FAST TRACK COMMUNICATIONS

Surface chemical bonding states and ferroelectricity of Ce-doped BiFeO₃ thin films prepared by sol–gel process

Zuci Quan · Hao Hu · Sheng Xu · Wei Liu · Guojia Fang · Meiya Li · Xingzhong Zhao

Received: 1 June 2008/Accepted: 4 August 2008/Published online: 25 August 2008 © Springer Science+Business Media, LLC 2008

Abstract $Bi_{1-x}Ce_xFeO_3$ (x = 0, 0.05, 0.1, 0.15 and 0.20) (BCFO) thin films were deposited on Pt/TiN/Si₃N₄/Si substrates by sol-gel technique. Crystal structures, surface chemical compositions and bonding states of BCFO films were investigated by X-ray diffraction and X-ray photoelectron spectroscopy (XPS), respectively. Compared to BiFeO₃ (BFO) counterparts, the fitted XPS narrow-scan spectra of Bi $4f_{7/2},$ Bi $4f_{5/2},$ Fe $2p_{3/2},$ Fe $2p_{1/2}$ and O 1s peaks for Bi_{0.8}Ce_{0.2}FeO₃ film shift towards higher binding energy regions by amounts of 0.33, 0.29, 0.43, 0.58 and 0.49 eV, respectively. Dielectric constants and loss tangents of the BCFO (x = 0, 0.1 and 0.2) film capacitors are 159, 131, 116, 0.048, 0.041 and 0.035 at 1 MHz, respectively. Bi_{0.8}Ce_{0.2}FeO₃ film has a higher remnant polarization ($P_r = 2.04 \ \mu C/cm^2$) than that of the BFO $(P_r = 1.08 \ \mu\text{C/cm}^2)$ at 388 kV/cm. Leakage current density of the Bi_{0.8}Ce_{0.2}FeO₃ capacitor is 1.47×10^{-4} A/cm² at 388 kV/cm, which is about two orders of magnitude lower than that of the BFO counterpart. Furthermore, Ce cations are feasibly substituted for Bi3+ in the Bi_{0.8}Ce_{0.2}FeO₃ matrix, possibly resulting in the enhanced ferroelectric properties for the decreased grain sizes and the reduced oxygen vacancies.

Keywords BiFeO₃ thin films · Ce-doping · Surface chemical bonding states · Dielectric and ferroelectric properties

1 Introduction

ABO₃-type perovskite structure BiFeO₃ (BFO) has attracted much attention because it is a known Pb-free and environmentally friendly material that simultaneously shows electric and magnetic ordering in the same phase [1-3]. Based on the specific characteristics, including magnetoelectric coupling effects [1], high Curie temperature ($T_{\rm C} \sim 1,103$ K) and Néel temperature ($T_{\rm N} \sim 643$ K) [4], and small optical bandgap ($E_{\rm g} \sim 2.5 \text{ eV}$) [5], BFO thin films have potential applications in magnetic/ferroelectric data storage media, spintronics, nonvolatile memories, photocatalytic compound and ultrafast optoelectronic devices [1, 3, 5, 6]. Though as promising as BiFeO₃ films are, there are still drawbacks, such as large dielectric loss, small remnant polarization, high coercive field, high leakage current and inhomogeneous magnetic spin structure [4, 7, 8], to be overcome before practical applications. One way of doing so is by doping BFO films with other elements. In order to enhance ferroelectric properties, reduce leakage current, modify its spatially inhomogeneous spin-modulated incommensurate structure, and intensify magnetoelectric interaction, several research groups have tried A-site, B-site or (A, B)-sites of BFO films substituting by La [2, 7], Tb [9], Gd [10], Sc [4], Ti [11], Cr [12], La and Ti [13], La and Mn [14] as well as La and Nb [15]. More recently, Shannigrahi et al. reported that impure phase appeared in 0.5 mol% Sc-doped BFO film, and leakage current, ferroelectric and magnetic properties were improved for 0.3 mol% Sc-doped BFO [4]. Wang et al. investigated that Ti⁴⁺ ions were used to substitute some of Fe cations to form $Ti_{Fe}^{\bullet\bullet}$ as substituting Fe^{2+} or Ti_{Fe}^{\bullet} as substituting Fe³⁺, which could eliminate oxygen vacancies in the Ti-doped BFO films because of the requirements of charge compensation [11]. Cheng et al. found that the

