

# **An Agmon estimate for Schrödinger operators on graphs**

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# **Abstract**

The Agmon estimate shows that eigenfunctions of Schrödinger operators,  $-\Delta\phi$  +  $V\phi = E\phi$ , decay exponentially in the 'classically forbidden' region where the potential exceeds the energy level  $\{x : V(x) > E\}$ . Moreover, the size of  $|\phi(x)|$  is bounded in terms of a weighted (Agmon) distance between *x* and the allowed region. We derive such a statement on graphs when  $-\Delta$  is replaced by the graph Laplacian  $L = D - A$ : we identify an explicit Agmon metric and prove a pointwise decay estimate in terms of the Agmon distance.

**Keywords** Agmon estimate · Agmon metric · Schrödinger operator · Graph

**Mathematics Subject Classification** 31B15 · 35J10 · 35R02

# **1 Introduction: Agmon estimates**

Let  $V : \mathbb{R}^n \to \mathbb{R}$  be a nonnegative potential growing at infinity, i.e.,  $V(x) \to \infty$  as *x*<sup> $||x||$ </sup> → ∞. Agmon estimates are concerned with eigenfunctions of the Schrödinger operator  $-\Delta + V$ : we study functions  $\phi : \mathbb{R}^n \to \mathbb{R}$  satisfying

$$
-\Delta\phi + V\phi = E\phi,
$$

where  $E \in \mathbb{R}$  is the eigenvalue. Multiplying with  $\phi$  and integrating by parts,

$$
\int_{\mathbb{R}^n} |\nabla \phi|^2 dx + \int_{\mathbb{R}^n} V(x) \cdot \phi(x)^2 dx = \int_{\mathbb{R}^n} E \cdot \phi(x)^2 dx.
$$

This identity implies that most of the *<sup>L</sup>*2−mass should be contained in the 'allowed' region  $\{x \in \mathbb{R}^n : V(x) \leq E\}$  and only very little mass can be in the 'forbidden' region

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 ${x \in \mathbb{R}^n : V(x) > E}$ . The celebrated Agmon estimate [\[1\]](#page-7-0) shows that this is indeed the case and, moreover, that  $\phi$  decays exponentially in terms of the distance from the allowed region for a suitable notion of distance. Agmon's estimate can be derived from an explicit integral identity. One way of motivating the estimate (taken from a summary of Deift  $[8]$  $[8]$ ) is as follows: if

$$
-\Delta\phi + V\phi = E\phi,
$$

then for any (sufficiently regular)  $g : \mathbb{R}^n \to \mathbb{R}$ 

$$
\int_{\mathbb{R}^n} \left| \nabla (e^g \phi) \right|^2 dx + \int_{\mathbb{R}^n} \left( V - E - |\nabla g|^2 \right) e^{2g} \phi^2 dx = 0.
$$

Ignoring the first (positive) term, this implies

$$
\int_{\mathbb{R}^n} \left( V - E - |\nabla g|^2 \right) e^{2g} \phi^2 dx \le 0.
$$

We note that  $e^{2g}$  and  $\phi^2$  are positive,  $V - E - |\nabla g|^2$  is negative in the allowed region and positive in the forbidden region provided  $|\nabla g|$  is sufficiently small. The inequality then naturally implies that there cannot be too much  $L^2$ −mass of  $\phi$  in the forbidden region except this is now coupled with an additional exponentially growing term  $e^{2g}$ . The statement becomes stronger, the larger we make  $g$ ; however, we want to maintain the nonnegativity of  $V - E - |\nabla g|^2$  in the forbidden region. This then suggests a way of defining *g*: the *Agmon metric* associated with the energy level *E* between two points  $x, y \in \mathbb{R}^n$  is given as the minimum energy taken

$$
\rho_E(x, y) = \inf_{\gamma} \int_0^1 \max \left( \sqrt{V(\gamma(t)) - E}, 0 \right) |\dot{\gamma}(t)| dt,
$$

