RESEARCH PAPER

Influence of Gd³⁺ ion doping on structural, optical, **and magnetic properties of (Mg–Ni–Co) nanoferrites**

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Abstract The characterization of (Mg–Ni–Co) nanoferrites doped with small diferent concentrations of gadolinium rare earth metal (Gd^{3+}) , prepared by the chemical co-precipitation method, is reported in this study. The XRD results reveal that the crystallite sizes, which are close to the particle size detected from TEM analysis, decrease upon adding Gd^{3+} . Furthermore, the lattice constant *a* increases from 8.351 to 8.367 \AA as the Gd^{3+} concentration (*x*) increases from 0.00 to 0.08.

Highlights

• (Mg–Ni–Co) nanoferrites doped with Gd^{3+} were synthesized by the coprecipitation method. • Doping the (Mg–Ni–Co) nanoferrites with Gd^{3+} increases the lattice constant and reduces the particle size and saturation magnetization. • Raman analysis has verifed the existence of active modes related to the spinel ferrite structure.

• Referring to the Mössbauer analysis, Gd^{3+} ions occupy the octahedral sites.

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HRTEM images reveal the presence of diferent grains with diferent lattice plane orientations. Furthermore, the appearance of concentric circles in the SAED patterns verifed the samples' polycrystalline nature. The presence of the active modes of the spinel structure was confrmed by the Raman analysis with a slight shift upon increasing the Gd^{3+} concentrations. Moreover, the photoluminescence (PL) analysis revealed the presence of a sharp peak in the UV region and structural defects in the prepared samples. Mössbauer spectroscopy and vibrating sample magnetometer (VSM) were used to examine the magnetic properties of the prepared samples. From the Mössbauer analysis, it was suggested that Gd^{3+} ions occupy the octahedral site. In addition, the saturation magnetization (M_s) and magnetocrystalline anisotropy (*K*) decreased from 31.87 to 15.76 emu/g and from 0.0057 to 0.0014 J/m³ with the increase of Gd^{3+} concentration from 0.00 to 0.08, respectively. Finally, diferent forms of the law of approach to saturation (LAS) for ferromagnetic materials were used for the ftting of M-H hysteresis loops at high applied felds. The obtained results show that the LAS model, mainly $M = M_s \left(1 - \frac{b}{H^2}\right) + \chi H$, can

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effectively describe the doping of Gd^{3+} for all concentrations. This model exhibits well-fitting with high R^2 values and provides relevant values of the magnetic parameters.

Keywords Spinel ferrites · Gd doping · Raman spectroscopy · Mössbauer spectroscopy · Law of approach to saturation · Magnetism · Nanocomposites

Introduction

The unique structural properties of pure spinel nanoferrites, characterized by the formula $MFe₂O₄$ (where M represents metal cations such as Mg^{2+} , Ni^{2+} , and $Co²⁺$, can be attributed to their special resilience orientation and valence changeability of metal cations with rich surface-active sites [\[1](#page-16-0)]. However, certain modifcations are required to improve the physical properties of spinel nanoferrites, including low saturation magnetization and chemical stability [\[2](#page-16-1)].

Among various types of ferrites, magnesium nanoferrite is highly signifcant. It is characterized as a soft magnetic n-type semiconducting material with a cubic normal spinel structure. It is widely used in magnetic technologies, electronic materials for nanodevices, storage devices, and catalysis [\[3](#page-16-2)]. Due to its typical ferromagnetic characteristics, abundance in nature, strong electrochemical stability, low conductivity, and catalytic action, the soft inverse nickel ferrite is one of the most useful and technologically signifcant nanoferrites [\[4](#page-17-0)]. Cobalt ferrite nanoparticles are a well-known moderate hard magnetic material with a small crystal size, moderate-high saturation magnetization, strong magneto-crystalline anisotropy, and moderate coercivity [\[5](#page-17-1)[–7](#page-17-2)]. It crystallizes as the inverse spinel structure, and is used in magnetic recording applications such as audio and video tape, high-density digital recording disks, and physical and chemical durability [[8\]](#page-17-3).

The magnetic properties of nickel-substituted cobalt ferrite demonstrated that the change in coercivity, saturation magnetization, remanent magnetization, and squareness ratio was correlated with the amount of nickel. Because nickel has a lower magnetic moment compared to cobalt, the magnetic properties gradually decreased as the amount of nickel substitution increased. The coercive feld decreased from 7134 to 279 Oe at 30 K as the nickel content increased from 0 to 1 in $Co_{1-x}Ni_xFe_2O_4$ nanoparticles [\[9](#page-17-4)]. Moreover, the increase in the saturation magnetization and remanent magnetization increased as the nickel content increased up to 0.4 was owed to the increase in crystallite size [\[9](#page-17-4), [10\]](#page-17-5). Doping cobalt ferrite with magnesium reduced the value of saturation magnetization from 37 to 29 emu/g as *x* increased from 0 to 0.25 in $Co_{1-x}Mg_xFe_2O_4$ nanoparticles [\[11](#page-17-6)]. Furthermore, the remanent magnetization, the coercivity, and the magnetic moment were diminished as well [[11\]](#page-17-6). The conversion of the magnetic characteristics of cobalt ferrite from moderately hard to soft magnetic materials upon doping it with magnesium is reported in a previous study [[12\]](#page-17-7). Magnesiumdoped nickel ferrite showed a growth of crystallite size and a decrease in the anisotropy. In addition, the value of the lattice constant increased from 8.309 to 8.456 Å with the increase in Mg^{2+} concentration [\[13](#page-17-8)]. Ashrafizadeh et al. [\[14](#page-17-9)] reported that the substitution of magnesium by copper in the $(Mg_0, Zn_0, 5)$ ferrites with the composition of $Cu_xMg_{0.5-x}Zn_0sFe_2O_4$ did not signifcantly change the average magnetic hyperfne feld. The saturation magnetization and magnetic susceptibility were improved by increasing copper concentrations. In addition, the coercive force was decreased by copper ions dopant. Suo et al. [[15\]](#page-17-10) synthesized $\text{Ni}_{0.2}\text{Mg}_{0.2}\text{Co}_{x}\text{Zn}_{0.6-x}\text{Fe}_{2}\text{O}_{4}$ (*x* $= 0, 0.15, 0.3, 0.45, 0.45, 0.6)$ nanoferrites and found that they had soft magnetic properties. The remanent magnetization, coercivity, and anisotropy constant have increased with the increase of $Co²⁺$ ion content. Suo et al. [[16\]](#page-17-11) declared that the Bi^{3+} doping reduced the magnetic properties in the ferromagnetic Ni_{0.2}Mg_{0.2}Co_{0.6}Fe_{2−*x*}Bi_{*x*}O₄ (*x* = 0, 0.025, 0.05, 0.075, and 0.1) nanoparticles.

