Stress analysis of steel beams strengthened with a bonded hygrothermal aged composite plate

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Abstract In this paper, the problem of interfacial stresses in steel beams strengthened with bonded hygrothermal aged composite laminates is analyzed using linear elastic theory. The analysis is based on the deformation compatibility approach developed by Tounsi (Int. J. Solids Struct. 43:4154-4174, 2006) where both the shear and normal stresses are assumed to be invariant across the adhesive layer thickness. The adopted model takes into account the adherend shear deformations by assuming a linear shear stress through the depth of the steel beam. This solution is intended for application to beams made of all kinds of materials bonded with a thin composite plate. For steel I-beam section, a geometrical coefficient ξ is determined to show the effect of the adherend shear deformations. This research is helpful for the understanding on mechanical behaviour of the interface and design of such structures.

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Nomenclature

A_{I}	The cross-sectional area of the steel
,	beam
A^{\prime}	Inverse of the extensional matrix A
$D^{'}$	Inverse of the flexural matrix D
E_{I}	The elastic modulus of the steel beam
E_a	The elastic modulus of adhesive
E_L	Longitudinal Young's modulus of
	FRP plate
E_T	Transversal Young's modulus of
	FRP plate
E_m	The elastic modulus of matrix
E_{f}	The elastic modulus of fibre
\check{G}_m	The transverse shear modulus of
	matrix
G_{f}	The transverse shear modulus of fibre
G_a	The transverse shear modulus of
	adhesive
G_I	The transverse shear modulus of the
	adherend 1
G_{LT}	The transverse shear modulus of FRP
	plate
Ι	The second moment of area
K_S	The shear stiffness of the adhesive
K_n	Normal stiffness of the adhesive per
	unit length
M(x)	The bending moment
	-

$M_T(x)$	The total applied moment		
$N_i (i = 1, 2)$	The longitudinal resultant force for		
	adherend "i"		
$U_1^N(x, y)$	Longitudinal displacements in steel		
	beam induced by the longitudinal		
	forces		
V(x)	Shear force		
V_m	Matrix volume fraction		
V_f C	Fibre volume fraction		
	Moisture		
Т	Temperature		
b_2	The width of the soffit plate		
q	The uniformly distributed load		
$t_i \ (i = 1, 2)$	The thickness of adherend "i"		
t_a	The thickness of adhesive		
u_1	The longitudinal displacement at the		
	base of adherend 1		
u_2	The longitudinal displacement at the		
	top of adherend 2		
u_I^N	The longitudinal displacement		
	induced by the longitudinal forces at		
	the interface between the upper		
	adherend and the adhesive		
$w_i \ (i = 1, 2)$	Vertical displacements of adherend " <i>i</i> "		
ε ₁	Strain at the base of adherend 1		
ε2	Strain at the top of adherend 2		
$arepsilon_2 \ arepsilon_M^M \ arepsilon_I$	Strains induced by the bending		
	moment at the steel beam		
ε_2^N	Strains induced by the longitudinal		
	forces at the steel beam		
$\sigma_{xy(1)}$	The shear stresses in steel beam		
σ_n	The normal stress in the adhesive		
γ1	The shear strain in steel beam		
$ au_{\mathrm{a}}$	The shear stresses through the		
	thickness of adhesive		
σ_{I}^{N}	Longitudinal normal stresses for steel		
	beam		
v_{LT}	Poisson's ratio of FRP plate		
v_m	Poisson's ratio of matrix		
v_f	Poisson's ratio of fibre		
ξ	Geometrical coefficient		

1 Introduction

Advanced composite materials have been used successfully for repairing metallic aircraft structures for a number of years (Baker 1984; Megueni et al. 2003, 2007). In Civil Engineering, composite plates have mainly been used to rehabilitate concrete structure, although the strengthening of metallic structures using FRP is gaining significant interesting in recent years. The main advantages of FRP are their high strength-to-weight ratio and their excellent resistance against corrosion and chemical attacks. An important topic arising in the study of plated steel beams is the evaluation of interactions at steel–FRP interface. These interactions, in fact, permit the transmission of stresses from the core to the plate; if they go over a limit value the premature failure of the strengthened beam can occur.

The determination of interfacial stresses has been researched for the last decade for steel or concrete beams bonded with either steel or advanced composite materials. In particular, several closed-form analytical solutions have been developed (Vilnay 1988; Roberts 1989; Roberts and Haji-Kazemi 1989; Stratford and Cadei 2006; Malek et al. 1998; Smith and Teng 2001; Benyoucef et al. 2006; Tounsi 2006). All these solutions are for linear elastic materials and employ the same key assumption that the adhesive is subject to normal and shear stresses that are constant across the thickness of the adhesive layer. It is this key assumption that enables relatively simple closedform solutions to be obtained. In the existing solutions, two different approaches have been employed. Roberts (1989) and Roberts and Haji-Kazemi (1989) used a staged analysis approach, while Vilnay (1988), Stratford and Cadei (2006), Malek et al. (1998) and Smith and Teng (2001) considered directly deformation compatibility conditions. Recently, Tounsi and his co-workers (Benyoucef et al. 2006; Tounsi 2006) developed theoretical solutions for interfacial stresses in concrete beams strengthened with FRP plate based also on deformation compatibility conditions. Benyoucef et al. (2006) present an alternative theoretical interfacial stress analysis, where, the FRP plate fibre orientation is considered and the flexural rigidity of the composite plate is estimated using lamination theory. Tounsi (2006) presented a new theoretical solution in which the adherend shear deformations have been included.