where  $\gamma : [0, 1] \to \mathbb{R}^n$  ranges over all paths from  $\gamma(0) = x$  to  $\gamma(1) = y$ . The integral identity can then be used (see for example Carmona and Simon [\[6](#page-7-1)]) to derive pointwise statements in the forbidden region along the lines of that for all  $\varepsilon > 0$ 

$$
|\phi(x)| \leq c_{\varepsilon} \sup_{\substack{y \in \mathbb{R}^n \\ V(y) \leq E}} e^{-(1-\varepsilon)\rho_E(x,y)}.
$$

We refer to Aizenman and Simon [\[2](#page-7-2)], Carmona [\[5\]](#page-7-3), Dimassi and Sjöstrand [\[9\]](#page-8-1), Helffer [\[13](#page-8-2)], Helffer and Sjöstrand [\[14,](#page-8-3) [15\]](#page-8-4), Hislop [\[16\]](#page-8-5), Simon [\[23](#page-8-6), [24\]](#page-8-7) and references therein for a more complete picture regarding Agmon's estimate in the continuous setting. Our paper is partially inspired by a recent probabilistic approach to obtain sharp pointwise Agmon estimates in the continuous setting [\[25](#page-8-8)].

#### **2 An Agmon estimate on graphs**

#### **2.1 Setup**

Let  $G = (V, E)$  be a finite, connected graph with  $V = \{v_1, \ldots, v_n\}$ . We introduce the diagonal matrix  $D \in \mathbb{R}^{n \times n}$  satisfying  $d_{ii} = \deg(v_i)$  and the adjacency matrix  $A \in \mathbb{R}^{n \times n}$  given by

$$
A_{ij} = \begin{cases} 1 & \text{if } (v_i, v_j) \in E \\ 0 & \text{otherwise.} \end{cases}
$$

There is a natural notion of a discrete Laplacian acting on functions  $f: V \to V$  given by the linear operator  $L = D - A \in \mathbb{R}^{n \times n}$ . We observe that *L* can be interpreted as the discrete analogue of  $-\Delta$  as both are positive semi-definite and allow for integration by parts: for  $f: V \to \mathbb{R}$ 

$$
\langle f, Lf \rangle = \sum_{(v_i, v_j) \in E} (f(v_i) - f(v_j))^2.
$$

Given an arbitrary potential  $W: V \to \mathbb{R}$ , our goal is to understand the behavior of eigenfunctions  $\phi: V \to \mathbb{R}$  satisfying

$$
L\phi + W\phi = E\phi
$$

for some eigenvalue  $E \in \mathbb{R}$ . Multiplying with  $\phi$  and integrating by parts

$$
\sum_{(v_i, v_j) \in E} (\phi(v_i) - \phi(v_j))^2 + \sum_{v \in V} W(v) \cdot \phi(v)^2 = \sum_{v \in V} E \cdot \phi(v)^2.
$$

This suggests, just as in the continuous case above, that there should be relatively little  $\ell^2$ −mass in the 'classically forbidden' region { $v \in V : W(v) > E$ }. The question is now whether, just as in the continuous case, one can expect exponential decay in the forbidden region and how this can be quantified.

#### **2.2 Main result**

We define a notion of Agmon distance  $\rho_E : V \to \mathbb{R}$  as the cost of the cheapest path starting in  $v \in V$  and ending in any arbitrary vertex in the allowed region where 'cheap' refers to an explicit cost function on *V* depending on the potential *W*, the energy *E* and the degree of the vertex. Formally,

$$
\rho_E(v) = \inf \left\{ \sum_{i=1}^{\ell} \log \left( 1 + \frac{(W(v_i) - E)_+}{\deg(v_i)} \right) : v = v_1 \to \cdots \to v_{\ell} \text{ and } W(v_{\ell}) \leq E \right\},\
$$

where the infimum is taken over all paths that start in v and end in a vertex  $v_\ell$  in the allowed region. As usual,  $(W(v_i) - E)_+ = \max \{W(v_i) - E, 0\}$ . Note that  $\rho_E \equiv 0$ in the allowed region and  $\rho_E > 0$  in the forbidden region.

