



Effectiveness of various shapes of Al₂O₃ nanoparticles on the MHD convective heat transportation in porous medium

CVFEM modelling

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Abstract

The influences of Al₂O₃ nanoparticles with various shapes on thermal characteristics of nanofluid within a permeable space concerning magnetic force have been simulated by means of CVFEM. To form the final PDEs, radiation term has been incorporated. Impacts of magnetic force, radiation constraint, Rayleigh number and shape factor on nanomaterial behaviour have been analysed. Results demonstrate that the higher values of shape factor lead to augmented convective heat transfer. By augmenting the magnetic strength, conductive heat transfer can be predominant than that of the convection.

Keywords MHD · Nanoparticle's shape · Darcy law · Radiation · Nanofluid · CVFEM

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Introduction

Because of simplicity, cost-effectiveness and low noise of free convection, this mode can affect thermal behaviour of a wide range of engineering equipment. The natural convection mechanism transpires under the impact of the magnetic field in multiple processes including metal casting, the liquid cooling blanket of the fusion reactors and crystal growth. However, the existence of Lorentz force imposes an adverse effect on the phenomenon and deteriorates the convective flow. In an investigation of free convection within a cavity under magnetic field influence, Rudraiah [1] observed that the stronger Hartmann destroys the convective heat transportation rate and this suppression is more characteristic in the regions of the low Grashof number. Kakarantzas et al. [2] investigated liquid metal MHD flow inside a container and reported that the implication of Hartmann flow deteriorates nanomaterial velocity. In a similar way, Selimefendigil and Oztop [3] and Sheikholeslami [4] demonstrated an inverse association between the Lorentz force and Nusselt number.

In the above-mentioned scenarios, the undesirable impact of Hartmann number on free convection can be compensated to some extent by replacing the conventional coolants with

the metallic nanofluids [5–9]. The term nanofluid is expressed for the colloidal solution of the traditional fluids and the nano-sized metallic or non-metallic particles exhibiting superior thermal behaviour than that of the hosting fluid and the suspended particles [10]. The magnetohydrodynamic (MHD) natural convection of the nanofluids for multiple operating conditions and nanofluid combinations has been presented by several studies. Sheremet et al. [11] analysed the MHD water–Cu nanofluid in a wavy enclosure under the influence of an isothermal corner heater. Kefayati [12, 13] implemented the FDLBM to scrutinize nanomaterial-free convection through a square cavity. Using the numerical approach, Sheikholeslami et al. [14–18] studied MHD and EHD convective transportation of water-based nanomaterials in a square enclosure, concentric annulus, semi-annulus and cubic cavity.

Fewer reports presented the MHD convective transportation within porous media. The permeable media possesses the advantages of the low density and large area for the optimal heat transfer. Selimefendigil and Oztop [19] scrutinized vented enclosure influenced by Lorentz forces to analyze mixed convection. Rashad et al. [20] explored the mutual impact of the internal heat generation and Lorentz forces on free convective flow of the copper–water within a rectangular tank. Finding new carrier fluid with greater thermal properties has been scrutinized by several researchers [21–33].

The objective of this investigation is to scrutinize study the impact of the nanoparticle shape variation on the MHD-free convection and the radiation of the Al₂O₃–water nanomaterial within a porous medium by implementing the CVFEM approach. The considered radiation parameter, shape factor, Hartmann number and particle volume fraction ranges are $0 \leq Rd \leq 0.8$, $3 \leq m \leq 5.7$, $0 \leq Ha \leq 20$ and $0\% \leq \phi \leq 5\%$ respectively.

