



# An unsteady thermal investigation of a wetted longitudinal porous fin of different profiles

G. Sowmya<sup>1</sup> · B. J. Gireesha<sup>1</sup> · Hamza Berrehal<sup>2</sup>

Received: 18 April 2020 / Accepted: 15 June 2020 / Published online: 2 July 2020  
© Akadémiai Kiadó, Budapest, Hungary 2020

## Abstract

The unsteady thermal behavior of a porous longitudinal fin in a fully wet circumstance in the existence of convection as well as radiation effect is modeled in the present analysis. The thickness of the fin is assumed to vary with the length of the fin. Therefore, different profiles of the fin such as rectangular, convex, as well as triangular-shaped fin have been considered. Darcy's model is imposed to study porous nature. The derived nonlinear partial differential equation is non-dimensionalized and solved numerically with the help of Maple software by the Finite Difference Method. The transient thermal response and fin efficiency for diverse values of the significant parameters have been discussed graphically. A comparative analysis of three different shaped fins has been performed. It is found that the temperature drop rate is faster in triangular fin and efficiency is higher in rectangular fin.

**Keywords** Porous fin · Different shape · Unsteady · Thermal analysis · Natural convection · Radiation

## List of symbols

$A_c$	Fin cross-sectional area
$K$	Permeability ( $m^2$ )
$L$	Length of the fin (m)
$Nc$	Convective parameter
$Nr$	Radiative parameter
$T$	Local fin temperature (K)
$T_a$	Ambient temperature (K)
$T_b$	Base temperature (K)
$b_2$	Variable parameter ( $K^{-1}$ )
$c_p$	Specific heat at constant pressure ( $J\ kg^{-1}\ K^{-1}$ )
$g$	Acceleration due to gravity ( $ms^{-2}$ )
$h$	Heat transfer coefficient ( $Wm^{-2}\ K^{-1}$ )
$h_a$	Heat transfer coefficient at temperature $T_a$ ( $Wm^{-2}\ K^{-1}$ )
$h_D$	Uniform mass transfer coefficient

$k_{eff}$	Effective thermal conductivity of the material ( $Wm^{-1}\ K^{-1}$ )
$m_0, m_1$	Constants and $m_2 = m_0 + m_1$
$m_2$	Wet porous parameter
$p$	Power index of heat transfer coefficient
$q$	Base heat transfer rate (W)
$t$	Fin thickness (m)
$t_b$	Fin base thickness (m)
$t^*$	Time (s)
$w$	Width (m)

## Greek symbols

$\alpha$	Thermal diffusivity ( $m^2\ s^{-1}$ )
$\rho$	Density ( $kg\ m^{-3}$ )
$\theta_a$	Dimensionless ambient temperature
$\tau$	Dimensionless time
$\omega$	Humidity ratio of the saturated air
$\omega_a$	Humidity ratio of the surrounding air
$\nu$	Kinematic viscosity ( $m^2\ s^{-1}$ )
$i_{fg}$	Latent heat of water evaporation ( $J\ kg^{-1}$ )
$\theta$	Non-dimensional temperature
$\phi$	Porosity
$\sigma$	Stefan–Boltzmann constant ( $Wm^{-2}\ K^{-4}$ )
$\varepsilon$	Surface emissivity of fin
$\beta$	Volumetric thermal expansion coefficient ( $K^{-1}$ )

✉ G. Sowmya  
g.sowmya34@gmail.com

B. J. Gireesha  
bjgireesu@kuvempu.ac.in

Hamza Berrehal  
hamza.berrehal@gmail.com

<sup>1</sup> Department of PG Studies and Research in Mathematics, Kuvempu University, Shankaraghatta, India

<sup>2</sup> Department of Physics, University of Brothers Mentouri Constantine 1, 25000 Constantine, Algeria

## Subscripts

a	Ambient
b	Base of fin

f Fluid  
s Solid

## Introduction

Many applications in the field of electrical, industrial, and mechanical engineering like heat sinks, gas turbines, super heaters, bike heads, vehicle radiators, automobile, and aircraft engines, which deal with the heat transfer phenomenon, will utilize the fin structure to provide additional surface area and thereby enhance the rate of heat transfer. The performance of permeable fin was modeled by Kiwan and Nimr [1], and they explored the significance of porous fin over the solid fin. In addition, the influence of convection on the permeable fin was discussed under different fin tip conditions by Kiwan [2]. The radiation effect along with natural convection on the permeable fin has been examined by Gorla and Bakier [3]. A similar analysis has been carried out via the Homotopy Analysis Method by Darvishi et al. [4]. Additionally, the nonlinear boundary condition and temperature-dependent thermal conductivity were explained by using the Differential Transform Method by Torabi and Yaghoobi [5]. Ghosemi et al. [6] have analytically discussed the significance of solid and permeable convective fins. Convection and radiation have a significant impact on the heat transmit rate. Hoseinzadeh et al. [7] have considered the numerical and analytical methods and did the comparative study for the temperature field of a porous fin.

Mainly, there are three types of the fin, namely longitudinal, radial, and pin fin. The longitudinal fin has different profiles depending upon the shape due to variation in the thickness. Some of the studies based on the different profiles of fins are as follows: the concave parabolic, trapezoidal, as well as rectangular profiled straight fin under fixed heat flux and fixed temperature boundary condition was presented by Aziz and Fang [8]. Moradi and Ahmadikia [9] considered the longitudinal fin of exponential, convex, and rectangular shapes under varying thermal conductivity conditions by using the Differential Transform Method. Torabi and Zhang [10] also discussed the efficiency of various fin profiles under nonlinear internal heat generation and surface emissivity. Torabi et al. [11] have comprehensively analyzed the performance of trapezoidal, rectangular, and concave parabolic profiled longitudinal fin.

The fully wet condition creates wet nature around the fin and aids in the heat transfer rate. In this regard, the mass and heat transfer study of a circular permeable fin in a fully wet circumstance has been performed by Hatami and Ganji [12]. The comparative study of the efficiency of straight versus exponential permeable fin in a completely wet situation was carried out by Turkyilmazoglu [13]. The thermal performance and efficiency of moving trapezoidal fin were modeled by

Turkyilmazoglu [14]. The influence of the occurrence of heat generation internally on the permeable fin has been studied via the Finite Element Method by Sowmya et al. [15]. Baslem et al. [16] scrutinized the behavior of wetted longitudinal fin in the presence of nanofluid. The functionally graded material has been considered by Sowmya et al. [17], and they analyzed the thermal nature of the fin under fully wetted condition and described the significance of the wet nature of a fin.

For some of the applications like high-speed solar energy system, heat exchangers, and aircraft, transient response is also an important aspect to be recognized. Therefore, the unsteady thermal analysis is a significant factor. The longitudinal convective fin of functionally graded material transient thermal performance was modeled by Khan and Aziz [18]. Darcy's model was utilized to analyze the thermal behavior in the permeable fin of rectangular shape in the existence of transient effect by Darvishi et al. [19]. The radial fin of diverse profiles was considered, and unsteady temperature response was examined by Mosayebidorcheh et al. [20]. Applying the Differential Transform technique, the heat transfer of longitudinal fin under unsteady conditions along with nonlinear thermal coefficients has been discussed by Pasha et al. [21].

