

Effect and feasibility analysis of the smoothening functions for clearance-type nonlinearity in a practical driveline system

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Abstract Discontinuous dynamic characteristics in gear backlash often cause convergence problems when a nonlinear torsional system is simulated. This clearance-type nonlinearity can be mathematically modeled by using smoothening functions which change the dynamic responses of the system under discontinuous ranges into continuous ones. However, the effect of the smoothening functions is not well known and difficult to anticipate under various nonlinear conditions. Thus, a new smoothening function is proposed. The effect and feasibility of the model were investigated with a practical vehicle driveline system. To examine the key factors of the smoothening function, the harmonic balance method was used with an ‘ n ’th order polynomial function and compared with hyperbolic-type smoothening functions. The harmonic balance method and numerical analysis were compared for nonlinear system responses that include much high order of the super-harmonic components with respect to impulsive contact motions to understand the limits of the method. The smoothening function is applicable for

simulating gear impact phenomena with limit conditions.

Keywords Harmonic balance method · Smoothening function · Clearance-type nonlinearity · Gear backlash · Gear mesh force · Torsional system · Driveline

1 Introduction

Clearance-type nonlinearities such as gear backlash are related to convergence problems when nonlinear dynamic responses in a practical system are simulated. The clearance itself contains discontinuities due to contact or non-contact behaviors [1,2]. Much prior research has been done to develop time-varying stiffness models or clearance-type nonlinear models with smoothening functions [3–12]. For example, Shen et al. [3] established a dynamical model of a spur gear pair by including the backlash, time-varying stiffness, and static transmission error. Rao et al. [4] studied the torsional instabilities in a two-stage gear system by considering the torsional flexibility of the shafts and the meshing time-varying stiffness of the gears.

Al-shayyab and Kahraman [5] investigated sub-harmonic and chaotic motions in a multi-mesh gear train using a nonlinear time-varying dynamic model. Raghothama and Narayanan [6] used the incremental harmonic balance method to obtain the periodic motions of a 3-degree-of-freedom (DOF) nonlinear model of a geared rotor system subjected to parametric

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excitation under sinusoidal excitation. Wong et al. [7] presented the non-linearities in the restoring force by employing the incremental harmonic balance method. Kim et al. [8] showed the differences in the nonlinear frequency response characteristics when selected smoothening functions were employed by different smoothening factors. Smoothening functions have also been used to simulate clearance or piecewise nonlinearities of the mechanical systems in many prior studies [9–12].

The examination by Kim et al. [8] of the effect of the smoothening functions was limited to a simple torsional model with a single DOF with idealized system parameters. In general, the parameters from a practical system can cause more difficulties in estimation, especially when the dynamic conditions of the system are affected by vibro-impacts, such as gear rattle in the wide-open-throttle (WOT) condition [1,2,9,10]. Under severe gear impact conditions, the gear mesh stiffness changes abruptly from 0 to $2.7 \times 10^8 \text{ N m}^{-1}$ (or from 2.7×10^8 to 0 N m^{-1}). In addition, the prior studies focused on specific smoothening functions using hyperbolic tangent or arctangent functions [1,2,8–12,21,24], which are not enough to overcome the convergence problems with respect to the gear impact behaviors in a practical system. Thus, the main objectives of this study contain: (1) suggesting a new smoothening function model using an n th order polynomial function by comparing with prior models [8]; (2) examining the key factors of new smoothening function model along with certain ranges; (3) and investigating the limits of the simulation based on the harmonic balance method (HBM) by focusing on gear impact conditions.

2 Modeling of the practical system and nonlinearity

2.1 Physical system and its parameters

Figure 1a depicts the physical system based on a front-engine and front-wheel driveline with a manual transmission [10]. Based on the practical system shown in Fig. 1a, a schematic diagram can be constructed as shown in Fig. 1b. All of the loaded gears are assumed to be lumped into the input shaft without changing the dynamic characteristics of the system, and focus is centered on only one unloaded gear pair to examine

the vibro-impact phenomenon [1,2]. In addition, the employed gear pairs are assumed to be geometrically ideal without any errors under the dynamic conditions. Thus, the reduced order of the lumped system model with 4 DOF will be used to investigate the dynamic behaviors with smoothening functions.

The symbols used and their parameter values shown in Fig. 1 are described in Tables 1 and 2, where the given properties are measured and given from the relevant industry with a practical driveline. Here, the drag torques are assumed to be constant values in the given engine operating condition, and the damping of the clutch damper c_f and the damping of the drive shaft c_{ve} are estimated based on the 5% modal damping ratio [1,2,10]. The scope of this study is limited to the 3rd gear engaged and 5th gear unloaded status under severe driving conditions such as the wide-open-throttle (WOT) condition.

2.2 Basic equations and their nonlinearities

Based on the system in Fig. 1, the matrix formulation of the basic equations is expressed as follows.

$$\underline{\mathbf{M}}\ddot{\underline{\boldsymbol{\theta}}}(\mathbf{t}) + \underline{\mathbf{C}}\dot{\underline{\boldsymbol{\theta}}}(\mathbf{t}) + \underline{\mathbf{K}}\underline{\boldsymbol{\theta}}(\mathbf{t}) + \underline{\mathbf{f}}_{\mathbf{n}}(\underline{\boldsymbol{\theta}}) = \underline{\mathbf{T}}_{\mathbf{E}}(\mathbf{t}), \tag{1a}$$

$$\underline{\mathbf{M}} = \text{diag} [I_f, I_{ie}, I_{ou}, I_{ve}], \tag{1b}$$

$$\underline{\mathbf{C}} = \begin{bmatrix} c_f & -c_f & 0 & 0 \\ -c_f & c_f + c_g R_{iu}^2 + c_{ve} & c_g R_{iu} R_{ou} & -c_{ve} \\ 0 & c_g R_{iu} R_{ou} & c_g R_{ou}^2 & 0 \\ 0 & -c_{ve} & 0 & c_{ve} \end{bmatrix},$$

$$\underline{\mathbf{K}} = \begin{bmatrix} k_f & -k_f & 0 & 0 \\ -k_f & k_f + k_{ve} & 0 & -k_{ve} \\ 0 & 0 & 0 & 0 \\ 0 & -k_{ve} & 0 & k_{ve} \end{bmatrix}, \tag{1c,d}$$

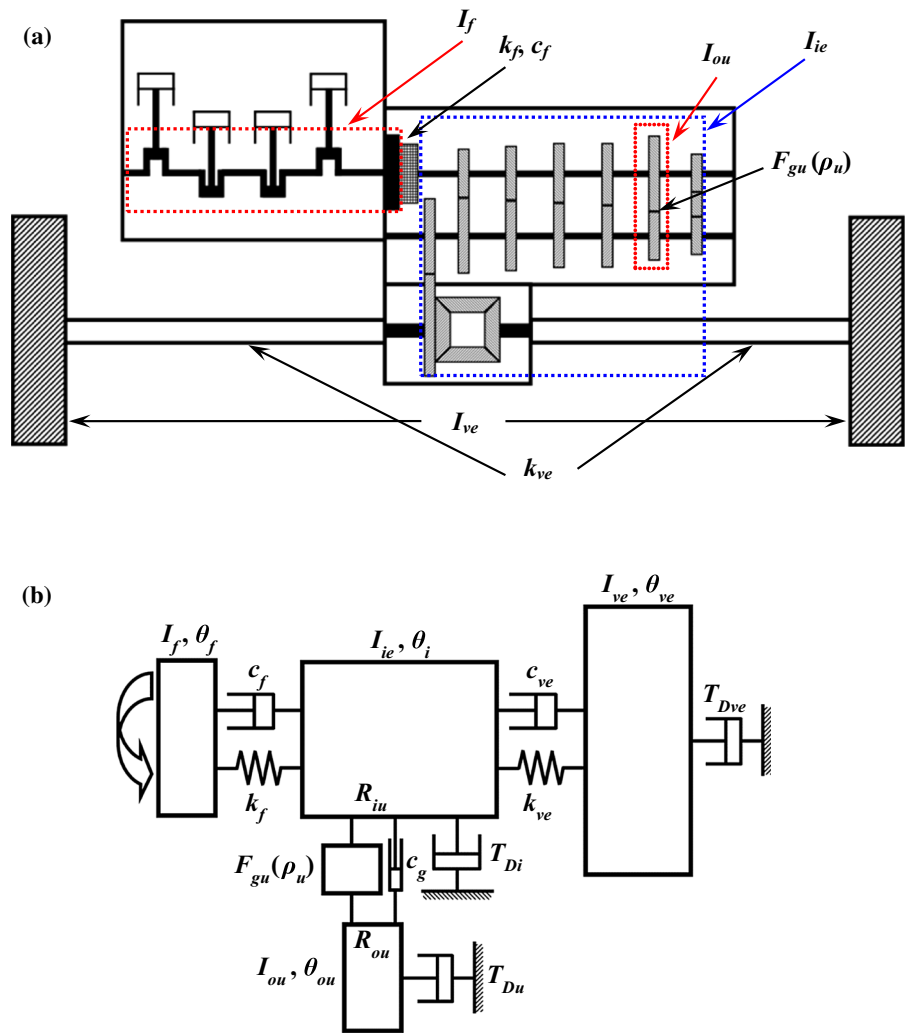
$$\underline{\mathbf{f}}_{\mathbf{n}}(\underline{\boldsymbol{\theta}}) = \begin{bmatrix} 0 \\ R_{iu} F_{gu}(\rho_u) + T_{Di} \\ R_{ou} F_{gu}(\rho_u) - T_{Du} \\ T_{Dve} \end{bmatrix},$$

