$$
(9)

2.6 Thrust of the main engine and its control torque

Figure 3 displays the main engine distribution of the submarine-launched missile. The control commands are four swing angles δ_1 , δ_2 , δ_3 and δ_4 , which are also given in Fig. 3.

Thus, the four swing angles can be converted into the equivalent three-channel swing angles independent of one another, such as δ_{φ} , δ_{ψ} and δ_{γ} . According to Fig. 3, it can be drawn that

$$\delta_{\varphi} = \frac{\delta_3 + \delta_4 - \delta_1 - \delta_2}{4}$$
$$\delta_{\psi} = \frac{\delta_2 + \delta_3 - \delta_1 - \delta_4}{4}$$
$$\delta_{\gamma} = \frac{\delta_1 + \delta_2 + \delta_3 + \delta_4}{4} \tag{10}$$

The thrust value of the four main engines is consistent. Therefore, the corresponding control forces and control torques are developed as follows.

$$\boldsymbol{P} = \begin{bmatrix} \boldsymbol{P} \\ \frac{\sqrt{2}}{2} \boldsymbol{P} \delta_{\varphi} \\ -\frac{\sqrt{2}}{2} \boldsymbol{P} \delta_{\psi} \end{bmatrix}, \quad \boldsymbol{M}_{\boldsymbol{P}} = \begin{bmatrix} -\boldsymbol{P} \delta_{\gamma} \boldsymbol{r} \\ -\frac{\sqrt{2}}{2} \boldsymbol{P} (\boldsymbol{x}_{R} - \boldsymbol{x}_{T}) \delta_{\psi} \\ -\frac{\sqrt{2}}{2} \boldsymbol{P} (\boldsymbol{x}_{R} - \boldsymbol{x}_{T}) \delta_{\varphi} \end{bmatrix}$$
(11)

where x_R is the distance between the position of main engines and vertex of the missile, x_T is the distance between the barycenter and vertex of the missile, and *P* denotes the total thrust. The four main engines are uniformly installed along a circle and the radius of the circle is *r*.

2.7 Multiple disturbances

(1) Sea wave

To facilitate the description of the motion for the submarine-launched missile, a new variable *l* is defined to describe the distance between the vertex of the missile and sea level in the Ox_1 direction.

$$l = \frac{Y + x_T \sin \theta - h}{\sin \theta} \tag{12}$$

where *Y* is the position of the submarine-launched missile in the ground coordinate system $O_e x_0 y_0 z_0$, *h* is the depth of the launching point, and θ is the trajectory inclination angle. The disturbance forces caused by the six levels of sea waves can be obtained as follows in [9].

$$R_{dy} = \begin{cases} 0.00391l^3 + 0.1436l^2 - 0.1892l - 89.68 & -15 \le l \le 0\\ 0.0003077l^3 - 0.4629l^2 + 12.72l - 89.12 & 0 < l \le 15 \end{cases}$$
(13)
$$R_{dz} = \begin{cases} 0.01512l^3 + 0.403l^2 + 0.2167l - 42.48 & -15 \le l < 0\\ -0.04748l^3 + 0.3908l^2 + 7.681l - 43.12 & 0 \le l \le 15 \end{cases}$$
(14)

where $R_{dx} = 0$, $\mathbf{R}_d = \begin{bmatrix} R_{dx} & R_{dy} & R_{dz} \end{bmatrix}^{\mathrm{T}}$ denotes the disturbance vector in the presence of sea waves and the unit is kN.

(2) Sea wind

The sea wind will result in disturbance force P_w . It can be described as

$$P_w = QC_{wx}K_z \tag{15}$$

where $Q = 0.5\rho_1 V_{wz}^2$, Q denotes the wind pressure, ρ_1 represents the air density, K_z is usually set to $K_z = 1$,

 V_{wz} represents the wind velocity and C_{wx} denotes the aerodynamic drag coefficient. The wind direction is denoted by ϕ . Hence, the disturbance forces produced by the sea wind can be indicated as

$$\begin{bmatrix} P_{wdx} \\ P_{wdy} \\ P_{wdz} \end{bmatrix} = \begin{bmatrix} P_w \cos \phi \\ 0 \\ -P_w \sin \phi \end{bmatrix}$$
(16)

Thus, the disturbance force in the ground coordinate system $O_e x_0 y_0 z_0$ can be transformed into that in the body coordinate system $O x_1 y_1 z_1$. The corresponding force P_{wt} and torque M_w will be obtained by the following equation.

$$\boldsymbol{P}_{wt} = \begin{bmatrix} P_{wtx} \\ P_{wty} \\ P_{wtz} \end{bmatrix} = \boldsymbol{B}_G \begin{bmatrix} P_w \cos\phi \\ 0 \\ -P_w \sin\phi \end{bmatrix}$$
(17)

$$\boldsymbol{M}_{w} = \begin{bmatrix} \boldsymbol{0} \\ P_{wtz}Dl\left(\boldsymbol{x}_{T} - \frac{l}{2}\right) \\ P_{wty}Dl\left(\boldsymbol{x}_{T} - \frac{l}{2}\right) \end{bmatrix}$$
(18)

where D is the diameter of the submarine-launched missile.

(3) Ocean current

The ocean current is similar to the wind in the atmosphere for the missile. It will result in the added angle of attack α_{wp} and added angle of sideslip β_{w_p} . As shown in [42], we have

$$\alpha_{wp} = -\arctan\frac{w_p \cos A \sin \theta}{V - w_p \cos A \cos \theta}$$

$$\beta_{w_p} = \arctan\frac{w_p \sin A}{V}$$
(19)

where w_p is the flow velocity of the seawater, A represents the wind direction, and V is the velocity of the missile. The total angle of attack α and the total angle of sideslip β are replaced by $\alpha + \alpha_{w_p}$ and $\beta + \beta_{w_p}$, respectively. With the added angle of attack and added angle of sideslip, there exist large disturbances in the fluid viscosity force \mathbf{R}_{α} and torque M_{α} .

2.8 Kinematics and dynamics model of the submarine-launched missile

The momentum Q and moment of momentum K of the submarine-launched missile are described as follows [43].

$$Q = m(V + \omega \times r_c)$$

$$K = J_0 \omega + r_c \times mV$$
(20)

where *m* is the mass of the submarine-launched missile, and $\mathbf{r}_c = \begin{bmatrix} x_c & y_c & z_c \end{bmatrix}^T$ is the position vector of buoyancy in the body coordinate system. $\mathbf{J}_0 = \text{diag}\left(\begin{bmatrix} J_x & J_y & J_z \end{bmatrix}\right)$ represents the rotational inertia matrix. By employing the momentum theorem and the moment of momentum theorem, the dynamics model of the submarine-launched missile is described as

$$\frac{\mathrm{d}Q}{\mathrm{d}t} + \omega \times Q = R + R_{\alpha} + R_{\omega} + N + G + P + R_{d} + P_{wt} \frac{\mathrm{d}K}{\mathrm{d}t} + \omega \times K + V \times Q = M + M_{\alpha} + M_{\omega} + M_{N} + M_{P} + M_{w}$$
(21)

Therefore, multiple disturbances are included in the dynamic model, such as R_{α} , R_{d} , P_{wt} , M_{α} and M_{w} .