The incorporation of rare earth metal ions through doping in nanoferrites is an efective method employed to enhance their magnetic, optical, and structural properties [\[17](#page-17-12)]. These metal ions exhibit a wide range of magnetic moments and can be arranged in the tetrahedral and octahedral sites in the spinel lattice, leading to improvements in the overall characteristics of the nanoferrites and enabling their utilization in various applications [\[18](#page-17-13), [19](#page-17-14)]. In the orbital shape of rare earth ions, the unpaired 4f electrons energy level that appeared is the key aspect for originating magnetic anisotropy [\[20](#page-17-15)]. The fndings of the research done on the substitution of rare earth elements in spinel ferrites with diferent concentrations revealed that low concentrations of Gd^{3+} could improve the properties of spinel ferrites and create a distinctive product [\[21](#page-17-16)]. Shao et al. [\[22](#page-17-17)] reported that $Ni_{0.2}Mg_{0.2}Co_{0.6}Fe₂O₄$ nanoferrites doped with rare-earth ions, mainly La^{3+} , Ce^{3+} , $Pr³⁺$, and $Sm³⁺$, were suitable for magnetic storage and magnetic recording materials. It had high coercivity, a large squareness ratio (M_1/M_S) and large crystal anisotropy compared to the pure sample. Yu et al. [[23\]](#page-17-18) reported that the magnetic properties of Ni_{0.2}Mg_{0.1}Co_{0.7}Fe_{2−*x*}Y_{*x*}O₄ nanoferrites can be optimized by adding a small amount of Y^{3+} ions. Consequently, the saturation magnetization increased to 31.09 emu/g with the increase of Y^{3+} concentrations to 0.02 and then decreased to reach 25.54 emu/g with a further increase in Y^{3+} up to 0.08.

Among spinel ferrite nanoparticles, $MgFe₂O₄$ is recognized as a soft magnetic semiconductor, featuring a normal spinel structure $[24]$ $[24]$. Similarly, NiFe₂O₄ is a soft magnetic semiconductor with an inverse spi-nel structure [[25\]](#page-17-20). In contrast, CoFe_2O_4 possesses an inverse spinel structure and is classifed as a semihard material $[26]$ $[26]$. Thus, exploring the properties of ferrite nanoparticles, obtained by combining Mg, Co, and Ni metals in equal proportions proves to be an interesting study. Knowing that combing the metals in equal proportions can lead to a synergistic efect, where the prepared nanoparticles can exhibit superior properties compared to MgFe₂O₄, NiFe₂O₄, and CoFe_2O_4 nanoparticles. It is worth mentioning that the choice of the nanoferrite preparation method depends on both the desired properties of the ferrite and its intended application [[27,](#page-17-22) [28\]](#page-18-0). Furthermore, the influence of Gd^{3+} ion doping on the structural, optical, and magnetic properties of (Mg–Co–Ni) nanoferrites represents a novel approach that has not been explored in previous research. Therefore, spinel Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ nanoferrites, with $0.00 \le x \le 0.08$, were prepared by the co-precipitation method. This synthesis process is known for its simplicity and low cost and allows the control of the grain size and particle size distribution. Several characterization techniques were used to investigate the properties of the prepared samples. To examine the structural properties, X-ray powder difraction (XRD), transmission electron microscope (TEM), high-resolution transmission electron microscope (HRTEM), selected area electron difraction (SAED), and Raman spectroscopy were employed. However,

the photoluminescence (PL) analysis provides information about the optical properties, whereas the identifcation of the magnetic properties was achieved through Mössbauer spectroscopy and M-H loop analysis.

Experimental methods

Materials

Highly pure (above 98%) chemical reagents, mainly magnesium chloride hexahydrate $(MgCl₂·6H₂O)$, nickel (II) chloride hexahydrate (NiCl₂·6H₂O), cobalt (II) chloride hydrate $(CoCl₂·H₂O)$, iron (III) chloride hexahydrate $(FeCl₃·6H₂O)$, gadolinium (III) chloride hexahydrate $(GdCl₃·6H₂O)$, and sodium hydroxide (NaOH) were used for the preparation of Mg0.33Ni0.33Co0.33Fe2−*x*Gd*x*O4 nanoferrites with 0.00 $\leq x \leq 0.08$.

Synthesis method

Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ nanoparticles were prepared by the co-precipitation method. For the preparation of 25 g of pure $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_2O_4$ nanoparticles, 2.21, 2.58, 2.59, and 17.62 g of $MgCl_2·6H_2O$, $NiCl₂·6H₂O$, $CoCl₂·H₂O$, and $FeCl₃·6H₂O$, respectively, were dissolved in deionized water. However, to synthesize Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ nanoparticles with $0.01 \le x \le 0.08$, the amount of FeCl₃·6H₂O and $GdCl₃·6H₂O$ ranges between 16.74 and 17.51 g and between 0.12 and 0.96 g, respectively. The solutions were then mixed and stirred magnetically for 30 min. The mixture was titrated with NaOH (3M) until it reached a pH of 12. The solution was then maintained at 80 °C for 2 h while stirring magnetically at a constant rate. The obtained precipitates were washed several times with a solution composed of 75% deionized water and 25% ethanol until the pH of the fltrate dropped to 7. The collected samples were dried at 100 °C for 18 h. Finally, the samples were annealed for 4 h at 550 $°C$.

