

FRACTIONAL WAVE EQUATION WITH CHANGING DIRECTION OF EVOLUTION

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We study the fractional wave equation with changing direction of evolution. The existence and uniqueness of a generalized solution are proved. Bibliography: 5 titles.

1 Introduction

Assume that $T > 0$ and $\Omega \subset R^m$ is a bounded domain with smooth boundary $\Gamma = \partial\Omega$. In the cylinder $Q = (0, T) \times \Omega$, $S = (0, T) \times \Gamma$, $0 < \nu < 1$, we consider the mixed problem for the model equation with fractional Gerasimov–Caputo derivative

$$\partial_t^\nu(k(t, x)u_t(t, x)) - \Delta u(t, x) + \gamma u_t(t, x) = f(t, x) \quad (1.1)$$

where the coefficient $k(t, x) \in C^1(\overline{Q})$ is of arbitrary sign.

The fractional diffusion equation was considered in [2], where the solvability of some boundary value problems was established by the method of a priori estimates of the form

$$\int_0^T \psi(v) D^\nu(kv) dt \geq C(\|v\|_B).$$

More general inequalities and their applications can be found in [1] (see the references in [2]).

In the present paper, we use a similar technique to study the fractional wave equation. The statement of the problem is similar to that for mixed type equations with usual derivative.

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Namely, under some conditions, the problem

$$\begin{aligned} (k(t, x)u_t(t, x))_t - \Delta u(t, x) + \gamma u_t(t, x) &= f(t, x), \\ u(t, x)|_S &= 0, \\ u(0, x) &= 0, \\ u_t(0, x) &= u_0(x), \quad x \in \Omega_0^+ = \{x | k(0, x) > 0\}, \\ u_t(T, x) &= u_1(x), \quad x \in \Omega_T^- = \{x | k(T, x) < 0\} \end{aligned}$$

is well-posed. A similar mixed problem is well-posed for Equation (1.1). We prove the existence of a generalized solution to this mixed problem. The uniqueness of a solution is established under additional conditions on the sign of $k(0, x)$ and $k(T, x)$.

2 Auxiliaries

We use the following definitions and properties of fractional derivatives (see [3]). Assume that $0 < \nu < 1$ and $t > 0$. The fractional integral of order ν with the origin at a point a is defined by

$$J_a^\nu y(t) = \frac{\text{sgn}(t-a)}{\Gamma(\nu)} \int_a^t \frac{y(s)}{|t-s|^{1-\nu}} ds.$$

The fractional Riemann–Liouville derivative of order ν with the origin at a is defined by

$$D_a^\nu y(t) = \frac{1}{\Gamma(1-\nu)} \frac{d}{dt} \int_a^t \frac{y(s)}{|t-s|^\nu} ds.$$

The fractional Gerasimov–Caputo derivative is defined by

$$\partial_a^\nu y(t) = \frac{1}{\Gamma(1-\nu)} \int_a^t \frac{y'(s)}{|t-s|^\nu} ds.$$

In the case $a = 0$, we write $J^\nu y(t)$, $D^\nu y(t)$, $\partial^\nu y(t)$. As known, for $\nu \neq 1/2$

$$C_0(\nu) \|y(t)\|_{W_2^\nu(0,T)}^2 \leq \int_0^T (y^2(t) + (D^\nu y(t))^2) dt \leq C_1(\nu) \|y(t)\|_{W_2^\nu(0,T)}^2.$$