Expanding Eq. (21) yields

$$(m + \lambda_{11})\dot{V}_{x} + m\dot{\omega}_{y}z_{c} - m\dot{\omega}_{z}y_{c} + m \left[V_{z}\omega_{y} - V_{y}\omega_{z} + y_{c}\omega_{y}\omega_{x} + z_{c}\omega_{x}\omega_{z} - x_{c} \left(\omega_{y}^{2} + \omega_{z}^{2}\right)\right] = -C_{x}\frac{1}{2}\rho V^{2}S + (N - G)\sin\varphi\cos\psi + P + P_{wtx} (22) (m + \lambda_{22})\dot{V}_{x} + (mx_{c} + \lambda_{26})\dot{\omega}_{z} - mz_{c}\dot{\omega}_{x} + m \left[V_{x}\omega_{z} - V_{z}\omega_{x} + x_{c}\omega_{y}\omega_{x} + z_{c}\omega_{y}\omega_{z} - y_{c} \left(\omega_{z}^{2} + \omega_{x}^{2}\right)\right] = \frac{1}{2}\rho V^{2}S \left(C_{y}^{\alpha}\alpha + C_{y}^{\tilde{\omega}_{z}}\bar{\omega}_{z}\right) + (N - G) (\sin\varphi\sin\psi\sin\gamma + \cos\varphi\cos\gamma) + \frac{\sqrt{2}}{2}P\delta_{\varphi} + R_{dy} + P_{wty} (23) (m + \lambda_{33})\dot{V}_{z} + my_{c}\dot{\omega}_{x} - (mx_{c} - \lambda_{35})\dot{\omega}_{y} + m \left[V_{y}\omega_{x} - V_{x}\omega_{y} + x_{c}\omega_{z}\omega_{x} + y_{c}\omega_{y}\omega_{z} - z_{c} \left(\omega_{x}^{2} + \omega_{y}^{2}\right)\right] = \frac{1}{2}\rho V^{2}S \left(C_{z}^{\beta}\beta + C_{z}^{\tilde{\omega}y}\omega_{y}\right) + (N - G) (\sin\varphi\sin\psi\cos\gamma - \cos\varphi\sin\gamma) - \frac{\sqrt{2}}{2}P\delta_{\psi} + R_{dz} + P_{wtz} (24) (J_{x} + \lambda_{44})\dot{\omega}_{x} + my_{c}\dot{V}_{z} - mz_{c}\dot{V}_{y} + my_{c} \left(V_{y}\omega_{x} - V_{x}\omega_{y}\right) + (M_{z} - J_{y})\omega_{y}\omega_{z} = \frac{1}{2}\rho V^{2}SL \left(m_{x}^{\beta}\beta + m_{x}^{\tilde{\omega}x}\omega_{x} + m_{x}^{\tilde{\omega}y}\omega_{y}\right)$$

$$+ N (\sin \varphi \sin \psi \sin \gamma + \cos \varphi \cos \gamma) z_{c}$$
$$-N (\sin \varphi \sin \psi \cos \gamma - \cos \varphi \sin \gamma) y_{c} - P \delta_{\gamma} r$$
(25)

$$(J_{y} + \lambda_{55})\dot{\omega}_{y} + mz_{c}V_{x} - (mx_{c} - \lambda_{35})V_{z} + mz_{c}(V_{z}\omega_{y} - V_{y}\omega_{z}) + mx_{c}(V_{x}\omega_{y} - V_{y}\omega_{x}) + (J_{x} - J_{z})\omega_{x}\omega_{z} = \frac{1}{2}\rho V^{2}SL \left(m_{y}^{\beta}\beta + m_{y}^{\bar{\omega}_{x}}\omega_{x} + m_{y}^{\bar{\omega}_{y}}\omega_{y}\right) -N \sin\varphi\cos\psi z_{c} + N (\sin\varphi\sin\psi\cos\gamma - \cos\varphi\sin\gamma)x_{c} -\frac{\sqrt{2}}{2}P(x_{R} - x_{T})\delta_{\psi} + M_{wy}$$
(26)
$$(J_{z} + \lambda_{66})\dot{\omega}_{z} + (mx_{c} + \lambda_{26})\dot{V}_{y} - my_{c}\dot{V}_{x} + mx_{c}(V_{x}\omega_{z} - V_{z}\omega_{x}) + my_{c}(V_{y}\omega_{z} - V_{z}\omega_{y}) + (J_{y} - J_{x})\omega_{x}\omega_{y} = \frac{1}{2}\rho V^{2}SL \left(m_{z}^{\alpha}\alpha + m_{z}^{\bar{\omega}_{z}}\omega_{z}\right) + N\sin\varphi\cos\psi y_{c}$$

$$-N\left(\sin\varphi\sin\psi\sin\gamma + \cos\varphi\cos\gamma\right)x_{c}$$
$$-\frac{\sqrt{2}}{2}P(x_{R} - x_{T})\delta_{\varphi} + M_{wz}$$
(27)

The attitude kinematic equations are presented as

$$\dot{\varphi} = \frac{1}{\cos\psi} \left(\omega_y \sin\gamma + \omega_z \cos\gamma \right)$$
$$\dot{\psi} = \omega_y \cos\gamma - \omega_z \sin\gamma$$
$$\dot{\gamma} = \omega_x - \omega_y \tan\psi \sin\gamma - \omega_z \tan\psi \cos\gamma$$
(28)

For the convenience of description of the added mass, some new variables are defined as

$$\begin{cases} K_{11} = \frac{\lambda_{11}}{\rho V_L}, K_{22} = \frac{\lambda_{22}}{\rho V_L}, K_{33} = \frac{\lambda_{33}}{\rho V_L} \\ K_{26} = K_{62} = \frac{\lambda_{26}}{\rho V_L^{4/3}} = \frac{\lambda_{62}}{\rho V_L^{4/3}} \\ K_{35} = K_{53} = \frac{\lambda_{35}}{\rho V_L^{4/3}} = \frac{\lambda_{53}}{\rho V_L^{4/3}}, K_{35} = -K_{26} \\ K_{44} = \frac{\lambda_{44}}{\rho V_L^{5/3}}, K_{55} = K_{66} = \frac{\lambda_{55}}{\rho V_L^{5/3}} = \frac{\lambda_{66}}{\rho V_L^{5/3}} \end{cases}$$
(29)

Moreover, the kinematic equations of the submarinelaunched missile are developed as

$$\begin{bmatrix} \dot{X} \\ \dot{Y} \\ \dot{Z} \end{bmatrix} = (\boldsymbol{B}_G)^{\mathrm{T}} \begin{bmatrix} V_x \\ V_y \\ V_z \end{bmatrix}$$
(30)

where X, Y and Z denote the position of the submarinelaunched missile. Thus, the trajectory inclination angle θ and trajectory deflection angle ψ_v are, respectively, calculated by

 $\theta = \arctan\left(\dot{Y}/\dot{X}\right)$

$$\psi_v = -\arcsin\left(\dot{Z}/V\right) \tag{31}$$

where $V = \sqrt{V_x^2 + V_y^2 + V_z^2}$. Furthermore, the mass satisfies

$$m = m_0 - \mathrm{d}m \cdot t \tag{32}$$

where m_0 is the initial mass of the missile, dm represents the consumption rate of the fuel and *t* denotes the running time of the submarine-launched missile.