Characterization techniques

The structure of the prepared samples was determined from the XRD analysis. Therefore, 1 g of each of the prepared nanoparticles was ground to create a uniform fne-grain powder. This powdered sample was then placed in a plate-shaped holder and subsequently analyzed using a Panalytical Co. Xpert Pro difractometer. The difractometer was operated at 45 kV and 40 mA, with Cu-K_{α} X-ray source of radiation ($\lambda = 1.5406$ Å) in the range $20^{\circ} \le 2\theta \le 80^{\circ}$ with the gradual increase by 0.02°. The obtained XRD patterns were refned by the Rietveld method, using the Material Analysis Using Difraction (MAUD) software. The morphology, particle size distribution, and crystallinity of synthesized nanoparticles were investigated by TEM, HRTEM, and SAED analysis via the JEOL JEM-2100 operated at 200 kV. This was done by dissolving 10 mg of the nanoparticles in a solution of 50% diluted hydrochloric acid (HCl). Then, the mixture was subjected to 20 min of sonication. Subsequently, a few drops of the suspension were deposited onto a carbon-coated copper grid and left to air-dry at room temperature for several hours. The particle size, distance between the lattice planes, and diameter of the rings were determined from the obtained TEM, HRTEM, and SAED images by using the ImageJ software. Then, Origin Lab software was used to draw the particle size distribution and then it was ftted using the Gaussian function. The Raman analysis was performed at room temperature using a dispersive Raman spectrometer (SENTERRA, Bruker Optics). The Raman spectra were obtained after depositing the powder samples in a silicon reference material under a laser excitation source at a wavelength of 785 nm. The PL spectra were recorded at room temperature by FP-8600 spectrofuorometer, equipped with a Xenon lamp, in the range 320–700 nm at an excitation wavelength of 290 nm. The samples were prepared for the PL analysis by dissolving 3 mg of the nanoparticles in 20 mL of diluted HCl solution. Then, a rectangular quartz cuvette of a path length of 1cm was used as a sample holder. Finally, the peaks of the PL spectra were deconvoluted via the Fityk program. Additionally, the Mössbauer spectra were recorded at 300 K and 77 K using 57 Co/Rh γ-ray source mounted on an electromagnetic transducer with a triangular velocity form. It is worth mentioning that the hyperfne structure was modeled by a least-square ftting method involving Zeeman sextets and quadrupolar doublets comprising Lorentzian lines. Finally, a vibrating sample magnetometer (VSM) (Lake Shore 7410) was used for estimating the magnetic properties. About 0.1 g of

the prepared samples is weighed and mounted on an insulated tube constructed of nonmagnetic material, a quartz tube, in the gap between an electromagnet. The M-H hysteresis loops were recorded in the feld range from20 to 20 kg.

Results and discussion

X-ray powder difraction analysis

Figure [1](#page-3-0) depicts Rietveld refinements of XRD that were performed to confrm the phase purity and crystallinity of the prepared $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ nanoparticles where $x = 0.00, 0.01, 0.02, 0.04,$ and 0.08. Following the standard sheet number JCPDS:77- 0013, the difraction peaks located at 2θ of 30.05°, 35.75°, 37.05°, 43.11°, 53.73°, 57.21°, and 62.77°

Fig. 1 XRD refinements of Mg $_{0.33}$ Ni $_{0.33}$ Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ where *x* = 0.00, 0.01, 0.02, 0.04, and 0.08

were identifed as the (220), (311), (222), (400), (422), (511), and (440) planes, respectively [\[29\]](#page-18-1). Thus, the prepared nanoferrites crystallize in the FCC crystal structure. A minor peak appeared at $2\theta \sim 33^{\circ}$ which assigns the presence of the $Fe₂O₃$ impurity phase which is usually accompanied by ferrite preparation.

The lattice constant *a* was measured for all the prepared samples and listed in Table [1](#page-4-0) according to the following equation [[30\]](#page-18-2):

$$
\frac{1}{d^2} = \frac{h^2 + k^2 + l^2}{a^2},
$$
\n(1)

where *h*, *k*, and *l* are Miller indices and *d* is the space between the planes. It is clear, from Table [1](#page-4-0), that the calculated *a* values show a remarkable increase with increasing the Gd^{3+} concentrations. This is well consistent with Vegard's law $[31]$ $[31]$. This finding is mainly due to the difference in ionic radii between Gd^{3+} (0.938 Å) and Fe³⁺ (0.645 Å) which could result in an enhancement in the unit cell parameter *a* of the prepared Mg–Ni–Co nanoferrites. The crystallite sizes of the produced Mg–Ni–Co doped with Gd^{3+} ion nanoferrites were estimated for the most intense peak (311) plane determined from the XRD data, using the Debye-Scherer's formula [[32\]](#page-18-4):

$$
D = \frac{\alpha \lambda}{\beta \cos \theta},\tag{2}
$$

where λ is the wavelength of the Cu target (1.5406 Å), β is the full width of half maximum (FWHM) of the most intensive diffraction peak (311) , and 2θ is the diffraction angle. It can be seen from Table [1](#page-4-0) that the crystallite sizes showed a signifcant decrease from 30.27 to 23.39 nm with increasing the Gd^{3+} concentration. Additionally, the observed broadening of the (311) peak with increasing substitution content is a good indication of Gd^{3+} incorporation in the Mg–Ni–Co nanoferrite lattice. This incorporation should result in a reduction in crystallite sizes [\[33](#page-18-5)]. The decrease of the particle size upon doping nanoferrites with Gd is attributed to the reduction in the surface energy which in turn prevents the growth of nanoparticles [\[34\]](#page-18-6). The dislocation density (δ), the X-ray density (ρ_x) , the specific surface area (S), strain (E), and polaron radius (γ_p), listed in Table [1](#page-4-0), were calculated using the following relations, respectively [\[35,](#page-18-7) [36\]](#page-18-8):

$$
\delta = \frac{1}{D^2},\tag{3}
$$

$$
\rho_x = \frac{8M}{N_A a^3},\tag{4}
$$

$$
S = \frac{6}{\rho_x D},\tag{5}
$$

$$
\varepsilon = \frac{\beta \cos \theta}{4},\tag{6}
$$

$$
\gamma_P = \frac{1}{2} \sqrt[3]{\frac{\pi \ a^3}{576}},\tag{7}
$$

where *M* is the molecular weight and N_A is Avogadro's number. The concentration of lattice imperfections,

x 0.00 0.01 0.02 0.04 0.08 *a* (A) 8.351(3) 8.361(9) 8.362(9) 8.364(9) 8.367(8) *D*_{XRD} (nm) 30.27 27.8 26.05 25.26 23.39 $\delta \times 10^{-3}$ (nm⁻²) 9.90 10.7 11.4 11.8 12.6 $ρ_x$ (g/cm³)) 4.955 5.092 5.118 5.155 5.242 $S(m^2/g)$ $\langle g \rangle$ 39.999(3) 42.383(6) 44.995(5) 46.069(9) 48.930(6) \mathcal{E} 0.166(8) 0.166(3) 0.166(4) 0.166(2) 0.165(2) $L_A(A)$ 3.616(2) 3.620(8) 3.621(2) 3.622(1) 3.623(3) *L*_B (A) 2.952(6) 2.956(3) 2.956(7) 2.957(4) 2.958(4) γ*P* (Å) 0.735(0) 0.735(9) 0.736(0) 0.736(2) 0.736(4) *D*_{TEM} (nm) 27.5 23.3 20.1 18.7 14.2