We use some assertions proved in [2]. Let $y(t) \in C^1(0, T)$. Then for some constant $C > 0$ we have (see [1, Lemma 4.1] and [2, formula (2.1)])

$$\int_0^T y(t) D^\nu y(t) dt \geq C \int_0^T y^2(t) \left(\frac{1}{t^\nu} + \frac{1}{(T-t)^\nu} \right) dt + C \|y(t)\|_{W_2^{\nu/2}(0,T)}^2; \quad (2.1)$$

moreover, if $y(0) = 0$, then (see [1, Lemma 5.4])

$$\int_0^T y'(t) D^\nu y(t) dt = \int_0^T D^{1-\nu}(D^\nu y(t)) D^\nu y(t) dt \geq C \|D^\nu y(t)\|_{W_2^{(1-\nu)/2}(0,T)}^2$$

and

$$\int_0^T y'(t) D^\nu y(t) dt \geq C \|y(t)\|_{W_2^{(1+\nu)/2}(0,T)}^2. \quad (2.2)$$

Further, for any smooth functions $f(t)$ and $g(t)$ we set

$$J(f, g) = \int_0^T f(t)g(t) dt.$$

Then for any $0 < \mu < 1/2$ (see [2, formula (2.2)])

$$|J(f, g)| \leq C_3(\mu, T) \|J^\mu f(t)\|_{L_2(0,T)} \|D^\mu g(t)\|_{L_2(0,T)}. \quad (2.3)$$

Assume that $0 < \mu < 1$ and $0 < \theta < 1$. We put

$$J_{\mu,\theta}y(t) = \frac{1}{\Gamma(\mu)} \int_0^t \frac{y(s)}{(t-s+\theta)^{1-\mu}} ds,$$

$$K_{\mu,\theta}y(t) = \frac{1}{\Gamma(1-\mu)} \frac{d}{dt} \int_0^t \frac{y(s)}{(t-s+\theta)^\mu} ds.$$

It is easy to see that

$$K_{\mu,\theta}v(t, x) = \frac{d}{dt} J_{1-\mu,\theta}.$$

For some constant $C_K = C_K(\mu, T) > 0$ independent of θ we have (see [2, formula (2.3)])

$$\int_0^T y(t) K_{\mu,\theta}y(t) dt \geq C_K \|y(t)\|_{L_2(0,T)}. \quad (2.4)$$

We set

$$D_{\mu,\theta}(y(t), z(t)) = \int_0^T (D^\mu y(t) - K_{\mu,\theta}y(t))z(t) dt.$$

Lemma 2.1 ([2, Lemma 2.1]). *There exists a constant $C(\mu, T) > 0$ independent of θ such that for any smooth functions $y(t)$ and $z(t)$ the following inequality holds:*

$$|D_{\mu,\theta}(y(t), z(t))| \leq C(\mu, T) \theta^{(1-\mu)/2} \|y(t)\|_{W_2^1(0,T)} \|z(t)\|_{W_2^1(0,T)}. \quad (2.5)$$

Corollary 2.1 ([2, Corollary 2.1]). *For any smooth function $y(t)$ the following estimate holds:*

$$\int_0^T y(t) K_{\mu,\theta}y(t) dt \geq C_{K1} \|y(t)\|_{W_2^{\mu/2}(0,T)}^2 - C_{K2} \theta^{(1-\mu)/2} \|y(t)\|_{W_2^1(0,T)}^2 \quad (2.6)$$

with some constants $C_{K1}(\mu, T) > 0$ and $C_{K2}(\mu, T) > 0$.

Corollary 2.2. For any smooth function $y(t)$ such that $y(0) = 0$

$$\int_0^T y'(t) K_{\mu, \theta} y(t) dt \geq C_{K3} \|y(t)\|_{W_2^{(1+\mu)/2}(0, T)}^2 - C_{K3} \theta^{(1-\mu)/2} \|y(t)\|_{W_2^2(0, T)}^2 \quad (2.7)$$

with some constants $C_{K3}(\mu, T) > 0$ and $C_{K4}(\mu, T) > 0$.