The relationships among the angle of attack α , the angle of sideslip β and the bank angle γ_v are depicted by [42].

$$\sin \beta = \cos (\theta - \varphi) \cos \psi_v \sin \psi \cos \gamma$$

+ $\sin (\varphi - \theta) \cos \psi_v \sin \gamma - \sin \psi_v \cos \psi \cos \gamma$
- $\sin \alpha \cos \beta = \cos (\theta - \varphi) \cos \psi_v \sin \psi \sin \gamma$
+ $\sin (\theta - \varphi) \cos \psi_v \cos \gamma - \sin \psi_v \cos \psi \sin \gamma$
 $\cos \psi_v \sin \gamma_v = \cos \alpha \cos \psi \sin \gamma - \sin \psi \sin \alpha$
(33)

Next, we will simplify the dynamic model to design the attitude control system. Define $X_1 = \begin{bmatrix} \varphi \ \psi \ \gamma \end{bmatrix}^T$, $X_2 = \begin{bmatrix} \omega_x \ \omega_y \ \omega_z \end{bmatrix}^T$ and $U = \begin{bmatrix} \delta_\varphi \ \delta_\psi \ \delta_\gamma \end{bmatrix}^T$. Combining Eqs. (25), (26) and (27) yields

$$\dot{\boldsymbol{X}}_1 = \boldsymbol{H}\boldsymbol{X}_2$$

$$\dot{\boldsymbol{X}}_2 = \boldsymbol{J}^{-1}\boldsymbol{B}\boldsymbol{U} + \boldsymbol{J}^{-1}\boldsymbol{D}_1$$
(34)

where J, H, \dot{H} , J and D_1 are given in "Appendix".

Differentiating Eq. (34) with respect to time obtains $\ddot{X}_1 = \dot{H}X_2 + H\dot{X}_2$

$$= \dot{\boldsymbol{H}}\boldsymbol{X}_2 + \boldsymbol{H}\left(\boldsymbol{J}^{-1}\boldsymbol{B}\boldsymbol{U} + \boldsymbol{J}^{-1}\boldsymbol{D}_1\right)$$
(35)

According to Eq. (34) and (35), the following second-order attitude tracking error dynamic system is presented as

$$\dot{\boldsymbol{e}}_1 = \boldsymbol{e}_2$$

$$\dot{\boldsymbol{e}}_2 = \dot{\boldsymbol{H}}\boldsymbol{X}_2 + \boldsymbol{H}\left(\boldsymbol{J}^{-1}\boldsymbol{B}\boldsymbol{U} + \boldsymbol{J}^{-1}\boldsymbol{D}_1\right) - \ddot{\boldsymbol{X}}_1^c \qquad (36)$$

where $X_1^c = \begin{bmatrix} \varphi^c \ \psi^c \ \gamma^c \end{bmatrix}^T$ is the attitude command vector, and $\dot{X}_1^c = \begin{bmatrix} \dot{\varphi}^c \ \dot{\psi}^c \ \dot{\gamma}^c \end{bmatrix}^T$ represents the derivative of the attitude command. $e_1 = X_1 - X_1^c$ and $e_2 = \dot{X}_1 - \dot{X}_1^c = HX_2 - \dot{X}_1^c$ denote the attitude tracking errors.

Simplifying Eq. (36) yields

$$\dot{\boldsymbol{e}}_1 = \boldsymbol{e}_2$$

$$\dot{\boldsymbol{e}}_2 = \boldsymbol{d} + \boldsymbol{B}_1 \boldsymbol{U}$$
(37)
where $\boldsymbol{d} = \dot{\boldsymbol{H}} \boldsymbol{X}_2 + \boldsymbol{H} \boldsymbol{J}^{-1} \boldsymbol{D}_1 - \ddot{\boldsymbol{X}}_1^c$ and $\boldsymbol{B}_1 = \boldsymbol{H} \boldsymbol{J}^{-1} \boldsymbol{B}$.

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Remark 1 According to Eq. (37), the attitude tracking control problem for a complicated dynamic model of the submarine-launched missile has been transformed into a simple second-order control system. After that, it is convenient to apply the corresponding modern control methods.

3 Adaptive non-singular fixed-time fast terminal sliding mode control

3.1 Definitions and Lemmas

Consider the following system

$$\dot{x} = f(x(t)), x(0) = 0, f(x(0)) = 0$$
 (38)

where $x \in \mathbb{R}^n$ and $f(x) : \mathbb{R}^n \to \mathbb{R}^n$ is a nonlinear function.

Definition 1 [44] If the origin of system (38) is Lyapunov stable, then system (38) is said to be finitetime stable and there exists a settling time function $T : \mathbb{R}^n \to \mathbb{R}^n$, such that for every $x_0 \in \mathbb{R}^n$, the solution $x (t, x_0)$ of the system satisfies $\lim_{t\to T(x_0)} x (t, x_0) = 0$.

Definition 2 [44] System (38) is said to be fixed-time stable if it is globally finite-time stable. Thus, it will converge to the origin within the bounded convergence time $T(x_0)$. In this case, there exists a positive constant T_{max} such that $T(x_0) < T_{\text{max}}$.

Lemma 1 [28] For system (38), suppose there exists a Lyapunov function V(x), parameters α_1 , β_1 , p, q, $k \in R^+$, pk < 1 and qk > 1, such that $\dot{V}(x) \leq -(\alpha_1 V(x)^p + \beta_1 V(x)^q)^k$. Then, the system is fixed-time stable. Furthermore, the upper bound of the convergence time is given as follows.

$$T \le \frac{1}{\alpha_1^k (1 - pk)} + \frac{1}{\beta_1^k (qk - 1)}$$
(39)

Lemma 2 [45] For any $x_i \in R, i = 1, 2..., n$, $\left(\sum_{i=1}^{n} |x_i|\right)^{v} \leq \sum_{i=1}^{n} |x_i|^{v}$, where $v \in R^+$ and $v \in (0, 1]$.

Lemma 3 [41] For any $x_i \in R, i = 1, 2..., n$, If $v \in R^+$ and $v_1 > 1$, we have $\left(\sum_{i=1}^n |x_i|\right)^{v_1} \le n^{v_1-1} \sum_{i=1}^n |x_i|^{v_1}$.

Notation 1 [46] In this paper, we use $\|\cdot\|$ for the Euclidean norm of vectors and induced norm of matrices. For a given vector $\mathbf{x} = [x_1, x_2, \dots, x_n]^T$, $\operatorname{sig}^p(\mathbf{x}) = [|x_1|^p \operatorname{sign}(x_1), |x_2|^p \operatorname{sign}(x_2), \dots, |x_n|^p \operatorname{sign}(x_n)]^T$, $\operatorname{sig}^p(x) = |x|^p \operatorname{sign}(x) \operatorname{sign}(\mathbf{x}) = [\operatorname{sign}(x_1), \operatorname{sign}(x_2), \dots, \operatorname{sign}(x_n)]^T$, where $\operatorname{sign}(.)$ denotes the sign function.

Assumption 1 [41] The disturbances d in Eq. (37) are assumed to be bounded and satisfy $||d|| \le d_m$, where d_m is a positive constant and d_m is the vector of the max value of the disturbances.

3.2 Control design

This section presents a novel adaptive non-singular fixed-time fast terminal sliding mode control. To guarantee that both e_1 and e_2 in Eq. (37) converge to zero within a fixed time regardless of the initial states, a novel non-singular fixed-time fast terminal sliding mode surface (NFFTSMS) is developed as follows [41].

$$\mathbf{S} = \mathbf{\Gamma} \left(\boldsymbol{e}_1 \right) \boldsymbol{e}_1 + \operatorname{sig}^{\gamma_2} \left(\boldsymbol{e}_2 \right) \tag{40}$$

where $e_1 \in R^3$, $e_2 \in R^3$ and $\Gamma(e_1)$ is a diagonal matrix. Moreover, the diagonal elements of $\Gamma(e_1)$ are given by

$$\Gamma(e_{1i}) = \left(\alpha_1 |e_{1i}|^{p-1/k\gamma_2} + \beta_1 |e_{1i}|^{q-1/k\gamma_2}\right)^{k\gamma_2}$$

$$\Gamma(e_1) = \operatorname{diag}\left(\left[\Gamma(e_{1x}) \Gamma(e_{1y}) \Gamma(e_{1z})\right]\right)$$
(41)

where $i = x, y, z, \alpha_1 > 0, \beta_1 > 0, k > 1, \gamma_2 > 1$ and p, q are positive parameters satisfying qk > 1 and $1/\gamma_2 < pk < 1$.