Formulation of problem and simulation

In the present problem, the impact of Hartmann flow on the behaviour of nanomaterial in a permeable geometry has been emphasized. CVFEM has been employed concerning triangular element. The imposed boundary conditions have been also depicted in Fig. 1. In order to model the porous terms, the Darcy law has been implemented. Shape factor influences on nanomaterial properties have been modelled. The problem under consideration has below equations:

$$\frac{\partial v}{\partial y} + \frac{\partial u}{\partial x} = 0 \tag{1}$$

$$\frac{\mu_{nf}}{K} u = -\frac{\partial P}{\partial x} + \sigma_{nf} B_0^2 [(\sin \gamma)v(\cos \gamma) - u(\sin \gamma)^2] \tag{2}$$

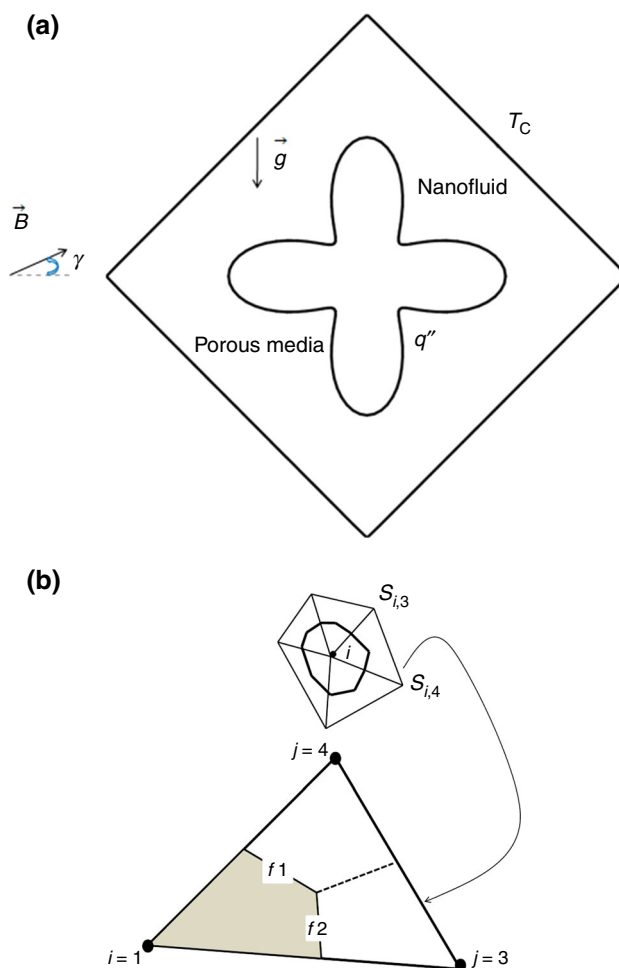


Fig. 1 a Current porous domain, b CVFEM element

$$\frac{\mu_{nf}}{K} v = -\frac{\partial P}{\partial y} + (T - T_c)g\rho_{nf}\beta_{nf} + \sigma_{nf} B_0^2 (\cos \gamma)[(\sin \gamma)u - (\cos \gamma)v] \tag{3}$$

$$\frac{1}{(\rho C_p)_{nf}} \frac{\partial q_r}{\partial y} + \left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = k_{nf} (\rho C_p)_{nf}^{-1} \left(\frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial x^2} \right), \tag{4}$$

$$\left[T^4 \cong 4T_c^3 T - 3T_c^4, q_r = -\frac{4\sigma_e}{3\beta_R} \frac{\partial T^4}{\partial y} \right]$$

The thermophysical properties of nanomaterial were estimated by the below equations

$$(\rho C_p)_{nf} / (\rho C_p)_f = PP, \quad PP = (1 - \phi) + \phi(\rho C_p)_s / (\rho C_p)_f \tag{5}$$

$$(\rho\beta)_{nf} = (\rho\beta)_f(1 - \phi) + \phi(\rho\beta)_s \tag{6}$$

$$\rho_{nf} - \phi\rho_s = (1 - \phi)\rho_f \tag{7}$$

$$1 + A = \sigma_s / \sigma_f, \frac{\sigma_{nf}}{\sigma_f} = B, \tag{8}$$

$$B - 1 = 3\phi \frac{A}{A + 3 - \phi A}$$

The effective viscosity μ_{nf} as involving Brownian motion:

$$\mu_{eff} = \left([1 - \phi]^{-2.5} + \frac{k_{Brownian}}{Pr_f k_f} \right) \mu_f$$