Table 1 Structural parameters valued from XRD and TEM analysis for $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ nanoparticles with 0.00 ≤ *x* ≤ 0.08

known as dislocation density δ, is infuenced by the concentration of Gd^{3+} and is inversely proportional to particle size. Consequently, as the Gd^{3+} content increases, the dislocation density δ increases. Similar results were reported upon doping $Ni_{0.5}Zn_{0.5}Fe₂O₄$ with ruthenium [\[37\]](#page-18-9). As listed in Table [1](#page-4-0), the linear increase in ρ _x from 4.955 to 5.242 g/cm³ by increasing the concentration of the Gd^{3+} is attributed to the increase in the molecular weight of the prepared samples [\[38\]](#page-18-10). Additionally, the surface area (S) increases from 39.99 to 43.93 m^2/g upon increasing the Gd^{3+} content from 0.00 to 0.08. This is owed to the inverse proportionality between the surface area and the crystallite size [\[39\]](#page-18-11). Upon doping nanoferrites with Gd^{3+} , the strain ε values, mentioned in Table [1](#page-4-0), are slightly reduced. This might be attributed to the diference in ionic radii of the Gd³⁺ (0.938 Å) and Fe³⁺ (0.645 Å) [\[40,](#page-18-12) [41](#page-18-13)]. When Gd^{3+} ions substitute for the host Fe^{3+} ions, the larger size of Gd^{3+} causes an expansion of the crystal lattice. This increase in the lattice parameter reduces the strain and subsequently decreases the microstrain [\[42\]](#page-18-14). The distance between magnetic ions (hopping length) in the tetrahedral and octahedral sites is given by [\[35\]](#page-18-7):

$$
L_A = a \frac{\sqrt{3}}{4} \tag{8}
$$

$$
L_B = a \frac{\sqrt{2}}{4}.
$$
\n(9)

As declared in Table [1,](#page-4-0) the values L_A and L_B increase with increasing the substitution of $Fe³⁺$ by Gd^{3+} ions and lattice parameter. This is owed to the fact that the hopping length referring to the A-B sites is longer than that of the A-A and B-B sites [\[36](#page-18-8)]. The polaron radius slightly increases from 0.735(0) to 0.736(4) Å by increasing the Gd^{3+} concentration from 0.00 to 0.08. This is owed to the electron-phonon interaction as the charge carriers are confned in the d-shell. This polaron defect appeared because of the moving of the neighboring atoms from their position immediately after the charge carriers get caught in the lattice site [[43\]](#page-18-15).

To compare the obtained results with those reported in previous studies, the structural parameters of pure and Gd-doped $MgFe₂O₄$, NiFe₂O₄, and $CoFe₂O₄$ $CoFe₂O₄$ $CoFe₂O₄$ are listed Table 2. It is worth mentioning that the preparation method and cationic distribution afect the structural parameters of ferrite nanoparticles. Smaller crystallite sizes are obtained in this study. However, comparable values of L_A and L_B were reported for NiFe₂O₄ and MgFe₂O₄ nanoparticles [[44,](#page-18-16) [45](#page-18-17)]. Additionally, an increasing trend in the lattice parameter accompanied by a shrinking trend in the crystallite size was observed upon increasing the Gd³⁺ content in CoFe_{2-x}Gd_xO₄, NiFe_{2-x}Gd_xO₄ and $Gd_xMgFe_{2-x}O_4$ [[47–](#page-18-18)[49\]](#page-18-19). This was attributed to the difference in the ionic radius between the Gd^{3+} ions and Fe^{3+} ions and the act of Gd^{3+} as a restraint to

Nanoparticle		$a(\AA)$		$D_{\rm XRD}$ (nm) $\delta \times 10^{-3}$ (nm ⁻²) ϵ		$L_{\rm A}$ (A)	$L_{\rm R}$ (A)	Ref.
NiFe ₂ O ₄		8.3402	54.57	0.34	0.064	3.6114	2.9482	[44]
MgFe ₂ O ₄		8.376	20.89	2.290	0.0951	3.627	2.961	[45]
CoFe ₂ O ₄		8.362	21	1.354	0.1583			[46]
$CoFe_{2-x}Gd_xO_4$	$x=0.00$	8.365	77.37					$[47]$
	$x=0.03$	8.371	32.56					
	$x=0.09$	8.388	23.26					
	$x=0.12$	8.376	22.13					
	$x=0.15$	8.377	20.82					
$NiFe_{2-x}Gd_{y}O_{4}$	$x=0.00$	8.365	30.12			3.6222	2.9575	[48]
	$x=0.025$	8.369	29.15			3.6239	2.9589	
	$x=0.05$	8.374	26.35			3.6261	2.9607	
	$x=0.075$	8.378	24.61			3.6278	2.9621	
	$x=0.1$	8.383	23.25			3.6299	2.9639	
Gd, $MgFe_{2-x}O_4$ $x=0.0$		8.4038	24					[49]
	$x=0.1$	8.4129	20					

Table 2 Structural parameters of pure and Gd-doped nanoferrites reported in the literature

the growth process. Therefore, consistent results are obtained in this study.

TEM, HRTEM, and SAED analysis

The morphology and average particle size of Mg0.33Ni0.33Co0.33Fe2−*x*Gd*x*O4 nanoparticles were determined by TEM analysis. The average particle size (D_{TEM}) , listed in Table [1,](#page-4-0) was examined from TEM images (Fig. $2(a_1-e_1)$ $2(a_1-e_1)$). In addition, the histograms, representing the particle size distribution, are obtained using the ImageJ software and illustrated in Fig. $2(a_2-e_2)$ $2(a_2-e_2)$. The prepared nanoparticles exhibit rounded-cube morphology with a particle size distribution that ranges between 0 to 60 nm. MgFe₂O₄, $NiFe₂O₄$, and $CoFe₂O₄$ have similar morphologies, according to earlier investigations [[50–](#page-18-22)[52\]](#page-19-0). The average particle size of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ nanoparticles is reduced from 27.4 to 14.3 nm as *x* increases from 0.00 to 0.08, respectively. This decrement in the particle size with increasing the Gd^{3+} concentration is attributed to the act of Gd^{3+} as a restraint to the growth process, causing the formation of smaller particle sizes. These results are in agreement with a previously published paper where the particle size of $Co_{0.6}Mn_{0.4}Gd_xFe_{2-x}O₄$ nanoparticles decreased from 47.88 to 40.97 nm as the Gd content increased from 0 to 0.1 [\[53](#page-19-1)]. The change in particle size, as estimated from TEM and XRD analysis, exhibits a consistent trend. However, the smaller values of D_{TEM} compared to D_{XRD} can be attributed to the aging in the reaction solution [\[54](#page-19-2)]. The obtained result is similar to a previous study where the particle size estimated from TEM was smaller than that valued from XRD analysis for Zn_0 ₅Co_{0.5}Fe_{2−*x*}Pr_{*x*}O₄ nanoparticles [\[55](#page-19-3)]. All the prepared samples show a degree of agglomeration that may be due to the magnetic interparticle interactions [\[53](#page-19-1), [56\]](#page-19-4). The HRTEM images, shown in Fig. $2(a_3-e_3)$ $2(a_3-e_3)$, reveal the presence of diferent grains with diferent lattice plane orientations. This observation is attributed to the polycrystalline nature of the samples. Additionally, the presence of lattice fringes confrms the crystalline nature of the samples. The values of distance between the lattice planes along with their corresponding miller indices for selected grains are displayed in Fig. $2(a_3-e_3)$ $2(a_3-e_3)$ and are comparable with the XRD results. The selected area of electron difraction (SAED) images, displayed in Fig. $2(a_4-e_4)$ $2(a_4-e_4)$, reveals the presence of concentric