Similarly, $sig^{\gamma_2}(e_2)$ can be denoted as

$$\operatorname{sig}^{\gamma_{2}}(\boldsymbol{e}_{2}) = \left[\operatorname{sig}^{\gamma_{2}}(\boldsymbol{e}_{2x}) \operatorname{sig}^{\gamma_{2}}(\boldsymbol{e}_{2y}) \operatorname{sig}^{\gamma_{2}}(\boldsymbol{e}_{2z})\right]^{\mathrm{T}}$$
(42)

Theorem 1 If the sliding mode surface satisfies S = 0, $e_1 = 0$ and $e_2 = 0$ can be reached within a fixed convergence time T_0 regardless of the initial states. Thus, we obtain the upper bound of the convergence time

$$T_0 \le \frac{1}{\alpha_1^k (1 - pk)} + \frac{1}{\beta_1^k (qk - 1)}$$
(43)

Proof When the sliding mode surface satisfies S = 0, we have

$$\dot{\boldsymbol{e}}_1 = -\operatorname{sig}^k\left(\alpha_1\operatorname{sig}^p(\boldsymbol{e}_1) + \beta_1\operatorname{sig}^q(\boldsymbol{e}_1)\right) \tag{44}$$

Next, consider a new variable $\Phi_i = |e_{1i}|^{1-pk}$, (*i* = *x*, *y*, *z*), Eq. (44) can be rewritten as

$$\dot{\Phi}_i = -(1-pk)\left(\alpha_1 + \beta_1 \Phi_i^{\mu}\right)^k \tag{45}$$

where $\mu = (q - p) / (1 - pk)$. By integrating Eq. (45), we obtain

$$T_{0} = \lim_{\Phi_{i} \to \infty} \frac{1}{1 - pk} \int_{0}^{\Phi_{i}} \frac{1}{(\alpha_{1} + \beta_{1} \Phi_{i}^{\mu})^{k}} d\Phi_{i}$$

$$= \lim_{\Phi_{i} \to \infty} \frac{1}{1 - pk} \left(\int_{1}^{\Phi_{i}} \frac{1}{(\alpha_{1} + \beta_{1} \Phi_{i}^{\mu})^{k}} d\Phi_{i} + \int_{0}^{1} \frac{1}{(\alpha_{1} + \beta_{1} \Phi_{i}^{\mu})^{k}} d\Phi_{i} \right)$$

$$\leq \lim_{\Phi_{i} \to \infty} \frac{1}{1 - pk} \left(\int_{1}^{\Phi_{i}} \frac{1}{\beta_{1} \Phi_{i}^{\mu k}} d\Phi_{i} + \int_{0}^{1} \frac{1}{\alpha_{1}^{k}} d\Phi_{i} \right)$$

$$= \frac{1}{1 - pk} \left(\frac{1}{\beta_{1}^{k} (\mu k - 1)} + \frac{1}{\alpha_{1}^{k}} \right)$$

$$= \frac{1}{\alpha_{1}^{k} (1 - pk)} + \frac{1}{\beta_{1}^{k} (qk - 1)}$$
(46)

This completes the proof of Theorem 1.

Remark 2 It is worth noting that this sliding mode surface is partly motivated by the reference [41]. Compared with other sliding mode surfaces, the proposed NFFTSMS can avoid singularity without adding additional parameters and achieve faster convergence and a more compact structure.

To design the non-singular fixed-time fast terminal sliding mode control, a nonlinear function is defined as

$$\eta(x) = \begin{cases} \sin\left(\frac{\pi}{2}\frac{x^2}{\tau^2}\right), \ |x| \le \tau\\ 1, \ |x| > \tau \end{cases}$$
(47)

where $\tau > 0$ is an unknown parameter that needs to be designed. $\eta(x)$ is a non-negative function, when $x \to 0$, there is $\eta(x)/x \to 0$.

In fact, the derivative of the attitude command X_1^c is usually unavailable in the process of real flight. Therefore, e_2 is unknown in Eq. (40). To solve this problem, a fixed-time differentiator can be developed based on the reference [47]. Accordingly, we have

$$\dot{v}_1 = v_2 - k_1 \operatorname{sig}^{a_1} (v_1 - \chi^c) - \kappa_1 \operatorname{sig}^{b_1} (v_1 - \chi^c)$$

$$\dot{v}_2 = -k_2 \operatorname{sig}^{a_3} (v_1 - \chi^c) - \kappa_2 \operatorname{sig}^{b_3} (v_1 - \chi^c)$$
(48)

where χ^c can be set as φ^c , ψ^c or γ^c , v_1 is the estimated value of χ^c , v_2 is the derivative of χ^c which corresponds to $\dot{\varphi}^c$, $\dot{\psi}^c$ or $\dot{\gamma}^c$. Therefore, \dot{X}_1^c is obtained and e_2 will be calculated. Moreover, the convergence time of the differentiator is set as T_d and the exact value can be found in Theorem 2 in the reference [47].

Subsequently, a fixed-time disturbance observer for Eq. (37) is established by [47].

$$\dot{\mathbf{Z}}_{1} = \mathbf{Z}_{2} - k_{1} \operatorname{sig}^{a_{1}} (\mathbf{Z}_{1} - \mathbf{e}) - \kappa_{1} \operatorname{sig}^{b_{1}} (\mathbf{Z}_{1} - \mathbf{e})$$

$$\dot{\mathbf{Z}}_{2} = \mathbf{Z}_{3} - k_{2} \operatorname{sig}^{a_{2}} (\mathbf{Z}_{1} - \mathbf{e}) - \kappa_{2} \operatorname{sig}^{b_{2}} (\mathbf{Z}_{1} - \mathbf{e}) + \mathbf{B}_{1} \mathbf{U}$$

$$\dot{\mathbf{Z}}_{3} = -k_{3} \operatorname{sig}^{a_{3}} (\mathbf{Z}_{1} - \mathbf{e}) - \kappa_{3} \operatorname{sig}^{b_{3}} (\mathbf{Z}_{1} - \mathbf{e})$$
(49)

where $\mathbf{Z}_1 \in \mathbb{R}^3$, $\mathbf{Z}_2 \in \mathbb{R}^3$, $\mathbf{Z}_3 \in \mathbb{R}^3$. $k_1, k_2, k_3, \kappa_1, \kappa_2$ and κ_3 are observer gains. $a_1 \in (1 - \varepsilon, 1)$, ε is a sufficiently small number, $a_2 = 2a_1 - 1$ and $a_3 = 3a_1 - 2$. Similarly, $b_1 \in (1, 1 + \varepsilon_1)$, $b_2 = 2b_1 - 1$, $b_3 = 3b_1 - 2$ and ε_1 is also sufficiently small. The resulting observer estimation error system for errors $\boldsymbol{\xi}_1 = \mathbf{Z}_1 - \boldsymbol{e}_1$, $\boldsymbol{\xi}_2 = \mathbf{Z}_2 - \dot{\boldsymbol{e}}_1$ and $\boldsymbol{\xi}_3 = \mathbf{Z}_3 - \boldsymbol{d}$ takes the following form

$$\dot{\boldsymbol{\xi}}_{1} = \boldsymbol{\xi}_{2} - k_{1} \operatorname{sig}^{a_{1}} (\boldsymbol{Z}_{1} - \boldsymbol{e}) - \kappa_{1} \operatorname{sig}^{b_{1}} (\boldsymbol{Z}_{1} - \boldsymbol{e})
\dot{\boldsymbol{\xi}}_{2} = \boldsymbol{\xi}_{3} - k_{2} \operatorname{sig}^{a_{2}} (\boldsymbol{Z}_{1} - \boldsymbol{e}) - \kappa_{2} \operatorname{sig}^{b_{2}} (\boldsymbol{Z}_{1} - \boldsymbol{e})
\dot{\boldsymbol{\xi}}_{3} = -k_{3} \operatorname{sig}^{a_{3}} (\boldsymbol{Z}_{1} - \boldsymbol{e}) - \kappa_{3} \operatorname{sig}^{b_{3}} (\boldsymbol{Z}_{1} - \boldsymbol{e})$$
(50)