$$\underline{\mathbf{T}}_{\mathbf{E}}(\mathbf{t}) = \begin{bmatrix} T_E(t) \\ 0 \\ 0 \\ 0 \end{bmatrix}. \tag{1e,f}$$

$$F_{gu}(\rho_u) \tag{1g,h}$$

$$= k_g \delta(\rho_u), \quad \rho_u = R_{iu} \theta_{ie}(t) + R_{ou} \theta_{ou}(t).$$

Fig. 1 A torsional system with a front-engine and front-wheel layout: **a** a practical driveline with a manual transmission; **b** a schematic diagram with gear mesh force



The angular displacements $\theta(t) = [\theta_f \ \theta_{ie} \ \theta_{ou} \ \theta_{ve}]^T$ are defined as the absolute motions of the flywheel, input shaft, unloaded gear, and vehicle, respectively. $F_{gu}(\rho_u)$ is defined as the gear mesh force between the input shaft and unloaded gear in terms of the translational relative displacement (ρ_u) of the unloaded gear pair. $\delta(\rho_u)$ reflects the discontinuous characteristics of the gear motions. To simulate the torsional system using the HBM, non-dimensionalization of the time scale was conducted by employing a new variable $\psi = \omega t$ and the time range $0 \leq t \leq \nu\tau$ for one period τ can be mapped onto $0 \leq \psi \leq 2\pi$ [9]. Here, ν indicates the sub-harmonic index [11, 12]. Also, when ω is parameterized, the normalized value $\bar{\omega} (\omega/\omega_N)$ is employed in order to avoid possible convergence problems. Here,

ω_N is the natural frequency relevant to vibro-impacts on the system shown in Fig. 1. In this study, 47.6 Hz is used for ω_N . Profound studies with respect to the relationship of simulation convergence to multiple choices of non-dimensionalization are beyond the scope of this research.

To calculate the system response, the input torque $T_E(t)$ is assumed as a sinusoidal excitation as follows.

$$T_E(t) = T_m + \sum_{i=1}^{N_{max}} T_{pi} \cos(i\omega_p t + \phi_{pi}). \quad (2)$$

The employed values of the input torques are as follows: the mean torque is $T_m = 168.9$ N m; the alternating torque is $T_{p1} = 251.53$ N m; and the phase is $\phi_{p1} = -1.93$. The system is assumed to be

Table 1 Employed parameters of the torsional system

Parameter	Value
Inertia (kg m ²)	
I_f (flywheel)	1.38×10^{-1}
I_{ie} (transmission)	1.49×10^{-2}
I_{ou} (unloaded gear)	5.23×10^{-4}
I_{ve} (wheels)	8.08×10^{-2}
Stiffness (N m rad ⁻¹)	
k_f (clutch damper)	9.20×10^2
k_{ve} (drive shaft)	3.95×10^2
Radius (mm)	
R_{iu} (unloaded gear on the input shaft)	45.9
R_{ou} (unloaded gear on the output shaft)	35.6

Table 2 Employed properties for the drag torques on the relevant sub-system

Drag torque	Value (Nm)
T_{Di} (drag torque on the input shaft)	1.22×10^2
T_{Du} (drag torque on the unloaded gear)	2.0
T_{Dve} (drag torque on the vehicle)	44.2

under steady state conditions with constant velocity. Thus, the drag torques of each sub-system shown in Fig. 1 are estimated by assuming that the summation of torques is equal to the mean torque value T_m as follows [1,2,9,10]:

$$T_m = T_{Di} + T_{Dve} + T_{Du} (R_{iu}/R_{ou}). \tag{3}$$

Figure 2a illustrates three different conditions of the gear pair for “contact on the driving side”, “no contact”, and “contact on the driven side”. When the dynamic behaviors of the gear pair show repetitive motions between “contact on the driving side” and “no contact” (or “contact on the driving side” and “contact on the driven side), the torsional system is severely affected by vibro-impacts such as gear rattle [1,2]. In general, this nonlinear dynamic effect is caused by the clearance between the driving and driven gears called backlash. The dynamic motions of the gear pairs are mathematically described as follows.

$$\delta(\rho_u) = \begin{cases} \rho_u + \frac{b}{2}, & \rho_u \leq -\frac{b}{2} \\ 0, & -\frac{b}{2} < \rho_u < \frac{b}{2} \\ \rho_u - \frac{b}{2}, & \rho_u \geq \frac{b}{2} \end{cases}. \tag{4}$$

The gear backlash is defined as b with a value of 0.1 mm. The backlash is a piecewise nonlinearity expressed using a unit step function $U(\rho_u)$ as follows:

$$\delta(\rho_u) = \left(\rho_u + \frac{b}{2}\right) + \left[\left(\rho_u - \frac{b}{2}\right) \times U\left(\rho_u - \frac{b}{2}\right) - \left(\rho_u + \frac{b}{2}\right) \times U\left(\rho_u + \frac{b}{2}\right)\right]. \tag{5}$$

Based on a prior study [8] and Eqs. (4, 5), the dynamic gear mesh forces are estimated using the Model I and II smoothening functions, which are mathematically described as follows. Here, the relative motions are defined as $\theta_{r1} = \rho_u - \frac{b}{2}$ and $\theta_{r2} = \rho_u + \frac{b}{2}$.

Model I

$$U_1(\theta_{r1}) \cong \frac{1}{2} [\tanh(\sigma_g \theta_{r1}) + 1], \tag{6a}$$

$$U_1(\theta_{r2}) \cong \frac{1}{2} [\tanh(\sigma_g \theta_{r2}) + 1], \tag{6b}$$

$$\delta_{gu1}(\rho_u) \tag{6c,d}$$

$$= \theta_{r2} + [\theta_{r1}U_1(\theta_{r1}) - \theta_{r2}U_1(\theta_{r2})],$$

$$F_{gu1}(\rho_u) = k_g \delta_{gu1}(\rho_u).$$

Model II

$$U_2(\theta_{r1}) \cong \frac{1}{2} \left[\frac{2}{\pi} \arctan(\sigma_g \theta_{r1}) + 1 \right], \tag{7a}$$

$$U_2(\theta_{r2}) \cong \frac{1}{2} \left[\frac{2}{\pi} \arctan(\sigma_g \theta_{r2}) + 1 \right], \tag{7b}$$

$$\delta_{gu2}(\rho_u) \tag{7c,d}$$

$$= \theta_{r2} + [\theta_{r1}U_2(\theta_{r1}) - \theta_{r2}U_2(\theta_{r2})],$$

$$F_{gu2}(\rho_u) = k_g \delta_{gu2}(\rho_u).$$

where $k_g (= 2.7 \times 10^8 \text{ N m}^{-1})$ is the gear mesh stiffness, and $b (= 0.1 \text{ mm})$ is the gear backlash. Also, 1×10^{10} is employed for the value of σ_g for both Models I and II in order to simulate the sudden change of gear mesh forces under vibro-impact conditions in a practical system [1,2]. Figure 3 compares the difference between two numerical models for unit step functions using $U_1(\theta_{r1})$ and $U_2(\theta_{r1})$. When those models are compared, $U_1(\theta_{r1})$ shows steeper change rather than $U_2(\theta_{r1})$. In order to compare two numerical unit step functions, 1×10^3 is employed for the value of σ_g .

