

# An Agmon estimate for Schrödinger operators on graphs

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# Abstract

The Agmon estimate shows that eigenfunctions of Schrödinger operators,  $-\Delta \phi + V\phi = E\phi$ , decay exponentially in the 'classically forbidden' region where the potential exceeds the energy level  $\{x : V(x) > E\}$ . Moreover, the size of  $|\phi(x)|$  is bounded in terms of a weighted (Agmon) distance between x and the allowed region. We derive such a statement on graphs when  $-\Delta$  is replaced by the graph Laplacian L = D - A: we identify an explicit Agmon metric and prove a pointwise decay estimate in terms of the Agmon distance.

Keywords Agmon estimate · Agmon metric · Schrödinger operator · Graph

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# 1 Introduction: Agmon estimates

Let  $V : \mathbb{R}^n \to \mathbb{R}$  be a nonnegative potential growing at infinity, i.e.,  $V(x) \to \infty$  as  $||x|| \to \infty$ . Agmon estimates are concerned with eigenfunctions of the Schrödinger operator  $-\Delta + V$ : we study functions  $\phi : \mathbb{R}^n \to \mathbb{R}$  satisfying

$$-\Delta\phi + V\phi = E\phi,$$

where  $E \in \mathbb{R}$  is the eigenvalue. Multiplying with  $\phi$  and integrating by parts,

$$\int_{\mathbb{R}^n} |\nabla \phi|^2 \mathrm{d}x + \int_{\mathbb{R}^n} V(x) \cdot \phi(x)^2 \mathrm{d}x = \int_{\mathbb{R}^n} E \cdot \phi(x)^2 \mathrm{d}x.$$

This identity implies that most of the  $L^2$ -mass should be contained in the 'allowed' region  $\{x \in \mathbb{R}^n : V(x) \le E\}$  and only very little mass can be in the 'forbidden' region

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 $\{x \in \mathbb{R}^n : V(x) > E\}$ . The celebrated Agmon estimate [1] shows that this is indeed the case and, moreover, that  $\phi$  decays exponentially in terms of the distance from the allowed region for a suitable notion of distance. Agmon's estimate can be derived from an explicit integral identity. One way of motivating the estimate (taken from a summary of Deift [8]) is as follows: if

$$-\Delta\phi + V\phi = E\phi,$$

then for any (sufficiently regular)  $g : \mathbb{R}^n \to \mathbb{R}$ 

$$\int_{\mathbb{R}^n} \left| \nabla (e^g \phi) \right|^2 \mathrm{d}x + \int_{\mathbb{R}^n} \left( V - E - \left| \nabla g \right|^2 \right) e^{2g} \phi^2 \mathrm{d}x = 0.$$

Ignoring the first (positive) term, this implies

$$\int_{\mathbb{R}^n} \left( V - E - |\nabla g|^2 \right) e^{2g} \phi^2 \mathrm{d}x \le 0.$$

We note that  $e^{2g}$  and  $\phi^2$  are positive,  $V - E - |\nabla g|^2$  is negative in the allowed region and positive in the forbidden region provided  $|\nabla g|$  is sufficiently small. The inequality then naturally implies that there cannot be too much  $L^2$ -mass of  $\phi$  in the forbidden region except this is now coupled with an additional exponentially growing term  $e^{2g}$ . The statement becomes stronger, the larger we make g; however, we want to maintain the nonnegativity of  $V - E - |\nabla g|^2$  in the forbidden region. This then suggests a way of defining g: the Agmon metric associated with the energy level E between two points  $x, y \in \mathbb{R}^n$  is given as the minimum energy taken

$$\rho_E(x, y) = \inf_{\gamma} \int_0^1 \max\left(\sqrt{V(\gamma(t)) - E}, 0\right) |\dot{\gamma}(t)| \mathrm{d}t,$$

where  $\gamma : [0, 1] \to \mathbb{R}^n$  ranges over all paths from  $\gamma(0) = x$  to  $\gamma(1) = y$ . The integral identity can then be used (see for example Carmona and Simon [6]) to derive pointwise statements in the forbidden region along the lines of that for all  $\varepsilon > 0$ 

$$|\phi(x)| \le c_{\varepsilon} \sup_{\substack{y \in \mathbb{R}^n \\ V(y) \le E}} e^{-(1-\varepsilon)\rho_E(x,y)}.$$

We refer to Aizenman and Simon [2], Carmona [5], Dimassi and Sjöstrand [9], Helffer [13], Helffer and Sjöstrand [14, 15], Hislop [16], Simon [23, 24] and references therein for a more complete picture regarding Agmon's estimate in the continuous setting. Our paper is partially inspired by a recent probabilistic approach to obtain sharp pointwise Agmon estimates in the continuous setting [25].