Proof. First of all, we note that

$$\int_0^T y'(t) K_{\mu, \theta} y(t) dt = \int_0^T y'(t) D^\mu y(t) dt - D_{\mu, \theta}(y(t), y'(t)).$$

Then we apply the inequality (2.2) and Lemma 2.1. □

3 Statement of the Problem and Existence Theorem

By technical reasons, we consider a bit more general problem

$$\partial^\nu(k(t, x)u_t(t, x)) - \Delta u(t, x) + \gamma_1 u_t(t, x) + \gamma_2 \partial^\nu u(t, x) = f(t, x), \quad (3.1)$$

$$u(t, x)|_S = 0, \quad (3.2)$$

$$u(0, x) = 0, \quad (3.3)$$

$$u_t(0, x) = 0, \quad x \in \Omega_0^+ = \{x | k(0, x) > 0\}, \quad (3.4)$$

$$u_t(T, x) = 0, \quad x \in \Omega_T^- = \{x | k(T, x) < 0\}. \quad (3.5)$$

A generalized solution to this problem is determined similarly to [2]. We denote

$$\chi_0(x) = k(0, x)u_t(0, x), \quad \chi_T(x) = k(T, x)u_t(T, x).$$

Note that (3.4) and (3.5) imply

$$\text{supp } \chi_0(x) \subseteq \Omega_0^-, \quad \text{supp } \chi_T(x) \subseteq \Omega_T^+, \quad (3.6)$$

where Ω_0^- and Ω_T^+ are defined as in (3.4) and (3.5). Formally applying the operator J^ν to Equation (3.1), we obtain the equality for $t \in [0, T]$

$$k(t, x)u_t(t, x) - J^\nu \Delta u(t, x) + \gamma_1 J^\nu u_t(t, x) + \gamma_2 u(t, x) = J^\nu f(t, x) + \chi_0(x) \quad (3.7)$$

and

$$\chi_T(x) - \chi_0(x) - J^\nu \Delta u(T, x) + \gamma_1 J^\nu u_t(T, x) + \gamma_2 u(T, x) = J^\nu f(T, x). \quad (3.8)$$

A function $u(t, x) \in L_2(Q)$ is called a *generalized solution* to the problem (3.1)–(3.5) if

$$u_t(t, x) \in L_2(Q), \quad u(0, x) = 0,$$

$$\partial^{1-\nu} u(t, x) \in C([0, T]; W_2^{-1}(\Omega)), \quad k(t, x)u_t(t, x) \in C([0, T], W_2^{-1}(\Omega)),$$

$$u(t, x) \in L_2(0, T; \mathring{W}_2^1(\Omega)), \quad J^\nu u(t, x) \in C([0, T]; \mathring{W}_2^1(\Omega));$$

moreover, (3.7), (3.8), and (3.6) hold for some functions $\chi_0(x), \chi_T(x) \in L_2(\Omega)$.

Everywhere below, we set $\mu = 1 - \nu$.

Theorem 3.1. *Assume that $\gamma_1 \geq 0$, $f(t, x) \in W_2^{\mu/2}(0, T; L_2(\Omega))$, and for some $\gamma_0 > 0$*

$$2C_K(\mu, T)\gamma_1 + 2\gamma_2 + k_t(t, x) \geq \gamma_0, \quad (t, x) \in Q, \quad (3.9)$$

where the constant $C_K(\mu, T)$ is taken from the inequality (2.4). Then the problem (3.1)–(3.5) has a generalized solution such that

$$u(t, x) \in W_2^{\frac{1+\mu}{2}}(0, T; \dot{W}_2^1(\Omega)), \quad u_t(t, x) \in W_2^{\mu/2}(0, T; L_2(\Omega)) \quad (3.10)$$

and

$$\|\chi_0\|_{L_2(\Omega)}^2 + \|\chi_T\|_{L_2(\Omega)}^2 + \|u_t\|_{W_2^{\mu/2}(0, T; L_2(\Omega))}^2 + \|u\|_{W_2^{\frac{1+\mu}{2}}(0, T; \dot{W}_2^1(\Omega))}^2 \leq C \|f\|_{W_2^{\mu/2}(0, T; L_2(\Omega))}^2.$$