Figure 2b indicates the expected gear mesh force with the relationship of $F_{gu}(\rho_u)$ versus ρ_u using Model I or II. Based on the expected gear mesh force shown

Fig. 2 Nonlinear behaviors of the gear pairs under the engine excitation condition: **a** gear contact conditions under vibro-impacts; **b** expected gear mesh force along with gear backlash regime

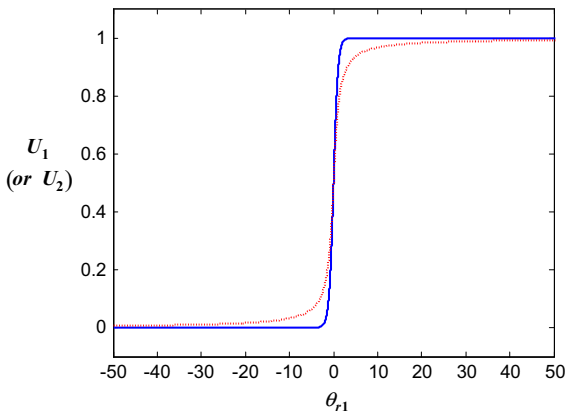
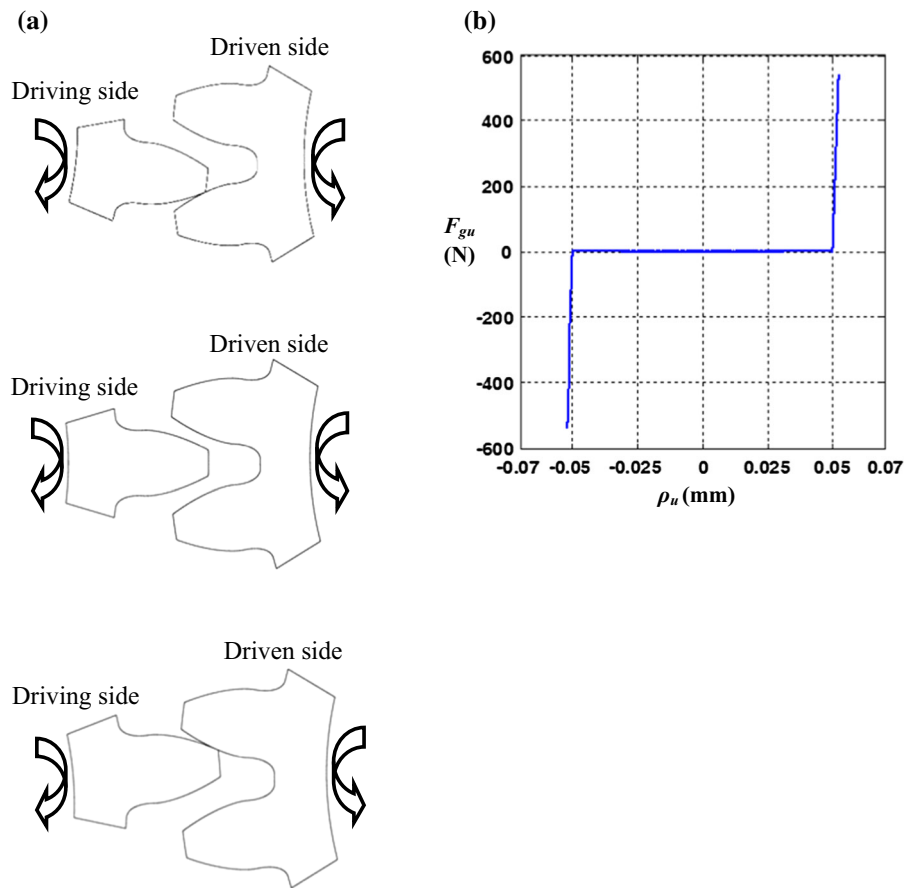


Fig. 3 Comparison of two numerical unit step functions. *Solid blue line* $U_1(\theta_{t1})$; *dotted red line* $U_2(\theta_{t1})$. (Color figure online)

in Fig. 2b, the dynamic conditions of the forces are abruptly changed at $b/2 (= 0.05 \text{ mm})$ or $-b/2$, which causes strong stiffness problems in the simulation with sudden changes of the values for the gear mesh force estimation. Differences between the two mod-

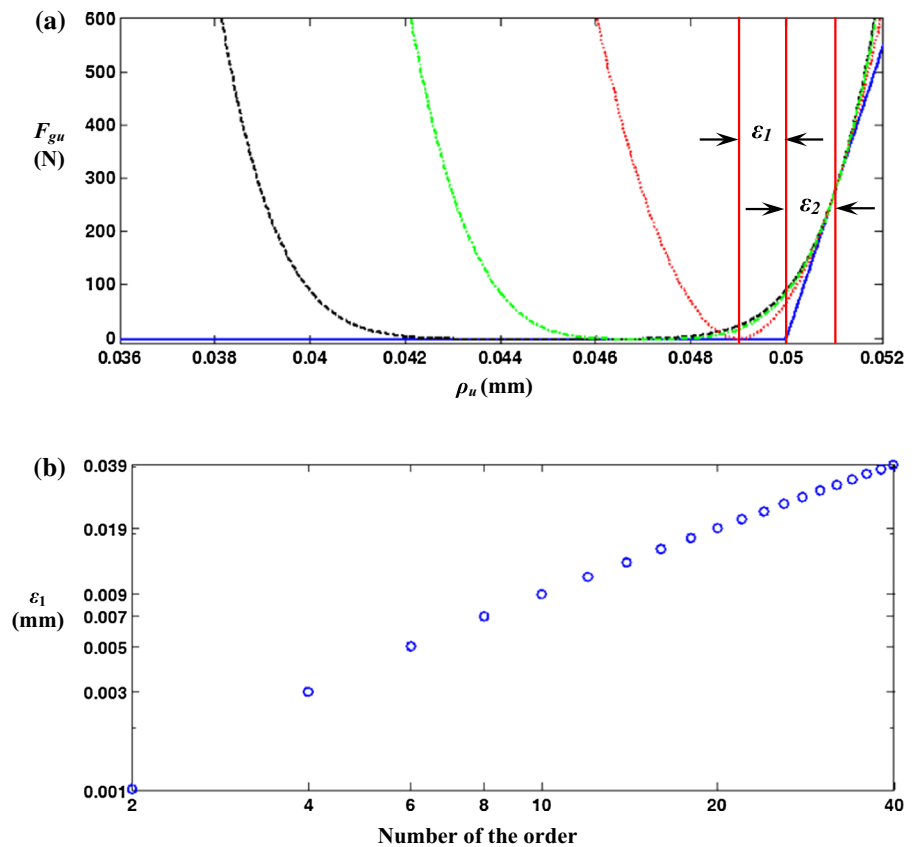
els will be explained with a new smoothening function model.