Fins have a wide range of applications in industries that are in the study of heat transfer analysis channels, etc. Xu et al. [22] and Xu [23] carried out a study of the heat transfer of thermal non-equilibrium in a metal-foam-filled solar collector and porous medium included micro heat exchangers, respectively. Hatami [24] studied the nanofluid flow through the wavy walled heated cylinder and discussed the RSM optimization. The heat transfer phenomenon in the microvessels and biomicrofluidics channel can be seen in Prakash et al. [25, 26]. MHD effect on the flexible fin in a cavity filled with nanofluid was analyzed by Selimefendigil et al. [27]. Asadollahi et al. [28] have scrutinized the channel with oblique fin coated with evacuating liquid. Some of the recent works can be observed in [29–32].

Based on all these observations, we intended to scrutinize the unsteady heat transfer through the rectangular, convex, and triangular-shaped longitudinal porous fin in a completely wet circumstance. The natural convection and radiations are also taken into consideration. The governed dimensionless partial differential equation has been solved numerically. The physical significance of the pertinent parameters has been illustrated through graphs and interpreted physically. The fin efficiency of all three profiles under diverse parameters is also plotted and discussed.

## Mathematical formulation

In the present study, we considered the one-dimensional heat transfer in the longitudinal porous fin of length  $L$  of diverse profiles like rectangular, triangular, as well as convex, and

these are as displayed in Fig. 1. The fin sectional area varies with  $t(x)$  function, where  $t(x)$  is nothing but the variation of the thickness of the fin with the length. The surface of the fin is attached to the primary surface maintained at temperature  $T_b$  and surrounded by a fully wetted condition with temperature  $T_a$ . Initially, the fin is in thermal equilibrium with the surroundings. At time  $t^* = 0$ , the fin base is subjected to an increase of temperature from  $T_a$  to  $T_b$ . Fin loses heat to the surrounding due to radiation and natural convection effect. Here, Darcy’s model is imposed to analyze the infiltration of ambient fluid through the pores of the fin. It is assumed that permeable fin is isotropic, homogeneous, and saturated with single-phase fluid. In addition, the radiant exchange through the surface is negligible.

The unsteady-state energy balance equation of a longitudinal permeable fin at a small element  $dx$  is defined as

$$\begin{aligned}
 & q(x) - q(x + dx) - 2\dot{m}c_p(T - T_a) - 2(1 - \phi)wh(T - T_a)dx \\
 & - 2\varepsilon\sigma wdx(T^4 - T_a^4) - 2h_D wdx i_{fg}(1 - \phi)(\omega - \omega_a) \\
 & = \rho c_p A_c \frac{\partial T}{\partial t^*}
 \end{aligned}
 \tag{1}$$

The first two terms in the LHS represent the heat transfer rate due to conduction, the second term stands for energy loss due to mass flow of immersed fluid, the third term is due to convection, the fourth term is the energy loss due to radiation, the fifth term signifies the heat lost because of the wet state, and the RHS is the unsteady term.

Here,  $\omega - \omega_a$  is the difference between the humidity ratio of saturated fluid and surrounding fluid, and it can be expressed as  $(\omega - \omega_a) = b_2(T - T_a)$  where  $b_2$  is the variable parameter.  $h_D$  is the uniform mass transfer coefficient,  $\phi$  is the porosity, and  $\dot{m}$  is the mass flow rate of fluid through porous fin and is defined as [4]

$$\dot{m} = \rho v(x) w dx,
 \tag{2}$$

where  $v(x)$  is the passage velocity, and according to Darcy’s model, it is as shown below [12]:

$$v(x) = \frac{gK\beta_f(T - T_a)}{v_f}.
 \tag{3}$$

Here, Fourier’s law of conduction takes the form as [11]

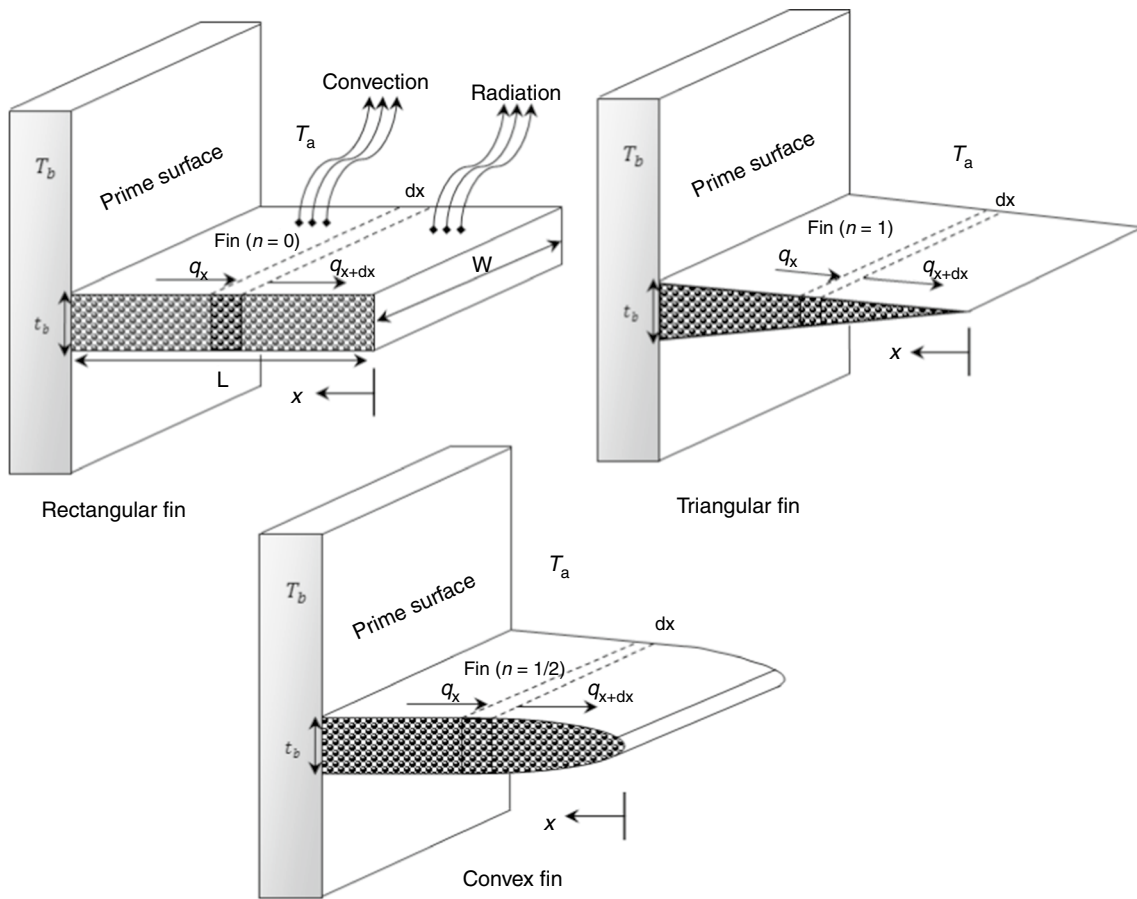


Fig. 1 Physical model of different profiles of the longitudinal porous fin

$$q = -k_{\text{eff}}wt(x)\frac{\partial T}{\partial x}, \tag{4}$$

where  $k_{\text{eff}}$  is the effective thermal conductivity and is represented as [12]

$$k_{\text{eff}} = k_f\phi + k_s(1 - \phi). \tag{5}$$

The local semi fin thickness  $t(x)$  is different for different cases. For case 1:  $t_b$  for rectangular profile, case 2:  $t_b\left(\frac{x}{L}\right)^{1/2}$  for convex profile, case 3:  $t_b\left(\frac{x}{L}\right)$  for triangular profile.