Moreover, if $D^\mu f(t, x) \in L_2(Q)$, then

$$\int_Q (C_K(\mu, T)\gamma_1 + \gamma_2 + k_t(t, x)/2) u_t^2(t, x) dQ \leq \int_Q D^\mu f(t, x) u_t(t, x) dQ. \quad (3.11)$$

Proof. As in [2], we use the regularization method proposed in [4]. Let $0 < \varepsilon < 1$. We introduce a family of smooth functions $f_\varepsilon(t, x)$ such that

$$f_\varepsilon(0, x) = 0, \quad (3.12)$$

$$\lim_{\varepsilon \rightarrow 0} \|f_\varepsilon(t, x) - f(t, x)\|_{W_2^{\mu/2}(0, T; L_2(\Omega))} = 0, \quad (3.13)$$

$$\lim_{\varepsilon \rightarrow 0} \varepsilon \|f_\varepsilon(t, x)\|_{W_2^1(0, T; L_2(\Omega))}^2 = 0. \quad (3.14)$$

Note that the condition (3.12) is compatible with the condition (3.13) since $\mu < 1$. Then we consider the problem (see [4])

$$-\varepsilon u_{ttt} + (k(t, x)u_t(t, x))_t - D^\mu \Delta u(t, x) + \gamma_1 D^\mu u_t(t, x) + \gamma_2 u_t(t, x) = D^\mu f_\varepsilon(t, x),$$

$$u(t, x)|_S = 0, \quad u(0, x) = 0,$$

$$-\varepsilon u_{tt}(0, x) + k^+(0, x)u_t(0, x) = 0,$$

$$-\varepsilon u_{tt}(T, x) + k^-(T, x)u_t(T, x) = 0.$$

As usual,

$$\eta^+ = \begin{cases} \eta, & \eta > 0, \\ 0, & \eta \leq 0, \end{cases} \quad \eta^- = \begin{cases} \eta, & \eta < 0, \\ 0, & \eta \geq 0. \end{cases}$$

To establish the solvability of the regularized problem, we use the special Galerkin method. Let $\{w_k(x)\}_{k \in \mathbb{N}}$ be the system of eigenfunctions of the problem

$$-\Delta w_k = \lambda_k w_k, \quad w_k(x)|_\Gamma = 0,$$

that are orthonormal in $L_2(\Omega)$. For any $n > 0$ we denote by $E_n \subset L_2(\Omega)$ the subspace of functions spanned by the vectors w_k , $k = \overline{1, n}$. We note that the space E_n is finite-dimensional and, consequently, for some constant $C(n)$ we have the inequality

$$\|h\|_{W_2^1(\Omega)}^2 \leq C(n) \|h\|_{L_2(\Omega)}^2, \quad \forall h \in E_n. \quad (3.15)$$

We denote by P_n the orthogonal projection in $L_2(\Omega)$ onto the space E_n . Assuming that $\theta = \theta(\varepsilon, n)$ (the exact value will be indicated later), we consider the problem

$$\begin{aligned} & -\varepsilon v_{nttt}(t, x) + P_n(k(t, x)v_{nt}(t, x))_t - K_{\mu, \theta} \Delta v_n(t, x) \\ & + \gamma_1 K_{\mu, \theta} v_{nt}(t, x) + \gamma_2 v_{nt} = K_{\mu, \theta} P_n f_\varepsilon(t, x), \end{aligned} \quad (3.16)$$

$$v_n(t, x)|_S = 0, \quad v_n(0, x) = 0, \quad (3.17)$$

$$-\varepsilon v_{ntt}(x, 0) + P_n(k^+(0, x)v_{nt}(0, x)) = 0, \quad (3.18)$$

$$-\varepsilon v_{ntt}(T, x) + P_n(k^-(T, x)v_{nt}(T, x)) = 0, \quad (3.19)$$

where $v_n(t, x) = \sum_{k=1}^n V_k(t)w_k(x)$. If no confusion arises, we omit the superscript n . It is easy to see that the solvability of this system follows from the uniqueness of a solution. Therefore, it suffices to derive a suitable a priori estimate for the solution.