3 Mathematical model of the smoothening function for gear mesh force

3.1 Development of the new smoothening function

Based on the prior smoothening function models, it is not possible to change the smoothened areas and values on the left (or right) side from the gear backlash $b/2$. The smoothening effect from the prior models given in Eqs. (4)–(7) are only dependent on the value of σ_g . Also, several problems have been observed: (1) unrealistic gear mesh forces by using Model I can be expected contrary to the measured ones which will be explained later; (2) computational problems using Model II can be caused with respect to the numerical convergence since fractional functions should be used when the Jacobin matrix is employed. Thus, the new smoothening func-

Fig. 4 Smoothing function Model III for different orders: **a** comparison of Model III with different orders; **b** ε_1 with different orders. Dotted red line $n = 2$; dashed with dot green line $n = 4$; dashed black line $n = 6$. (Color figure online)



tion model should be considered to avoid those problems observe in the prior studies as well as to have more flexibility to adjust the smoothing areas with asymmetrical way for both left and right sides from $b/2$. For satisfying this conditions, this study employs an n th order polynomial function since it is easy to adapt more flexible smoothing effects than the models in Eqs. (4)–(7). Here, small amounts of the displacements ε_1 and ε_2 at $b/2$ or $-b/2$ can be defined for gear mesh force estimation under smoothing conditions. These factors also lead to the smoothing change of $F_{gu}(\rho_u)$ in the stiff regimes, as illustrated in Fig. 4a. Model III is mathematically described as follows:

$$\delta(\rho_u) = \begin{cases} \rho_u + \frac{b}{2}, & \rho_u < -B_2 \\ -\alpha(\rho_u + B_1)^n, & -B_2 \leq \rho_u < -B_1 \\ 0, & -B_1 \leq \rho_u < B_1 \\ \alpha(\rho_u - B_1)^n, & B_1 \leq \rho_u < B_2 \\ \rho_u - \frac{b}{2}, & \rho_u \geq B_2 \end{cases}, \tag{8a}$$

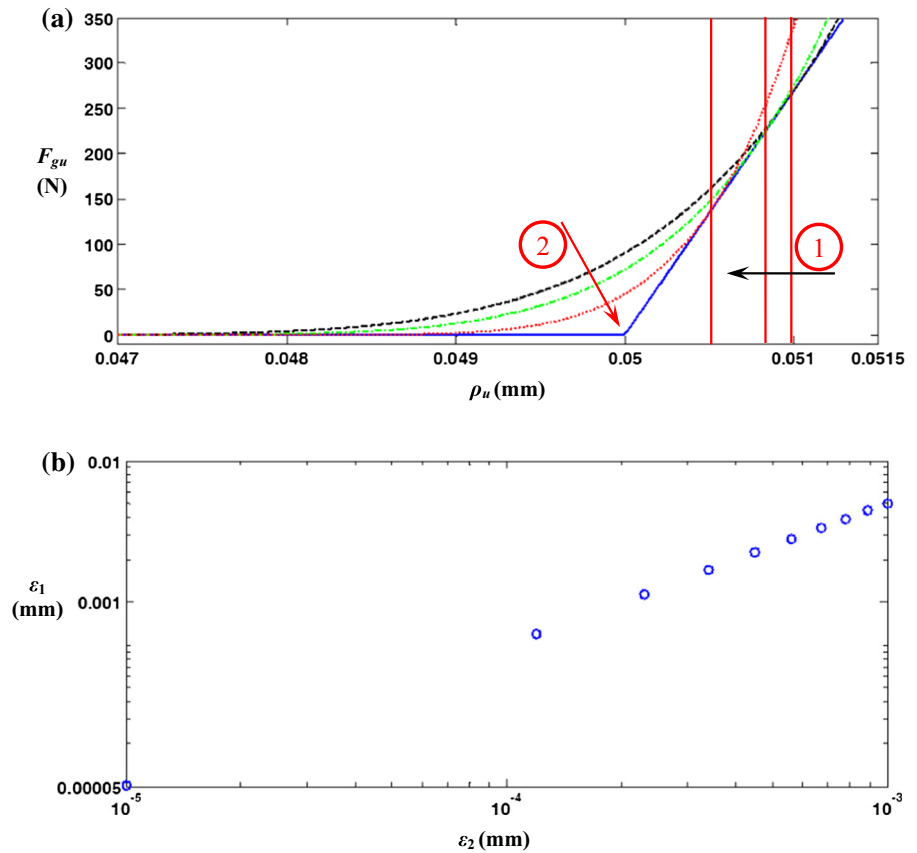
$$B_2 = \frac{b}{2} + \varepsilon_2, \quad B_1 = B_2 - n\varepsilon_2, \quad \varepsilon_1 = \frac{b}{2} - B_1, \tag{8b-e}$$

$$\alpha = \frac{\varepsilon_2}{(B_2 - B_1)^n}.$$

where ε_1 and ε_2 are the distances from $b/2$ or $-b/2$, as shown in Fig. 4a. To create the smoothing changes in the area of $b/2$ or $-b/2$, the n th order polynomial functions $\alpha(\rho_u - B_1)^n$ or $-\alpha(\rho_u + B_1)^n$ are used, respectively.

As described in Eq. (8b), B_2 is estimated simply by adding ε_2 to $b/2$, where ε_2 can be determined arbitrarily. Thus, if ε_2 is small, ε_1 becomes small and makes the gear contact motions stiff. In Eqs. (8a)–(8e), the properties B_1 and α are determined based on several conditions. For example, the polynomial function should first contact the linear line of $k_g(\rho_u - \frac{b}{2})$ (or $k_g(\rho_u + \frac{b}{2})$) at the range of $\rho_u \geq B_2$ (or $\rho_u < -B_2$) tangentially. Thus, the derivative value at the contact point is the same as the slope of tangent line k_g , which is the gear mesh stiffness. Second, the polynomial function $\alpha\rho_u^n$ (or $-\alpha\rho_u^n$) is shifted with the amount of B_1 (or $-B_1$) on the ρ_u axis. Third, the order n must be even, since an odd number can cause $F_{gu}(\rho_u)$ to be estimated below 0N in the area

Fig. 5 Smoothening function Model III with ε_2 : **a** comparison of Model III with different values of ε_2 ; **b** ε_1 versus ε_2 . Dotted red line $\varepsilon_2 = 5 \times 10^{-4}$; dashed with dot green line $\varepsilon_2 = 8 \times 10^{-4}$; dashed black line $\varepsilon_2 = 1 \times 10^{-3}$. (Color figure online)



between $-b/2$ and $b/2$, which does not occur in a practical system. Thus, from Eqs. (8a)–(8e), the gear mesh force $F_{gu}(\rho_u)$ is derived as follows. Here, the relative motions are defined as $\theta_{r1} = \rho_u - \frac{b}{2}$, $\theta_{r2} = \rho_u + \frac{b}{2}$, $\theta_{p1} = \rho_u - B_1$, $\theta_{p2} = \rho_u - B_2$, $\theta_{n1} = \rho_u + B_1$ and $\theta_{n2} = \rho_u + B_2$.

$$\begin{aligned}
 U_3(\theta_{n1}) &\cong \frac{1}{2} [\tanh(\sigma_g \theta_{n1}) + 1], \\
 U_3(\theta_{n2}) &\cong \frac{1}{2} [\tanh(\sigma_g \theta_{n2}) + 1],
 \end{aligned} \tag{9a}$$

$$\begin{aligned}
 U_3(\theta_{p1}) &\cong \frac{1}{2} [\tanh(\sigma_g \theta_{p1}) + 1], \\
 U_3(\theta_{p2}) &\cong \frac{1}{2} [\tanh(\sigma_g \theta_{p2}) + 1],
 \end{aligned} \tag{9b}$$