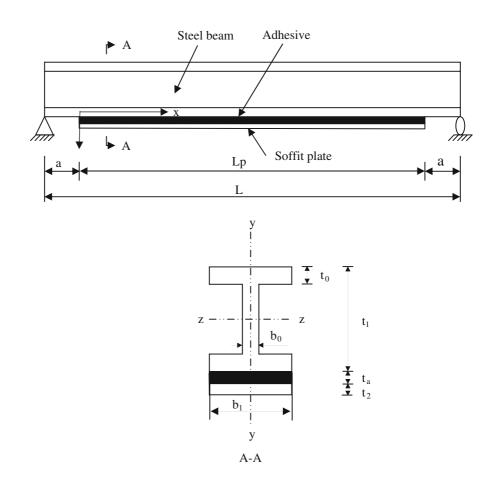
In this paper, we present an alternative theoretical interfacial analysis for simply supported steel beams bonded with a thin hygrothermal aged FRP plate. It Fig. 1 Simply supported

beam strengthened with bonded composite plate

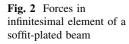
is well known that during the operational life, the variation of temperature and moisture reduces the elastic moduli and degrades the strength of composite material (Megueni et al. 2003, 2007; Shen and Springer 1981; Adams and Miller 1977; Bowles and Tompkins 1989; Shen 2001; Tounsi and Amara 2005; Amara et al. 2005; Tounsi et al. 2005). The material properties of the bonded composite plate are assumed to be functions of temperature and moisture. This effect is studied here, due to technological considerations. For example, when we use these plates to reinforce beams, they can be already exposed to the environmental conditions. Thus, these plates are aged with time and the mechanical properties will be reduced. In this study, we aim to show if the reduction in stiffness due to hygrothermal ageing has an effect on the interfacial stresses. Both ambient temperature and moisture are assumed to have a uniform distribution. The plate is fully saturated such that the variation of temperature and moisture are independent of time and position. The adopted model describes better the actual response of the FRP-steel hybrid beam and permits the evaluation of the interfacial stresses, the knowledge of which is very important in the design of such structures.

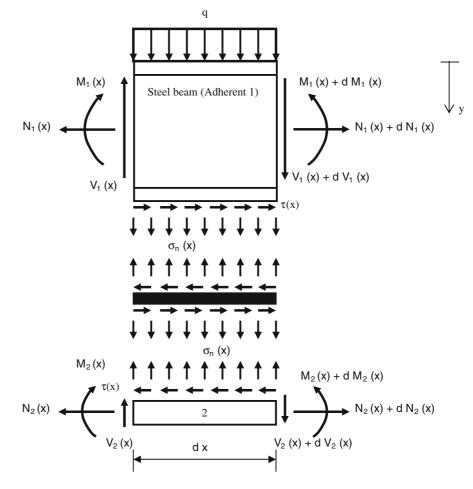
2 The method of solution

A differential section dx, can be cut out from the FRP strengthened steel beam (Fig. 1), as shown in Fig. 2. The composite beam is made from three materials: steel, adhesive layer and FRP reinforcement. In the present analysis, linear elastic behaviour is regarded to be for all the materials; the adhesive is assumed to only play a role in transferring the stresses from the steel to the FRP reinforcement and the stresses in the adhesive layer do not change through the direction of the thickness.



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2.1 Basic equation of elasticity

The strain $\varepsilon_1(x)$ in the steel beam near adhesive interface can be expressed as

$$\varepsilon_1(x) = \frac{\mathrm{d}u_1(x)}{\mathrm{d}x} = \varepsilon_1^M(x) + \varepsilon_1^N(x) \tag{1}$$

where $u_I(x)$ is the longitudinal displacement at the base of steel beam. $\varepsilon_1^M(x)$ is the strain induced by the bending moment at the adherend 1 and it is written as follow:

$$\varepsilon_1^M(x) = \frac{t_1}{2E_1 I_1} M_1(x) \tag{2}$$

where $M_1(x)$, is the bending moment applied in the steel beam at its mid-height; E_1 is Young's moduli of the steel; I_1 is the second moment area; t_1 is the depth of steel beam. $\varepsilon_1^N(x)$ is the unknown longitudinal strain of the steel beam, at the adhesive interface and it is due to the longitudinal forces. This strain is given as follow:

$$\varepsilon_1^N(x) = \frac{\mathrm{d}u_1^N(x)}{\mathrm{d}x} \tag{3}$$

where u_1^N represents the longitudinal force-induced adhesive displacement at the interface between the steel beam and the adhesive.

To determine the unknown longitudinal strain $\varepsilon_1^N(x)$, shear deformation of the steel beam is incorporated in this analysis. It is reasonable to assume that the shear stresses, which develop in the adhesive, are continuous across the adhesive–adherend interface. In addition, equilibrium requires the shear stress to be zero at the free surface. Using the same methodology developed by Tounsi (2006), this effect is taken into account. The importance of including shear-lag effect of the adherents was also shown by Tsai et al. (1998). In the present analysis, shear deformations of the FRP plate is ignored. However, a linear shear stress through the depth of the steel beam is assumed.

$$\sigma_{xy(1)} = \frac{\tau(x)}{t_1} y \tag{4}$$

Equation 4 is based on zero shear stresses at the top surface of the upper adherend (i.e. at y = 0) and $\sigma_{xy(1)} = \tau_a$ at $y = t_1$. Then with a linear material constitutive relationship the adherend shear strain γ_1 for the adherend 1 (steel beam) is written as:

$$\gamma_{xy(1)} = \gamma_1 = \frac{\tau_a}{G_1 t_1} y \tag{5}$$

 G_I is the transverse shear moduli of the adherend 1 (steel beam), and $\tau_a = \tau(x)$.

The longitudinal displacement functions U_1^N for the upper adherend, due to the longitudinal forces, is given by

$$U_1^N(y) = U_1^N(0) + \int_0^y \gamma_1(y) dy = U_1^N(0) + \frac{\tau_a}{2G_1 t_1} y^2 \quad (6)$$

where $U_1^N(0)$ represents the displacement at the top surface of the upper adherend (due to the longitudinal forces).