**Theorem** *We have, for all*  $v \in V$ ,

 $|\phi(v)| < e^{-\rho_E(v)} \cdot ||\phi||_{\ell^{\infty}}.$ 

The maximum principle shows that  $|\phi|$  assumes its maximum in the allowed region and thus  $\|\phi\|_{\ell^{\infty}} = \|\phi\|_{\ell^{\infty}(W(v)\leq E)}$ . Since  $\rho_E \equiv 0$  in the allowed region, the inequality is sharp in the maximum and the implicit constant 1 in front cannot be improved any further. In terms of the exponential decay, there are graphs where the inequality is asymptotically optimal: such examples are constructed in §3.2.

#### **2.3 Related results**

There is relatively little work regarding Agmon estimates on graphs. However, we emphasize one recent result which is close in spirit to our result. Filoche, Mayboroda and Tao [\[11](#page-8-9)] study eigenvector localization for a fairly general class of matrices  $A \in$  $\mathbb{R}^{n \times n}$ . They obtain an integrated exponential estimate in terms of an explicit Agmontype distance. Considering  $A = D - A + W$  and  $u = (1, 1, \ldots, 1)$  in their approach, one arrives at a notion of distance

$$
\rho(v, w) = \inf \left\{ \sum_{i=1}^{\ell} \log \left( 1 + \sqrt[4]{(W(v_i) - E)_+(W(v_{i+1}) - E)_+} \right) \right\},\,
$$

where the infimum ranges over all paths that start in  $v = v_1$  and end in  $w = v_{\ell+1}$  (and, as in our approach, traveling through the allowed region is free which we suppress in the equation above for simplicity of exposition). This is very similar in flavor to our distance above: using this, they then obtain an integrated estimate also involving a landscape-type potential  $A^{-1}$ **1** [\[11](#page-8-9), Theorem 2.5] as well as more general integrated estimates [\[11,](#page-8-9) Theorem 2.7]. A main difference is the dependency on the degree of a vertex which is locally built into our distance while arising in the integrated estimates of [\[11\]](#page-8-9) more globally (somewhat unsurprisingly: integrated estimates themselves are global). Both our estimate and the estimates in  $[11]$  are complementary: which one ends up being better will depend (among other things) on whether there is a lot of variation in the degrees of the vertices.

There is also a recent work of Keller and Pogorzelski [\[18\]](#page-8-10) who study Agmon estimates in the more general setting of weighted, infinite graphs where the Agmon distance is given in terms of Hardy weights. There is a philosophical overlap with work of Dodziuk [\[10](#page-8-11)]. We also note the work of Akduman and Pankov [\[3](#page-7-4), [4\]](#page-7-5) on metric graphs, the work of Harrell and Maltsev [\[12](#page-8-12)] on quantum graphs, Damanik, Fillman and Sukhtaiev [\[7\]](#page-8-13) on tree graphs, results of Hua and Lu [\[17](#page-8-14)] and Wojciechowski [\[27](#page-8-15)] as well as work of Klein and Rosenberger  $[19, 20]$  $[19, 20]$  $[19, 20]$  $[19, 20]$ , Mandich  $[22]$  and Wang and Zhang  $[26]$  on  $\mathbb{Z}^d$ .