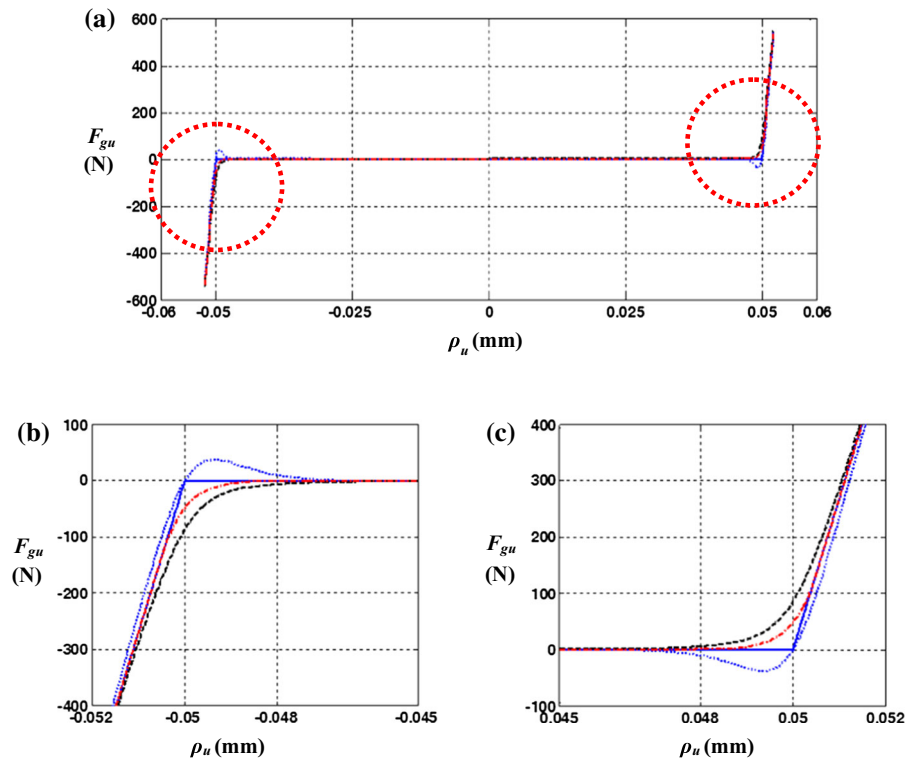
$$\begin{aligned}
 \delta_{gu3}(\rho_u) &= \theta_{r2} + [\theta_{r1} \cdot U(\theta_{p2}) - \theta_{r2} \cdot U(\theta_{n2})] \\
 &\quad - \alpha \theta_{n1}^n [U(\theta_{n2}) - U(\theta_{n1})] + \alpha \theta_{p1}^n [U(\theta_{p1}) \\
 &\quad - U(\theta_{p2})],
 \end{aligned} \tag{9c}$$

$$\begin{aligned}
 F_{gu3}(\rho_u) &= k_g \delta_{gu3}(\rho_u).
 \end{aligned} \tag{9d}$$

Figure 4 compares the smoothening effects with regard to n . As shown in Fig. 4a, when n is increased to a certain value such as 6, ε_1 becomes large, as shown with a dashed line in Fig. 4a. This indicates that the smoothening effect starts more gradually than when n is reduced to 2, as shown with the dotted line. Thus, increasing the smoothening factor ε_1 leads to improved numerical convergence conditions by avoiding sudden changes of the gear mesh stiffness from 0 N m^{-1} to $2.7 \times 10^8 \text{ N m}^{-1}$.

Accordingly, Fig. 4b shows the relationship between ε_1 and n . When n becomes large, the distance ε_1 from $b/2$ or $-b/2$ increases. As explained previously, ε_2 also affects the smoothening change of the gear contact motion. Figure 5 illustrates the relationship of ε_2 to the smoothening effect. Figure 5a shows that $F_{gu}(\rho_u)$

Fig. 6 Comparison of three smoothing functions: **a** comparison of the smoothing function Models I, II, and III; **b** smoothing effect with Models I, II, and III in the area of $-b/2$; **c** smoothing effect with Models I, II, and III in the area of the $b/2$. Dotted blue line Model I; dashed black line Model II; dashed with dot red line Model III. (Color figure online)



becomes stiff, as indicated by ②, when the value of the ε_2 is reduced, indicated by ①. When ε_2 is a small value such as 1×10^{-5} , the dynamic characteristics of the gear contact are closer to the practical motions than when ε_2 is a large value such as 1×10^{-3} . Also, ε_1 is estimated based on ε_2 as described in Eqs. (8b)–(8e), from which the smoothing effect can be mathematically determined.

3.2 Comparison of the smoothing function models

Figure 6 shows the differences between Models I, II, and III. First, Model I uses a hyperbolic tangent function, as described in Eq. (6). This function shows a negative or positive overshoot in the area of $b/2$ or $-b/2$. Thus, Model I includes severe error since the gear mesh forces never fall into the negative or positive values before the gear contact occurs at $b/2$ or $-b/2$. Models II and III do not show the same problem as Model I, which is shown in Fig. 6. However, Models II and III slightly differ with respect to the areas where the smoothing changes of $F_{gu}(\rho_u)$ occur. For example, the smoothed $F_{gu}(\rho_u)$ in Model II is determined only by the factor σ_g described in Eq. (7), but

it is difficult to anticipate the location of ε_1 and ε_2 . In contrast, ε_1 with Model III is exactly estimated as long as ε_2 is defined by Eq. (8). Moreover, ε_1 is determined based on the order n from Eqs. (8c) and (8d). Thus, the dynamic characteristics with the smoothing effect in Model III are more definite than in Models I and II.

HBM was used to investigate the effect of the smoothing functions with different models [3, 5–25]. The development and basic process of the HBM for systems with one or more DOF have been introduced in prior studies [9, 10]. Figure 7 compares the HBM results using Models I, II, and III. The system responses in the figure were estimated by the maximum number of harmonics $N_{max} = 6$. The value of σ_g is 1×10^{10} for all of the models as described in Sect. 2.2, and ε_2 and n for Model III are 1×10^{-5} and 20, respectively. Here, $\delta_2(t) = R_{iu}\theta_{ie}(t) + R_{ou}\theta_{ou}(t)$ is the relative displacement between the input shaft and unloaded gear, and $\delta_{2(max)}$, $\delta_{2(mean)}$, and $\delta_{2(min)}$ are the maximum, mean, and minimum values of $\delta_2(t)$, respectively, in one period of time responses, as shown in Fig. 7 with frequency sweeping conditions.

The results for all the models correlated well with each other. Also, single-sided and double-sided vibro-

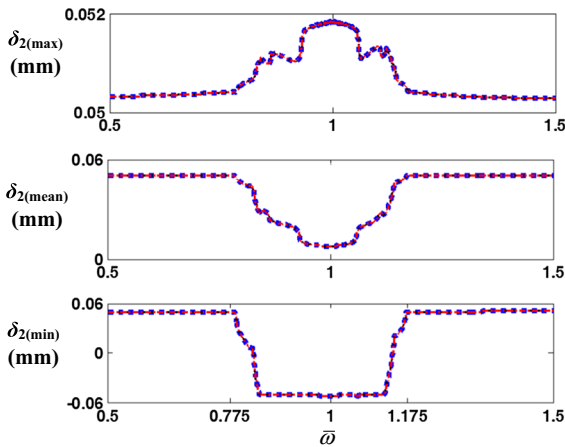


Fig. 7 Comparison of the HBM ($N_{\max} = 6$) using smoothening function Models I, II, and III. *Dotted blue line* Model I; *dashed black line* Model II; *dashed with dot red line* Model III. (Color figure online)

impacts were observed clearly between $\bar{\omega} = 0.775$ and $\bar{\omega} = 1.175$. $\bar{\omega}$ is the frequency normalized using the natural frequency $\omega_N = 47.6$ Hz. When the number of harmonics is increased, the simulation using HBM shows many differences between Models I, II, and III.

Table 3 describes the feasibility of the simulation using each smoothening function model. The nonlinear responses with $N_{\max} = 6$ can be successfully simulated by using HBM. However, when the number of harmonics increased to more than 6, the simulation never accomplished calculating the nonlinear responses using Models I and II with any values of σ_g , as indicated in Table 3. However, the simulation with Model III was successfully conducted with all numbers of harmonics with some discrepancies. Based on the results described in Table 3, the simulation cases of Model III are more feasible for adaption with different harmonic conditions than Models I and II, since Model

Table 3 Feasibility analysis of the smoothening functions with the number of harmonics

Number of harmonics	Smoothening function		
	Model I	Model II	Model III
6	O	O	O
8	×	×	O
10	×	×	O
12	×	×	O

III includes more factors to manage the smoothening conditions, such as ε_2 , n , and σ_g based on Eq. (8).