Z. Quan \cdot H. Hu \cdot S. Xu \cdot W. Liu \cdot G. Fang \cdot M. Li \cdot X. Zhao (\boxtimes)

Key Laboratory of Acoustic and Photonic Materials and Devices of Ministry of Education, Department of Physics, Wuhan University, Wuhan 430072, People's Republic of China e-mail: xzzhao@whu.edu.cn

ferroelectric properties were significantly improved, and the ferromagnetic moments were kept at a high level for the (La, Nb)-codoped BFO films [15]. Being a lanthanide element, however, Ce has remained relatively unexplored in doping BFO films.

In this study, crystal structures, surface chemical compositions and bonding states of pure and Ce-doped BFO films were characterized by X-ray diffraction (XRD) and X-ray photoelectron spectroscopy (XPS), respectively. Chemical shifts of Bi $4f_{7/2}$, Bi $4f_{5/2}$, Fe $2p_{3/2}$, Fe $2p_{1/2}$, O 1s, Ce $3d_{5/2}$ and Ce $3d_{3/2}$ peaks for pure BFO and Bi_{0.8}Ce_{0.2}FeO₃ films were analyzed by fitting the corresponding XPS narrow-scan spectra. Dielectric and ferroelectric properties of the pure and Ce-doped BFO film capacitors were also discussed.

2 Experimental details

 $[Bi(NO_3)_3 \cdot 5H_2O],$ **Bismuth** nitrate iron nitrate $[Fe(NO_3)_3 \cdot 9H_2O]$ and cerium nitrate $[Ce(NO_3)_3 \cdot 6H_2O]$ were suitably dissolved in the mixtures of 2-methoxyethanol (CH₃OCH₂CH₂OH) and glacial acetic acid (CH₃COOH), while Ce contents were determined by composition formulae $Bi_{1-x}Ce_xFeO_3$ (abbreviated as BCFO), where x is equal to 0 (for pure BFO), 0.05, 0.10, 0.15 and 0.20, respectively. Meanwhile, acetylacetone (CH₃COCH₂COCH₃) and N-Ndimethyl formamide [(CH₃)₂NOCH] in 1 to 1 molar ratio with respect to $Fe(NO_3)_3 \cdot 9H_2O$ were used to stabilize BCFO solutions. The stock solutions were stirred at 318 K for 240 min to obtain homogeneous BCFO precursors. BCFO sols were spin-coated onto Pt (111)/TiN/Si₃N₄/Si (100) substrates at 3,000 rpm for 20 s. After each spincoating, the green BCFO films were dried at 393 K for 10 min, and then prebaked at 693 K for 30 min in the air to remove organic contaminations. This procedure was repeated several times to obtain appropriate thickness. Finally, BCFO films were annealed in a quartz tube furnace at 873 K for 60 min under a N2 ambience, and then cooled down slowly to room temperature.