Note that due to the perfect bonding of the joints, the displacements are continuous at the interfaces between the adhesive and adherends. As a result, the u_1^N (the adhesive displacement at the interface between the adhesive and upper adherend) should be the same as the upper adherend displacement at the interface. Based on Eq. 6 the u_1^N can be expressed as:

$$u_1^N = U_1^N(y = t_1) = U_1^N(0) + \frac{\tau_a t_1}{2G_1}$$
(7)

Using Eq. 7, Eq. 6 can be rewritten as:

$$U_1^N(y) = u_1^N + \frac{\tau_a}{2G_1 t_1} y^2 - \frac{\tau_a t_1}{2G_1}$$
(8)

The longitudinal resultant force, N_I for the upper adherends having I section (Fig. 1), is:

$$N_{1} = b_{1} \int_{0}^{t_{0}} \sigma_{1}^{N}(y) dy + b_{0} \int_{t_{0}}^{t_{1}-t_{0}} \sigma_{1}^{N}(y) dy + b_{1} \int_{t_{1}-t_{0}}^{t_{1}} \sigma_{1}^{N}(y) dy$$

$$(9)$$

where σ_1^N is longitudinal normal stress for the upper adherend. By changing this stress into function of displacement and substituting Eq. 8 into the displacement, Eq. 9 can be rewritten as:

$$N_{1} = E_{1}b_{1}\int_{0}^{t_{0}} \frac{\mathrm{d}U_{1}^{N}}{\mathrm{d}x}\mathrm{d}y + E_{1}b_{0}\int_{t_{0}}^{t_{1}-t_{0}} \frac{\mathrm{d}U_{1}^{N}}{\mathrm{d}x}\mathrm{d}y + E_{1}b_{1}\int_{t_{1}-t_{0}}^{t_{1}} \frac{\mathrm{d}U_{1}^{N}}{\mathrm{d}x}\mathrm{d}y$$
(10)
$$N_{1} = E_{1}A_{1}\left[\frac{\mathrm{d}u_{1}^{N}}{\mathrm{d}x} - \frac{1}{6GtA_{1}}\left(b_{1}\left((t_{1}-t_{0})^{3}\right)\right)\right]$$

Hence, the longitudinal strains induced by the longitudinal forces Eq. 3 can be expressed as

$$\varepsilon_{1}^{N}(x) = \frac{\mathrm{d}u_{1}^{N}}{\mathrm{d}x} = \frac{N_{1}}{E_{1}A_{1}} + \frac{1}{6G_{1}t_{1}A_{1}} \left(b_{1} \left((t_{1} - t_{0})^{3} - t_{0}^{3} - t_{1}^{3} + 6t_{1}^{2}t_{0} \right) + b_{0} \left(3t_{1}^{2}(t_{1} - 2t_{0}) - (t_{1} - t_{0})^{3} + t_{0}^{3} \right) \right) \frac{\mathrm{d}\tau(x)}{\mathrm{d}x}$$
(12)

Substituting Eqs. 2 and 12 into Eq. 1, this latter becomes:

$$\varepsilon_{1}(x) = \frac{du_{1}(x)}{dx} = \frac{t_{1}}{2E_{1}I_{1}}M_{1}(x) + \frac{N_{1}(x)}{E_{1}A_{1}} + \frac{b_{1}\left((t_{1}-t_{0})^{3}-t_{0}^{3}-t_{1}^{3}+6t_{1}^{2}t_{0}\right) + b_{0}\left(3t_{1}^{2}(t_{1}-2t_{0})-(t_{1}-t_{0})^{3}+t_{0}^{3}\right)}{6G_{1}t_{1}A_{1}}\frac{d\tau(x)}{dx}$$

$$(13)$$

 A_1 is the cross-sectional area.

On the other hand, the laminate theory is used to determine the stress and strain of the externally bonded composite plate in order to investigate the whole mechanical performance of the compositestrengthened structure. The effective moduli of the composite laminate are varied by the orientation of the fibre directions and arrangements of the laminate patterns. The classical laminate theory (CLT) is used to estimate the strain of the composite plate (Hera-kovich 1998), i.e.

$$\begin{cases} \varepsilon^{0} \\ k \end{cases} = \begin{bmatrix} A' & B' \\ C' & D' \end{bmatrix} \begin{cases} N \\ M \end{cases}$$
 (14)

where $[A'] = [A]^{-1} + [A]^{-1}[B][D^x]^{-1}[B][A]^{-1}, [B'] = -[A^{-1}][B][D^x]^{-1}, [C'] = -[D^x]^{-1}[B][A]^{-1}, [D'] = ([D] - [B][A]^{-1}[B])^{-1}, [D^x] = [D] - [B][A]^{-1}[B]$ and Extensional matrix:

$$A_{ij} = \sum_{k=1}^{NN} \overline{Q}_{ij}^{(k)} ((y_2)_k - (y_2)_{k-1})$$
(15)

Extensional-bending coupled matrix:

$$B_{ij} = \frac{1}{2} \sum_{k=1}^{NN} \overline{Q}_{ij}^{(k)} \left((y_2^2)_k - (y_2^2)_{k-1} \right)$$
(16)

Flexural matrix:

$$D_{ij} = \frac{1}{3} \sum_{k=1}^{NN} \overline{Q}_{ij}^{(k)} \left((y_2^3)_k - (y_2^3)_{k-1} \right)$$
(17)