4 Results

4.1 Comparison of the numerical analysis and HBM using the new smoothening function model

Figure 8 shows the comparison of the simulated results using numerical analysis (NS) and HBM in the frequency domain. When the computational time is compared, NS needs 2h for both up-(or down-)frequency sweeping. On the other hand, the calculation time with the HBM is 20 min. Both results reflect the vibro-impacts well. However, the HBM does not follow the NS results for the gear impact conditions. When the system has double-sided impacts between $\bar{\omega} = 0.8403$ and $\bar{\omega} = 1.122$, as indicated with a solid line in Fig. 8, severe differences between NS and HBM at $\delta_{2(\max)}$ are observed. This is caused by the limits of the HBM, since the number of harmonics is 6, which was examined in a prior study [9]. Since the vibro-impacts contain inherently impulsive responses, they include multiple numbers of harmonics and higher super-harmonic components than $N_{\max} = 6$ or 12 [2]. Also, the limit of the number of harmonics with HBM causes a discrepancy of the phase in the Fourier components in comparison with the FFT results from NS. These behaviors are shown in Figs. 9, 10, 11 and 12.

Figure 9 compares the relative displacements from HBM and NS in the time domain. Figure 9a shows

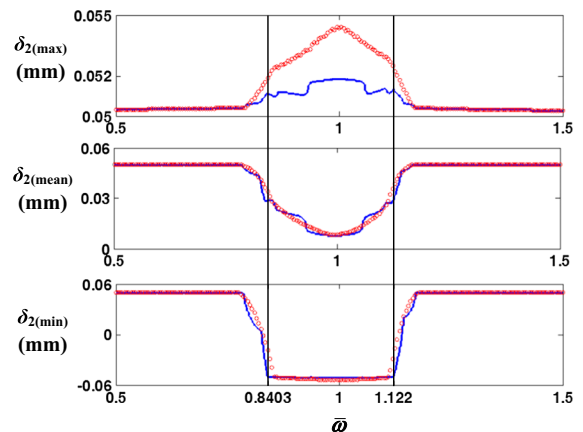


Fig. 8 Comparison of the HBM and NS results. *Solid blue line* HBM with Model III; *circled red line* NS. (Color figure online)

Fig. 9 Comparison of the relative displacements with HBM and NS under different excitation conditions in the time domain: **a** time histories of $\delta_1(t)$ and $\delta_2(t)$ at $\bar{\omega} = 0.825$; **b** time histories of $\delta_1(t)$ and $\delta_2(t)$ at $\bar{\omega} = 1.0$. *Solid blue line* NS; *dashed red line* HBM. (Color figure online)

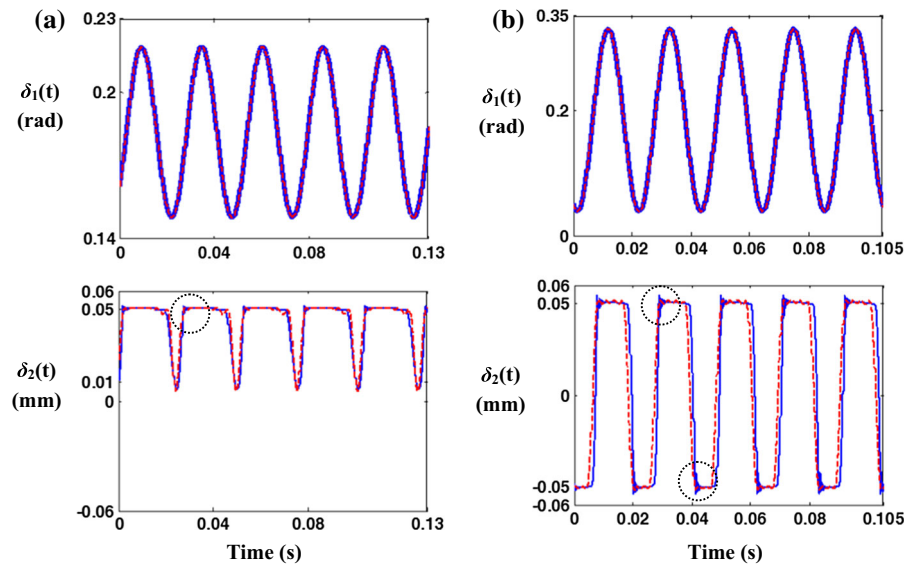
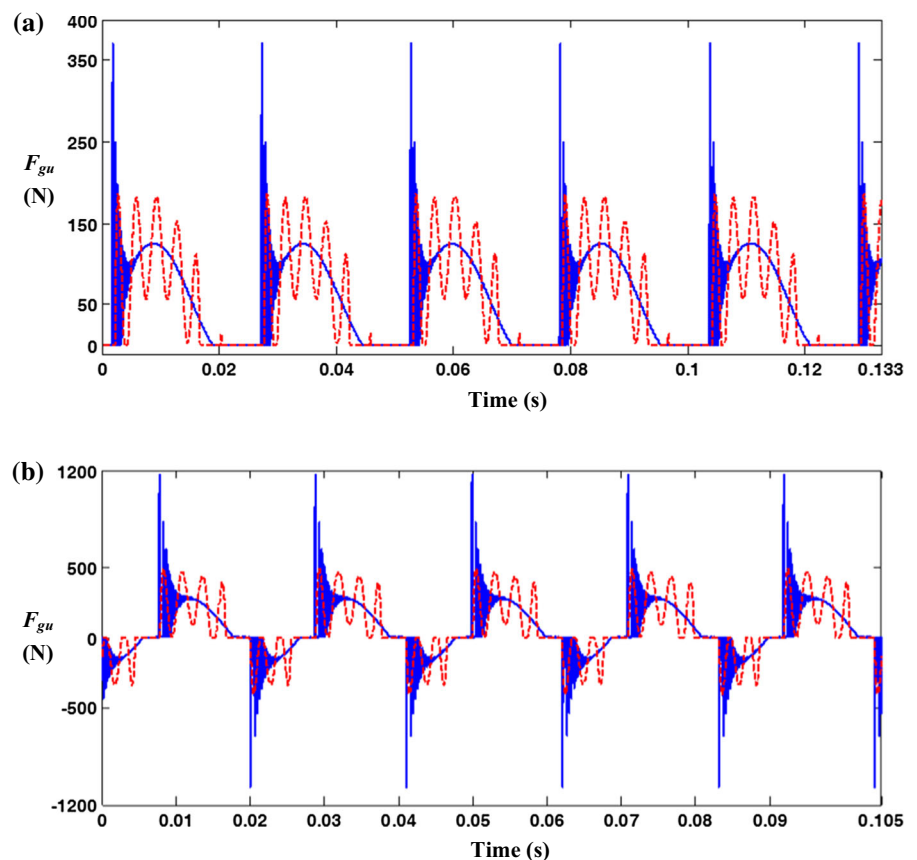


Fig. 10 Comparison of the gear mesh forces with HBM and NS under different excitation conditions in the time domain: **a** time histories of $F_{gu}(t)$ at $\bar{\omega} = 0.825$; **b** time histories of $F_{gu}(t)$ at $\bar{\omega} = 1.0$. *Solid blue line* NS; *dashed red line* HBM. (Color figure online)



the results estimated at the excitation frequency of $\bar{\omega} = 0.825$, and Fig. 9b shows the results at $\bar{\omega} = 1$. Here, $\delta_1(t) = \theta_f(t) - \theta_{ie}(t)$ is the relative displacement between the flywheel and input shaft. Figure 9a shows

the comparisons of time histories under single-sided impact. The results of $\delta_1(t)$ from HBM and NS correlated well with each other. However, $\delta_2(t)$ shows minor differences, especially when the gear impact occurs

Fig. 11 Comparison of the relative displacements with HBM and NS under different excitation conditions in the frequency domain: **a** FFT results of $\delta_2(t)$ at $\bar{\omega} = 0.825$; **b** FFT results of $\delta_2(t)$ at $\bar{\omega} = 1.0$. Line with square red line NS; line with circle blue line HBM. (Color figure online)