#### 2 An Agmon estimate on graphs

#### 2.1 Setup

Let G = (V, E) be a finite, connected graph with  $V = \{v_1, \ldots, v_n\}$ . We introduce the diagonal matrix  $D \in \mathbb{R}^{n \times n}$  satisfying  $d_{ii} = \deg(v_i)$  and the adjacency matrix  $A \in \mathbb{R}^{n \times n}$  given by

$$A_{ij} = \begin{cases} 1 & \text{if } (v_i, v_j) \in E \\ 0 & \text{otherwise.} \end{cases}$$

There is a natural notion of a discrete Laplacian acting on functions  $f: V \to V$  given by the linear operator  $L = D - A \in \mathbb{R}^{n \times n}$ . We observe that *L* can be interpreted as the discrete analogue of  $-\Delta$  as both are positive semi-definite and allow for integration by parts: for  $f: V \to \mathbb{R}$ 

$$\langle f, Lf \rangle = \sum_{(v_i, v_j) \in E} (f(v_i) - f(v_j))^2.$$

Given an arbitrary potential  $W : V \to \mathbb{R}$ , our goal is to understand the behavior of eigenfunctions  $\phi : V \to \mathbb{R}$  satisfying

$$L\phi + W\phi = E\phi$$

for some eigenvalue  $E \in \mathbb{R}$ . Multiplying with  $\phi$  and integrating by parts

$$\sum_{(v_i,v_j)\in E} (\phi(v_i) - \phi(v_j))^2 + \sum_{v\in V} W(v) \cdot \phi(v)^2 = \sum_{v\in V} E \cdot \phi(v)^2.$$

This suggests, just as in the continuous case above, that there should be relatively little  $\ell^2$ -mass in the 'classically forbidden' region { $v \in V : W(v) > E$ }. The question is now whether, just as in the continuous case, one can expect exponential decay in the forbidden region and how this can be quantified.

#### 2.2 Main result

We define a notion of Agmon distance  $\rho_E : V \to \mathbb{R}$  as the cost of the cheapest path starting in  $v \in V$  and ending in any arbitrary vertex in the allowed region where 'cheap' refers to an explicit cost function on V depending on the potential W, the energy E and the degree of the vertex. Formally,

$$\rho_E(v) = \inf\left\{\sum_{i=1}^{\ell} \log\left(1 + \frac{(W(v_i) - E)_+}{\deg(v_i)}\right) : v = v_1 \to \dots \to v_{\ell} \text{ and } W(v_{\ell}) \le E\right\},\$$

where the infimum is taken over all paths that start in v and end in a vertex  $v_{\ell}$  in the allowed region. As usual,  $(W(v_i) - E)_+ = \max \{W(v_i) - E, 0\}$ . Note that  $\rho_E \equiv 0$  in the allowed region and  $\rho_E > 0$  in the forbidden region.

**Theorem** We have, for all  $v \in V$ ,

 $|\phi(v)| \le e^{-\rho_E(v)} \cdot \|\phi\|_{\ell^{\infty}}.$ 

The maximum principle shows that  $|\phi|$  assumes its maximum in the allowed region and thus  $\|\phi\|_{\ell^{\infty}} = \|\phi\|_{\ell^{\infty}(W(v) \le E)}$ . Since  $\rho_E \equiv 0$  in the allowed region, the inequality is sharp in the maximum and the implicit constant 1 in front cannot be improved any further. In terms of the exponential decay, there are graphs where the inequality is asymptotically optimal: such examples are constructed in §3.2.

#### 2.3 Related results

There is relatively little work regarding Agmon estimates on graphs. However, we emphasize one recent result which is close in spirit to our result. Filoche, Mayboroda and Tao [11] study eigenvector localization for a fairly general class of matrices  $A \in \mathbb{R}^{n \times n}$ . They obtain an integrated exponential estimate in terms of an explicit Agmontype distance. Considering A = D - A + W and u = (1, 1, ..., 1) in their approach, one arrives at a notion of distance

$$\rho(v, w) = \inf \left\{ \sum_{i=1}^{\ell} \log \left( 1 + \sqrt[4]{(W(v_i) - E)_+ (W(v_{i+1}) - E)_+} \right) \right\},\$$

where the infimum ranges over all paths that start in  $v = v_1$  and end in  $w = v_{\ell+1}$  (and, as in our approach, traveling through the allowed region is free which we suppress in the equation above for simplicity of exposition). This is very similar in flavor to our distance above: using this, they then obtain an integrated estimate also involving a landscape-type potential  $A^{-1}\mathbf{1}$  [11, Theorem 2.5] as well as more general integrated estimates [11, Theorem 2.7]. A main difference is the dependency on the degree of a vertex which is locally built into our distance while arising in the integrated estimates of [11] more globally (somewhat unsurprisingly: integrated estimates themselves are global). Both our estimate and the estimates in [11] are complementary: which one ends up being better will depend (among other things) on whether there is a lot of variation in the degrees of the vertices.