The subscript *NN* represents the number of laminate layers of the FRP plate. Parameter \overline{Q}_{ij} can be estimated by using the off-axis orthotropic plate theory (Herakovich 1998). Assume that the ply arrangement of the plate is symmetrical with respect to the mid-plane axis $y_2 = 0$. A great simplification in laminate analysis then occurs by assuming that the coupling matrix *B* is identically zero (Herakovich 1998). Therefore Eqs. 14–17 can be simplified to the following matrix form for a plate with a width of b_2 :

$$\left\{\varepsilon^{0}\right\} = \left[A'\right]\left\{N\right\}_{2} \text{ and } \left\{k\right\} = \left[D'\right]\left\{M\right\}_{2}$$
 (18)

where:

$$\left\{ \varepsilon^{0} \right\}_{2} = \left\{ \begin{array}{c} \varepsilon^{0}_{x} \\ \varepsilon^{0}_{y} \\ \gamma^{0}_{xy} \end{array} \right\}; \left\{ N \right\}_{2} = \left\{ \begin{array}{c} N_{x} \\ N_{y} \\ N_{xy} \end{array} \right\}_{2}; \left\{ k \right\}$$
$$= \left\{ \begin{array}{c} k_{x} \\ k_{y} \\ k_{xy} \end{array} \right\} \text{ and } \left\{ M \right\}_{2} = \left\{ \begin{array}{c} M_{x} \\ M_{y} \\ M_{xy} \end{array} \right\}_{2}$$
(19)

In the present study, only an axial load N_x and the bending moment M_x in the beam's longitudinal axis are considered, i.e. $N_y = N_{xy} = 0$ and $M_y =$ $M_{xy} = 0$. Therefore, Eq. 18 can be simplified to:

$$\varepsilon_x^0 = A'_{11}N_x \frac{1}{b_2}$$
 and $k_x = D'_{11}M_x \frac{1}{b_2}$ (20)

Using CLT, the strain at the top of the FRP plate 2 is given as:

$$\varepsilon_2(x) = \varepsilon_x^0 - k_x \frac{t_2}{2} \tag{21}$$

Substituting Eq. 20 in 21 gives the following equation:

$$\varepsilon_{2}(x) = \frac{\mathrm{d}u_{2}(x)}{\mathrm{d}x} = -D_{11}^{'}\frac{t_{2}}{2b_{2}}M_{2}(x) + A_{11}^{'}\frac{N_{2}(x)}{b_{2}} \qquad (22)$$

where:

$$N_2(x) = N_x \quad \text{and} \quad M_2(x) = M_x \tag{23}$$

The subscripts 1 and 2 denote adherends 1 and 2, respectively. M(x), N(x) and V(x) are the bending moment, axial and shear forces in each adherend.

It is well known in many studies (Shen and Springer 1981; Adams and Miller 1977; Bowles and Tompkins 1989; Shen 2001; Tounsi and Amara 2005; Amara et al. 2005; Tounsi and Amara 2005) that the material properties of composite are function of temperature and moisture. In terms of a micromechanical model of laminate, the material properties may be written as (Tsai and Hahn 1980)

$$E_L = V_f E_f + V_m E_m \tag{24}$$

$$\frac{1}{E_T} = \frac{V_f}{E_f} + \frac{V_m}{E_m} - V_f V_m \frac{v_f^2 \left(\frac{E_m}{E_f}\right) + v_m^2 \left(\frac{E_f}{E_m}\right) - 2v_f v_m}{V_f E_f + V_m E_m}$$
(25)

$$\frac{1}{G_{LT}} = \frac{V_f}{G_f} + \frac{V_m}{G_m} \tag{26}$$

$$v_{LT} = V_f v_f + V_m v_m \tag{27}$$

In the above equations, V_f and V_m are the fibre and matrix volume fractions and are related by

$$V_f + V_m = 1 \tag{28}$$

 E_f , G_f and v_f are the Young's modulus, shear modulus and Poisson's ratio, respectively, of the fibre, and E_m , G_m and v_m are corresponding properties for the matrix. It is assumed that E_m is a function of temperature and moisture, as is shown in Sect. 3, then E_L , E_T and G_{LT} are also functions of temperature and moisture. By adopting the equilibrium conditions of the steel beam, we have: Along *x*-direction:

$$\frac{\mathrm{d}N_1(x)}{\mathrm{d}x} = -\tau(x)b_2\tag{29}$$

where $\tau(x)$ is shear stress in the adhesive layer. Along *y*-direction:

$$\frac{\mathrm{d}V_1(x)}{\mathrm{d}x} = -[\sigma_n(x)b_2 + q] \tag{30}$$

where $V_1(x)$ is shear force applied in the steel beam; $\sigma_n(x)$ is normal stress in the adhesive layer; *q* is the uniformly distributed load and b_1 is width of steel beam. Moment equilibrium:

$$\frac{\mathrm{d}M_1(x)}{\mathrm{d}x} = V_1(x) - \tau(x)b_2\frac{t_1}{2} \tag{31}$$

The equilibrium of the external FRP reinforcement along x-, y-direction and moment equilibrium can be also written as: Along x-direction:

$$\frac{\mathrm{d}N_2(x)}{\mathrm{d}x} = \tau(x)b_2\tag{32}$$

Along y-direction:

$$\frac{\mathrm{d}V_2(x)}{\mathrm{d}x} = \sigma_n(x)b_2\tag{33}$$

Moment equilibrium:

$$\frac{\mathrm{d}M_2(x)}{\mathrm{d}x} = V_2(x) - \tau(x)b_2\frac{t_2}{2} \tag{34}$$

where $V_2(x)$ is shear force applied in the external FRP reinforcement.