The convective heat transfer coefficient ( $h$ ) is taken as a power function of temperature and is given by

$$h = h_a\left[\frac{T - T_a}{T_b - T_a}\right]^p = h_Dc_pLe^{\frac{2}{3}}. \tag{6}$$

Equation (1) is simplified on applying Eqs. (2)–(6) in it and get reduced as follows

Case 1: For rectangular profile

$$\begin{aligned} \frac{\partial^2 T}{\partial x^2} - \frac{2\rho g\beta_f Kc_p}{v_f k_{\text{eff}}t_b}(T - T_a)^2 \\ - \frac{2\sigma\varepsilon}{k_{\text{eff}}t_b}(T^4 - T_a^4) - \frac{2(1 - \phi)h_a(T - T_a)^{p+1}}{k_{\text{eff}}t_b(T_b - T_a)^p} \\ - \frac{2(1 - \phi)h_a i_{fg} b_2(T - T_a)^{p+1}}{c_p Le^{2/3} k_{\text{eff}}t_b(T_b - T_a)^p} = \frac{\rho c_p}{k_{\text{eff}}} \frac{\partial T}{\partial t^*}, \end{aligned} \tag{7a}$$

Case 2: For convex profile

$$\begin{aligned} \frac{\partial}{\partial x}\left(\left(\frac{x}{L}\right)^{1/2}\frac{\partial T}{\partial x}\right) - \frac{2\rho g\beta_f Kc_p}{v_f k_{\text{eff}}t_b}(T - T_a)^2 \\ - \frac{2\sigma\varepsilon}{k_{\text{eff}}t_b}(T^4 - T_a^4) - \frac{2(1 - \phi)h_a(T - T_a)^{p+1}}{k_{\text{eff}}t_b(T_b - T_a)^p} \\ - \frac{2(1 - \phi)h_a i_{fg} b_2(T - T_a)^{p+1}}{c_p Le^{2/3} k_{\text{eff}}t_b(T_b - T_a)^p} = \frac{\rho c_p}{k_{\text{eff}}} \frac{\partial T}{\partial t^*} \end{aligned} \tag{7b}$$

Case 3: For triangular profile

$$\begin{aligned} \frac{\partial}{\partial x}\left(\frac{x}{L}\frac{\partial T}{\partial x}\right) - \frac{2\rho g\beta_f Kc_p}{v_f k_{\text{eff}}t_b}(T - T_a)^2 \\ - \frac{2\sigma\varepsilon}{k_{\text{eff}}t_b}(T^4 - T_a^4) - \frac{2(1 - \phi)h_a(T - T_a)^{p+1}}{k_{\text{eff}}t_b(T_b - T_a)^p} \\ - \frac{2(1 - \phi)h_a i_{fg} b_2(T - T_a)^{p+1}}{c_p Le^{2/3} k_{\text{eff}}t_b(T_b - T_a)^p} = \frac{\rho c_p}{k_{\text{eff}}} \frac{\partial T}{\partial t^*}. \end{aligned} \tag{7c}$$

The following dimensionless terms are introduced to non-dimensionalize the above equations [18].

$$\theta = \frac{T}{T_b}, \theta_a = \frac{T_a}{T_b}, X = \frac{x}{L}, \tau = \frac{\alpha t^*}{L^2}. \tag{8}$$

On applying Eq. (8) in Eq. (7), it is modified to the dimensionless form as follows:

Case 1: For rectangular profile

$$\begin{aligned} \frac{\partial \theta}{\partial \tau} = \frac{\partial^2 \theta}{\partial X^2} - Nc(\theta - \theta_a)^2 \\ - Nr(\theta^4 - \theta_a^4) - \frac{m_2(\theta - \theta_a)^{p+1}}{(1 - \theta_a)^p}. \end{aligned} \tag{9a}$$

Case 2: For convex profile

$$\begin{aligned} \frac{\partial \theta}{\partial \tau} = \frac{\partial}{\partial X}\left(X^{1/2}\frac{\partial \theta}{\partial X}\right) - Nc(\theta - \theta_a)^2 \\ - Nr(\theta^4 - \theta_a^4) - \frac{m_2(\theta - \theta_a)^{p+1}}{(1 - \theta_a)^p}. \end{aligned} \tag{9b}$$

Case 3: For triangular profile

$$\begin{aligned} \frac{\partial \theta}{\partial \tau} = \frac{\partial}{\partial X}\left(X\frac{\partial \theta}{\partial X}\right) - Nc(\theta - \theta_a)^2 \\ - Nr(\theta^4 - \theta_a^4) - \frac{m_2(\theta - \theta_a)^{p+1}}{(1 - \theta_a)^p}. \end{aligned} \tag{9c}$$

Corresponding initial and boundary conditions will be

$$\theta(X, 0) = 0, \theta(1, \tau) = 1, \frac{\partial \theta}{\partial X}(0, \tau) = 0. \tag{10}$$

In Eq. (9),  $\theta$  is the dimensionless temperature,  $X$  is the dimensionless length,  $\tau$  is the dimensionless time,

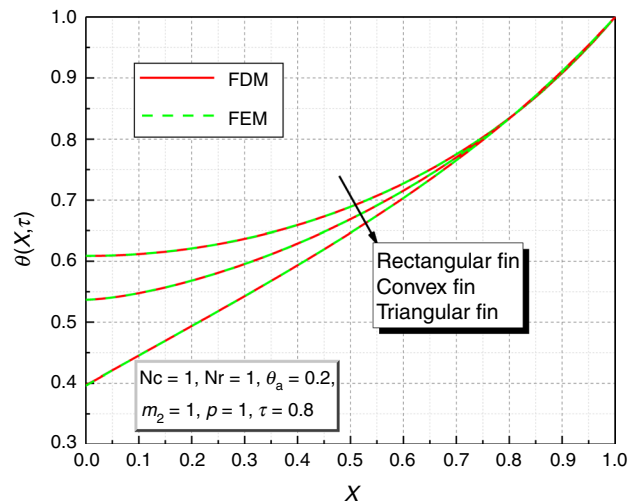
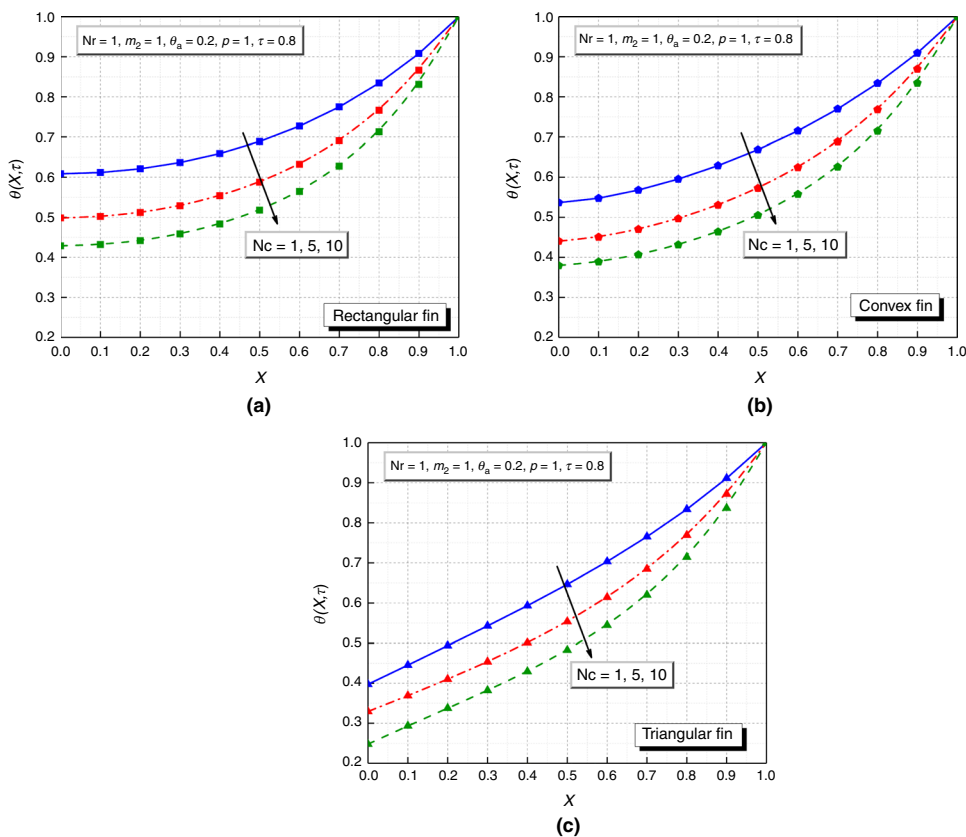