## **3 Proof**

#### **3.1 Proof of the Theorem**

*Proof* Note first that the eigenfunction satisfies

$$
(D-A)\phi = (E-W)\phi.
$$

Considering this linear system of equations in a fixed vertex  $u \in V$  one obtains

$$
deg(u)\phi(u) - \sum_{(u,w)\in E} \phi(w) = (E - W(u))\phi(u).
$$

This equation can be rewritten as

$$
\[1 + \frac{W(u) - E}{\deg(u)}\]phi(u) = \frac{1}{\deg(u)} \sum_{(u,w) \in E} \phi(w).
$$

We observe that if  $\phi(u) = 0$ , then the theorem is trivially true in *u*. It thus suffices to prove it for vertices  $u \in V$  where  $\phi(u) \neq 0$ . Note, moreover, that in the forbidden region  $\{u \in V : W(u) > E\}$ , one trivially has

$$
1 + \frac{W(u) - E}{\deg(u)} \ge 1
$$

and thus, for  $u \in V$  in the forbidden region, it is possible to divide and

$$
\phi(u) = \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \frac{1}{\deg(u)} \sum_{(u,w) \in E} \phi(w).
$$

Taking absolute values on both sides, we have

$$
|\phi(u)| \le \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \frac{1}{\deg(u)} \sum_{(u,w) \in E} |\phi(w)|
$$
  
 
$$
\le \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \max_{(u,w) \in E} |\phi(w)|.
$$

Since  $\phi(u) \neq 0$ , we deduce

$$
\max_{(u,w)\in E} |\phi(w)| > |\phi(u)|.
$$

We can now move from *u* to its neighbor w maximizing  $|\phi(w)|$  and then apply the very same argument again in  $w$ . The argument can be applied iteratively as long as the new vertex is still in the forbidden region. Note that  $|\phi|$  is increasing along the way  $u = v_1 \rightarrow v_2 \rightarrow \cdots \rightarrow v_m \rightarrow v_{m+1}$  where  $W(v_{m+1}) \leq E$ 

since otherwise the path could be further extended. Collecting all the factors

$$
|\phi(u)| \le \left(\prod_{i=1}^m \left[1 + \frac{W(v_i) - E}{\deg(v_i)}\right]^{-1}\right) \cdot |\phi(v_{m+1})|
$$
  

$$
\le \left(\prod_{i=1}^m \left[1 + \frac{W(v_i) - E}{\deg(v_i)}\right]^{-1}\right) \cdot \|\phi\|_{\ell^\infty}
$$

Note that

$$
\prod_{i=1}^m \left[1+\frac{W(v_i)-E}{\deg(v_i)}\right]^{-1} = \exp\left(-\sum_{i=1}^m \log\left(1+\frac{W(v_i)-E}{\deg(v_i)}\right)\right).
$$

By definition of  $\rho_E$ , we have

$$
\exp\left(-\sum_{i=1}^{m}\log\left(1+\frac{W(v_i)-E}{\deg(v_i)}\right)\right)\leq \exp\left(-\rho_E(v_1)\right)
$$

and this concludes the proof.

*Remark* We note that the final estimate in the argument implies

$$
|\phi(u)| \le \exp\left(-\sum_{i=1}^m \log\left(1 + \frac{W(v_i) - E}{\deg(v_i)}\right)\right) \cdot |\phi(w_{m+1})|
$$

for any path starting in  $u = v_1$  and ending in the vertex  $v_{m+1}$  in the allowed region. This would imply a slightly refined estimate where one is not only interested in minimizing the Agmon metric but also wants to end up in a vertex in the allowed region such that  $|\phi(w_{m+1})|$  is not too small.