$$g'(d_p, \phi, T) = gg$$

$$k_{Brownian} = 5 \times 10^4 c_{p,f} gg \rho_f \phi \sqrt{\frac{\kappa_b T}{\rho_p d_p}}$$

$$gg = \text{Ln}(T) (a_1 + a_3 \text{Ln}(\phi) + a_4 \text{Ln}(d_p) \text{Ln}(\phi) + a_2 \text{Ln}(d_p) + a_5 \text{Ln}(d_p)^2) + (a_6 + \text{Ln}(d_p) a_{10} \text{Ln}(d_p) + a_8 \text{Ln}(\phi) + a_9 \text{Ln}(d_p) \text{Ln}(\phi) + a_7 \text{Ln}(d_p)) \tag{9}$$

The impact of the nanoparticle shape factor has been involved in the estimation of the effective thermal conductivity k_{nf} as:

$$k_p - k_f = kr$$

$$\frac{k_{nf}}{k_f} = \frac{kr m \phi + k_f + kr \phi + m k_f + k_p}{k_p + m k_f + kr \phi + k_f} \tag{10}$$

The various shape factors, related coefficient and properties are enlisted in Tables 1–3 [34].

Equation (11) has been considered to attain the dimensionless form.

Table 1 The coefficient values of Al₂O₃–water nanofluid

Coefficient values	Al ₂ O ₃ –water
<i>a</i> ₁	52.813488759
<i>a</i> ₂	6.115637295
<i>a</i> ₃	0.6955745084
<i>a</i> ₄	4.17455552786E–02
<i>a</i> ₅	0.176919300241
<i>a</i> ₆	– 298.19819084
<i>a</i> ₇	– 34.532716906
<i>a</i> ₈	– 3.9225289283
<i>a</i> ₉	– 0.2354329626
<i>a</i> ₁₀	– 0.999063481

Table 2 Thermophysical properties of water and nanoparticles

	$\rho/\text{kg m}^{-3}$	$C_p/\text{J kgK}^{-1}$	$k/\text{W m K}^{-1}$	$\beta \times 10^5/\text{K}^{-1}$	$\sigma/\Omega \text{m}^{-1}$
Water	997.1	4179	0.613	21	0.05
Al ₂ O ₃	3970	765	25	0.85×10^{-5}	1×10^{-10}

$$(y, x) = (Y, X)L, \Psi = \psi / \alpha_{nf}, \theta = \frac{T - T_c}{\Delta T}, \Delta T k_f L^{-1} = q'' \tag{11}$$

The final expression can be written as:

$$\frac{\partial^2 \Psi}{\partial X^2} + \frac{\partial^2 \Psi}{\partial Y^2}$$

$$= -\text{Ha} \frac{A_6}{A_5} \left[(\cos \gamma) \frac{\partial^2 \Psi}{\partial X^2} (\cos \gamma) + 2 \frac{\partial^2 \Psi}{\partial X \partial Y} (\sin \gamma) (\cos \gamma) + \frac{\partial^2 \Psi}{\partial Y^2} (\sin^2 \gamma) \right] - \frac{A_3 A_2}{A_4 A_5} \frac{\partial \theta}{\partial X} \text{Ra}$$

$$\left(\frac{\partial^2 \theta}{\partial X^2} \right) + \left(1 + \frac{4}{3} \left(\frac{k_{nf}}{k_f} \right)^{-1} \text{Rd} \right) \frac{\partial^2 \theta}{\partial Y^2} = - \frac{\partial \Psi}{\partial X} \frac{\partial \theta}{\partial Y} + \frac{\partial \theta}{\partial X} \frac{\partial \Psi}{\partial Y} \tag{12}$$

$$\left(\frac{\partial^2 \theta}{\partial X^2} \right) + \left(1 + \frac{4}{3} \left(\frac{k_{nf}}{k_f} \right)^{-1} \text{Rd} \right) \frac{\partial^2 \theta}{\partial Y^2} = - \frac{\partial \Psi}{\partial X} \frac{\partial \theta}{\partial Y} + \frac{\partial \theta}{\partial X} \frac{\partial \Psi}{\partial Y} \tag{13}$$