Crystal structures of BCFO films were analyzed by XRD (D8 Advance, Bruker Corp.) with Cu K_{α} radiation. Thicknesses of BCFO films were measured by scanning electron microscope (SEM) (Sirion FEG, FEI Co.). Surface chemical compositions and binding states of the corresponding elements for BFO and Bi_{0.8}Ce_{0.2}FeO₃ films were investigated by XPS (VG MultiLab 2000, Thermo Electron Corp.). The XPS Al K_{α} source operating at 300 W provided nonmono-chromatic X-rays at 1,486.6 eV. Photoelectric peak of C 1s located at 284.6 eV was assigned to carbon from adventitious contaminations, and it was used as the criterion to rectify binding energies of XPS spectra. Software of XPSPEAK Version 4.1 was used to fit the narrow-scan

spectra of Bi 4f, Fe 2p, O 1s and Ce 3d after Shirley-type background subtraction [16, 17]. For comparison, 120-nmthick Pt top electrodes with diameter of 0.3 mm were sputtered on BCFO films through a metal shadow mask. To understand Ce-doped influence on dielectric and ferroelectric properties of BFO film, capacitance-frequency and dielectric loss-frequency curves for the BCFO (x = 0, 0.1and 0.2) film capacitors were measured by a precision impedance analyzer (Agilent 4294A, Agilent Technologies Ltd.). Polarization–electric field (P-E) hysteresis loops and leakage current density–electric field (J-E) curves for the same BCFO capacitors were measured by a precision materials analyzer (RT6000, Radiant Technology Inc.). All measurements were carried out at room temperature.

3 Results and discussion

Figure 1 shows XRD patterns of BCFO films annealed in 873 K for 60 min under a N₂ ambience. It is found that all the BCFO films are polycrystalline and exhibit a singlephase perovskite structure. Non-perovskite phase such as Bi₂Fe₄O₉, Bi₃₆Fe₂O₅₇ and Bi₁₂(Bi_{0.5}Fe_{0.5})O_{19.5} [6, 18] are not detected in the XRD patterns. The obvious peaksplitting XRD pattern for the BCFO film with x = 0indicate that the BFO film is rhombohedral phase (space group *R3c*) [10]. As shown in the broken square of Fig. 1, (104) and (110) peaks of BCFO films move closer together along with broadening, indicating a structural change [19, 20] with increasing Ce content. Since ionic radii of Ce³⁺ (1.18 Å) and Ce⁴⁺ (1.02 Å) are smaller than that of Bi³⁺ (1.20 Å) [21], Ce cations are partially substituted for Bi³⁺



Fig. 1 XRD patterns of $\text{Bi}_{1-x}\text{Ce}_x\text{FeO}_3$ (x = 0, 0.05, 0.1, 0.15 and 0.2) films deposited on Pt/TiN/Si₃N₄/Si substrates. As shown in the broken square of Fig. 1, (104) and (110) peaks merge when x increases from 0 to 0.2

with increasing Ce dopant, resulting in smaller grain sizes of the BCFO films compared to those of the BFO counterparts.

Figure 2 shows XPS survey spectra of BFO and $Bi_{0.8}Ce_{0.2}TiO_3$ films. There are XPS photoelectron peaks and the corresponding Auger lines of Bi, Fe, O and C elements on the BFO film surface. Herein, C 1s peak located at 284.6 eV is used as the criterion to rectify the binding energy of XPS spectra. Characteristic peaks for Bi 4f_{7/2} (158.56 eV), Bi 4f_{5/2} (163.90 eV), Fe 2p_{3/2} (710.33 eV), Fe 2p_{1/2} (723.69 eV), O 1s (529.52 eV) and Auger peaks for Fe L₃M₄₅M₄₅ (783.86 eV), Fe L₃M₂₃M₄₅ (838.24 eV) and Fe LM₂₃M₂₃ (888.29 eV) [17] are identified in Fig. 2. Undoubtedly, there is no signal of Ce element in the pure BFO film. As for Bi_{0.8}Ce_{0.2}FeO₃ film, there are XPS signals of Bi, Ce, Fe, O and C elements, where Ce $4d_{5/2}$ (108.71 eV), Ce $4d_{3/2}$ (112.43 eV), Ce 3d_{5/2} (882.32 eV) and Ce 3d_{3/2} (900.64 eV) and Auger peak for Ce $M_{45}N_{45}N_{45}$ (832.87 eV) [17] are assigned in Fig. 2. It is concluded that some Ce-O bonds are present in Bi_{0.8}Ce_{0.2}FeO₃ matrix.