Fig. 2 Comparison of present solution via FDM with FEM solution

**Fig. 3** Impact of convective parameter on the unsteady thermal profile for **a** rectangular fin, **b** convex fin, and **c** triangular fin



$Nc = \frac{2\rho g\beta_f Kc_p T_b L^2}{v_f k_{eff} t_b}$  is the convective parameter,  $Nr = \frac{2\varepsilon\sigma L^2 T_b^3}{k_{eff} t_b}$  is the radiative parameter,  $p$  is the power index associated with the convective heat transfer coefficient,  $\theta_a$  is the dimensionless ambient temperature,  $X$  is the dimensionless length,  $m_2$  is the wet porous parameter and is the sum of  $m_0 = \frac{2h_a L^2 (1-\phi)}{k_{eff} t_b}$  and  $m_1 = \frac{2h_a i_{fg} (1-\phi) b_2 L^2}{k_{eff} t_b c_p L e^{2/3}}$ .

**Table 1** Influence of dimensionless time on the unsteady thermal profile of the different profiled fin at  $Nc = 1, Nr = 1, p = 1, m_2 = 1, \theta_a = 0.2$

$\theta(0, \tau)$	Rectangular fin	Convex fin	Triangular fin
$\tau = 1.0$	0.62058515	0.56074330	0.45296611
$\tau = 1.5$	0.62584218	0.57435767	0.49617832
$\tau = 2.0$	0.62613929	0.57555652	0.50256936
$\tau = 2.5$	0.62615603	0.57566117	0.50349330
$\tau = 3.0$	0.62615697	0.57567030	0.50362650
$\tau = 3.5$	0.62615702	0.57567110	0.50364570
$\tau = 4.0$	0.62615702	0.57567117	0.50364847
$\tau = 4.5$	0.62615702	0.57567118	0.50364887
$\tau = 5.0$	0.62615702	0.57567118	0.50364892

**Fin efficiency**

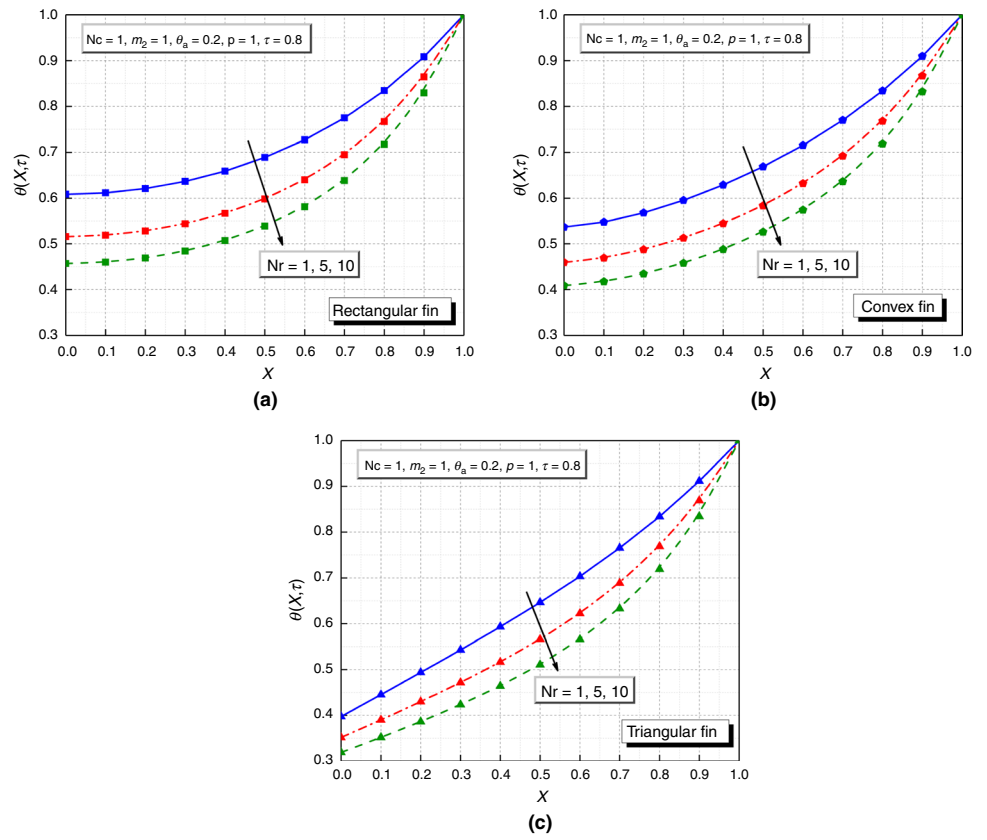
The total transfer of heat from the fin surface is the sum of heat transfer due to porosity, convection, as well as radiation, and it is expressed as follows:

$$Q_f = \int_0^L \left[ \frac{2\rho g w \beta_f K c_p}{v_f} (T - T_a)^2 + 2\sigma \varepsilon w (T^4 - T_a^4) + \left( 2h_a (1 - \phi) w + \frac{2h_a i_{fg} (1 - \phi) b_2 w}{c_p L e^{2/3}} \right) \frac{(T - T_a)^{p+1}}{(T_b - T_a)^p} \right] dx \tag{11}$$

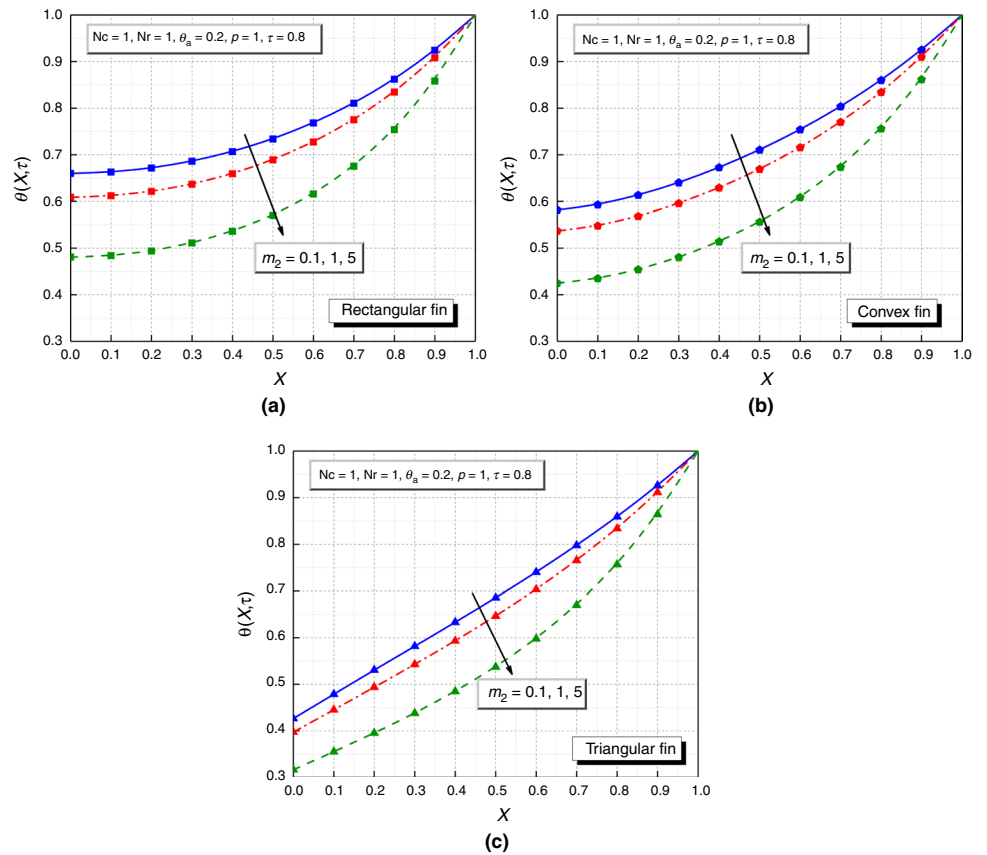
The ideal transfer of heat from the fin means the heat transfer from the fin surface when it is kept at a base temperature  $T_b$ . The ideal case is as follows:

$$Q_{ideal} = \frac{2\rho g w \beta_f K c_p L}{v_f} (T_b - T_a)^2 + 2\sigma \varepsilon w L (T_b^4 - T_a^4) + \left( 2h_a (1 - \phi) w L + \frac{2h_a i_{fg} (1 - \phi) b_2 w L}{c_p L e^{2/3}} \right) (T_b - T_a) \tag{12}$$

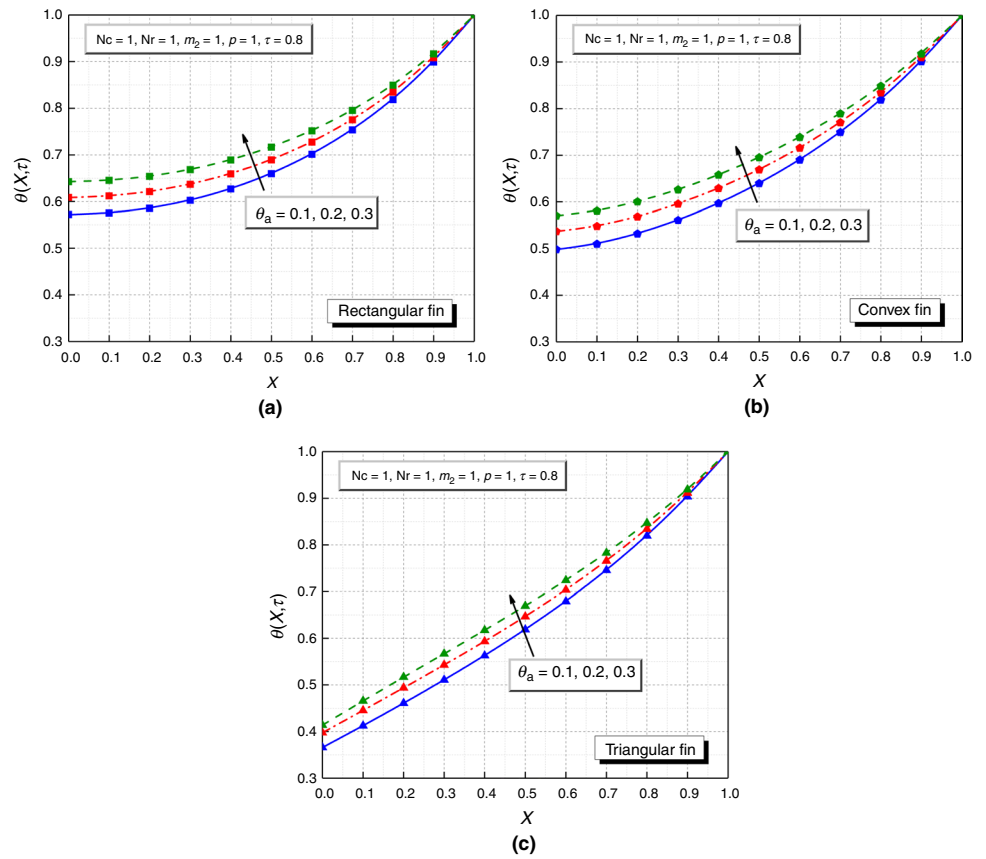
**Fig. 4** Impact of radiative parameter on the unsteady thermal profile for **a** rectangular fin, **b** convex fin, and **c** triangular fin



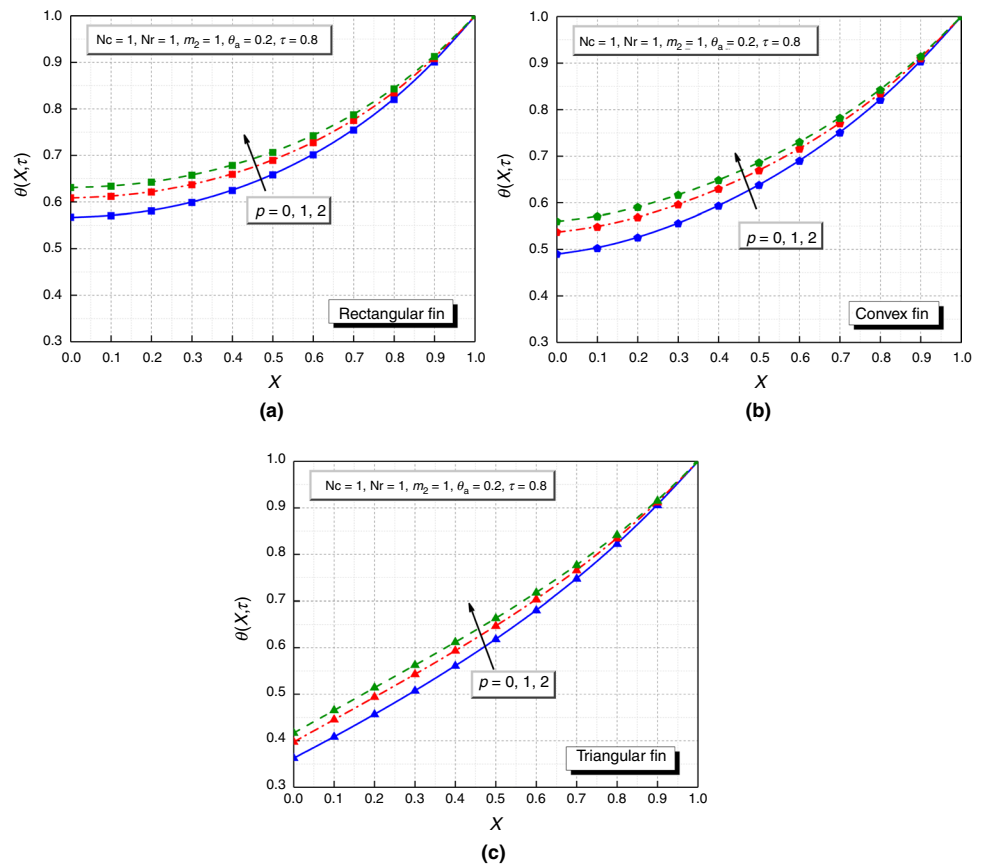
**Fig. 5** Impact of a wet porous parameter on the unsteady thermal profile for **a** rectangular fin, **b** convex fin, and **c** triangular fin