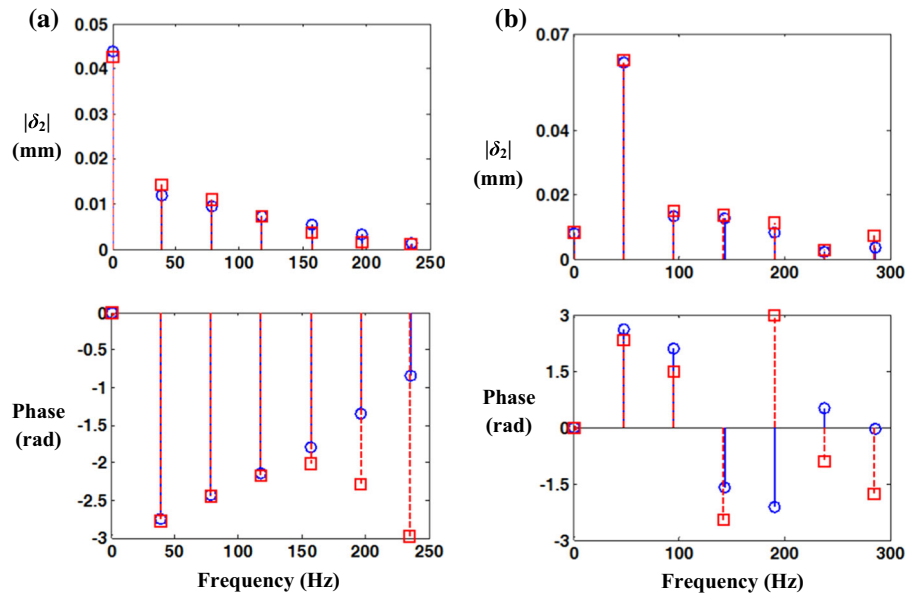
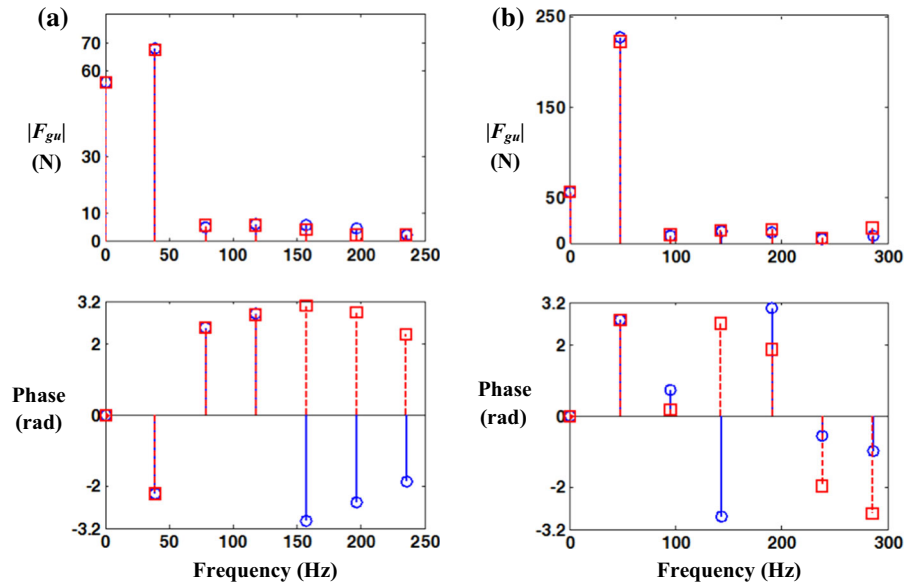


Fig. 12 Comparison of the gear mesh forces with HBM and NS under different excitation conditions in the frequency domain: **a** FFT results of $F_{gu}(t)$ at $\bar{\omega} = 0.825$; **b** FFT results of $F_{gu}(t)$ at $\bar{\omega} = 1.0$. Line with square red line NS; line with circle blue line HBM. (Color figure online)



at $b/2$, as indicated by the dotted circle. In general, NS includes the impulsive responses well, since it is based on time domain analysis. This discrepancy of HBM is clearly observed when the system has double-sided impacts, as shown in Fig. 9b. Since the double-sided impact conditions are worse than the single-sided impacts, the impulsive responses are more dominant than the single-sided impacts, as indicated with dotted lines.

Figure 10 compares the gear mesh forces $F_{gu}(t)$ at two different excitation conditions at $\bar{\omega} = 0.825$ and $\bar{\omega} = 1$. Figure 10a, b simulates the single-sided and

double-sided impacts, respectively. For example, the dynamic behaviors of $F_{gu}(t)$ shown in Fig. 10a follow the zero values at the first time, and then they go up to high values around 370N on only the positive side. However, the NS results show more density of the impulsive force responses than the HBM results. The result of $F_{gu}(t)$ shown in Fig. 10b depicts the double-sided impacts. In a similar manner to the single-sided impacts in Fig. 10b, $F_{gu}(t)$ from NS shows more density at the moment when the gear impacts occur. The dynamic behaviors of $F_{gu}(t)$ for the double-sided impacts show repetitive motions on both the positive

and negative sides. Also, the peak values for the double-sided impacts are higher than those for the single-sided impacts. Each local maxima is 1200 N, and the minima are -1100 N. Overall, both comparisons between HBM and NS show that NS results contain more super-harmonic components than the HBM results. These discrepancies between the different simulations were found based on the FFT results as well.

Figures 11 and 12 compare the FFT results of $\delta_2(t)$ and $F_{gu}(t)$ with 6 harmonics. Figure 11 shows the comparison of the FFT results from $\delta_2(t)$ based on HBM and NS. When the single-sided impact occurs, severe discrepancies between the two simulations are not observed, as shown in Fig. 11a. However, the phase values at the 4th, 5th, and 6th harmonics do not correlate well. The FFT results for the double-sided impact show severe differences in the phases between HBM and NS for all ranges of harmonic components, which is clearly observed in Fig. 11b.

Figure 12 compares the FFT results of $F_{gu}(t)$ for both single- and double-sided impacts. In a similar manner to the FFT results of $\delta_2(t)$, the FFT results of $F_{gu}(t)$ for the single-sided impact in Fig. 12a show that the phase of the 4th, 5th, and 6th harmonics from HBM is not correlated well with those of the NS. Also, the harmonic spectra of the phases for the double-sided impact for HBM show severe differences from the results of the NS, except for the phase component at the 1st harmonic. Thus, based on the results in Figs. 11 and 12, the limit number of harmonics employed for HBM causes discrepancies for the time response estimations in which high super-harmonic components such as those beyond 500 Hz are not included [2]. However, the HBM can still be used to assess the vibro-impact phenomena in frequency sweeping conditions in both frequency and time domains, as shown in Figs. 8 and 9.

4.2 Effect of the ϵ_2 values in HBM with respect to vibro-impacts

In general, clearance nonlinearity inherently contains discontinuity, which makes gear mesh forces suddenly change from zero to high stiffness levels, or vice versa. These dynamic characteristics can cause convergence problems or consume calculation time. Thus, the smoothening functions are employed to resolve these problems. Based on Eqs. (8) and (9), smoothen-

ing effects can be adapted using different levels of ϵ_2 . For example, as the level of ϵ_2 is increased, the relative displacements $\delta_2(t)$ or gear mesh forces $F_{gu}(\rho_u)$ of the unloaded gear pair are estimated more smoothly than when ϵ_2 is reduced.

Figure 13 compares the dynamic behaviors of $\delta_2(t)$ estimated with different ϵ_2 . The simulations with $\epsilon_2 = 5 \times 10^{-5}$ and 1×10^{-4} show discrepancies with the results of $\epsilon_2 = 1 \times 10^{-5}$. The HBM results with $\epsilon_2 = 1 \times 10^{-5}$ are assumed to simulate the vibro-impacts reasonably except for the impulsive responses, as described in Sect. 4.1. As shown in Fig. 13, the effect of ϵ_2 is observed clearly in both the $\delta_{2(max)}$ and $\delta_{2(RMS)}$ results, where $\delta_{2(RMS)} = \sqrt{\frac{1}{m} \{ \delta_2^2(t_0) + \delta_2^2(t_1) + \dots + \delta_2^2(t_{m-1}) \}}$, and m is the number of steps in one period of the time history.