Figure 3 shows the fitted narrow-scan spectra of Bi 4f, Fe 2p, O 1s and Ce 3d peaks for the BFO and Bi_{0.8}Ce_{0.2}FeO₃ films. Figure 3a(1) shows that pure Bi 4f doublet consists of two peaks at 158.56 and 163.90 eV, which are mainly identified as a signal from Bi–O bonds [17]. Spin-orbit splitting energy (Δ) of the pure Bi 4f doublet is 5.34 eV, which is comparable to theoretical value ($\Delta_{Bi \ 4f}$) of 5.31 eV [16]. Two fitted subpeaks located at 158.28 and 163.61 eV are ascribed to Bi (4f_{7/2})–O and Bi (4f_{5/2})–O bonds, while other subpeaks located at 158.84 and 164.18 eV may be related to Bi–O–Fe bonds in oxygen octahedron and/or relaxed Bi phase, probably caused by



Fig. 2 XPS survey spectra of BFO and $Bi_{0.8}Ce_{0.2}FeO_3$ films annealed at 873 K for 60 min under a N_2 ambience

oxygen vacancies and cation defects [17, 22]. Figure 3a(2)shows that compared to the pure Bi 4f doublet, Bi $4f_{7/2}$ (158.89 eV) and Bi $4f_{5/2}$ (164.19 eV) peaks of the Bi_{0.8}Ce_{0.2}FeO₃ film shift towards higher binding energy regions by amounts of 0.33 and 0.29 eV, respectively. Herein, two subpeaks located at 158.56 and 163.86 eV are ascribed to Bi (4f_{7/2})–O and Bi (4f_{5/2})–O bonds, while other subpeaks located at 159.22 and 164.52 eV may be related to Bi-O-Fe and Bi-O-Ce bonds in oxygen octahedron and/ or relaxed Bi phase. It is conceivable that Ce cations are substituted for Bi³⁺ at A-site in the Bi_{0.8}Ce_{0.2}FeO₃ matrix because of the decreased Bi 4f peak area in comparison with the pure counterpart. As reported by Gao et al., the N-doped Ba 3d_{5/2} and Ba 3d_{3/2} peaks shift towards the lower binding energy regions because the bonding energies of Ba-N bonds are lower than that of Ba-O bonds considering the slightly lower electronegativity of N element (3.04) than that of O element (3.44) [21]. So the chemical shifts of Bi 4f, Fe 2p and O 1s peaks for the Bi_{0.8}Ce_{0.2}FeO₃ film may be related to the various electronegativity values for Bi (2.02), Ce (1.12), Fe (1.83) and O (3.44) elements [16, 21, 23]. To estimate ionicity of Bi-O, Ce-O and Fe-O bonds for the Bi_{0.8}Ce_{0.2}FeO₃ film, fraction of ionicity (F_i) [24] is given by:

$$F_i = 1 - \exp\left[-\frac{\left(\Delta EN\right)^2}{4}\right] \tag{1}$$

where ΔEN is difference in the anion and cation electronegativities. According to Eq. (1), the fraction of ionicity $[F_i(Ce-O)]$ (0.74) is much larger than $F_i(Bi-O)$ (0.40), and bonding energy of (Bi,Ce)–O bond in oxygen octahedron may be larger than that of single Bi–O bond, suggesting that the Bi 4f doublet of Bi_{0.8}Ce_{0.2}FeO₃ film shifts to the higher binding energy region.