In two recent equations, the definition of new parameters is:

$$\text{Ha} = \frac{\sigma_f K B_0^2}{\mu_f}, \text{Ra} = \frac{g K (\rho \beta)_f L \Delta T}{\mu_f \alpha_f}, \text{Rd} = 4 \sigma_e T_c^3 / (\beta_R k_f)$$

$$A_1 = \frac{\rho_{nf}}{\rho_f}, A_2 = \frac{(\rho C_p)_{nf}}{(\rho C_p)_f}, A_5 = \frac{\mu_{nf}}{\mu_f},$$

$$A_3 = \frac{(\rho \beta)_{nf}}{(\rho \beta)_f}, A_6 = \frac{\sigma_{nf}}{\sigma_f}, A_4 = \frac{k_{nf}}{k_f} \tag{14}$$

In addition, the boundaries are summarized as:

$$\theta = 0.0 \quad \text{on outer surfaces}$$

$$\Psi = 0.0 \quad \text{all boundaries}$$

$$\frac{\partial \theta}{\partial n} = 1.0 \quad \text{on inner surface} \tag{15}$$

To calculate Nu_{loc} and Nu_{ave} , the below equations were employed:

Table 3 Different values of *m*


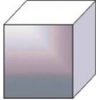


Platelet	Brick	Cylinder	Spherical	Shape
				<i>m</i>
5.7	3.7	4.8	3	

Table 4 Variation of Nu_{ave} with change of mesh size at $Ra = 600, Rd = 0.8, Ha = 20$ and $\phi = 0.04$.

Mesh size in radical direction \times angular direction				
51×151	61×181	71×211	81×241	91×271
1.777101	1.7797347	1.783917	1.784450	1.7860465

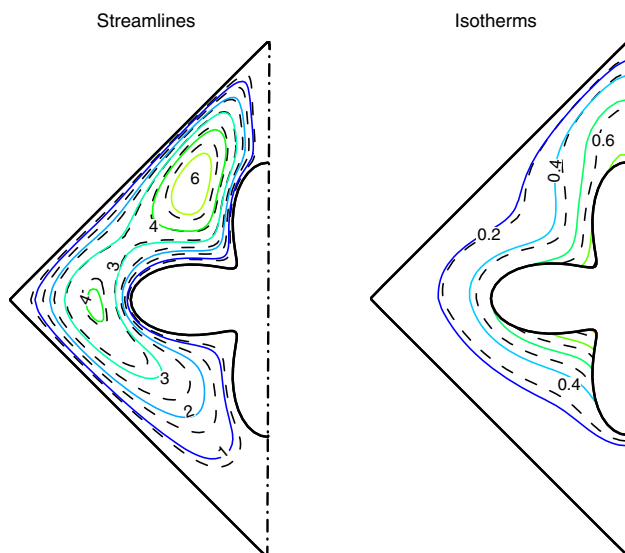
Table 5 Variation of Nu_{ave} at $Pr = 0.733$

Ha	Gr = 2×10^5	
	Present	Rudraiah et al. [1]
50	2.67911	2.8442
10	4.9047	4.8053

$$Nu_{loc} = \frac{1}{\theta} \left(1 + \frac{4}{3} \left(\frac{k_{nf}}{k_f} \right)^{-1} Rd \right) \left(\frac{k_{nf}}{k_f} \right) \quad (16)$$

$$Nu_{ave} = \frac{1}{S} \int_0^s Nu_{loc} ds \quad (17)$$

Sheikholeslami [34] is the pioneer of CVFEM. This technique combines FEM and FVM and uses the benefits of both approaches. In the final step, the Gauss–Seidel technique has been applied to find the values of scalars in each corner of triangular element. Various improvements in numerical approaches have been reported in recent decade [35–47].