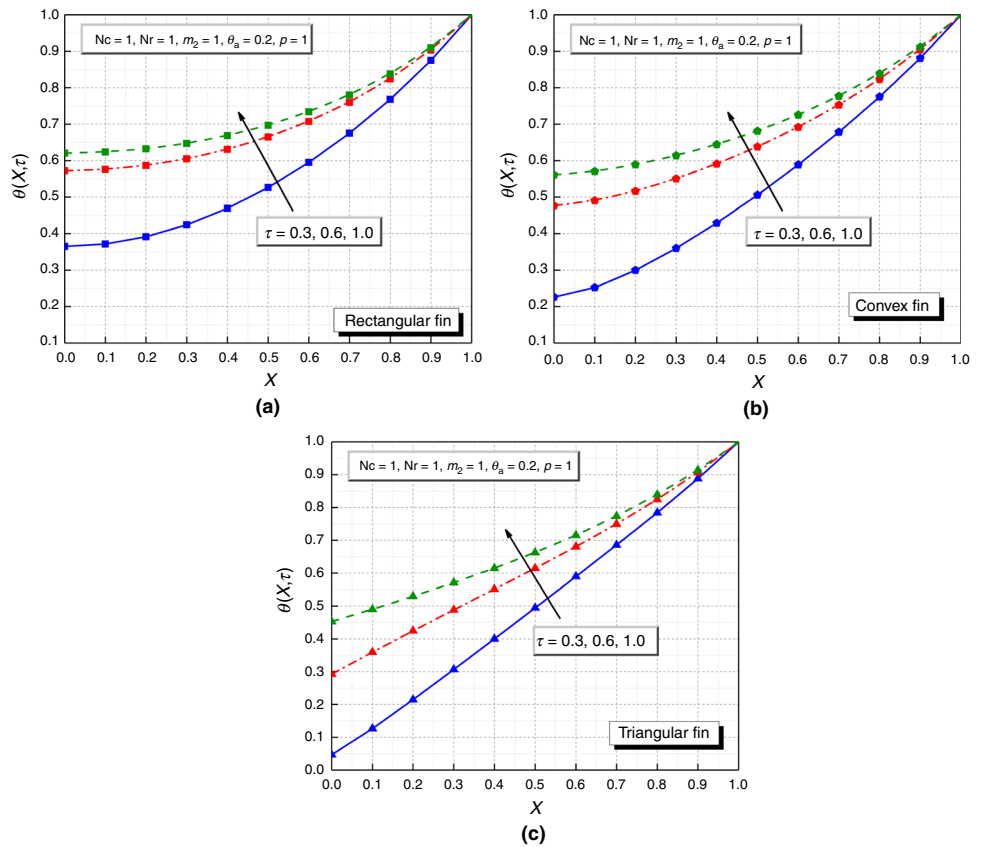
**Fig. 6** Impact of ambient temperature on the unsteady thermal profile for **a** rectangular fin, **b** convex fin, and **c** triangular fin



**Fig. 7** Impact of power index on the unsteady thermal profile for **a** rectangular fin, **b** convex fin, and **c** triangular fin



**Fig. 8** Impact of dimensionless time on the unsteady thermal profile for **a** rectangular fin, **b** convex fin, and **c** triangular fin



The fin efficiency is described as the ratio of the actual dissipation of heat from the fin to the heat that would be transferred if the entire fin surface is kept at the base temperature. That is

$$\eta = \frac{Q_f}{Q_{ideal}} \tag{13}$$

The dimensionless form of Eq. (13) is

$$\eta = \frac{\int_0^1 \left[ Nc(\theta - \theta_a)^2 + Nr(\theta^4 - \theta_a^4) + \frac{m_2(\theta - \theta_a)^{p+1}}{(1 - \theta_a)^p} \right] dX}{Nc(1 - \theta_a)^2 + Nr(1 - \theta_a^4) + m_2(1 - \theta_a)} \tag{14}$$

**Solution procedure**

The parabolic partial differential Eq. (9) with initial and boundary conditions (10) was solved numerically via the pdsolve command in MAPLE 18. This is based on the algorithm of the Finite Difference Method (FDM) with

the discretization of the differential equation by the center implicit scheme. This command extracts numerical data from the solution obtained in the form of a module. The space step and time step are chosen as  $\Delta x = \Delta t = 0.001$ .

The validation of the current solution is done by comparing the obtained solution with the finite element method and is displayed in Fig. 2. The solution found in FDM via Maple is in good agreement and accurate with the FEM solution for rectangular, convex, and triangular profiled fins considered in this study.

**Results and discussion**

The parametric study of the obtained solution has been conducted to determine the prominence of dimensionless significant constraints on the unsteady thermal distribution as well as efficiency for the longitudinal fin of triangular, convex, and rectangular profiles. The results have been shown in Figs. 3–11 and in Table 1.

Figure 3a–c displays the combined effect of variation in convective parameter and different fin profiles on the thermal field along the surface of the porous fin. Here, an increase



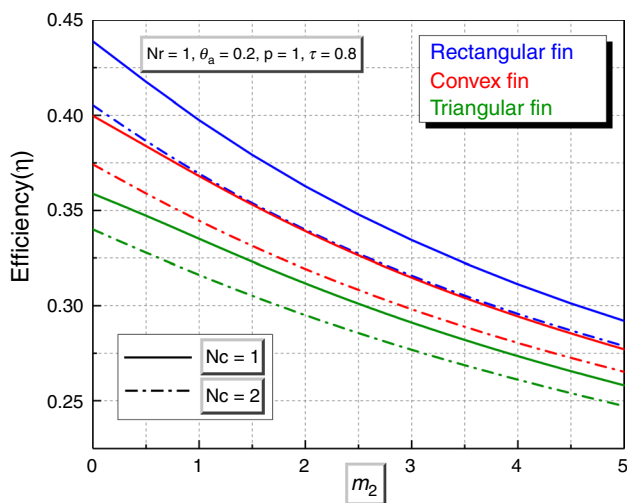


Fig. 9 Fin efficiency versus  $m_2$  for diverse  $N_c$

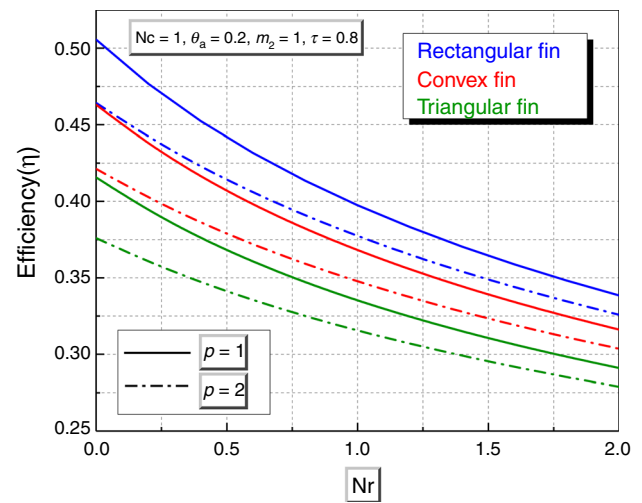


Fig. 10 Fin efficiency versus  $N_r$  for diverse  $p$

in convective parameter decreases the temperature profile. This is because, as the convection effect increases, more heat from the fin surface transform to the surrounding. Therefore, the rate of transfer of heat increases and the temperature decreases. The same nature is noticeable in rectangular, convex, as well as triangular profiled fin structure. On the other hand, the thermal drop rate is more for triangular fin followed by convex and rectangular fin.