Figure 3b shows the fitted Fe 2p narrow-scan spectra for the BFO and Bi_{0.8}Ce_{0.2}FeO₃ films. In Fig. 3b(1), the pure Fe 2p doublet consists of two wide peaks of Fe $2p_{3/2}$ (710.33 eV) and Fe $2p_{1/2}$ (723.69 eV), which are mainly ascribed to Fe–O bonds [17]. Spin-orbit splitting energy of the pure Fe 2p doublet is equal to 13.36 eV, which is compared to theoretical value ($\Delta_{\text{Fe 2p}}$) of 13.6 eV for Fe₂O₃ [17]. Figure 3b(2) shows that compared to the pure Fe 2p doublet, Fe $2p_{3/2}$ (710.76 eV) and Fe $2p_{1/2}$ (724.27 eV) peaks of the Bi_{0.8}Ce_{0.2}FeO₃ film shift towards higher binding energy regions by amounts of 0.43 and 0.58 eV, respectively. Herein, two subpeaks located at 711.51 and 725.15 eV are ascribed to (Fe 2p_{3/2})₂–O₃ and (Fe 2p_{1/2})₂–O₃ bonds, while the subpeaks located at 710.00 and 723.38 eV may be related to Fe-O bonds for Fe²⁺, Fe-O-Bi bonds, Fe-O-Ce bonds in oxygen octahedra and/or other relaxed Fe phase [16, 25]. It is the reason why the Fe $2p_{3/2}$ and Fe $2p_{1/2}$ peaks of the Bi_{0.8}Ce_{0.2}FeO₃ film shift towards the higher binding energy

Fig. 3 The fitted XPS narrowscan spectra of (a) Bi 4f, (b) Fe 2p, (c) O 1 s and (d) Ce 3d peaks for (1) BFO and (2) Bi_{0.8}Ce_{0.2}FeO₃ films. The black solid lines represent experimental result, red short dots represent the fitting results after Shirley-type background subtraction, the dashed peaks with colors of blue, magenta, violet and green are fitted subpeaks. Iterations are performed until a minimum for square of the difference (χ^2) between the experimental curve and the fitted curve is reached for each fitted sum curve



regions in comparison with the Fe 2p doublet of the BFO film. According to ratio of the fitted peak areas for Fe³⁺ and Fe²⁺, concentration ratios of Fe³⁺ and Fe²⁺ in the BFO and Bi_{0.8}Ce_{0.2}FeO₃ films are 67:33 and 70:30, respectively, indicating that the presence of Fe²⁺ ions are less in the Bi_{0.8}Ce_{0.2}FeO₃ film compared to that of the BFO film. This could be an evidence for the decreased oxygen vacancies in the Bi_{0.8}Ce_{0.2}FeO₃ film.

Figure 3c shows the fitted O 1s narrow-scan spectra for the BFO and $Bi_{0.8}Ce_{0.2}FeO_3$ films. Figure 3c(1) shows that a broad O 1s peak (529.52 eV) for the BFO film consists of three subpeaks located at 528.77, 529.34 and 530.45 eV. Since the BFO film consists of three components (Bi₂O₃, Fe₂O₃ and FeO) in BFO solid solution, the subpeaks are mainly ascribed to Fe₂-(O 1s)₃ (529.6 eV), Fe-(O 1s) (529.8 eV), Bi₂-(O 1s)₃ (530.0 eV) bonds [17] and relaxed O phase, especially for the fitted subpeak located at the higher binding energy region. It may be mainly attributed to the absorbed oxygen (531.0 eV) [17] associated with oxygen vacancies and surface species, such as H₂O and CO_2 absorbed from the air during the sol-gel process [16, 17]. Figure 3c(2) shows that compared to the counterpart of the BFO film, O 1s peak (530.01 eV) for the Bi_{0.8}Ce_{0.2}FeO₃ film shifts towards higher binding energy region by amount of 0.49 eV. These subpeaks are mainly ascribed to Fe₂-(O 1s)₃ (529.6 eV for the O 1s peak), Fe-(O 1s) (529.8 eV), Bi₂-(O 1s)₃ (530.0 eV) and Ce₂-(O 1s)₃ (530.3 eV) bonds [17]. Considering that Ce^{3+} ions are partially substituted for Bi^{3+} ions, and binding energy of O 1s in Ce_2 -(O 1s)₃ bond is larger than that of Bi_2 -(O 1s)₃ bond, suggesting that the O 1s peak of the $Bi_{0.8}Ce_{0.2}FeO_3$ film shifts toward the higher binding energy regions in comparison with the BFO counterpart.