**Fig. 2** Impacts of ϕ on nanofluid behaviour ($\phi = 0.04$ (—) and $\phi = 0$ (- -)) when $Ra = 600, Ha = 0, m = 5.7, Rd = 0.8$

In order to attain mesh insensitive results, a mesh analysis has been conducted for all the states. Table 4 is an example which demonstrates the results of various mesh sizes for a certain case. Furthermore, to ensure the correctness of the written code, the outputs have been compared with the previously published studies [1] which employed the same code. Table 5 validates that the present results are in reasonable accordance with the past literature. In addition, other validation exists in Refs. [48, 49].

Results and discussion

In the present investigation, the role of nanoparticle shape on the transportation of nanomaterial within porous tank with employing buoyancy and radiation parameters was depicted. To control the velocity, magnetic force was involved. The results have been analysed to predict the impact of the radiation parameter ($0 \leq Rd \leq 0.8$), shape factor ($3 \leq m \leq 5.7$), Hartmann number ($0 \leq Ha \leq 20$) and concentration ($0\% \leq \phi \leq 5\%$) of the alumina nanoparticles.

Figure 2 illustrates the impact of platelet-shaped ($m = 5.7$) nanoparticles addition ($\phi = 4\%$) on the streamlines and the isotherm profiles of the porous medium at $Ra = 600$ and $Rd = 0.8$ without magnetic field effect. It is evident that the convective coefficient reduces while the Ψ_{max} augments with the introduction of the nanomaterial within hosting fluid because k_{nf} is greater than k_f . The impact of the Hartmann number on the nanofluid migration at $Rd = 0.8, m = 5.7, \phi = 4\%$ for the $Ra = 100$ and $Ra = 600$ is illustrated in Figs. 3 and 4. The impact of the Lorentz forces results in the augmentation of temperature surface. The fluid convection incites a clockwise recirculating eddy which is subdivided into the bottom and top sections of the cavity, and this division is more perceptible at lower Ra and higher Ha . Lorentz forces results in the augment of Ψ_{max} . The mode of heat transfer is conductive at lower Ra , while it is relatively convective at the higher Ra .

The dependence of Nu_{avg} on the shape factor (m), (Rd), (Ha) have been shown in Fig. 5. Results show that Nu_{avg} observes the direct dependence on the m . The platelet shapes ($m = 5.7$) demonstrated the best performance of all the selected nanoparticle shapes followed by cylindrical ($m = 4.8$)-, brick ($m = 3.7$)- and spherical ($m = 3$)-shaped particles. Like m , the Rd also demonstrates a direct relationship with the Nu_{avg} . However, compared to the m , Ra and the Ha , the results of Nu_{avg} are more sensitive towards the variation of the Rd . The trend of the Nu_{avg} is also

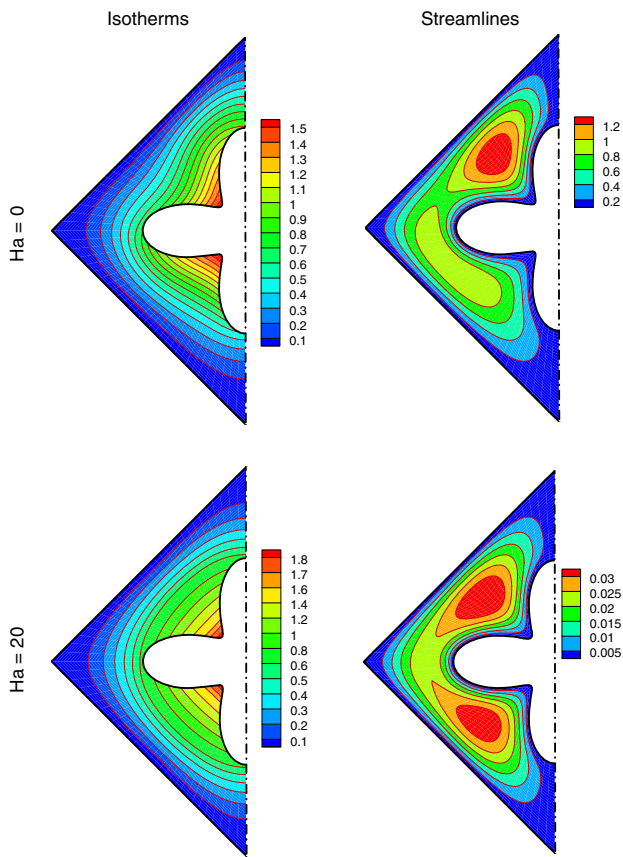


Fig. 3 Impacts of Ha on nanofluid flow when $Ra = 100, Rd = 0.8, m = 5.7, \phi = 0.04$