Theorem 2 Consider the observer (49) and the corresponding estimation error system (50). The errors ξ_1, ξ_2 and ξ_3 converge to the origin and the observer states Z_1, Z_2, Z_3 converge to the state variables e, \dot{e}_1 and d, respectively, within a fixed time T_1 .

$$T_1 \le \frac{\lambda_{\max}^{\rho_1}\left(P_1\right)}{r_1\rho_1} + \frac{1}{r_2\sigma\gamma^{\sigma}} \tag{51}$$

where $\rho_1 = 1 - a_1$, $\sigma = b_1 - 1$, $\lambda_{\min}(Q_1) > 0$ is the minimum eigenvalue of the matrix Q_1 , $\lambda_{\max}(P_1) > 0$ is the maximum eigenvalue of the matrix P_1 , $r_1 = \lambda_{\min}(Q_1) / \lambda_{\max}(P_1)$, $r_2 = \lambda_{\min}(Q_2) / \lambda_{\max}(P_2)$, and $\gamma \leq \lambda_{\min}(P_2)$ is a positive number. Q_1 and Q_2 are positive definite matrices and the matrices P_1 , P_2 satisfy

$$P_1 A_1 + A_1^{\mathrm{T}} P_1 = -Q_1$$

$$P_2 A_2 + A_2^{\mathrm{T}} P_2 = -Q_2$$
(52)

where A_1 and A_2 are defined as

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$$A_{1} = \begin{bmatrix} -k_{1} & 1 & 0 \\ -k_{2} & 0 & 1 \\ -k_{3} & 0 & 0 \end{bmatrix}, A_{2} = \begin{bmatrix} -\kappa_{1} & 1 & 0 \\ -\kappa_{2} & 0 & 1 \\ -\kappa_{3} & 0 & 0 \end{bmatrix}$$
(53)

Theorem 2 can be proved by the reference [47].

Next, by employing the proposed NFFTSMS in Eq. (40), an adaptive non-singular fixed-time fast terminal sliding mode control (ANFFTSMC) is designed. Thus we have

$$\boldsymbol{u} = -\boldsymbol{B}_{1}^{-1} \left[\frac{1}{\gamma_{2}} \left(\bar{\boldsymbol{\Gamma}} \left(\boldsymbol{e}_{1} \right) + \boldsymbol{\Gamma} \left(\boldsymbol{e}_{1} \right) \right) \operatorname{sig}^{2-\gamma_{2}} \left(\boldsymbol{e}_{2} \right) + \frac{1}{\gamma_{2}} \operatorname{diag} \left\{ \eta \left(|\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right) \right\} \operatorname{diag} \left\{ \left(|\boldsymbol{e}_{2}|^{1-\gamma_{2}} \right) \right\} \operatorname{sig}^{\gamma_{5}} \left(\sigma_{1} \operatorname{sig}^{\gamma_{3}} \left(\boldsymbol{S} \right) + \sigma_{2} \operatorname{sig}^{\gamma_{4}} \left(\boldsymbol{S} \right) \right) + \operatorname{diag} \left(\boldsymbol{\Upsilon} \left| \boldsymbol{Z}_{3} \right|_{\max} \right) \operatorname{sign} \left(\boldsymbol{S} \right) \right]$$
(54)

where $\gamma_5 > 1, \gamma_3\gamma_5 < 1, \Upsilon > 1$ and $\gamma_4\gamma_5 > 1$. σ_1 and σ_2 are two positive constants. \mathbf{Z}_3 represents the estimate value of the multiple disturbances in Eq. (37), which has the form of $\mathbf{Z}_3 = \mathbf{d}$ within a fixed time T_1 . $|\mathbf{Z}_3|_{\text{max}} = [|\mathbf{Z}_{3x}|_{\text{max}} |\mathbf{Z}_{3y}|_{\text{max}} |\mathbf{Z}_{3z}|_{\text{max}}]^{\text{T}}$ denotes the max value of the disturbances, which can be updated online. $\mathbf{\bar{\Gamma}}(\mathbf{e}_1)$ is a diagonal matrix which has the diagonal elements as

$$\bar{\mathbf{\Gamma}}(e_{1i}) = k\gamma_2 \left(\alpha_1 |e_{1i}|^{p-1/k\gamma_2} + \beta_1 |e_{1i}|^{q-1/k\gamma_2} \right)^{k\gamma_2 - 1} \Lambda$$
(55)

where $\Lambda = \alpha_1 \left(p - \frac{1}{k\gamma_2} \right) |e_{1i}|^{p-1/k\gamma_2} + \beta_1 \left(q - \frac{1}{k\gamma_2} \right) |e_{1i}|^{q-1/k\gamma_2}$ and i = x, y, z.

The following theorem presents the fixed-time stability of the proposed ANFFTSMC in Eq. (54) for the second-order system (37).

Theorem 3 Consider system (37) and the controller (54), the state variables e_1 and e_2 will converge to the origin within a fixed time and the upper bound of the convergence time is given as follows.

$$T < T_{\max} = T_0 + T_1 + T_2 + T_d + T_L$$
(56)

where

$$T_{2} = \frac{2}{2^{(1+\gamma_{3}\gamma_{5})/2} \sigma_{1}^{\gamma_{5}} 3^{\gamma_{5}-1} (1-\gamma_{3}\gamma_{5})} + \frac{2}{2^{(1+\gamma_{4}\gamma_{5})/2} 3^{\gamma_{5}-1} (\gamma_{4}\gamma_{5}-1)} T_{L} = \left(\tau^{1/(\gamma_{2}-1)}\right) / (\|\Upsilon\|\mathbf{Z}_{3}\|_{\max} - d_{m}\|)$$

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Proof Consider the following Lyapunov function candidate

$$\boldsymbol{V}_{\boldsymbol{S}} = \frac{1}{2} \boldsymbol{S}^{\mathrm{T}} \boldsymbol{S} \tag{57}$$

The derivative of V_S is

$$\begin{split} \dot{\boldsymbol{V}}_{S} &= \boldsymbol{S}^{\mathrm{T}} \dot{\boldsymbol{S}} \\ &= \boldsymbol{S}^{\mathrm{T}} \left(\frac{\mathrm{d}\boldsymbol{\Gamma} \left(\boldsymbol{e}_{1} \right)}{\mathrm{d}t} \boldsymbol{e}_{1} + \boldsymbol{\Gamma} \left(\boldsymbol{e}_{1} \right) \boldsymbol{e}_{2} + \gamma_{2} \mathrm{diag} \left\{ |\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right\} \dot{\boldsymbol{e}}_{2} \right) \\ &= \boldsymbol{S}^{\mathrm{T}} \left(\bar{\boldsymbol{\Gamma}} \left(\boldsymbol{e}_{1} \right) \boldsymbol{e}_{1} + \boldsymbol{\Gamma} \left(\boldsymbol{e}_{1} \right) \boldsymbol{e}_{2} + \gamma_{2} \mathrm{diag} \left\{ |\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right\} \\ &\quad (\boldsymbol{d} + \boldsymbol{B}_{1} \boldsymbol{U})) \end{split}$$
(58)