Figure 3d shows the fitted Ce 3d_{5/2} peak located at 882.24 eV and Ce 3d_{3/2} peak located at 900.61 eV. Both of them are mainly assigned to Ce-O bonds. Spin-orbit splitting energy of the Ce 3d doublet is equal to 18.37 eV, which is compared to theoretical value ($\Delta_{Ce 3d}$) of 18.10 eV for Ce or 18.30 eV for CeO₂ [17]. The fitted area ratio for Ce $3d_{5/2}$ and Ce $3d_{3/2}$ peaks is about 1.47, which is comparative to theoretical value of 1.50 for Ce 3d doublet [17]. Figure 3d also shows that each of the Ce 3d peak is accompanied by two satellite lines, called 3d⁹4f⁰ and 3d⁹4f² [17, 26]. Furthermore, each of the Ce $3d_{5/2}$ and Ce $3d_{3/2}$ peak could be fitted by two subpeaks separated by 1.33 ± 0.01 eV. The subpeaks located at 882.91 and 901.27 eV are assigned to Ce 3d_{5/2}–O and Ce 3d_{3/2}–O bonds for Ce³⁺. Meanwhile, the subpeaks located at 881.57 and 899.95 eV may be associated with Ce $3d_{5/2}$ –O and Ce $3d_{3/2}$ –O bonds for Ce⁴⁺ [17], (Bi, Ce)-O bonds and/or relaxed Ce phase, probably caused by oxygen vacancies and other defects.

Figure 4 shows typical dielectric properties of the BCFO (x = 0, 0.1 and 0.2) film capacitors at the measured frequency ranging from 40 Hz to 1 MHz. The



Fig. 4 Dielectric constants and loss tangents of $Bi_{1-x}Ce_xFeO_3$ (x = 0, 0.1 and 0.2) film capacitors as a function of frequency ranging from 40 Hz to 1 MHz

corresponding dielectric constant (ε_r) and loss tangent (tan δ) for the BFO, Bi_{0.9}Ce_{0.1}FeO₃ and Bi_{0.8}Ce_{0.2}FeO₃ capacitors are 159, 131, 116, 0.048, 0.041 and 0.035 at 1 MHz, respectively. It means that compared to the BFO counterparts, both ε_r and tan δ of the Bi_{0.8}Ce_{0.2}FeO₃ capacitor decrease when 20 mol% Ce doped into the BFO matrix. Moreover, this result could be comparable to the values of PLD-derived Bi_{0.8}La_{0.2}FeO₃ ($\varepsilon_r = 86$, tan $\delta = 0.69$) and Bi_{0.8}La_{0.2}Nb_{0.01}Fe_{0.99}O₃ ($\varepsilon_r = 184$, tan $\delta = 0.038$) films [15] deposited on Pt/Ti/SiO₂/Si substrates at 100 kHz as reported by Cheng et al., respectively.

Figure 5 shows P-E hysteresis loops for the BFO, Bi_{0.9}Ce_{0.1}FeO₃ and Bi_{0.8}Ce_{0.2}FeO₃ film capacitors. These capacitors have typical ferroelectric characteristics, and the corresponding remnant polarization $[P_r = (P_r^+ - P_r^-)/2]$