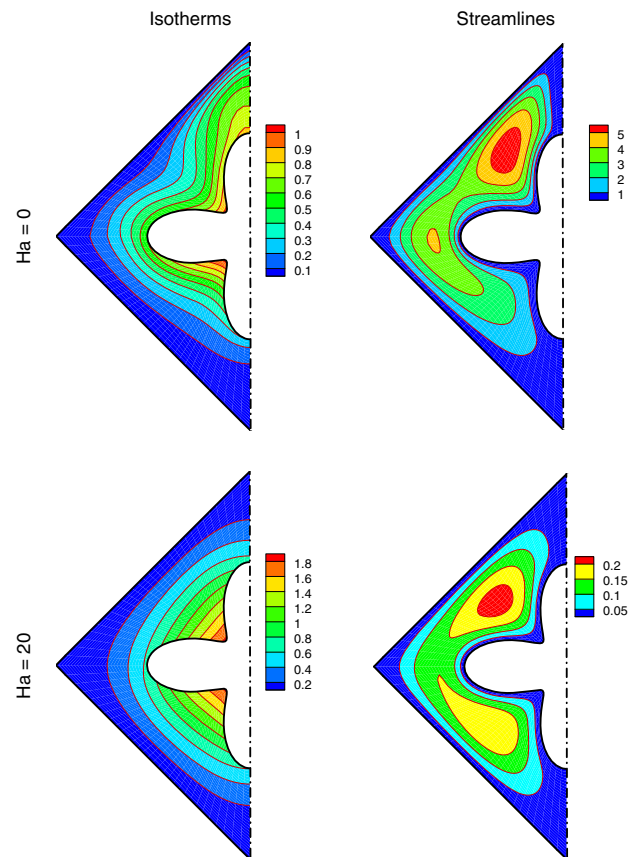


Fig. 4 Impacts of Ha on nanofluid flow when $Ra = 600, Rd = 0.8, m = 5.7, \phi = 0.04$

ascending as a function of the Ha and Ra; however, for the conditions of constant a Rd , m and φ, the Nu_{avg} demonstrates declining trend with the increasing values of Ha and Ra. Based on the results, a correlation for the estimation of the Nu_{avg} as a function of the considered parameter can be predicted as;

$$\begin{aligned}
 Nu_{ave} = & 1.86 + 0.042 m + 0.64 Rd + 0.3 Ra - 0.25 Ha \\
 & + 7 \times 10^{-3} m Ha - 0.16 Rd Ha - 0.3 Ra Ha \\
 & + 1.4 \times 10^{-4} m^2
 \end{aligned}
 \tag{18}$$

The validity of the proposed correlation is applicable for the studied radiation parameter ($0 \leq Rd \leq 0.8$), shape factor ($3 \leq m \leq 5.7$), nanomaterial concentration ($0\% \leq \phi \leq 5\%$), Hartmann number ($0 \leq Ha \leq 20$) and ($100 \leq Re \leq 600$).

Conclusions

The manuscript investigates the influence of the nanomaterial shape variation on the MHD-free convection and the radiative heat transfers of the Al₂O₃–water nanomaterial within a porous medium by employing the CVFEM approach. The considered radiation parameter, shape factor, Hartmann number and particle volume fraction ranges are $0 \leq Rd \leq 0.8$, $3 \leq m \leq 5.7$, $0 \leq Ha \leq 20$ and $0\% \leq \phi \leq 5\%$, respectively. The findings of the study can be concluded as:

- The higher shape factor augments the Nu_{ave}. The platelet shapes ($m = 5.7$) demonstrated the highest Nu_{avg} followed by cylindrical ($m = 4.8$)-, brick ($m = 3.7$)- and spherical ($m = 3$)-shaped particles.
- Compared to the m, Ra and the Ha, the results of Nu_{ave} are more sensitive towards the variation of the Rd. The convective mode is predominating when Ha = 0.