Substituting the ANFFTSMC into Eq. (58) yields

$$\dot{V}_{S} = -S^{T} \operatorname{diag} \left\{ \eta \left(|\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right) \right\}$$

$$\operatorname{sig}^{\gamma_{5}} \left(\sigma_{1} \operatorname{sig}^{\gamma_{3}} \left(S \right) + \sigma_{2} \operatorname{sig}^{\gamma_{4}} \left(S \right) \right)$$

$$-\gamma_{2} S^{T} \operatorname{diag} \left\{ \left(|\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right) \right\}$$

$$\left(\operatorname{diag} \left(|\boldsymbol{Z}_{3}| \right) \operatorname{sign} \left(S \right) - \boldsymbol{d} \right)$$

$$\leq -\sum_{i=1}^{3} \left(\sigma_{1} \eta^{1/\gamma_{5}} \left(|\boldsymbol{e}_{2i}|^{\gamma_{2}-1} \right) |\boldsymbol{S}_{i}|^{1/\gamma_{5}+\gamma_{3}} + \sigma_{2} \eta^{1/\gamma_{5}} \left(|\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right) |\boldsymbol{S}_{i}|^{1/\gamma_{5}+\gamma_{4}} \right)^{\gamma_{5}}$$

$$-\gamma_{2} \sum_{i=1}^{3} |\boldsymbol{e}_{2i}|^{\gamma_{2}-1} |\boldsymbol{S}_{i}| \left(\Upsilon |\boldsymbol{Z}_{3}|_{\max} - |\boldsymbol{d}_{i}| \right) (59)$$

With the proposed fixed-time disturbance observer and Theorem 2, we have $Z_3 = d$ within a fixed time T_1 and $\Upsilon |Z_3|_{\text{max}} - |d_i| > 0$. Meanwhile, according to Lemma 2 and Lemma 3, Eq. (59) can be written as

$$\begin{split} \dot{V}_{S} &\leq -\sum_{i=1}^{3} \left(\sigma_{1} \eta^{1/\gamma_{5}} \left(|\boldsymbol{e}_{2i}|^{\gamma_{2}-1} \right) |\boldsymbol{S}_{i}|^{1/\gamma_{5}+\gamma_{3}} \right. \\ &+ \sigma_{2} \eta^{1/\gamma_{5}} \left(|\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right) |\boldsymbol{S}_{i}|^{1/\gamma_{5}+\gamma_{4}} \right)^{\gamma_{5}} \\ &\leq -3^{1-\gamma_{5}} \left(\sigma_{1} \min \left(\eta^{1/\gamma_{5}} \left(|\boldsymbol{e}_{2i}|^{\gamma_{2}-1} \right) \right) \right) \\ &\sum_{i=1}^{3} \left(|\boldsymbol{S}_{i}|^{2} \right)^{(1+\gamma_{3}\gamma_{5})/(2\gamma_{5})} \\ &+ \sigma_{2} \min \left(\eta^{1/\gamma_{5}} \left(|\boldsymbol{e}_{2i}|^{\gamma_{2}-1} \right) \right) \right) \\ &\sum_{i=1}^{3} \left(|\boldsymbol{S}_{i}|^{2} \right)^{(1+\gamma_{4}\gamma_{5})/(2\gamma_{5})} \right)^{\gamma_{5}} \\ &\leq -3^{1-\gamma_{5}} \\ &\left(\tau_{m} \boldsymbol{V}_{S}^{(1+\gamma_{3}\gamma_{5})/(2\gamma_{5})} + \upsilon_{m} \boldsymbol{V}_{S}^{(1+\gamma_{4}\gamma_{5})/(2\gamma_{5})} \right)^{\gamma_{5}} \tag{60} \end{split}$$

where $\tau_m = 2^{(1+\gamma_3\gamma_5)/(2\gamma_5)}\sigma_1 \min(\eta^{1/\gamma_5}(|\boldsymbol{e}_{2i}|^{\gamma_2-1})),$ $\upsilon_m = 2^{(1+\gamma_4\gamma_5)/(2\gamma_5)}\sigma_2 \min(\eta^{1/\gamma_5}(|\boldsymbol{e}_{2i}|^{\gamma_2-1})).$ If $\boldsymbol{e}_{2i} \neq 0$, then $\eta(|\boldsymbol{e}_{2i}|^{\gamma_2-1}) > 0$, (i = 1, 2, 3). Two new variables are defined as follows for the convenience of the proof of the proposed ANFFTSMC.

$$\Omega_{1} = \left\{ (\boldsymbol{e}_{1i}, \boldsymbol{e}_{2i}) | |\boldsymbol{e}_{2i}|^{\gamma_{2}-1} \ge \tau, \forall i = x, y, z \right\}$$

$$\Omega_{2} = \left\{ (\boldsymbol{e}_{1i}, \boldsymbol{e}_{2i}) | |\boldsymbol{e}_{2i}|^{\gamma_{2}-1} < \tau, \forall i = x, y, z \right\}$$
(61)

If the error system states e_{1i} and e_{2i} reach the region Ω_1 , we have $\eta (|e_{2i}|^{\gamma_2-1}) = 1$ based on Eq. (47). Therefore, τ_m and υ_m are positive constants regardless of the system states. They can be simplified as

$$\tau_m = 2^{(1+\gamma_3\gamma_5)/(2\gamma_5)} \sigma_1$$

$$\upsilon_m = 2^{(1+\gamma_4\gamma_5)/(2\gamma_5)} \sigma_2$$
(62)

Equation (60) can be converted into

$$\dot{V}_{S} \leq -3^{1-\gamma_{5}} \left(2^{(1+\gamma_{3}\gamma_{5})/(2\gamma_{5})} \sigma_{1} V_{S}^{(1+\gamma_{3}\gamma_{5})/(2\gamma_{5})} + 2^{(1+\gamma_{4}\gamma_{5})/(2\gamma_{5})} \sigma_{2} V_{S}^{(1+\gamma_{4}\gamma_{5})/(2\gamma_{5})} \right)^{\gamma_{5}}$$
(63)

Based on Lemma 1 and Eq. (63), the system states e_{1i} and e_{2i} will be driven on the sliding mode surface S = 0 within a bounded time T_2 .

$$T_{2} = \frac{2}{2^{(1+\gamma_{3}\gamma_{5})/2} \sigma_{1}^{\gamma_{5}} 3^{\gamma_{5}-1} (1-\gamma_{3}\gamma_{5})} + \frac{2}{2^{(1+\gamma_{4}\gamma_{5})/2} 3^{\gamma_{5}-1} (\gamma_{4}\gamma_{5}-1)}$$
(64)

Furthermore, if $e_{2i} \neq 0$, we have $0 < |e_{2i}|^{\gamma_2 - 1} < \tau$ in the region Ω_2 . It can be similarly verified according to the definition of τ_m and υ_m that $\tau_m > 0$ and $\upsilon_m > 0$. Thus, based on Lemma 1 and Eq. (60), the sliding mode surface $S_i = 0$ is still an attractor in the region Ω_2 . Next, we will verify that the region Ω_2 is not an attractor area except for the origin. Substituting the ANFFTSMC into Eq. (37) yields

$$\dot{\boldsymbol{e}}_{2} = -\frac{1}{\gamma_{2}} \left(\bar{\boldsymbol{\Gamma}} \left(\boldsymbol{e}_{1} \right) + \boldsymbol{\Gamma} \left(\boldsymbol{e}_{1} \right) \right) \operatorname{sig}^{2-\gamma_{2}} \left(\boldsymbol{e}_{2} \right) - \frac{1}{\gamma_{2}} \operatorname{diag} \left\{ \eta \left(|\boldsymbol{e}_{2}|^{\gamma_{2}-1} \right) \right\} \operatorname{diag} \left\{ \left(|\boldsymbol{e}_{2}|^{1-\gamma_{2}} \right) \right\} \operatorname{sig}^{\gamma_{5}} \left(\sigma_{1} \operatorname{sig}^{\gamma_{3}} \left(\boldsymbol{S} \right) + \sigma_{2} \operatorname{sig}^{\gamma_{4}} \left(\boldsymbol{S} \right) \right) - \operatorname{diag} \left(\boldsymbol{\Upsilon} \left| \boldsymbol{Z}_{3} \right|_{\max} \right) \operatorname{sign} \left(\boldsymbol{S} \right) + \boldsymbol{d}$$
(65)