As indicated with the solid line in Fig. 13a, b, the vibro-impacts occur between $\bar{\omega} = 0.775$ and $\bar{\omega} = 1.175$. However, the simulation results with $\epsilon_2 = 1 \times 10^{-4}$ show severe discrepancy, since $\delta_{2(max)}$ has values below $b/2 = 0.05$ mm in the frequency ranges with $\bar{\omega} < 0.775$ and $\bar{\omega} > 1.175$. The results of $\delta_{2(RMS)}$ with both $\epsilon_2 = 5 \times 10^{-5}$ and $\epsilon_2 = 1 \times 10^{-4}$ are very different from those with $\epsilon_2 = 1 \times 10^{-5}$ at $\bar{\omega} < 0.775$ and $\bar{\omega} > 1.175$. In the simulated result of $\delta_{2(max)}$ with $\epsilon_2 = 1 \times 10^{-4}$ shown in Fig. 13a, $\delta_{2(max)}$ levels below $b/2 = 0.05$ mm reflect that the gear pairs never contact each other, which is not practically reasonable for the physical system. Thus, to achieve reasonable simulation results, ϵ_2 should be determined within a certain range. For this study, all of the HBM results are estimated with $\epsilon_2 = 1 \times 10^{-5}$.

Figure 14 and Table 4 show the relationship of ϵ_2 values considered with the number of harmonics. Figure 14a shows the HBM results with different numbers of harmonics N_{max} . As N_{max} is increased, the discrepancies become severe compared with the HBM results with $N_{max} = 6$. The reason why the accuracy of the

Table 4 Employed minimum values of ϵ_2 for the smoothening function Model III with the number of harmonics

Number of harmonics	ϵ_2
6	1×10^{-7}
8	1×10^{-3}
10	2×10^{-3}
12	3×10^{-3}

Fig. 13 Comparison of the HBM results under different values of ε_2 with the smoothening function Model III: **a** max, mean, and min values of $\delta_2(t)$ with ε_2 ; **b** RMS values of $\delta_2(t)$ with ε_2 . *Solid blue line* $\varepsilon_2 = 1 \times 10^{-5}$; *dotted red line* $\varepsilon_2 = 5 \times 10^{-5}$; *dashed black line* $\varepsilon_2 = 1 \times 10^{-4}$. (Color figure online)

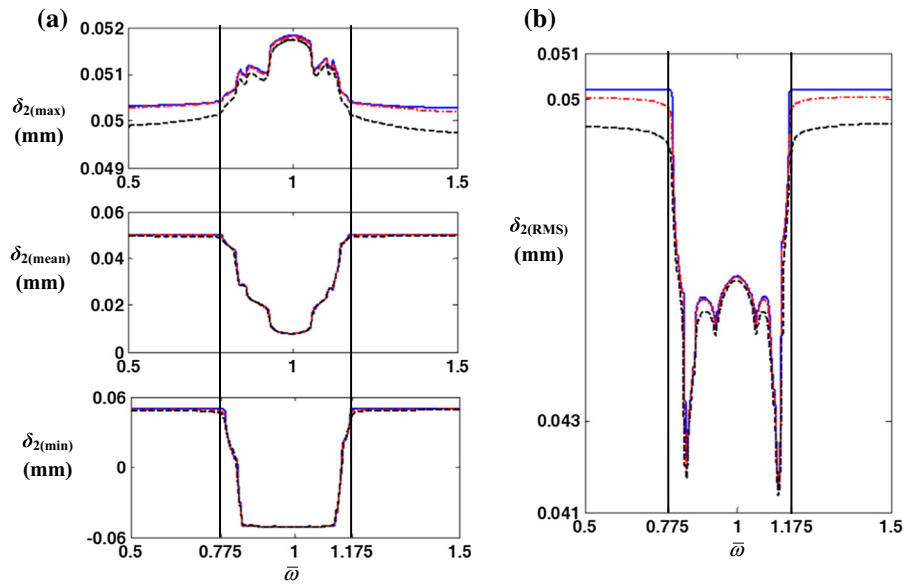
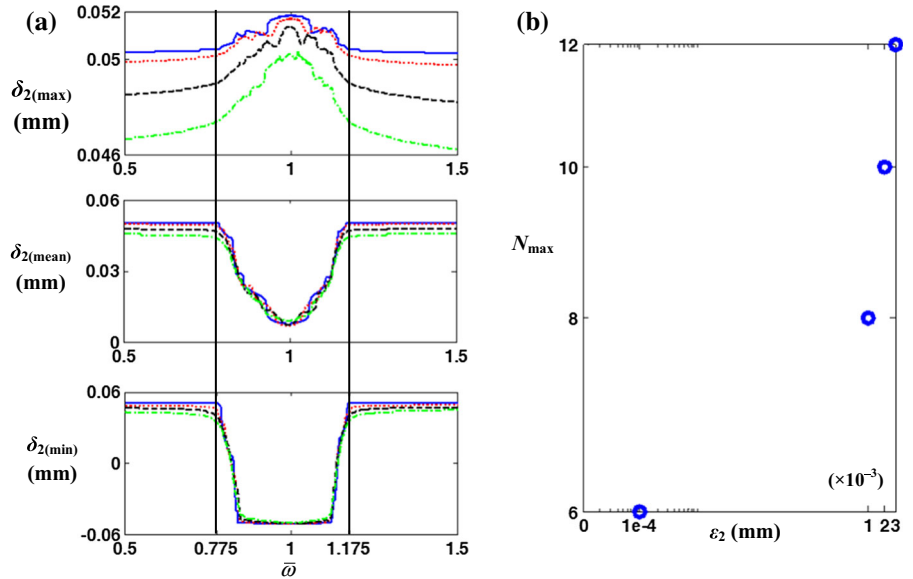


Fig. 14 Effect of ε_2 with number of harmonics: **a** comparison of the max, mean, and min values of $\delta_1(t)$ with different ε_2 values and numbers of harmonics; **b** number of harmonics versus available minimum values of ε_2 . *Solid blue line* $N_{\max} = 6$; *dotted red curve* $N_{\max} = 8$; *dashed black line* $N_{\max} = 10$; *dashed with dot green line* $N_{\max} = 12$. (Color figure online)



HBM results dropped is the ε_2 value. The ε_2 values versus N_{\max} are indicated in Fig. 14b and described in Table 4. To overcome the convergence problems as N_{\max} is increased, ε_2 must be larger than 1×10^{-5} with $N_{\max} = 6$. For example, the minimum value of ε_2 considered for $N_{\max} = 6$ is 1×10^{-7} . However, ε_2 should be equal to or larger than 1×10^{-3} , when N_{\max} is increased to more than 6, as described in Table 4. Thus, due to the increase in ε_2 , the convergence problems are resolved. On the other hand, large ε_2 causes the discrepancies shown in Fig. 14a.

5 Conclusion

New smoothening functions have been suggested and compared with other models in prior studies [8]. Using this model, the limits of HBM have been examined with respect to the impulsive responses in the time histories of $\delta_2(t)$, $F_{\text{gu}}(t)$, and the FFT results. The specific contributions of this study are summarized as follows. First, the new smoothening function model has been suggested by using an n th order polynomial function. This model could achieve smoothening changes

of the gear mesh force in the discontinuous regime, which leads to solving the convergence problems. Second, the feasibility of the smoothening factor ε_2 has been examined. Increasing the value of ε_2 assured that the convergence problems were removed. However, ε_2 beyond certain values caused severe errors, especially in the frequency ranges where the vibro-impacts do not occur. Third, the limit of the HBM with respect to the impact phenomena such as gear rattle has been investigated. In the comparison of the HBM with NS, the vibro-impact behaviors simulated using HBM were still revealed, even though all of the impulsive response components were not included, which is directly due to the limited number of harmonics.

Based on the study of N_{\max} versus ε_2 , the Model III smoothening function still has a discrepancy, even though it successfully overcomes the convergence problems. Thus, the development of multiple non-dimensionalization methods or the investigation of analytical formulations for clearance-type nonlinearities for HBM is the potential subjects for further study.

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