*Remark* We quickly note a part in the derivation where the argument can be lossy: the main inequality is

$$
|\phi(u)| \le \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \frac{1}{\deg(u)} \sum_{w \sim u} |\phi(w)|
$$
  
 
$$
\le \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \max_{(u, w) \in E} |\phi(w)|.
$$

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This inequality could also be interpreted as a type of martingale inequality and can be exploited in this sense. Let us define a sequence of random vertices given by  $X_0 = u$ and such that  $X_{k+1}$  is a randomly chosen neighbor of  $X_k$  and that this random walk is continued until  $W(X_k) \leq E$ . We will denote the smallest such k by the stopping time  $\tau$ . Assume furthermore that

$$
\frac{W(u) - E}{\deg(u)} \ge \delta \quad \text{for all } u \in V \text{ in the forbidden region.}
$$

An iterative application of the inequality then implies

$$
|\phi(u)| \leq \left(\sum_{\ell=0}^{\infty} \frac{\mathbb{P}(\tau=\ell)}{(1+\delta)^{\ell}}\right) \cdot \|\phi\|_{\ell^{\infty}}.
$$

We note that this inequality can lead to improved results in settings where a random walk needs a very long time before arriving in the allowed region. Observe that the sum can be interpreted as an exponential moment  $\mathbb{E}$  exp ( $\tau/(\delta + 1)$ ) of the stopping time  $\tau$  which is a well-studied object. We refer to [\[25](#page-8-8)] for the derivation of Agmon estimates via this more stochastic perspective in the continuous setting.

#### **3.2 An Example**

The purpose of this section is to construct a graph where the inequality is nearly sharp. One example of such graphs is given by *q*−regular trees of a certain depth where the final layer of vertices is then additionally connected to another vertex  $v_*$  (see Fig. [1,](#page-6-0) for an example). We consider the potential given by  $W(v_*) = 0$  and, for all other vertices  $v \neq v_*,$  we choose the potential to be constant and  $W(v) = W \gg q^k \gg 1$ for some very large constant  $W \in \mathbb{R}$  where k is the depth of the tree. The function we will consider is the first eigenfunction of  $L + W$ .

By Rayleigh–Ritz, the smallest eigenvalue of  $-\Delta + W$  satisfies

$$
\lambda_1 = \inf_{f:V \to \mathbb{R}} \frac{\sum_{(v_i, v_j) \in E} (f(v_i) - f(v_j))^2 + \sum_{v \in V} W(v) f(v)^2}{\sum_{v \in V} f(v)^2}.
$$

<span id="page-6-0"></span>**Fig. 1** A *q*−regular tree (here,  $q = 3$ ) of depth *k* (here,  $k = 2$ ) with final layer being connected to a single additional vertex  $v<sub>*</sub>$ 



Taking  $f: V \to \mathbb{R}$  given by  $f(v_*) = 1$  and  $f(v) = 0$  for all  $v \neq v^*$ , we deduce that

$$
E = \lambda_1 \le q^k \quad \text{independently of } W.
$$

We can now use the equation

$$
\phi(u) = \frac{E - W(u)}{\deg(u)} \phi(u) + \frac{1}{\deg(u)} \sum_{w \sim u} \phi(w)
$$

and we shall restrict its use to vertices in the *q*−regular tree. The value of  $\phi$  then only depends on the level. We shall therefore write  $\phi(v) = \phi_i$  whenever the vertex v is in the *i*−th level where  $1 \le i \le k - 1$  (the case  $i = 0$  and  $i = k$  will be ignored since the algebra is slightly different). The equation then simplifies to

$$
\phi_i = \frac{\lambda_1 - W}{q + 1} \phi_i + \frac{1}{q + 1} (q \cdot \phi_{i+1} + \phi_{i-1}).
$$

This can be rewritten as

$$
(W - \lambda_1 + q + 1) \cdot \phi_i = q \cdot \phi_{i+1} + \phi_{i-1}.
$$

For fixed *q* and  $W \gg q^k \ge E$ , this implies that approximately  $\phi_i \sim (q/W) \cdot \phi_{i+1}$  as  $W \rightarrow \infty$  which implies exponential decay. Conversely, we have

$$
\log\left(1+\frac{W-E}{\deg(v)}\right) \sim \log\left(\frac{W}{q}\right)
$$

which implies, to leading order, the same kind of decay.

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#### **Declarations**

**Conflict of interest** There are no associated data produced and there are no conflicts of interest to declare.

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