We multiply Equation (3.16) by $2v_t(t, x)$ and integrate over the cylinder Q

$$\begin{aligned} & \int_{\Omega} (|k(0, x)|v_t^2(0, x) + |k(T, x)|v_t^2(T, x)) + 2\varepsilon \int_Q v_{tt}^2 dQ + \int_Q (2\gamma_2 + k_t)v_t^2 dQ \\ & + 2 \int_Q ((K_{\mu, \theta} \nabla v, \nabla v_t) + \gamma_1 v_t K_{\mu, \theta} v_t) dQ = 2 \int_Q v_t K_{\mu, \theta} f_\varepsilon dQ. \end{aligned} \quad (3.20)$$

Let $\delta > 0$. Using (2.4), (2.6), (2.7), (3.15), and (2.3), we get

$$\begin{aligned} & \int_{\Omega} (|k(0, x)|v_t^2(0, x) + |k(T, x)|v_t^2(T, x)) + 2\varepsilon \int_Q v_{tt}^2 dQ + 2\delta C_{K1} \|v_t\|_{W_2^{\mu/2}(0, T; L_2(\Omega))}^2 \\ & + \int_Q (2C_K(\gamma_1 - \delta) + 2\gamma_2 + k_t)v_t^2 dQ + C_{K3} \|\nabla v\|_{W_2^{(1+\mu)/2}(0, T; L_2(\Omega))} \\ & \leq 2\theta^{(1-\mu)/2} \|v_{tt}\|_{L_2(Q)}^2 (\delta C_{K2} + C(n)C_{K4}) + C \|f\|_{W_2^{\mu/2}(0, T; L_2(\Omega))} \|u_t\|_{W_2^{\mu/2}(0, T; L_2(\Omega))}. \end{aligned}$$

Choosing δ and θ sufficiently small, we get the required estimate

$$\begin{aligned} & \int_{\Omega} (|k(0, x)|v_t^2(0, x) + |k(T, x)|v_t^2(T, x)) + \varepsilon \int_Q v_{tt}^2 dQ \\ & + \|v_t\|_{W_2^{\mu/2}(0, T; L_2(\Omega))}^2 + \|\nabla v\|_{W_2^{(1+\mu)/2}(0, T; L_2(\Omega))} \leq C \|f\|_{W_2^{\mu/2}(0, T; L_2(\Omega))}^2. \end{aligned} \quad (3.21)$$

This estimate guarantees that the system (3.16)–(3.19) is uniquely solvable. Note that for all $(t, x) \in Q$ we have the equality

$$(-\varepsilon v_{ntt} + P_n(kv_{nt}))|_0^t = J_{\nu, \theta}(P_n f_\varepsilon + \Delta v_n - \gamma_1 v_{nt}) - \gamma_2 v_n$$

and, due to the conditions (3.17), (3.18), and (3.19),

$$P_n(k^+(T, x)v_{nt}(T, x)) - P_n(k^-(0, x)v_{nt}(0, x)) = J_{\nu, \theta}(P_n f_\varepsilon + \Delta v_n - \gamma_1 v_{nt})(T, x) - \gamma_2 v_n(T, x).$$