circles which confrms the polycrystalline nature of the samples. The identifed planes mainly (220), (311), (222), (400), (422), (551), and (440) are in agreement with those obtained from XRD measurements. Thus, the SAED images with the identifed planes confrm the formation of the desired spinel ferrite nanoparticles.

Raman analysis

Raman spectra of $Mg_{0,33}Ni_{0,33}Co_{0,33}Fe_{2-x}Ga_xO_4$ nanoparticles $(x = 0.00, 0.02,$ and (0.08) were recorded in the wavenumber range of 125–800 cm^{-1} . The prepared nanoferrites belong to the Fd-3m space group, which has an alternating A-B sublattice made up of 56 atoms in the entire unit cell and 14 atoms in the smallest Bravais lattice [[57\]](#page-19-5). The obtained spectra comprise the Raman active modes associated with the formation of the spinel ferrite nanoparticles. The analysis of group theory of the lattice vibrations reveals Raman active modes where *A*, *E*, and *T* represent one-, two-, and three-dimensional Raman modes, respectively [[58\]](#page-19-6). Clearly, all the samples exhibit more than fve Raman active modes. The appearance of additional modes is ascribed to the vibrations of various cations, mainly Fe, Mg, Ni, Co, and Gd, each with distinct properties [\[59](#page-19-7)].

The Raman spectra were deconvoluted to get more specifc details about the lattice structure and the vibrational modes by the Gaussian superposition ftting method using OriginLab software. The resulting deconvoluted spectra are shown in Fig. [3.](#page-8-0) In addition, Table [3](#page-8-1) summarizes the Raman mode's peak position of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ for $x = 0.00, 0.02,$ and 0.08. The detected Raman modes are attributed to the diferent cations that can occupy an identical position with the same lattice structure [\[60](#page-19-8)]. Furthermore, the broadening of certain peaks and the presence of weak modes are ascribed to the formation of small-sized ferrite nanoparticles [\[61](#page-19-9)]. The represented active modes confrm the spinel phase formation and are in good agreement with previously published reports [[62,](#page-19-10) [63\]](#page-19-11). The variation of Raman peak position upon doping nanoferrites with Gd^{3+} gives information about the deformation and the rearrangement of cations and anions in the lattice structure [\[64](#page-19-12)].

Vibrational modes observed above 600 cm−1 in the Raman spectrum were assigned to vibrations within the

Fig. 2 TEM images (a_1-e_1) , histograms of the particle size distribution (a_2-e_2) , HRTEM image (a_3-e_3) , and SAED images (a_4-d_e) of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Gd_xO_4$ where 0.00 ≤ *x* ≤ 0.08

tetrahedral complex $(AO₄)$. Conversely, those below 600 cm−1 were associated with oxygen vibrations within the octahedral complex (BO_6) . The A_{1g} mode corresponds to the symmetric stretch of oxygen anion in the tetrahedral arrangement [[65](#page-19-13)]. The band at A_{1g} mode was deconvoluted into three peaks, which are assigned to Fe–O and M–O where M=Ni and Co [\[66\]](#page-19-14). The diferent cations in the samples are responsible for the presence of the

Fig. 3 Deconvoluted Raman spectra of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ where **a** $x = 0.00$, **b** 0.02, and **c** 0.08

Table 3 Raman peaks of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Gd_xO_4$ where $x = 0.00, 0.02,$ and 0.08

\boldsymbol{x}	0.00	0.02	0.08
$T_{2g}(1)$	198	196	192
$\rm E_g$	311	310	301
$T_{2g}(2)$	464	463	462
$T_{2g}(3)$	542	549	543
A_{1g}	600	601	614
	667	649	648
	688	686	683

shouldered peak in A_{1g} mode because of those movements to the tetrahedral site at the same position [\[66\]](#page-19-14). Furthermore, the A_{1g} mode, located between 600 and 750 cm⁻¹, is red-shifted with increasing the Gd^{3+} concentration. This is owed to the dynamic motion of $Fe³⁺$ in the direction of the A sites $[67]$ $[67]$ $[67]$. Furthermore, the Raman modes labeled $T_{2g}(3)$ and $T_{2g}(2)$ arise from the symmetric and asymmetric bending vibrations of the oxygen atom within the M–O bond at the octahedral site. The E_g Raman mode corresponds to the symmetric bending motion of the oxygen atom relative to the metal ion located at the octahedral site [\[68](#page-19-16)]. Finally, the $T_{2g}(1)$

mode, located in the range of 192–198 cm−1 corresponds to the translational movement of the entire MO_4 tetrahe-dral units [\[69\]](#page-19-17). The redshift in the peak related to the A_{1g} band T_2 g (2) might be attributed to the alteration in the population of $Fe³⁺$ ions at the tetrahedral and octahedral sites [[70](#page-19-18)]. Interestingly, the peak located $<$ 200 cm⁻¹ was associated with the A_{1g} mode and two peaks of E_g mode, positioned in the range of 275–305 cm⁻¹ and 367–380 cm−1, correspond to the vibrational modes of hematite Fe₂O₃ [\[71](#page-19-19)[–73\]](#page-19-20). The recorded Raman spectra show a decrease in peak position by increasing the Gd^{3+} concentrations for all samples as shown in Table [2.](#page-5-0) This slight redshift may be due to the large diference in reduced mass of Fe³⁺ (55.85 g) and Gd³⁺ (157.25 g) ions causing changes in the M–O (M= Fe, Gd) bonds $[41]$ $[41]$ $[41]$.