2.2 Shear stress distribution along the FRP–steel interface

The shear stress in the adhesive can be expressed as follows:

$$\tau(x) = K_s \Delta u(x) = K_s [u_2(x) - u_1(x)]$$
(35)

where K_s is shear stiffness of the adhesive per unit length and can be deduced as:

$$K_s = \frac{\tau(x)}{\Delta u(x)} = \frac{\tau(x)}{\Delta u(x)/t_a} \frac{1}{t_a} = \frac{G_a}{t_a}$$
(36)

 $\Delta u(x)$ is relative horizontal displacement at the adhesive interface; G_a is the shear modulus in the adhesive and t_a is the thickness of the adhesive.Differentiating Eqs. 13, 22 and 35 with respect to x, respectively:

$$\frac{d\tau(x)}{dx} = K_{S} \left(A'_{11} \frac{N_{2}(x)}{b_{2}} - D'_{11} \frac{t_{2}}{2b_{2}} M_{2}(x) - \frac{t_{1}}{2E_{1}I_{1}} M_{1}(x) - \frac{N_{1}(x)}{E_{1}A_{1}} \right) - \frac{K_{S}}{6G_{1}t_{1}A_{1}} \left[b_{1} \left((t_{1} - t_{0})^{3} - t_{0}^{3} - t_{1}^{3} + 6t_{1}^{2}t_{0} \right) + b_{0} \left(3t_{1}^{2}(t_{1} - 2t_{0}) - (t_{1} - t_{0})^{3} + t_{0}^{3} \right) \right] \frac{d\tau(x)}{dx}$$
(37)

Assuming equal curvature in the beam and the FRP plate, the relationship between the moments in the two adherends can be expressed as:

$$M_1(x) = RM_2(x) \tag{38}$$

with

$$R = \frac{E_1 I_1 D_{11}}{b_2} \tag{39}$$

Note that Eq. 38 is an approximation which was also used by previous researchers studying interfacial stresses between a beam and a strengthening plate, and was found to lead to accurate results (Smith and Teng 2001).

Moment equilibrium of the differential segment of the plated beam in Fig. 2 gives:

$$M_T(x) = M_1(x) + M_2(x) + N(x) \left[\frac{t_1 + t_2}{2} + t_a \right] \quad (40)$$

where, $M_T(x)$ is the total applied moment and from Eqs. 29 and 32, the axial forces are given as:

$$N_{1}(x) = -N(x) = -b_{2} \int_{0}^{x} \tau(x) \text{ and}$$
$$N_{2}(x) = N(x) = b_{2} \int_{0}^{x} \tau(x)$$
(41)

The bending moment in each adherend, expressed as a function of the total applied moment and the interfacial shear stress, is given as

$$M_{1}(x) = \frac{R}{R+1} \left[M_{T}(x) - b_{2} \int_{0}^{x} \tau(x) \left(\frac{t_{1} + t_{2}}{2} + t_{a} \right) \mathrm{d}x \right]$$
(42)

and

$$M_{2}(x) = \frac{1}{R+1} \left[M_{T}(x) - b_{2} \int_{0}^{x} \tau(x) \left(\frac{t_{1} + t_{2}}{2} + t_{a} \right) dx \right]$$
(43)

The first derivative of the bending moment in each adherend gives:

$$\frac{\mathrm{d}M_1(x)}{\mathrm{d}x} = \frac{R}{R+1} \left[V_T(x) - b_2 \tau(x) \left(\frac{t_1 + t_2}{2} + t_a \right) \right]$$
(44)

and

$$\frac{\mathrm{d}M_2(x)}{\mathrm{d}x} = \frac{1}{R+1} \left[V_T(x) - b_2 \tau(x) \left(\frac{t_1 + t_2}{2} + t_a \right) \right]$$
(45)

Differentiating Eq. 37:

$$\frac{d^{2}\tau(x)}{dx^{2}} = K_{S} \left(\frac{A_{11}'}{b_{2}} \frac{dN_{2}(x)}{dx} - D_{11}' \frac{t_{2}}{2b_{2}} \frac{dM_{2}(x)}{dx} - \frac{t_{1}}{2E_{1}I_{1}} \frac{dM_{1}(x)}{dx} - \frac{1}{E_{1}A_{1}} \frac{dN_{1}(x)}{dx} \right) - \frac{K_{S}}{6G_{1}t_{1}(2b_{1}t_{0} + b_{0}(t_{1} - 2t_{0}))} \times \left[b_{1} \left((t_{1} - t_{0})^{3} - t_{0}^{3} - t_{1}^{3} + 6t_{1}^{2}t_{0} \right) + b_{0} \left(3t_{1}^{2}(t_{1} - 2t_{0}) - (t_{1} - t_{0})^{3} + t_{0}^{3} \right) \right] \frac{d^{2}\tau(x)}{dx^{2}}$$
(46)

Substitution of the shear forces (Eqs. 44 and 45) and axial forces Eq. 41 into Eq. 46 gives the following

governing differential equation for the interfacial shear stress.

$$\frac{d^{2}\tau(x)}{dx^{2}} - K_{1}\left(A_{11}^{'} + \frac{b_{2}}{E_{1}A_{1}} + \frac{(t_{1}+t_{2})(t_{1}+t_{2}}{2} + t_{a})}{E_{1}I_{1}D_{11}^{'} + b_{2}}b_{2}D_{11}^{'}\right)\tau(x) + \frac{K_{1}}{2}\left(\frac{t_{1}+t_{2}}{E_{1}I_{1}D_{11}^{'} + b_{2}}D_{11}^{'}\right)V_{T}(x) = 0$$
(47)

where

$$K_1 = \frac{1}{\left(\frac{t_a}{G_a} + \frac{t_1}{3G_1}\xi\right)} \tag{48}$$

and ξ is a geometrical coefficient which is given as

$$\xi = \frac{1}{2A_1t_1^2} \Big[b_1 \Big(-t_0^3 + 6t_0t_1^2 - t_1^3 + (t_1 - t_0)^3 \Big) \\ + b_0 \Big(3t_1^2(t_1 - 2t_0) - (t_1 - t_0)^3 + t_0^3 \Big) \Big]$$
(49)