Now, we can pass to the limit as $n \rightarrow \infty$. Passing, if necessary, to a subsequence, we assume that for some functions $z_{\varepsilon 0}(x)$, $z_{\varepsilon T}(x)$, $u_{\varepsilon}(t, x)$ the following convergences take place:

$$\begin{aligned} \sqrt{|k(0, x)|}v_{nt}(0, x) &\rightharpoonup z_{\varepsilon 0}(x), && \text{weakly in } L_2(\Omega), \\ \sqrt{|k(T, x)|}v_{nt}(T, x) &\rightharpoonup z_{\varepsilon T}(x), && \text{weakly in } L_2(\Omega), \\ v_n(t, x) &\rightharpoonup u_{\varepsilon}(t, x), && \text{weakly in } W_2^{(1+\mu)/2}(0, T; W_2^1(\Omega)), \\ v_{nt}(t, x) &\rightharpoonup u_{\varepsilon t}(t, x), && \text{weakly in } W_2^{\mu/2}(0, T; L_2(\Omega)) \end{aligned}$$

as $n \rightarrow \infty$. In this case, due to the estimate (3.21),

$$\varepsilon \|u_{\varepsilon tt}\|_{L_2(Q)}^2 + \|J^{\nu}u_{\varepsilon t}\|_{W_2^{\frac{1+\nu}{2}}(0, T; W_2^1(\Omega))} \leq C. \quad (3.22)$$

In particular, $J^{\nu}u_{\varepsilon t}(t, x) \in C([0, T]; W_2^1(\Omega))$ and $J^{\nu}u_{\varepsilon t}(0, x) = 0$. Next, we denote

$$\chi_{0\varepsilon}(x) = -\sqrt{|k^-(0, x)|}z_{\varepsilon 0}(x), \quad \chi_{T\varepsilon}(x) = \sqrt{|k^+(T, x)|}z_{\varepsilon T}(x).$$

It is clear that these functions satisfy the conditions (3.6) and, by virtue of (3.18) and (3.19), the following equalities hold:

$$\begin{aligned} -\varepsilon u_{\varepsilon tt} + k u_{\varepsilon t} - \chi_{0\varepsilon} &= J_{\nu}(f_{\varepsilon} + \Delta u_{\varepsilon} - \gamma_1 u_{\varepsilon t}) - \gamma_2 u_{\varepsilon t}, \\ \chi_{T\varepsilon}(x) - \chi_{0\varepsilon}(x) &= J_{\nu}(f_{\varepsilon} + \Delta u_{\varepsilon} - \gamma_1 u_{\varepsilon t})(T, x) - \gamma_2 u_{\varepsilon t}(T, x). \end{aligned}$$

Now, we can pass to the limit as $\varepsilon \rightarrow 0$. All actions are standard, and we omit them.

Now, assume that $D^{\mu}f(t, x) \in L_2(Q)$. Passing to the limit as $n \rightarrow \infty$ and then as $\varepsilon \rightarrow 0$ in (3.20), we get (3.11). \square

4 Uniqueness Theorems

In a more general case, the study of the uniqueness of such a generalized solution is difficult even in the case of ordinary derivatives (see [4]). We consider a simpler case where the functions $k(0, x)$ and $k(T, x)$ do not change the sign.

Theorem 4.1. *Assume that $k(0, x) \geq 0$, $k(T, x) \geq 0$, $\gamma_1 > 0$, and $\gamma_2 > 0$ is large enough. Then a generalized solution to the problem (3.1)–(3.5) is unique.*

Proof. Consider a solution $u(t, x)$ to the problem (3.1)–(3.5) with $f(t, x) \equiv 0$. Then

$$k(t, x)u_t(t, x) - J^{\nu}\Delta u(t, x) + \gamma_1\partial^{1-\nu}u(t, x) + \gamma_2u(t, x) = \chi_0(x). \quad (4.1)$$

By (3.6) and the inequality $k(0, x) \geq 0$, we have $\chi_0(x) \equiv 0$. Multiplying (4.1) by $2u(t, x)$ and integrating over Q , we have

$$\begin{aligned} \int_{\Omega} k(T, x)u^2(T, x) dx + 2 \int_Q (J^{\nu}\nabla u(t, x), \nabla u(t, x)) dQ \\ + 2\gamma_1 \int_Q u(t, x)\partial^{1-\nu}u(t, x) dQ + \int_Q (2\gamma_2 - k_t(t, x))u^2(t, x) dQ = 0. \end{aligned}$$

Thus, $\int_Q (2\gamma_2 - k_t(t, x))u^2(t, x) dQ \leq 0$ and $u(t, x) \equiv 0$ provided that $2\gamma_2 - k_t(t, x) \geq \bar{\gamma} > 0$. \square

We note that the above proof is similar to that of the uniqueness theorem for the wave equation [5, Theorem 3.1].