PL examination

The optical properties of the synthesized Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ nanoparticles, where *x*=0.00, 0.01, 0.02, 0.04, and 0.08, were investigated via PL spectroscopy performed at room temperature. PL is used to analyze the structural defects in terms of the energy state of sub-energy level defects and the process of emission as well as the excitation. The recorded PL spectra of the prepared samples are shown in Fig. [4](#page-10-0). The particle size, surface defects, and oxygen functionalities play a crucial role in the photoluminescent property of the prepared samples [\[74](#page-19-21)]. The intensity of the PL signal provides information on the quality of surfaces and interfaces [\[75](#page-19-22)]. A sharp peak with the highest intensity of PL emission was detected near 334 nm in the prepared NPs. This is due to the recombination of electronhole pairs at the defect site due to electron transitions from the valence band to the conduction band. With the increase of Gd^{3+} concentration, the intensity of the PL peak located near 335 nm decreases, along with the appearance of small humps. Furthermore, the slowest recombination rate was detected for $Mg_{0,33}Ni_{0,33}Co_{0,33}Fe_{2-x}Gd_xO_4$ nanoparticles with $x=0.08$. Violet emission, observed at 410 nm, is attributed to the emission from some deep level in the band gap [[76\]](#page-20-0). Furthermore, the blue emission is detected in the range of 435–470 nm and is allocated to the recombination of electrons from vacancies to the holes in the valence band. This phenomenon occurs due to quantum confnement, resulting from the presence of radiative defects located at the surface grain boundaries of spinel ferrites. In addition, the blue emission band is due to the $3d^5 \rightarrow 3d^4$ 4s transition of Fe^{3+} at tetrahedral sites [[77\]](#page-20-1). The green emission band that appears near 550 nm in the prepared NPs is related to oxygen vacancies or other defects within the crystal lattice [\[37](#page-18-9)].

Mössbauer spectroscopy

Figures [5](#page-11-0) and [6](#page-12-0) display the Mössbauer spectra recorded at 300 and 77 K without an external applied field of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ nanoparticles where $x = 0.00, 0.02,$ and 0.08. The obtained Mössbauer parameters, including the mean values of the isomer shift (IS), quadrupole shift (QS), the hyperfine field (B_{hvo}) , and the area percentage (*S*) are listed in Table [3.](#page-8-1) Being less than 0.5, the IS values range between 0.305–0.308 mm/s and 0.379–0.489 mm/s at 300 and 77 K, respectively, indicating the existence of Fe^{3+} ions and lack of Fe^{2+} ions [[78,](#page-20-2) [79\]](#page-20-3).

The Mössbauer spectra recorded at 300 K for $x =$ 0.00, displayed in Fig. [5](#page-11-0)a, show a clear magnetic sextet comprising asymmetrical lines. This sextet can be described by at least two magnetic sextets, each having diferent isomer shift and hyperfne feld values based on the iron coordination at the tetrahedral A and octahedral B sites. However, due to the limited resolution of the recorded spectra, it is not possible to achieve a more precise ftting. Therefore, the ftting is performed using a discrete hyperfne feld distribution, which allows for the determination of average hyperfne parameter values. As Gd concentration increases from 0.00 to 0.08, an additional central paramagnetic doublet is observed alongside the relaxed sextet as shown in Fig. [5c](#page-11-0). This is mainly attributed to the existence of reduced particle sizes in Gd-doped nanoparticles [[80](#page-20-4)]. As *x* increases from 0.00 to 0.08, average particle size decreases from 27.5 to 14.2 nm along with the decrease in B_{hvo} values. This might be attributed to an increase in the lattice distortion upon doping nanoferrites with Gd^{3+} ions.

The Mössbauer spectra obtained at 77 K, represented in Fig. [6,](#page-12-0) reveal diferent hyperfne patterns characterized by well-defned magnetic sextets. Notably, each spectrum can be separated into two prominent sextets. The absence of paramagnetic doublet in Gd-doped nanoferrite $(x = 0.08)$ Mössbauer spectrum recorded at 77 K, shown in Fig. [6c](#page-12-0), verifes that the doublet observed at 300 K was a result of thermal superparamagnetic

Fig. 4 Deconvoluted PL spectra of $Mg_{0,33}Ni_{0,33}Co_{0,33}Fe_{2-x}Gd_xO_4$ where (a) $x = 0.00$, (b) $x = 0.01$, (c) $x = 0.02$, (d) $x = 0.04$, and (e) $x = 0.08$

Fig. 5 Mössbauer spectra of Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ (x = 0.00, 0.02, and 0.08) samples recorded at 300 K

fluctuations [\[81\]](#page-20-5). Furthermore, it is observed that the IS values at the B site are greater than those of the A-site. This phenomenon is attributed to the larger bond separation between Fe^{3+} – O^{2-} for the octahedral ions compared to the tetrahedral ions [[82](#page-20-6)]. Knowing that S_B and S_A represent area percentages for B and A sites respectively, the S_B/S_A ratio was determined and listed in Table [4](#page-12-1). As Gd concentration increases from 0.00 to 0.08, the S_B/S_A ratio is reduced from 0.82 to 0.43. This can be attributed to the settling of Gd^{3+} dopant ions at the B site causing the migration of $Fe³⁺$ ions from the octahedral site to the tetrahedral site.

The occurrence of chemical disorder in spinel ferrite structures leads to the generation of an electric feld gradient (EFG) which is responsible for the quadrupole splitting (QS) observed in the Mössbauer spectra [[83](#page-20-7)]. The QS values, listed in Table [4](#page-12-1), show that the values at examined 300 K are less than that at 77 K. This is owed to the fact that as the temperature of the nanoferrites rises, the thermal energy causes increased atomic movement within the nanoparticle. This enhanced atomic motion leads to a broader distribution of electric charges, resulting in a reduction in the EFG at the nucleus. Consequently, the quadrupole splitting (QS) observed in the Mössbauer spectra decreases. In comparison to tetrahedral symmetry, octahedral symmetry in crystal structures exhibits a higher degree of asymmetry. Consequently, the larger EFG leads to a larger quadrupole splitting (QS) observed at B sites compared to those at A sites [\[84](#page-20-8)]. The small values of QS observed in both the A and B suggest that the local symmetry of the obtained ferrite nanoparticles is closely approaching cubic symmetry [[85\]](#page-20-9).