There is also a recent work of Keller and Pogorzelski [18] who study Agmon estimates in the more general setting of weighted, infinite graphs where the Agmon distance is given in terms of Hardy weights. There is a philosophical overlap with work of Dodziuk [10]. We also note the work of Akduman and Pankov [3, 4] on metric graphs, the work of Harrell and Maltsev [12] on quantum graphs, Damanik, Fillman and Sukhtaiev [7] on tree graphs, results of Hua and Lu [17] and Wojciechowski [27] as well as work of Klein and Rosenberger [19, 20], Mandich [22] and Wang and Zhang [26] on  $\mathbb{Z}^d$ .

## 3 Proof

#### 3.1 Proof of the Theorem

**Proof** Note first that the eigenfunction satisfies

$$(D-A)\phi = (E-W)\phi.$$

Considering this linear system of equations in a fixed vertex  $u \in V$  one obtains

$$\deg(u)\phi(u) - \sum_{(u,w)\in E} \phi(w) = (E - W(u))\phi(u).$$

This equation can be rewritten as

$$\left[1 + \frac{W(u) - E}{\deg(u)}\right]\phi(u) = \frac{1}{\deg(u)}\sum_{(u,w)\in E}\phi(w).$$

We observe that if  $\phi(u) = 0$ , then the theorem is trivially true in *u*. It thus suffices to prove it for vertices  $u \in V$  where  $\phi(u) \neq 0$ . Note, moreover, that in the forbidden region  $\{u \in V : W(u) > E\}$ , one trivially has

$$1 + \frac{W(u) - E}{\deg(u)} \ge 1$$

and thus, for  $u \in V$  in the forbidden region, it is possible to divide and

$$\phi(u) = \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \frac{1}{\deg(u)} \sum_{(u,w) \in E} \phi(w).$$

Taking absolute values on both sides, we have

$$\begin{aligned} |\phi(u)| &\leq \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \frac{1}{\deg(u)} \sum_{(u,w) \in E} |\phi(w)| \\ &\leq \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \max_{(u,w) \in E} |\phi(w)|. \end{aligned}$$

Since  $\phi(u) \neq 0$ , we deduce

$$\max_{(u,w)\in E} |\phi(w)| > |\phi(u)|.$$

We can now move from u to its neighbor w maximizing  $|\phi(w)|$  and then apply the very same argument again in w. The argument can be applied iteratively as long as the new vertex is still in the forbidden region. Note that  $|\phi|$  is increasing along the way

which implies that the arising path can never cross itself and must eventually end up in the allowed region  $\{v \in V : W(v) \le E\}$ . Altogether, this results in a path

$$u = v_1 \rightarrow v_2 \rightarrow \cdots \rightarrow v_m \rightarrow v_{m+1}$$
 where  $W(v_{m+1}) \leq E$ 

since otherwise the path could be further extended. Collecting all the factors

$$\begin{aligned} |\phi(u)| &\leq \left(\prod_{i=1}^{m} \left[1 + \frac{W(v_i) - E}{\deg(v_i)}\right]^{-1}\right) \cdot |\phi(v_{m+1})| \\ &\leq \left(\prod_{i=1}^{m} \left[1 + \frac{W(v_i) - E}{\deg(v_i)}\right]^{-1}\right) \cdot \|\phi\|_{\ell^{\infty}} \end{aligned}$$

Note that

$$\prod_{i=1}^{m} \left[ 1 + \frac{W(v_i) - E}{\deg(v_i)} \right]^{-1} = \exp\left( -\sum_{i=1}^{m} \log\left( 1 + \frac{W(v_i) - E}{\deg(v_i)} \right) \right).$$

By definition of  $\rho_E$ , we have

$$\exp\left(-\sum_{i=1}^{m}\log\left(1+\frac{W(v_i)-E}{\deg(v_i)}\right)\right) \le \exp\left(-\rho_E(v_1)\right)$$

and this concludes the proof.