Corollary 4.1. *Let the assumptions of Theorem 3.1 hold. Assume that $k(0, x) \geq 0$ and $k(T, x) \geq 0$. Then a generalized solution to the problem (3.1)–(3.5) is unique.*

Proof. Consider a solution $u(t, x)$ to the problem (3.1)–(3.5) with $f(t, x) \equiv 0$. It is clear that for any λ

$$\partial^\nu(k(t, x)u_t(t, x)) - \Delta u(t, x) + \gamma_1 u_t(t, x) + (\gamma_2 + \lambda)\partial^\nu u(t, x) = \lambda\partial^\nu u(t, x). \quad (4.2)$$

We note that $D^\mu \partial^\nu u(t, x) = u_t(t, x) \in L_2(Q)$. By Theorem 4.1, the solution to problem (4.2), (3.2)–(3.5) is unique, and, by Theorem 3.1 (see (3.11))

$$\int_Q (C_K(\mu, T)\gamma_1 + \gamma_2 + \lambda + k_t(t, x)/2)u_t^2(t, x) dQ \leq \int_Q \lambda u_t^2(t, x) dQ.$$

Thus, $\gamma_0 \int_Q u_t^2(t, x) dQ \leq 0$ and $u(t, x) \equiv 0$. □

A simple idea proposed in Theorem 4.1 does not work if $k(0, x) \leq 0$ or $k(T, x) \leq 0$. In such cases, we use convolution mollifiers. The arguments are not difficult, but a rigorous proof requires a lot of tedious technical work.

For any a and b we set

$$Q_{a,b} = (a, b) \times \Omega.$$

First of all, we need to extend the functions $k(t, x)$ and $u(t, x)$ inside $Q_{-T, 2T}$. We set

$$u(t, x) = \begin{cases} 0, & t < 0, \\ u(2T - t, x), & t > T. \end{cases}$$

The function $k(t, x)$ is extended in such a way $k(t, x) \in C^1(\overline{Q_{-T, 2T}})$, and the condition (3.9) is satisfied in $\overline{Q_{-T, 2T}}$. Let $\rho(t) \geq 0$ be a smooth even function such that $\text{supp } \rho(t) \subset (-1, 1)$, and let

$$\text{sgn}(t)\rho'(t) \leq 0, \quad \int_{-1}^1 \rho(t) dt = 1. \quad (4.3)$$

We note that

$$\int_{-1}^1 t\rho'(t) dt = -1. \quad (4.4)$$

For any $\delta > 0$ we denote

$$\rho_\delta(t) = \frac{1}{\delta}\rho(t/\delta), \quad u_\delta(t, x) = \rho_\delta(t) * u(t, x).$$

Lemma 4.1. *Assume that $a < b$, $w(t, x) \in L_2(Q_{a,b})$, $0 < 2\delta < b - a$. Denote $\tilde{Q}_\delta = Q_{a+\delta, b-\delta}$ and $F(\delta, t, x) = k(t, x)\rho_\delta(t) * w(t, x) - \rho_\delta(t) * (k(t, x)w(t, x))$. Then $F_t(\delta, t, x) \in L_2(\tilde{Q}_\delta)$ and $\|F_t(\delta, t, x)\|_{L_2(\tilde{Q}_\delta)} \rightarrow 0$ as $\delta \rightarrow 0$.*

Proof. Let $(t, x) \in \tilde{Q}_\delta$. It is easy to see that

$$F_t(\delta, t, x) = \int_a^b (\rho_\delta(t-s)k'(t, x) - \rho'_\delta(t-s)(k(s, x) - k(t, x)))w(s, x) ds.$$