M-H hysteresis loops

The magnetic properties of Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ nanoparticles were evaluated at room temperature using VSM with a maximum applied feld of 20 kG. The obtained M-H hysteresis loops, shown in Fig. [7](#page-13-0), were used to estimate the magnetic parameters like saturation magnetization (M_s) , retentivity (M_r) , coercive force (H_c) ,

Fig. 6 Mössbauer spectra of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ ($x = 0.00, 0.02,$ and 0.08) samples recorded at 77 K

the squareness ratio ($SQR = M_l/M_s$), and magnetocrystalline anisotropy constant (K). The values of the magnetic parameters are listed in Table [5.](#page-13-1)

The values of the M_s and H_c for pure $Mg_{0,33}Ni_{0,33}Co_{0,33}Fe_2O_4$ nanoparticles are 31.870 emu/g and 284.89 G, respectively. Naaz et al. [\[86\]](#page-20-10) reported lower values for M_s (24.520 emu/g) and H_c (93.78 Oe) for $MgFe₂O₄$ nanoparticles. However, CoFe₂O₄ nanoparticles exhibited superior M_s (52.57) emu/g) and H_c (377 Oe) values [[87](#page-20-11)]. Improved M_s value (41.67 emu/g) and lower H_c (0.027 Oe) were reported for $NiFe₂O₄$ nanoparticles in a previous study

Fig. 7 M-H hysteresis loops of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Ga_xO_4$ nanoparticles where $0.00 \le x \le 0.08$

Table 5 Magnetic parameters of $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Gd_xO_4$ where $0.00 \le x \le 0.08$

\mathcal{X}		M_s (emu/g) M_r (emu/g) H_c (G) SQR K (J/m ³)			
0.00	31.870	8.306	284.89	0.260	0.0057
0.01	27.357	6.357	274.59	0.225	0.0047
0.02	26.524	6.359	276.11	0.239	0.0046
0.04	23.126	5.265	307.05	0.227	0.0044
0.08	15.760	1.137	149.17	0.072	0.0014

[\[88\]](#page-20-12). It is worth mentioning that all the mentioned references used the coprecipitation method in the synthesis of MgFe₂O₄, CoFe₂O₄, and NiFe₂O₄ nanoparticles. However, the difference in the values of M_s and H_c can be attributed to the cation distribution [[89\]](#page-20-13). Thus, the modifcation in the nature and concentration of ions in A and B sites causes a variation in the resultant magnetization.

Doping Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ with different Gd concentrations initiates a significant effect on the magnetic properties of ferrite nanoparticles. The saturation magnetization of the prepared samples decreases from 31.870 to 15.760 emu/g as the concentration of Gd^{3+} increases from 0.00 to 0.08, respectively. The decrease of M_s can be explained by the 4f-3d electron coupling which refers to the interaction between the localized 4f electrons of the rare earth ions and the 3d electrons of the transition metal ions [\[48,](#page-18-21) [90](#page-20-14)]. Moreover, the weak $Gd^{3+} - Gd^{3+}$ interactions, mainly arising from the indirect 4f-5d-4f mechanism, lead to a reduction in magnetization [\[91\]](#page-20-15). Similarly in a previous study, the substitution of magnetic $Fe³⁺$ ions with nonmagnetic Gd^{3+} ions in the B site leads to a reduction in the magnetic moment of the B sublattice and caused a decrease in M_s values in Ni_{0.9}Cd_{0.1}Gd_xFe_{2−*x*}O₄ nanoparticles [\[92\]](#page-20-16). Peng et al. [\[93](#page-20-17)] reported a decrease in M_s values with the increase of $Gd³⁺$ concentrations in $CoFe₂,Gd_xO₄$ nanoparticles. In addition, the diminishing trend of the M_s was attributed to the variation in the cation distribution upon varying *x* in the Co_xFe_{3−*x*}O₄/ $SiO₂$ system [\[7\]](#page-17-2). The magnetocrystalline anisotropy constant K was calculated from the equation $[94, 95]$ $[94, 95]$ $[94, 95]$ $[94, 95]$:

$$
K = \frac{\mu_0 H_c M_s}{2} \tag{9}
$$

The calculated values of *K*, listed in Table [5,](#page-13-1) showed a decrease with increasing the Gd^{3+} concentration. This may be due to the large magnetic moment Gd³⁺ (7.94 μ_B) ions that replace the smaller Fe³⁺ (5.92 μ_B) ions at B sites [[53](#page-19-1), [96](#page-20-20)]. The squareness ratio (SQR= M_t/M_s) values are less than 0.5 revealing that the prepared samples are formed of single magnetic domain nanoparticles as stated by the Stoner Wohlfarth model [\[97,](#page-20-21) [98\]](#page-20-22). Furthermore, surface effects led to the uniaxial anisotropy in the prepared magnetic nanoferrites [[99\]](#page-20-23). Among the prepared samples, $Mg_{0,33}Ni_{0,33}Co_{0,33}Fe_{2-x}Ga_xO_4$ nanoparticles with $x =$ 0.04 exhibited the highest value of coercivity where H_c = 307.05 G. This might be ascribed to a higher level of spin-orbit coupling. Also, the enhanced H_c value may be due to the single ion anisotropy of rare-earth ions in the crystal lattice and the surface efects that alter the magnetic structure of nanoparticles [[100](#page-20-24), [101\]](#page-20-25). The reduced coercivity value ($H_c = 149.17$ G) revealed by nanoferrites with $x = 0.08$ is related to the small particle size, being below the critical size, as revealed from TEM analysis ($D_{\text{TEM}} = 14.2$ nm). Besides, the centered doublet, observed in the Mössbauer spectrum of $Mg_{0,33}Ni_{0,33}Co_{0,33}Fe_{1,92}Gd_{0,08}O_4$ at 300 K, can be related to superparamagnetic behavior of some particles of sizes below the critical size. The small value of squareness ratio, $SQR = 0.072$ for $x = 0.08$ demonstrates that the nanoparticles are relaxing so fast and act as superparamagnet at room temperature even in the absence of an external magnetic field $[102]$ $[102]$ $[102]$.

The law of approach to saturation (LAS) plays a vital role in studying the magnetization processes and their variation with the magnetic feld in highfeld regions. To describe the behavior near saturation of the prepared $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Gd_xO_4$, where *x* = 0.00, 0.01, 0.02, 0.04, and 0.08, the LAS was applied. To determine the most reliable model for obtaining information about magnetic parameters such as M_s and K , a study was conducted on three different LAS models applicable to various structures of spinel ferrites. LAS models, displayed in Fig. [8,](#page-15-0) are described in the following equations [[103\]](#page-21-1):

Model 1 :
$$
M = M_s \left(1 - \frac{c}{H} - \frac{b}{H^2} \right) + \chi H
$$
 (10)

Model 2 :
$$
M = M_s \left(1 - \frac{c}{H} - \frac{b}{H^2} \right)
$$
 (11)

$$
\text{Model 3}: M = M_s \left(1 - \frac{b}{H^2} \right) + \chi H \tag{12}
$$

where γH demonstrates the forced magnetization that grows up in magnetization with the magnetic feld. The ftting parameter *b* indicates the variation of magnetocrystalline anisotropy during the increase of the magnetic feld. The term *c*/*H* shows an indication of the deformation in the lattice structure and *c* indicates the magnetic hardness and is related to the sample microstructure. Although these models could be applied for spinel ferrites, the particular type of magnetic material, the applied magnetic feld range, and experimental conditions could afect the extracted results of magnetic parameters [[103\]](#page-21-1). At low regions of the applied magnetic field, the magnetization includes domain wall motion whereas in highfeld regions the magnetization process consists of the rotation of domains [[104\]](#page-21-2). To ensure that the magnetization process is dominated by domain rotations rather than other factors such as domain wall motion and to obtain accurate results, the ftting of LAS models is often performed in the high-feld regions, typically specified as $2000 < H < 20000$ [\[105](#page-21-3)].