For a rectangular section $(b_1 = b_0)$, $\xi = 1$ which corresponds to the same expression given by Tounsi (2006) by neglecting shear deformations of the FRP plate. However, for I-beam section (the present case) we have $\xi < 1$.For simplicity, the general solutions presented below are limited to loading which is either concentrated or uniformly distributed over part or the whole span of the beam, or both. For such loading, $d^2V_T(x)/dx^2 = 0$, and the general solution to Eq. 47 is given by:

$$\tau(x) = B_1 \cosh(\lambda x) + B_2 \sinh(\lambda x) + m_1 V_T(x) \qquad (50)$$

where

$$\ell^{2} = K_{1} \left(A_{11}^{'} + \frac{b_{2}}{E_{1}A_{1}} + \frac{\left(\frac{t_{1}+t_{2}}{2}\right)\left(\frac{t_{1}+t_{2}}{2} + t_{a}\right)}{E_{1}I_{1}D_{11}^{'} + b_{2}} b_{2}D_{11}^{'} \right)$$
(51)

and

$$m_1 = \frac{K_1}{2\lambda^2} \left(\frac{t_1 + t_2}{E_1 I_1 D_{11}' + b_2} D_{11}' \right)$$
(52)

 B_1 and B_2 are constant coefficients determined from the boundary conditions. In the present study, a simply supported beam is investigated which is subjected to a uniformly distributed load as shown in Fig. 1. The interfacial shear stress for this load case at any point is written as:

$$\tau(x) = \left[\frac{m_2 a}{2}(L-a) - m_1\right] \frac{q e^{-\lambda x}}{\lambda} + m_1 q \left(\frac{L}{2} - a - x\right) \quad 0 \le x \le L_p$$
(53)

where *q* is the uniformly distributed load and *x*, *a*, *L* and L_p are defined in Fig. 1. And m_2 is given as follow:

$$m_2 = \frac{K_1 t_1}{2E_1 I_1} \tag{54}$$

In the case where the RC beam is subjected to a two symmetric point loads as shown in Fig. 3, the general solution for the interfacial shear stress is given by the following expressions Tounsi (2006) a < b

$$\tau(x) = \begin{cases} \frac{m_2}{\lambda} Pae^{-\lambda x} + m_1 P \cosh(\lambda x) e^{-k} & 0 \le x \le (b-a) \\ \frac{m_2}{\lambda} Pae^{-\lambda x} + m_1 P \sinh(k) e^{-\lambda x} & (b-a) \le x \le \frac{L_p}{2} \end{cases}$$
(55)

Fig. 3 Comparison of interfacial shear stress of the steel plated RC beam with the experimental results

$$\tau(x) = \frac{m_2}{\lambda} P b e^{-\lambda x} \quad 0 \le x \le L_p \tag{56}$$

where *P* is the concentrated load and $k = \lambda(b - a)$. The expression of m_1 and m_2 takes into considerations the shear deformation of adherends.

2.3 Normal stress distribution along the FRP-steel interface

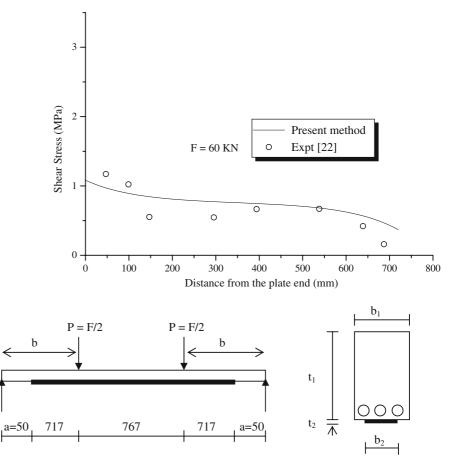
The normal stress in the adhesive can be expressed as follows:

$$\sigma_n(x) = K_n \Delta w(x) = K_n[w_2(x) - w_1(x)]$$
(57)

where K_n is normal stiffness of the adhesive per unit length and can be deduced as:

$$K_n = \frac{\sigma_n(x)}{\Delta w(x)} = \frac{\sigma_n(x)}{\Delta w(x)/t_a} \left(\frac{1}{t_a}\right) = \frac{E_a}{t_a}$$
(58)

 $w_1(x)$ and $w_2(x)$ are the vertical displacements of



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adherend 1 and 2, respectively.Differentiating Eq. 50 twice results in:

$$\frac{d^2\sigma_n(x)}{dx^2} = K_n \left[\frac{d^2w_2(x)}{dx^2} - \frac{d^2w_1(x)}{dx^2} \right]$$
(59)

Considering the moment–curvature relationships for the beam to be strengthened and the external reinforcement, respectively:

$$\frac{\mathrm{d}^2 w_1(x)}{\mathrm{d}x^2} = -\frac{M_1(x)}{E_1 I_1}, \quad \frac{\mathrm{d}^2 w_2(x)}{\mathrm{d}x^2} = -\frac{D_{11}' M_2(x)}{b_2} \tag{60}$$

Based on the equilibrium Eqs. 29–34, the governing differential equations for the deflection of adherends 1 and 2, expressed in terms of the interfacial shear and normal stresses, are given as follows:Adherend 1:

$$\frac{\mathrm{d}^4 w_1(x)}{\mathrm{d}x^4} = \frac{1}{E_1 I_1} b_2 \sigma_n(x) + \frac{t_1}{2E_1 I_1} b_2 \frac{\mathrm{d}\tau(x)}{\mathrm{d}x} + \frac{q}{E_1 I_1}$$
(61)

Adherend 2:

$$\frac{\mathrm{d}^4 w_2(x)}{\mathrm{d}x^4} = -D_{11}^{'}\sigma_n(x) + D_{11}^{'}\frac{t_2}{2}\frac{\mathrm{d}\tau(x)}{\mathrm{d}x} \tag{62}$$

Substitution of Eqs. 59 and 60 into the fourth derivation of the interfacial normal stress obtainable from Eq. 55 gives the following governing differential equation for the interfacial normal stress:

$$\frac{d^{4}\sigma_{n}(x)}{dx^{4}} + \frac{E_{a}}{t_{a}} \left(D_{11}^{'} + \frac{b_{2}}{E_{1}I_{1}} \right) \sigma_{n}(x) - \frac{E_{a}}{t_{a}} \left(D_{11}^{'}y_{2} - \frac{t_{1}b_{2}}{2E_{1}I_{1}} \right) \frac{d\tau(x)}{dx} + \frac{qE_{a}}{t_{a}E_{1}I_{1}} = 0 \quad (63)$$