The significance of the radiative parameter on the temperature behavior of the permeable fin of different profiles is as shown in Fig. 4a–c. The radiation effect positively impacts on the heat transfer purpose. Therefore, the temperature of the fin surface declines with amplification in the radiative parameter because the radiation effect carries away the excess heat from the fin surface. This nature is similar for all the three shaped profiles considered in this study.

The wet nature around the fin will aid in absorbing extra heat from the surface of the fin and therefore decreases the fin temperature. This is evident in Fig. 5a–c, where the effect of the wet porous parameter on the different profiles as well as the thermal feature of the fin is plotted. It is seen that, as the wet porous parameter enhances, the temperature profile decreases. The wet porous parameter shows a significant effect on the temperature field of all the three shapes of fin considered in the present analysis. In addition, the temperature decrease rate is faster in the triangular fin compared to the convex and rectangular fins.

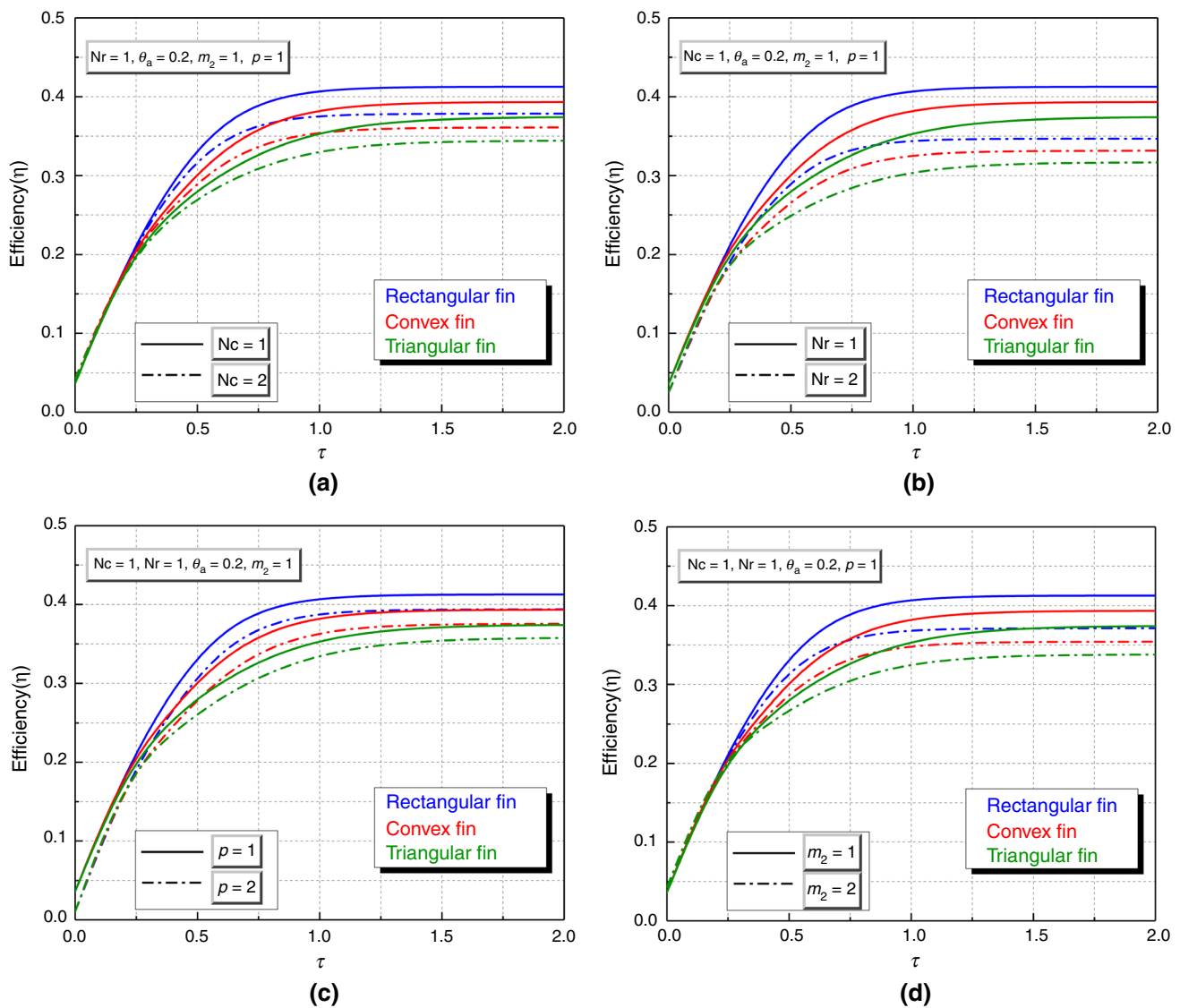
The impact of dimensionless ambient temperature parameter on the temperature attribute of the rectangular, convex, as well as triangular fin is represented in Fig. 6a–c. Here, as  $\theta_a$  rises, the thermal profile is more. This is due to the rise in the ratio of the surrounding temperature to the base

temperature of the fin. This implies that when the surrounding temperature amplifies, the thermal variation between the surface of the fin and the surrounding decreases. According to Newton’s law of cooling, the cooling effect decreases due to a reduction in thermal difference between surface and surrounding.

The effect of power index  $p$  incorporated in a convective heat transfer coefficient on the temperature field of the rectangular, convex, and triangular profiled fin has been captured in Fig. 7a–c. As the power index rises, the temperature profile increases, due to the nonlinearity. Therefore, we notice less heat transfer rate with the rise in the power index in all the three different fin structures.

The transient effect on the temperature feature of the porous fin of the different profile is displayed in Fig. 8a–c. One can see that as time increases the temperature profile also rises to a certain limit. It is very interesting to note that, after a certain time, the thermal profile becomes constant. The transient effect is visible initially and after some time a steady condition is reached. This is shown in Table 1. Clearly, we saw that the unsteady thermal behavior value increases with time but as time increases it reaches a steady state. On the other hand, the unsteady thermal value of the triangular fin is less than the convex and rectangular profiled fin.

The efficiency of rectangular, convex, and triangular porous fin versus a convective parameter and radiative parameter for diverse values of ambient temperature and power index is as shown in Figs. 9 and 10, respectively. It is observed that as the convective and radiative parameters rise, efficiency shows decreasing nature. This is because an increase in convection effect reduces the temperature of the



**Fig. 11** Efficiency of fin versus dimensionless time for diverse values of **a**  $N_c$ , **b**  $N_r$ , **c**  $p$ , and **d**  $m_2$

fin surface, which leads to the decrease in the ratio of actual heat transport to transfer of heat from the fin surface kept at base temperature. This result can be qualitatively compared with that of Torabi et al. [11].

The transient efficiency of fin of rectangular, convex and triangular profile for different values of  $N_c$ ,  $N_r$ ,  $p$ ,  $m_2$  is presented in Fig. 11a–d, respectively. Initially, when there is no heat flow from the base, the fin efficiency is zero and the surface of the fin is in thermal equilibrium with the surrounding but as the heat flow through the fin, the efficiency rises sharply and attains constant value as the steady state is reached. In the transient region, it is noticed that, as the dimensionless time rises, efficiency also enhances. Khan and Aziz [18] attained the same conclusion for unsteady

convective fin case. It is observed that steady state is reached earlier for higher  $N_c$ ,  $N_r$ ,  $p$ , and  $m_2$  and as these parameters rise, efficiency decreases. Additionally, it is observed that the efficiency of the rectangular fin is higher and the triangular fin is lower than the convex fin.