To choose the suitable model for describing the dependency of the magnetic feld at appropriate highlevel regions and extracting acceptable values of magnetic parameters, a Mathematica 13.3 Wolfram coding package has been created. The values measured from the three models are listed in Table [5.](#page-13-1) From the ftting parameters, the magnetocrystalline anisotropy K can be calculated from equation $[106]$ $[106]$:

$$
K = \mu_0 M_S \sqrt{\frac{105 b}{8}}\tag{13}
$$

From Table [5](#page-13-1), the decrease in the M_s values, calculated using LAS models, is comparable to the data extracted from hysteresis loops. The increase in parameter *c* with increasing the concentration of Gd^{3+} ions, as observed in Table [6,](#page-15-1) could be attributed to the substitution of Fe^{3+} ions (with an ionic radius of 0.645 Å) by Gd^{3+} ions (with a larger ionic radius of 0.938 Å). This substitution introduces structural defects within the material, leading to the observed increase in the c parameter. Moreover, the introduction of Gd^{3+} dopant in the crystal lattice can cause dislocations, which disrupt the alignment of magnetic domains and result in an increase in the magnetic hardness [\[107\]](#page-21-5). All the fitted LAS models exhibited high coefficient of determination $(R²)$ values. However, the LAS models, mainly $M = M_s \left(1 - \frac{c}{H} - \frac{b}{H^2}\right) + \chi H$ and $M = M_s \left(1 - \frac{c}{H} - \frac{b}{H^2}\right)$, reveal negative values of *b* in all the prepared samples. Thus, these models are not valid for calculating the

Fig. 8 Fitting plots of LAS models for $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Gd_xO_4$ where $\mathbf{a} \times \mathbf{a} = 0.00$, $\mathbf{b} \times \mathbf{b} = 0.02$, and $\mathbf{c} \times \mathbf{b} = 0.08$

Table 6 Parameters obtained from the fitting of LAS models in the field region 2000 *G* <*H* < 20000 *G* for $Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2-x}Gd_xO_4$ where $x = 0.00, 0.02,$ and 0.08

Model	Equation	Parameter	\boldsymbol{x}			
			0.00	0.02	0.08	
Model 1	$M = M_{\rm s} \left(1 - \frac{c}{H} - \frac{b}{H^2} \right) + \chi H$	M_{s} (emu/g)	30.747	24.74	12.678	
		c(G)	776.772	911.486	1651.096	
		$b \times 10^6$ (G ²)	-0.548	-0.696	-1.41	
		χ	0.00012	0.00015	0.00021	
Model 2	$M = M_s \left(1 - \frac{c}{H} - \frac{b}{H^2} \right)$	$M_{\rm s}$ (emu/g)	33.347	28.153	17.397	
		c(G)	1140.77	1446.88	2628.86	
		$b \times 10^6$ (G ²)	-1.08	-1.49	-2.94	
Model 3	$M = M_s \left(1 - \frac{b}{H^2} \right) + \chi H$	M_s (emu/g)	26.662	20.899	9.0724	
		$b \times 10^6$ (G ²)	0.655	0.736	1.297	
		χ	0.00029	0.00031	0.00036	
		K	0.099	0.082	0.047	

magnetocrystalline anisotropy *K*. By excluding the term *c*/*H*, the model $M = M_s \left(1 - \frac{b}{H^2}\right) + \chi H$ reveals reliable values of M_s , b , and χ that describe the decrease of M_s and K upon increasing the Gd^{3+} concentration which in turn matches the obtained experimental results.

Conclusion

In this study, the influence of Gd^{3+} doping on the structural, optical, and magnetic properties of Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ was examined. The structural properties of the samples were examined by XRD, TEM, and HRTEM. As Gd^{3+} concentration increases from 0.00 to 0.08, the lattice parameter increases from 8.315 to 8.367 Å whereas the average particle size decreases from 27.5 to 14.2 nm, respectively. This was attributed to the diference in the ionic radii between the host $Fe³⁺$ (0.645) Å) and the dopant Gd^{3+} (0.938 Å) ions. Agglomeration among the prepared nanoparticles, revealed from TEM analysis, was owed to the magnetic interparticle interactions. Compared with previously published reports, smaller crystallite sizes are obtained upon doping nanoferrites with Gd^{3+} . Smaller nanoparticles have a higher surface area-to-volume ratio, which makes them good candidates for catalysis, adsorption, and surface modifcation applications. The optical properties studied by PL showed evidence of the existence of structural defects. The slowest recombination rate of electron-hole pair was revealed by $Mg_{0,33}Ni_{0,33}Co_{0,33}Fe_{2-x}Ga_xO_4$ nanoparticles with $x=0.08$. Thus, these nanoparticles can be used as photocatalysts for the degradation of pollutants and wastewater treatment. A shift in the peak position in the Raman modes revealed upon increasing Gd^{3+} concentrations was due to the large difference in reduced mass between the Fe^{3+} (55.85 g) and the Gd³⁺ (157.25 g) ions. The settling of Gd³⁺ dopant ions at the B site caused the migration of $Fe³⁺$ ions from the octahedral site to the tetrahedral site as revealed by Mössbauer spectra recorded at 77 K. In addition, the M_s values were reduced from 31.870 to 15.760 emu/g with the increase of Gd^{3+} concentration from 0.00 to 0.08, respectively. The small value of the squareness ratio, $SQR = 0.072$ for $x = 0.08$, demonstrates that the nanoparticles act as superparamagnets at room temperature. The relationship between

magnetization and the applied feld was investigated by ftting the experimental data into diferent forms of the law of approach to saturation. The prepared Mg_{0.33}Ni_{0.33}Co_{0.33}Fe_{2−*x*}Gd_{*x*}O₄ nanoparticles exhibited unique structural, optical, and magnetic properties which could have a futuristic signifcance in industrial applications such as sensors, light-emitting diodes (LEDs), photocatalysts, and electric generators.

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Data availability The data supporting the reported results are available upon request.

Compliance with ethical standard

Competing interests The authors declare no competing interests.

Ethics approval and consent to participate Not applicable.

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