The general solution to this fourth-order differential equation is:

$$\sigma_n(x) = e^{-\beta x} [C_1 \cos(\beta x) + C_2 \sin(\beta x)]$$
$$+ e^{\beta x} [C_3 \cos(\beta x) + C_4 \sin(\beta x)] - n_1 \frac{\mathrm{d}\tau(x)}{\mathrm{d}x} - n_2 q(64)$$

For large values of x it is assumed that the normal stress approaches zero, and as a result $C_3 = C_4 = 0$. The general solution therefore becomes:

$$\sigma_n(x) = e^{-\beta x} [C_1 \cos(\beta x) + C_2 \sin(\beta x)] - n_1 \frac{\mathrm{d}\tau(x)}{\mathrm{d}x} - n_2 q$$
(65)

where

$$\beta = \sqrt{4} \frac{E_a}{4t_a} \left(\frac{b_2}{E_1 I_1} + D'_{11} \right) \tag{66}$$

$$n_{1} = \left(\frac{t_{1}b_{2} - D_{11}^{'}E_{1}I_{1}t_{2}}{2(D_{11}^{'}E_{1}I_{1} + b_{2})}\right)$$
(67)

and

$$n_2 = \frac{1}{D'_{11}E_1I_1 + b_2} \tag{68}$$

The constants C_1 and C_2 in Eq. 63 are written as follow:

$$C_{1} = \frac{E_{a}}{2\beta^{3}t_{a}E_{1}I_{1}}[V_{T}(0) + \beta M_{T}(0)] - \frac{n_{3}}{2\beta^{3}}\tau(0) + \frac{n_{1}}{2\beta^{3}}\left(\frac{d^{4}\tau(0)}{dx^{4}} + \beta\frac{d^{3}\tau(0)}{dx^{3}}\right)$$
(69)

$$C_2 = -\frac{E_a}{2\beta^2 t_a E_1 I_1} M_T(0) - \frac{n_1}{2\beta^2} \frac{d^3 \tau(0)}{dx^3}$$
(70)

where

$$n_3 = \frac{E_a b_2}{t_a} \left(\frac{t_1}{2E_1 I_1} - \frac{D_{11}' t_2}{2b_2} \right) \tag{71}$$

The above expressions for the constants C_1 and C_2 have been left in terms of the bending moment $M_T(0)$ and shear force $V_T(0)$ at the end of the soffit plate.

3 Results and discussion

A computer code based on the preceding equations was written to compute the interfacial stresses in a steel beam bonded with a hygrothermal aged FRP plate.

Graphite/epoxy composite material was selected in the present examples as a bonded plate. However, the analysis is equally applicable to other types of composite material. For these examples the thickness of each ply is 0.125 mm, the width is $b_2 = 150$ mm and the material properties adopted are (Adams and Miller 1977; Bowles and Tompkins 1989; Shen 2001; Tounsi and Amara 2005; Amara et al. 2005; Tounsi and Amara 2005): E_f = 230.0 GPa, G_f = 9.0 GPa, $v_f = 0.203$, $v_m = 0.34$ and $E_m = (3.51 - 0.003T - 0.142C)$ GPa, in which $T = T_0 + \Delta T$ and $T_0 = 25^{\circ}$ C (room temperature), and $C = C_0 + \Delta C$ and $C_0 =$ 0 wt% H₂O. Using the formula of the Young's modulus E_m and the relations (24)–(28), we calculate E_L , E_T , G_{LT} and v_{LT} .

3.1 Comparison with experimental results

To validate the present method, a rectangular section $(\xi = 1)$ is used here. One of the tested beams bonded with steel plate by Jones et al. (1988), beam F31, is analysed here using the present improved solution. The beam is simply supported and subjected to four-point bending, each at the third point. The geometry and materials properties of the specimen are summarized in Table 1.

The interfacial shear stress distributions in the beam bonded with a soffit steel plate under the applied load 60 kN, i.e. P = 30 kN in Fig. 3, are compared between the experimental results and those obtained by the present method. As it can be seen from Fig. 3, the predicted theoretical results are in reasonable agreement with the experimental results.

3.2 Theoretical parametric study

In this section, numerical results of the present solutions are presented to study the effect of various parameters on the distributions of the interfacial stresses in a steel beam bonded with an FRP plate. These results are intended to demonstrate the main characteristics of interfacial stress distributions in these strengthened beams.

The steel beam is simply supported and subjected to a uniformly distributed load q = 50 kN/m. The Young's modulus and Poisson's ratio are respectively: $E_I = 210$ GPa and $v_1 = 0.3$. The span of the steel beam is 3,000 mm; the width is $b_I = 150$ mm; the total depth is $t_I = 300$ mm; the depth of the flange is $t_0 = 10.7$ mm; the thickness of the web is $b_0 = 7.1$ mm; the distance from the support to the end of the plate is 300 mm.

The geometric and material properties of the adhesive layer are: the thickness is $t_a = 4$ mm; the Young's modulus and Poisson's ratio are respectively, $E_a = 3$ GPa and $v_a = 0.35$.