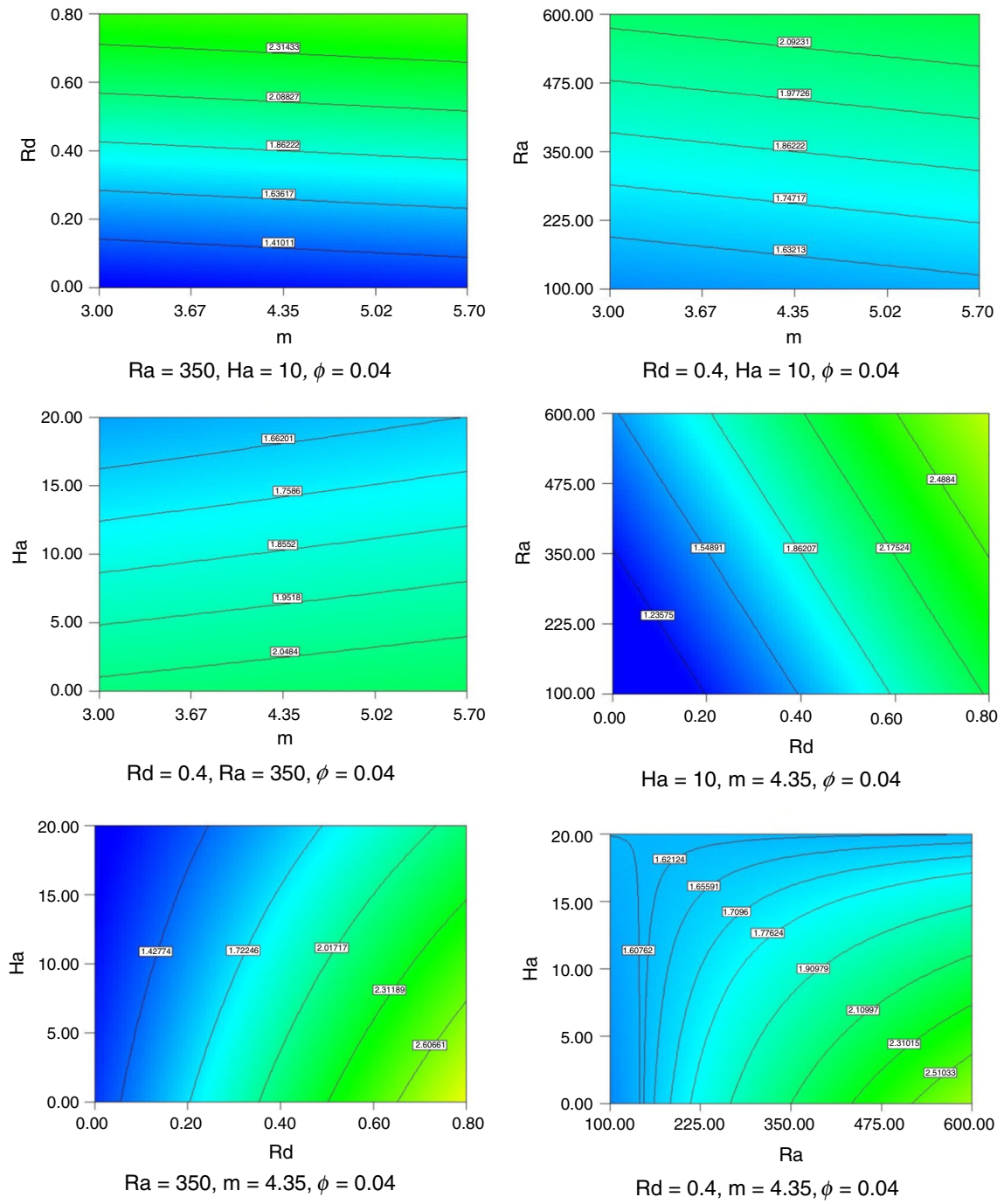


Fig. 5 Impacts of Ra, Ha, Rd, m, ϕ on Nu_{ave}