If $e_{2i} \rightarrow 0$, Eq. (65) can be simplified as

$$\dot{\boldsymbol{e}}_2 = -\text{diag}\left(\Upsilon \left| \boldsymbol{Z}_3 \right|_{\text{max}}\right) \text{sign}\left(\boldsymbol{S}\right) + \boldsymbol{d} \tag{66}$$

If $S_i > 0$, then $\dot{e}_2 \leq d_m - \Upsilon |Z_3|_{max} < 0$, and if $S_i < 0$, then $\dot{e}_2 > \Upsilon |Z_3|_{max} - d_m > 0$, where d_m is the vector of the max value of the disturbances. Therefore, the system states will transgress the region Ω_2 into Ω_1 monotonically within the time $T_L = (\tau^{1/(\gamma_2 - 1)}) / (||\Upsilon |Z_3|_{max} - d_m||)$. It is obvious that the system states from anywhere will be driven on the proposed sliding mode surface S = 0. Moreover, according to Theorem 1, the system states e_{1i} and e_{2i} will converge to the origin within a fixed time $T < T_0 + T_1 + T_2 + T_d + T_L$. This completes the proof of Theorem 3.

Remark 3 The proposed controller is developed based on the NFFTSMS and the total disturbances that are estimated by the fixed-time disturbance observer. Compared with other non-singular fixed-time fast terminal sliding mode controllers, the proposed ANFFTSMC does not require the bounded disturbances in advance, and the control parameters are not required to be equal to or greater than the disturbances.

4 Simulations and results

In this section, to prove the effectiveness of the proposed adaptive non-singular fixed-time fast terminal sliding mode control (ANFFTSMC) in Eq. (54), various results of the numerical simulation are presented. The fluid coefficients and parameters of the submarinelaunched missile are listed in Tables 1 and 2 [43].

4.1 Simulation results for the vertical launching mode

The proposed control algorithm will be first tested for the vertical launching mode, and the initial parameters are set to be $x = 0, y = 0, z = 0, \varphi = 89^{\circ}, \psi =$

Table 1 Fluid coefficients of the submarine-launched missile

Parameters	Value	Parameters	Value	Parameters	Value
C_x	0.141	m_y^{β}	0.69	<i>K</i> ₁₁	0.0222
m_x^β	0.00152	$m_y^{\omega_y}$ -	- 0.61	<i>K</i> ₂₂	1.1096
$m_x^{\omega_x}$ -	- 0.0044	C_z^β -	- 2.32	K ₃₃	1.1096
$m_x^{\omega_y}$	0.0008	$C_z^{\omega_y}$ -	- 1.17	K_{44}	0.1406
C_y^{α}	2.32	m_z^{α}	0.69	K_{66}	3.8129
$C_y^{\omega_z}$	1.17	$m_z^{\omega_z}$ –	- 0.61	K ₂₆	- 0.363

Table 2 Parameters of the submarine-launched missile

Parameters	Value	Parameters	Value
ρ (kg/m ³)	1024	x_c (m)	0.02
$V_L(m^3)$	37.68	y_c (m)	-0.005
<i>m</i> ⁰ (kg)	40,000	z_c (m)	0.008
$g (m/s^2)$	9.80665	J_x (kg.m ²)	2.0729e+04
<i>L</i> (m)	12.5	J_y (kg m ²)	3.9213e+05
$S(m^2)$	3.1416	$J_z (\mathrm{kg}\mathrm{m}^2)$	3.9213e+05
P(kN)	560	dm(kg/s)	100
r (m)	0.6	x_R (m)	12.0
<i>D</i> (m)	2.0	x_T (m)	8.3276

1°, $\gamma = 1^\circ$, $V_x = 0.5$, $V_y = 0$, $V_z = 0$ and $\omega_x = 0.001$, $\omega_y = 0.001$, $\omega_z = 0.001$. The expected attitude commands are chosen as $\varphi^c = 90^\circ$, $\psi^c = 0^\circ$ and $\gamma^c = 0^\circ$. Meanwhile, the depth of the launching point for the submarine-launched missile is h = 30 m, and the simulation sampling step is 5ms. The control command is limited within $[-15^\circ, 15^\circ]$. Moreover, the PID controller parameters for both the pitch channel and the yaw channel are set as (10, 2, 6), and the roll channel is (5.0, 1.7). Furthermore, the control parameters

of ANFFTSMC are selected to be $\alpha_1 = 1.5$, $\beta_1 = 1.5$, p = 0.4, q = 0.7, k = 2, $\gamma_2 = 5/3$, $\gamma_3 = 1/3$, $\gamma_4 = 3/4$, $\gamma_5 = 1.85$, $\tau = 0.005$, $\sigma_1 = 1.0$, $\Upsilon = 2$ and $\sigma_2 = 1.0$. The parameters of the fixed-time disturbance observer are selected to be $a_1 = 0.9$, $a_2 = 2a_1 - 1$, $a_3 = 3a_1 - 2$, $b_1 = 1.8$, $b_1 = 2b_1 - 1$, $b_1 = 3b_1 - 2$ and the observer gains are set as $k_1 = \kappa_1 = 24$, $k_2 = \kappa_2 = 216$ and $k_3 = \kappa_3 = 864$ [47]. As for the parameter uncertainty, its bias is set to be 20% for the density of the seawater, and 10% for the fluid coefficients of the submarine-launched missile. The sea current velocity is $w_p = 1.5$ m/s; the air density is $\rho_1 = 1.225$; the sea wind velocity is $V_{wz} = 30$ m/s; the aerodynamic drag coefficient is $C_{wx} = 0.5$ and the wind direction is $\phi = 60^{\circ}$.

To reduce the chattering of the control system, the switching function sign(x) will be replaced by a hyperbolic tangent function $tanh(x) = (e^x - e^{-x})/(e^x + e^{-x})$. The simulation results of the submarine-launched missile with six degrees of freedom (6DOF) in the water-exit process are shown in Figs. 4, 5, 6, 7, 8, 9.

From Figs. 4, 5, 6, it can be seen that the proposed ANFFTSMC has a more satisfactory performance than



Fig. 4 Attitude tracking results under the PID controller and ANFFTSMC





Fig. 6 Angular velocity results under the PID controller and ANFFTSMC





Fig. 8 Control commands under the PID controller and ANFFTSMC



Fig. 9 Fixed-time disturbance observer results

the PID controller in spite of multiple disturbances including sea winds, sea waves and ocean currents, etc. The attitude tracking error of the three channels for the ANFFTSMC is $< 0.003^{\circ}$, while the attitude tracking error of the PID controller is 0.56°. Hence, the simulation results indicate that the ANFFTSMC has less overshoot and smaller steady-state error than the PID control law. Meanwhile, the total angle of attack and the total angle of sideslip are shown in Fig. 7, which indicate large fluid viscosity forces and torques. The responses of control commands are shown in Fig. 8, where the ANFFTSMC's control commands are larger than the PID controller's. Also, the ANFFTSMC has stronger robustness in the presence of multiple disturbances especially near the sea level. Moreover, the control command is more exquisite during the water-exit process especially from 2.0 to 2.5 s, which indicates large disturbances for the submarine-launched missile. From Fig. 9, it can be seen that, for multiple disturbances and un-modeled dynamics, the proposed fixedtime disturbance observer has excellent estimated performance.