Table 1 Dimensions and material properties

Concrete	$b_I = 155 \text{ mm}$	$t_I = 225 \text{ mm}$	$E_I = 31,000 \text{ MPa}$
Steel	$b_2 = 125 \text{ mm}$	$t_2 = 6 \text{ mm}$	$E_2 = 200,000 \text{ MPa}$
Adhesive	$b_a = 123 \text{ mm}$	$t_a = 1.5 \text{ mm}$	$E_a = 280 \text{ MPa}$
			$G_a = 108$ MPa

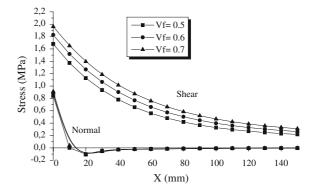


Fig. 4 The effect of fiber volume fraction on the variation of both shear and normal adhesive stresses in steel beam bonded with composite plate [016]_S. $\Delta T = 0^{\circ}$ C, $\Delta C = 0\%$

3.2.1 Fiber volume fractions effect

Figure 4 shows, the effect of fiber volume fractions V_f (= 0.5, 0.6 and 0.7) on the variation of shear and normal adhesive stresses. It can be seen that the interfacial shear stresses are reduced with decreases in fiber volume fraction. However, almost no effect is observed on the variation of interfacial normal stresses.

3.2.2 Hygrothermal effect on the adhesive stresses

This effect is studied here, due to technological considerations. For example, when we use these plates to strengthen beams, they can be already exposed to the environmental conditions. Thus, these plates are aged with time and the mechanical properties will be reduced. In this study, we aim to show if the reduction in stiffness of the bonded composite plate due to hygrothermal ageing has an effect on the interfacial stresses.

From results presented in Fig. 5 we can conclude that the hygrothermal ageing of the bonded composite plate has no effect on the variation of adhesive stresses.

3.2.3 Effect of the adhesive layer thickness

Figure 6 shows the effects of the thickness of the adhesive layer on the interfacial stresses. It is seen that increasing the thickness of the adhesive layer leads to significant reduction in the peak interfacial stresses. Thus using thick adhesive layer, especially in the vicinity of the edge, is recommended.

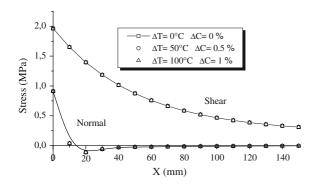


Fig. 5 Hygrothermal effect on the variation of both shear and normal adhesive stresses in steel beam bonded with composite plate $[016]_{S}$. (Vf = 0.7)

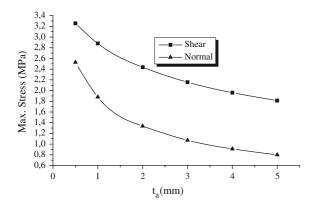


Fig. 6 Effect of adhesive layer thickness on interfacial stresses in steel beam bonded with composite plate [016]_S. (Vf = 0.7). $\Delta T = 0^{\circ}$ C, $\Delta C = 0\%$

3.2.4 Effect of FRP plate thickness

Peak shear and peeling stresses for various thicknesses of the FRP plate within the range of 1–6 mm appear in Fig. 7. The results reveal that thickness of the FRP plates significantly increases the edge peeling and shear stresses. This relation is an outcome of the local bending effects in the FRP plate governed by the flexural rigidity of the plate. Thus any increase in the flexural rigidity leads to an increase in the magnitude of the edge stresses.

3.2.5 Effect on Plate Length of the strengthened beam region L_p

The influence of length of the strengthened beam region L_p appears in Fig. 8. It is seen that, as the plate terminates further away from the supports, the

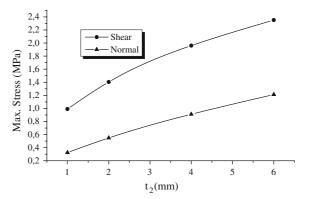


Fig. 7 Influence of thickness of frp plate on edge stresses in steel beam bonded with composite plate $[016]_{s}$. (Vf = 0.7). $\Delta T = 0^{\circ}$ C, $\Delta C = 0\%$

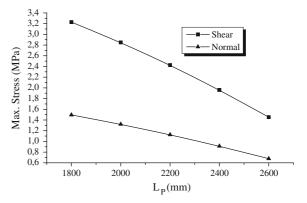


Fig. 8 Effect of plate length on edge stresses interfacial stresses in steel beam bonded with composite plate $[016]_{\rm S}$. (Vf = 0.7). $\Delta T = 0^{\circ}$ C, $\Delta C = 0\%$

interfacial stresses increase significantly. This result reveals that in any case of strengthening, including cases where retrofitting is required in a limited zone of maximum bending moments at midspan, it is recommended to extend the strengthening strip as close as possible to the support lines.

4 Conclusion

A systematic rigorous general approach for the analysis of interfacial stresses in steel beams strengthened with externally bonded hygrothermal aged FRP plate has been presented. This approach is based on elastic foundation model in which the adherend shear deformations have been included by assuming a linear shear stress through the depth of the steel beam. By comparing with experimental results, the present closed-solution provides satisfactory predictions to the interfacial shear stress in the plated beams.

The material properties of FRP plate were considered to be dependent on temperature and moisture, which are given explicitly in terms of the fibre and matrix properties and the fibre–volume ratio. The conclusions from this research can be outlined as follows.

- There are stress concentrations at the end of the FRP plate. The normal stress concentration is a tensile stress but quickly towards to zero through a small oscillation. The initial delamination of the FRP plate from steel beam results from joints effects of the shear and normal stress at the end of the FRP plate.
- The interfacial shear stresses are reduced with decreases in fiber volume fraction. However, almost no effect is observed on the variation of interfacial normal stresses.
- No effect on the variation of adhesive stresses is observed when we use hygrothermal aged FRP plate to strengthen the steel beam.
- The interfacial stresses are influenced by the geometry parameters such as thickness of the adhesive layer and FRP plate in range of the different degrees. It is shown that the edge stresses and levels increase obviously with the increase of the thickness of the FRP plate. However, it is seen that increasing the thickness of the adhesive layer leads to significant reduction in the peak interfacial stresses.
- Another outcome based on the parametric study indicates that extending the FRP strip as close as possible to the support reduces the stresses at the edge.

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