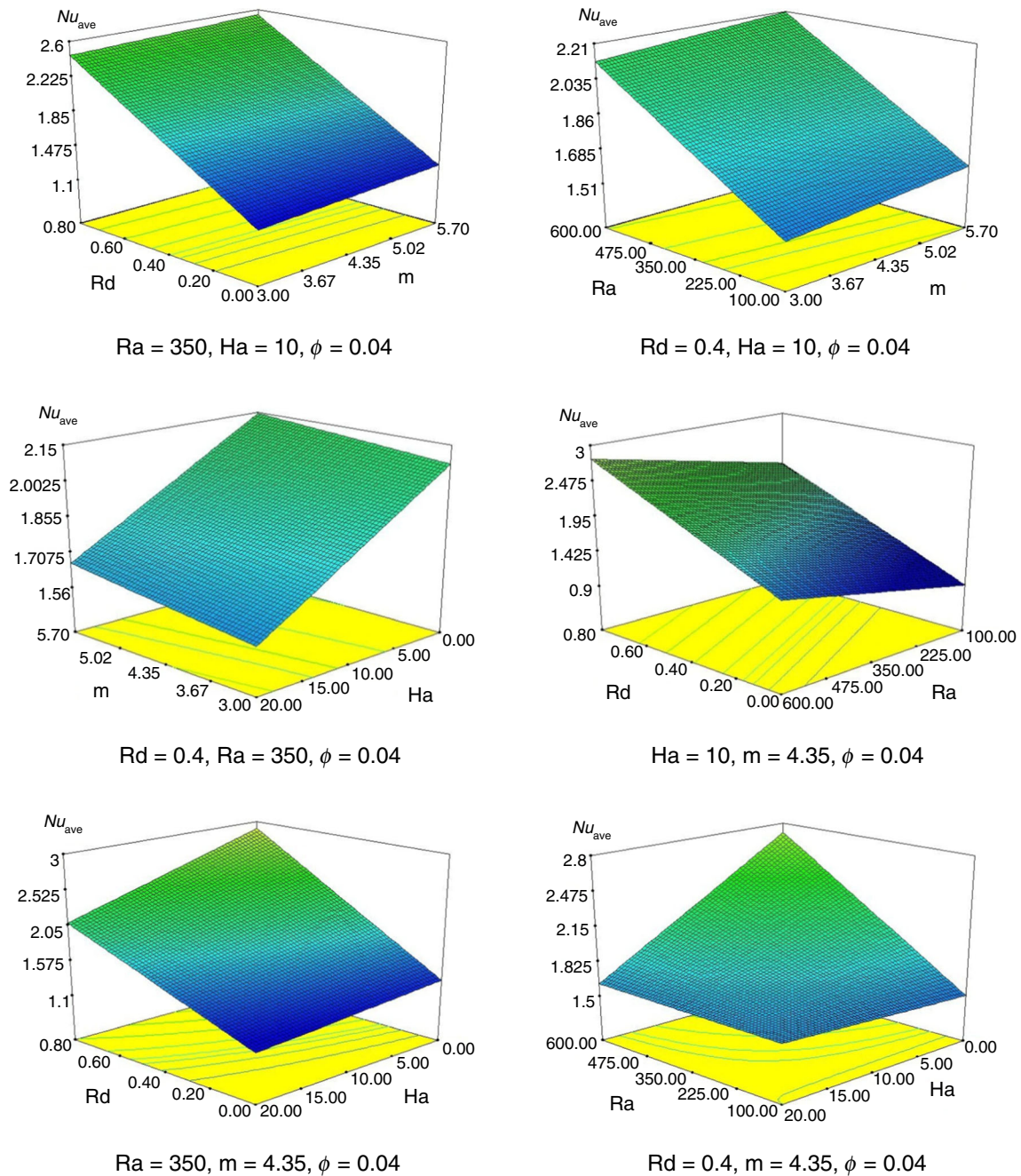


Fig. 5 continued

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