4.2 Simulation results for the oblique launching mode

Next, the simulation is carried out on the missile launched in oblique launching mode. The initial parameters are set to $\varphi = 46^{\circ}$, $\psi = 0.5^{\circ}$, $\gamma = 0.5^{\circ}$ and other parameters are the same as those in Sect. 4.1. Moreover, the expected attitude commands are designed as follows

$$\begin{cases} \varphi^{c} = 45 + 5t \\ \psi^{c} = 2\sin(\pi t/6) \\ \gamma^{c} = \sin(\pi t/6) \end{cases}$$
(67)

The simulation results of the oblique launching mode are shown in Figs. 10, 11, 12, 13 and 14.

From Figs. 10, 11, it can be seen that the simulation results indicate the effectiveness of the proposed





Fig. 11 Attitude tracking error results for the oblique launching mode





Fig. 13 Simulation results of the differentiator



Fig. 14 Simulation results of the fixed-time disturbance observer for the oblique launching mode

ANFFTSMC with respect to the oblique launching mode for the submarine-launched missile. Moreover, the tracking error for the three channels is $< 0.03^{\circ}$. Meanwhile, the curves of the control commands are smooth in Fig. 12. As can be seen in Figs. 13, 14, both the differentiator and the proposed fixed-time disturbance observer show satisfactory performances for the oblique launching mode, which exhibits strong robustness of the proposed control scheme.

5 Conclusion

In this work, a novel fixed-time attitude tracking control named as ANFFTSMC is presented for a submarineposed ANFFTSMC is designed by employing a NFFTSMS, a fixed-time disturbance observer and a fixed-time differentiator. Compared with other sliding mode surfaces, the proposed NFFTSMS can avoid singularity without adding additional parameters and achieve faster convergence and a more compact structure. Moreover, the proposed ANFFTSMC does not require the bounded disturbances in advance, and the control parameters do not need to be equal to or greater than the disturbances. Simulations are carried out to demonstrate the strong robustness and effectiveness of the proposed control scheme in comparison with the PID control law. The proposed fixed-time distur-

launched missile with multiple disturbances. The pro-

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bance observer and the differentiator exhibit satisfactory performances in the 6DOF simulations. Moreover, the proposed control scheme achieves an attitude tracking error of 0.03° , which satisfies the requirement of the control system design. Acknowledgements The authors would like to thank the financial supports by the National Nature Science Fund of China (Grant No. 61403100), the open Fund of National Defense Key Discipline Laboratory of Micro-Spacecraft Technology (Grant No. HIT.KLOF.MST.201704), and the Fundamental Research Funds for the Central Universities (Grant No. HIT.NSRIF.2015.037).

Compliance with ethical standards

Conflict of interests The authors declare that there is no conflict of interest regarding the publication of this paper.

Appendix

$$\begin{split} \mathbf{D}_{1} &= \begin{bmatrix} -m_{V}c\dot{\mathbf{v}} + m_{z}c\dot{\mathbf{v}} - m_{y}c\left(\mathbf{V}_{y}\omega_{x} - \mathbf{V}_{x}\omega_{y}\right) - m_{z}c\left(\mathbf{V}_{z}\omega_{x} - \mathbf{V}_{x}\omega_{z}\right) - (J_{z} - J_{y})\omega_{y}\omega_{z} \\ &+ \frac{1}{2}\rho\mathbf{V}^{2}SL\left(m_{x}^{\beta}\beta + m_{x}^{\tilde{\omega}_{x}}\bar{\omega}_{x} + m_{x}^{\tilde{\omega}_{y}}\bar{\omega}_{y}\right) + B\left(\sin\varphi\sin\psi\sin\gamma + \cos\varphi\cos\gamma\right)z_{c} \\ &- B\left(\sin\varphi\sin\psi\cos\gamma - \cos\varphi\sin\gamma\right)y_{c} \\ -m_{z}c\dot{\mathbf{v}}_{x} + (m_{x}c - \lambda_{35})\dot{\mathbf{v}}_{z} - m_{z}c(\mathbf{V}_{z}\omega_{y} - \mathbf{V}_{y}\omega_{z}) - m_{x}c(\mathbf{V}_{x}\omega_{y} - \mathbf{V}_{y}\omega_{x}) - (J_{x} - J_{z})\omega_{x}\omega_{z} \\ &+ \frac{1}{2}\rho\mathbf{V}^{2}SL\left(m_{y}^{\beta}\beta + m_{y}^{\tilde{\omega}_{x}}\bar{\omega}_{x} + m_{y}^{\tilde{\omega}_{y}}\bar{\omega}_{y}\right) - B\sin\varphi\cos\psi_{z}c + B\left(\sin\varphi\sin\psi\cos\gamma - \cos\varphi\sin\gamma\right)x_{c} + M_{wy} \\ &- (m_{x}c + \lambda_{26})\dot{\mathbf{v}}_{y} + m_{y}c\dot{\mathbf{v}}_{x} - m_{x}c(\mathbf{V}_{x}\omega_{z} - \mathbf{V}_{z}\omega_{x}) - m_{y}c(\mathbf{V}_{y}\omega_{z} - \mathbf{V}_{z}\omega_{y}) - (J_{y} - J_{x})\omega_{x}\omega_{y} \\ &+ \frac{1}{2}\rho\mathbf{V}^{2}SL\left(m_{z}^{\alpha}\alpha + m_{z}^{\tilde{\omega}_{z}}\bar{\omega}_{z}\right) + B\sin\varphi\cos\psi_{z}c - B\left(sim\varphi\sin\psi\sin\gamma + \cos\varphi\cos\gamma\right)x_{c} + M_{wz} \\ \mathbf{J} = \begin{bmatrix} J_{x} + \lambda_{44} & 0 & 0 \\ 0 & J_{y} + \lambda_{55} & 0 \\ 0 & 0 & J_{z} + \lambda_{66} \end{bmatrix}, \\ \mathbf{H} = \begin{bmatrix} 0 \sin\gamma/\cos\psi\cos\gamma/\cos\psi \\ 0 & \cos\gamma &- \sin\gamma \\ 1 - \tan\psi\sin\gamma - \tan\psi\cos\gamma \end{bmatrix} \\ \mathbf{B} = \begin{bmatrix} -Pr & 0 & 0 \\ 0 & -\frac{\sqrt{2}}{2}P(x_{R} - x_{T}) & 0 \\ 0 & 0 & -\frac{\sqrt{2}}{2}P(x_{R} - x_{T}) \end{bmatrix} \\ \dot{\mathbf{H}} = \begin{bmatrix} 0\dot{\gamma}\cos\gamma\sec\psi + \dot{\psi}\sin\gamma\tan\psi\sin\psi\sin\gamma + \sin\psi\psi\sin\gamma\sin\psi\sin\gamma + \sin\psi\psi\psi\psi - \dot{\gamma}\sin\gamma\psi\psi\psi\psi \\ 0 &-\dot{\gamma}\sin\gamma &-\dot{\gamma}\cos\gamma \\ 0 &-\dot{\gamma}\sin\psi\cos\gamma - \dot{\psi}\sec^{2}\psi\sin\gamma &\dot{\gamma}\sin\psi\sin\gamma - \dot{\psi}\sec^{2}\psi\cos\gamma \end{bmatrix} \end{aligned}$$

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