Fig. 5 Polarization-electric field hysteresis loops of $Bi_{1-x}Ce_xFeO_3$ (x = 0, 0.1 and 0.2) capacitors

of 1.08, 1.43 and 2.04 μ C/cm² and coercive field [$E_c =$ $(E_{\rm c}^{+} - E_{\rm c}^{-})/2$] of 120, 112 and 123 kV/cm, respectively, are obtained under the applied electric field of 388 kV/cm. $P_{\rm r}$ value of the Bi_{0.8}Ce_{0.2}FeO₃ capacitor is larger than that of the BFO film, or those of the BiFe_{1-x}Sc_xO₃ (x = 0.1 and 0.3 mol%) films [4] or the Ti⁴⁺ and Ni²⁺ doped BFO films prepared by PLD process [27], but it is much small than that of the Gd-doped BFO film prepared by metal organic decomposition process [10]. According to Wang et al., ferroelectricity of BFO film originated from the displacements of Bi with respect to the FeO_6 cages along (111) plane [22]. Compared to the BFO counterpart, the increased P_r of the Bi_{0.8}Ce_{0.2}FeO₃ film might be attributed to the structural change in the oxygen octahedron, since Ce cations are feasibly substituted for Bi³⁺ as evidenced by the XPS spectra.

Figure 6 shows J–E curves of the BCFO (x = 0, 0.1and 0.2) capacitors. Leakage current densities of BFO, $Bi_{0.9}Ce_{0.1}FeO_3$ and $Bi_{0.8}Ce_{0.2}FeO_3$ capacitors are 1.38 × 10^{-2} , 1.06×10^{-3} and 1.47×10^{-4} A/cm² at 388 kV/cm, respectively. It means that the leakage current density of the Bi_{0.8}Ce_{0.2}FeO₃ capacitor is about two orders of magnitude lower than that of the BFO counterpart. One reason for this is that an appropriate amount of Ce doping reduces grain sizes of the BCFO film as evidenced by XRD result, leading to an increase in the density of grain boundaries, which makes contribution to the decreased leakage current density [9]. Another reason is that the decreased Bi volatilization and the restrained reduction of Fe³⁺ to Fe²⁺ are realized by virtue of the reducible oxygen vacancies and cation defects after Ce-doping, as proofed by XPS data.



Fig. 6 Leakage current density-electric field curves of $Bi_{1-x}Ce_xFeO_3$ (x = 0, 0.1 and 0.2) capacitors when electric filed ranges from -388 to 388 kV/cm

4 Conclusions

BCFO films were deposited on Pt/TiN/Si₃N₄/Si substrates by sol-gel technique. Compared to the BFO counterparts, the fitted Bi $4f_{7/2}$, Bi $4f_{5/2}$, Fe $2p_{3/2}$, Fe $2p_{1/2}$ and O 1s peaks for the Bi_{0.8}Ce_{0.2}FeO₃ film shift towards the higher binding energy regions by amounts of 0.33, 0.29, 0.43, 0.58 and 0.49 eV, respectively. The corresponding ε_r and tan δ of the BFO, Bi_{0.9}Ce_{0.1}FeO₃ and Bi_{0.8}Ce_{0.2}FeO₃ capacitors are 159, 131, 116, 0.048, 0.041 and 0.035 at 1 MHz, respectively. The corresponding P_r of 1.08, 1.43 and 2.04 μ C/cm² and E_c of 120, 112 and 123 kV/cm are obtained under the electric field of 388 kV/cm, respectively. The leakage current density of the Bi_{0.8}Ce_{0.2}FeO₃ capacitor is about two orders of magnitude lower than that of the BFO counterpart. Compared to the BFO counterparts, the increased P_r and the decreased leakage current density of the $Bi_{0.8}Ce_{0.2}FeO_3$ film might be attributed to the decreased Bi volatilization and the restrained reduction of Fe³⁺ to Fe²⁺ by virtue of the reducible oxygen vacancies and cation defects after Ce-doping, as proved by the XPS spectra.

Acknowledgements The authors gratefully acknowledge the financial supports from Hi-tech Plan of Ministry of Science and Technology (Grant No. 2006AA03Z347) and National Nature Science Foundation of People's Republic of China (Grant No. 50125309).