## Conclusions

The transient thermal performance and efficiency of the longitudinal porous fin of different profiles like rectangular, triangular, and convex shapes under a fully wet circumstance has been scrutinized numerically. The main findings of the present investigation are listed below:

- The radiative parameter, as well as a convective parameter, induces a pronounced effect on the fin cooling.
- The higher value of the wet porous parameter is also attributed to the augmentation of the heat dissipation from the fin.
- The unsteady thermal profile rises for a rise in power index and ambient temperature, which means that it reduces the heat transfer rate.
- The different profile of the fin surface has a prominent impact on the rate of heat transfer.
- The efficiency of the fin is affected by the convective and radiative parameters.
- The efficiency rises with the rise in dimensionless time and becomes constant as a steady state is reached.
- The temperature drop rate is more for a triangular fin profile compared to rectangular and convex fin profiles.
- More fin efficiency is observed for a rectangular fin profile followed by convex and triangular fins.

## References

1. Kiwan S, Al-Nimr MA. Using porous fins for heat transfer enhancement. *J Heat Transf.* 2001;123(4):790–5.
2. Kiwan S. Thermal analysis of natural convection porous fins. *Transp Porous Med.* 2007;67(1):17.
3. Gorla RS, Bakier AY. Thermal analysis of natural convection and radiation in porous fins. *Int Commun Heat Mass Transf.* 2011;38(5):638–45.
4. Darvishi MT, Gorla RS, Khani F. Natural convection and radiation in porous fins. *Int J Numer Methods Heat Fluid Flow.* 2013;23(8):1406–20.
5. Torabi M, Yaghoobi H. Accurate solution for convective–radiative fin with variable thermal conductivity and nonlinear boundary condition by DTM. *Arab J Sci Eng.* 2013;38(12):3575–85.
6. Ghasemi SE, Valipour P, Hatami M, Ganji DD. Heat transfer study on solid and porous convective fins with temperature-dependent heat generation using efficient analytical method. *J Central South Univ.* 2014;21(12):4592–8.
7. Hoseinzadeh S, Heyns PS, Chamkha AJ, Shirkhani A. Thermal analysis of porous fins enclosure with the comparison of analytical and numerical methods. *J Therm Anal Calorim.* 2019;138(1):727–35.
8. Aziz A, Fang T. Alternative solutions for longitudinal fins of rectangular, trapezoidal, and concave parabolic profiles. *Energy Convers Manag.* 2010;51(11):2188–94.
9. Moradi A, Ahmadikia H. Analytical solution for different profiles of fin with temperature-dependent thermal conductivity. *Math Prob Eng.* 2010;2010.
10. Torabi M, Bao Zhang Q. Analytical solution for evaluating the thermal performance and efficiency of convective–radiative straight fins with various profiles and considering all nonlinearities. *Energy Convers Manag.* 2013;66:199–210.
11. Torabi M, Aziz A, Zhang K. A comparative study of longitudinal fins of rectangular, trapezoidal and concave parabolic profiles with multiple nonlinearities. *Energy.* 2013;51:243–56.
12. Hatami M, Ganji DD. Investigation of refrigeration efficiency for fully wet circular porous fins with variable sections by combined heat and mass transfer analysis. *Int J Refrig.* 2014;40:140–51.
13. Turkyilmazoglu M. Efficiency of heat and mass transfer in fully wet porous fins: exponential fins versus straight fins. *Int J Refrig.* 2014;46:158–64.
14. Turkyilmazoglu M. Efficiency of the longitudinal fins of trapezoidal profile in motion. *J Heat Transf.* 2017;139(9):094501.
15. Sowmya G, Gireesha BJ, Madhu M. Analysis of a fully wetted moving fin with temperature-dependent internal heat generation using the finite element method. *Heat Transf.* 2020;49(4):1939–54.
16. Baslem A, Sowmya G, Gireesha BJ, Prasannakumara BC, Rahimi-Gorji M, Hoang NM. Analysis of thermal behavior of a porous fin fully wetted with nanofluids: convection and radiation. *J Mol Liq.* 2020;307:112920.
17. Sowmya G, Gireesha BJ, Khan MI, Momani S, Hayat T. Thermal investigation of fully wet longitudinal porous fin of functionally graded material. *Int J Numer Methods Heat Fluid Flow.* 2020. <https://doi.org/10.1108/HFF-12-2019-0908>.
18. Khan WA, Aziz A. Transient heat transfer in a functionally graded convecting longitudinal fin. *Heat Mass Transf.* 2012;48(10):1745–53.
19. Darvishi MT, Gorla RS, Khani F. Unsteady thermal response of a porous fin under the influence of natural convection and radiation. *Heat Mass Transf.* 2014;50(9):1311–7.
20. Mosayebidorcheh S, Rahimi-Gorji M, Ganji DD, Moayebidorcheh T, Pourmehran O, Biglarian M. Transient thermal behavior of radial fins of rectangular, triangular and hyperbolic profiles with temperature-dependent properties using DTM-FDM. *J Central South Univ.* 2017;24(3):675–82.
21. Pasha AV, Jalili P, Ganji DD. Analysis of unsteady heat transfer of specific longitudinal fins with temperature-dependent thermal coefficients by DTM. *Alex Eng J.* 2018;57(4):3509–21.
22. Xu H, Gong L, Huang S, Xu M. Non-equilibrium heat transfer in metal-foam solar collector with no-slip boundary condition. *Int J Heat Mass Transf.* 2014;76:357–65.
23. Xu H. Performance evaluation of multi-layered porous-medium micro heat exchangers with effects of slip condition and thermal non-equilibrium. *Appl Therm Eng.* 2017;116:516–27.
24. Hatami M. Nanoparticles migration around the heated cylinder during the RSM optimization of a wavy-wall enclosure. *Adv Powder Technol.* 2017;28(3):890–9.
25. Prakash J, Ramesh K, Tripathi D, Kumar R. Numerical simulation of heat transfer in blood flow altered by electroosmosis through tapered micro-vessels. *Microvasc Res.* 2018;118:162–72.
26. Prakash J, Sharma A, Tripathi D. Thermal radiation effects on electroosmosis modulated peristaltic transport of ionic nanoliquids in biomicrofluidics channel. *J Mol Liq.* 2018;249:843–55.
27. Selimefendigil F, Oztop HF, Chamkha AJ. MHD mixed convection in a nanofluid filled vertical lid-driven cavity having a flexible fin attached to its upper wall. *J Therm Anal Calorim.* 2019;135(1):325–40.
28. Asadollahi A, Esfahani JA, Ellahi R. Evacuating liquid coatings from a diffusive oblique fin in micro-/mini-channels. *J Therm Anal Calorim.* 2019;138(1):255–63.
29. Noreen S, Kausar T, Tripathi D, Ain QU, Lu DC. Heat transfer analysis on creeping flow Carreau fluid driven by peristaltic pumping in an inclined asymmetric channel. *Therm Sci Eng Progress.* 2020;17:100486.
30. Xu HJ. Thermal transport in microchannels partially filled with micro-porous media involving flow inertia, flow/thermal slips, thermal non-equilibrium and thermal asymmetry. *Int Commun Heat Mass Transf.* 2020;110:104404.

31. Jing D, Hu S, Hatami M, Xiao Y, Jia J. Thermal analysis on a nanofluid-filled rectangular cavity with heated fins of different geometries under magnetic field effects. *J Therm Anal Calorim.* 2020;139(6):3577–88.
32. Nazari S, Ellahi R, Sarafraz MM, Safaei MR, Asgari A, Akbari OA. Numerical study on mixed convection of a non-Newtonian nanofluid with porous media in a two lid-driven square cavity. *J Therm Anal Calorim.* 2020;140:1121–45.

**Publisher's Note** Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.