**Remark** We note that the final estimate in the argument implies

$$|\phi(u)| \le \exp\left(-\sum_{i=1}^{m} \log\left(1 + \frac{W(v_i) - E}{\deg(v_i)}\right)\right) \cdot |\phi(w_{m+1})|$$

for any path starting in  $u = v_1$  and ending in the vertex  $v_{m+1}$  in the allowed region. This would imply a slightly refined estimate where one is not only interested in minimizing the Agmon metric but also wants to end up in a vertex in the allowed region such that  $|\phi(w_{m+1})|$  is not too small.

*Remark* We quickly note a part in the derivation where the argument can be lossy: the main inequality is

$$\begin{aligned} |\phi(u)| &\leq \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \frac{1}{\deg(u)} \sum_{w \sim u} |\phi(w)| \\ &\leq \left[1 + \frac{W(u) - E}{\deg(u)}\right]^{-1} \max_{(u,w) \in E} |\phi(w)|. \end{aligned}$$

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This inequality could also be interpreted as a type of martingale inequality and can be exploited in this sense. Let us define a sequence of random vertices given by  $X_0 = u$  and such that  $X_{k+1}$  is a randomly chosen neighbor of  $X_k$  and that this random walk is continued until  $W(X_k) \le E$ . We will denote the smallest such k by the stopping time  $\tau$ . Assume furthermore that

$$\frac{W(u) - E}{\deg(u)} \ge \delta \quad \text{for all } u \in V \text{ in the forbidden region.}$$

An iterative application of the inequality then implies

$$|\phi(u)| \leq \left(\sum_{\ell=0}^\infty \frac{\mathbb{P}\left(\tau=\ell\right)}{(1+\delta)^\ell}\right) \cdot \|\phi\|_{\ell^\infty}.$$

We note that this inequality can lead to improved results in settings where a random walk needs a very long time before arriving in the allowed region. Observe that the sum can be interpreted as an exponential moment  $\mathbb{E} \exp(\tau/(\delta + 1))$  of the stopping time  $\tau$  which is a well-studied object. We refer to [25] for the derivation of Agmon estimates via this more stochastic perspective in the continuous setting.

#### 3.2 An Example

The purpose of this section is to construct a graph where the inequality is nearly sharp. One example of such graphs is given by q-regular trees of a certain depth where the final layer of vertices is then additionally connected to another vertex  $v_*$  (see Fig. 1, for an example). We consider the potential given by  $W(v_*) = 0$  and, for all other vertices  $v \neq v_*$ , we choose the potential to be constant and  $W(v) = W \gg q^k \gg 1$  for some very large constant  $W \in \mathbb{R}$  where k is the depth of the tree. The function we will consider is the first eigenfunction of L + W.

By Rayleigh–Ritz, the smallest eigenvalue of  $-\Delta + W$  satisfies

$$\lambda_1 = \inf_{f:V \to \mathbb{R}} \frac{\sum_{(v_i, v_j) \in E} (f(v_i) - f(v_j))^2 + \sum_{v \in V} W(v) f(v)^2}{\sum_{v \in V} f(v)^2}$$

**Fig. 1** A *q*-regular tree (here, q = 3) of depth *k* (here, k = 2) with final layer being connected to a single additional vertex  $v_*$ 



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$$E = \lambda_1 \le q^k$$
 independently of  $W$ .

We can now use the equation

$$\phi(u) = \frac{E - W(u)}{\deg(u)}\phi(u) + \frac{1}{\deg(u)}\sum_{w \sim u}\phi(w)$$

and we shall restrict its use to vertices in the q-regular tree. The value of  $\phi$  then only depends on the level. We shall therefore write  $\phi(v) = \phi_i$  whenever the vertex v is in the *i*-th level where  $1 \le i \le k - 1$  (the case i = 0 and i = k will be ignored since the algebra is slightly different). The equation then simplifies to

$$\phi_i = \frac{\lambda_1 - W}{q+1}\phi_i + \frac{1}{q+1}(q \cdot \phi_{i+1} + \phi_{i-1}).$$

This can be rewritten as

$$(W - \lambda_1 + q + 1) \cdot \phi_i = q \cdot \phi_{i+1} + \phi_{i-1}.$$

For fixed q and  $W \gg q^k \ge E$ , this implies that approximately  $\phi_i \sim (q/W) \cdot \phi_{i+1}$  as  $W \to \infty$  which implies exponential decay. Conversely, we have

$$\log\left(1 + \frac{W - E}{\deg(v)}\right) \sim \log\left(\frac{W}{q}\right)$$

which implies, to leading order, the same kind of decay.

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### Declarations

Conflict of interest There are no associated data produced and there are no conflicts of interest to declare.

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