References

- 1. Zhao T, Scholl A, Zavaliche F, Lee K, Barry M, Doran A, Cruz MP, Chu YH, Ederer C, Spaldin NA, Das RR, Kim DM, Baek SH, Eom CB, Ramesh R (2006) Nat Mater 5:823
- Chu YH, Zhan Q, Yang CH, Cruz MP, Martin LW, Zhao T, Yu P, Ramesh R, Joseph PT, Lin IN, Tian W, Schlom DG (2008) Appl Phys Lett 92:102909
- Kumar A, Murari NM, Katiyar RS (2008) Appl Phys Lett 92:152907
- 4. Shannigrahi SR, Huang A, Chandrasekhar N, Tripathy D, Adeyeye AO (2007) Appl Phys Lett 90:022901

- 5. Takahashi K, Kida N, Tonouchi M (2006) Phys Rev Lett 96:117402
- Gao F, Chen XY, Yin KB, Dong S, Ren ZF, Yuan F, Yu T, Zou ZG (2007) Adv Mater 19:2889
- Simões AZ, Cavalcante LS, Riccardi CS, Varela JA, Longo E (2008) Curr Appl Phys. doi:10.1016/j.cap.2008.05.001
- 8. Ederer C, Spaldin NA (2005) Phys Rev B 71:060401
- 9. Wang Y, Nan CW (2008) J Appl Phys 103:024103
- 10. Hu GD, Cheng X, Wu WB, Yang CH (2007) Appl Phys Lett 91:232909
- 11. Wang Y, Nan CW (2006) Appl Phys Lett 89:052903
- Singh SK, Sato K, Maruyama K, Ishiwara H (2006) Jpn J Appl Phys Part 2 45:L1087
- 13. Lee CC, Wu JM (2007) Electrochem Solid State 10:G58
- Kartopu G, Lahmar A, Habouti S, Solterbeck CL, Elouadi B, Souni ME (2008) Appl Phys Lett 92:151910
- Cheng ZX, Wang XL, Dou SX, Kimura H, Ozawa K (2008) Phys Rev B 77:092101
- Zhang BS, Quan ZC, Zhang TJ, Guo T, Mo SB (2007) J Appl Phys 101:014107
- Moulder JF, Stickle WF, Sobol PE, Bomben KD (1992) Handbook of X-ray photoelectron spectroscopy. Perkin-Elmer Corporation, Minnesota
- Wang YP, Zhou L, Zhang MF, Chen XY, Liu JM, Liu ZG (2004) Appl Phys Lett 84:1731
- Brinkman K, Iijima T, Nishida K, Katoda T, Funakubo H (2007) Ferroelectrics 357:599
- Brinkman K, Iijima T, Takamura H (2007) Jpn J Appl Phys Part 2 46:93
- Schaffer JP, Saxena A, Antolovich SD, Sanders TH, Warner JSB (1999) The science and design of engineering materials, 2nd edn. McGraw-Hill Companies, New York
- 22. Wang J, Neaton JB, Zheng H, Nagarajan V, Ogale SB, Liu B, Viehland D, Vaithyanathan V, Schlom DG, Waghmare UV, Spaldin NA, Rabe KM, Wuttig M, Ramesh R (2003) Science 299:1719
- Gao YH, Shen H, Ma JH, Xue JQ, Sun JL, Meng XJ, Chu JH, Wang PN (2007) J Appl Phys 102:064106
- 24. Askeland DR (1994) The science and engineering of materials, 3rd edn. PWS Publishing Company, Boston
- 25. Wang Y, Jiang QH, He HC, Nan CW (2006) Appl Phys Lett 88:142503
- 26. Talik E, Guzik A (2003) Phys Stat Sol (a) 196:332
- 27. Qi XD, Dho J, Tomov R, Blamire MG, MacManus-Driscoll JL (2005) Appl Phys Lett 86:062903