By the Taylor formula $k(s, x) = k(t, x) + k'(t, x)(s-t) + o(|s-t|)$, it follows that

$$F_t(\delta, t, x)(t, x) = k'(t, x)(\rho_\delta(t) + t\rho'_\delta(t)) * w(t, x) + o(\delta) \int_a^b |\rho'_\delta(t-s)||w(s, x)| ds.$$

Taking into account (4.3) and (4.4), we obtain the required statement. \square

Lemma 4.2. *Assume that $a < b$, $0 < \alpha < 1/2$, $w(t, x) \in W_2^\alpha(a, b; L_2(\Omega))$, $0 < 2\delta < b - a$. Then for some $\beta = \beta(\alpha, a, b) > 0$ and any $t_0 \in [a + \delta, b - \delta]$*

$$\|w_\delta(t_0, x)\|_{L_2(\Omega)}^2 \leq C\delta^{\beta-1} \|w(t, x)\|_{W_2^\alpha(a, b; L_2(\Omega))}^2.$$

Proof. Let $1/q = 1/2 - \alpha$ and $1/p = 1 - 1/q$. By the embedding theorem,

$$|w_\delta(t_0, x)| \leq C \|\rho_\delta(t)\|_{L_p(-\delta, \delta)} \|w(t, x)\|_{W_2^\alpha(a, b)} \leq C\delta^{(1-p)/p} \|w(t, x)\|_{W_2^\alpha(a, b)}.$$

Note that $1 < p < 2$. So, we put $\beta = (2-p)/p$ and get

$$\|w_\delta(t_0, x)\|_{L_2(\Omega)}^2 \leq C\delta^{\beta-1} \|w(t, x)\|_{W_2^\alpha(a, b; L_2(\Omega))}^2.$$

The lemma is proved. \square

Lemma 4.3. *Assume that $a < b$, $w(t, x) \in L_2(Q_{a, b})$, $\tilde{k}(t, x) \in C^1(Q_{a, b})$, $\tilde{k}(a, x) = 0$. $0 < 2\delta < b - a$. Extend by zero $\tilde{k}(t, x)$, $w(t, x)$ for $t < a$ and consider $(\tilde{k}w)_\delta(t, x)$. Then*

$$\|(\tilde{k}w)_\delta(t, x)\|_{L_2(Q_{a-\delta, a})} \leq C \|w(t, x)\|_{L_2(Q_{a, a+\delta})}.$$

Proof. It is easy to see that for $t < a$

$$(\tilde{k}w)_\delta(t, x) = \tilde{k}(t, x)\rho'_\delta * w(t, x) + \int_a^{a+\delta} \rho'_\delta(t-s)(\tilde{k}(s, x) - \tilde{k}(t, x))w(s, x) ds.$$

It remains to note that $\tilde{k}(s, x) - \tilde{k}(t, x) = O(|s-t|)$ and $\int_{-\delta}^\delta |\rho'_\delta(t)| dt \leq C/\delta$. \square

Lemma 4.4. *Assume that $a < b$, $0 < \alpha < 1$, $w(t) \in W_2^\alpha(a, b)$, and (an even continuation) $w(t) = w(2b-t)$ for $t > b$. Then $w(t) \in W_2^\alpha(a, 2b-a)$.*

Proof. Without loss of generality we can assume that $a = 0$ and $b = 1$. According to the definition, we consider the integral

$$J = \int_0^2 dt \int_0^2 \frac{(w(t) - w(\tau))^2}{|t - \tau|^{1+2\alpha}} d\tau.$$

After trivial transformations we get

$$J = 2 \int_0^1 dt \int_0^1 \frac{(w(t) - w(\tau))^2}{|t - \tau|^{1+2\alpha}} d\tau + 2 \int_0^1 dt \int_0^1 \frac{(w(\tau) - w(t))^2}{|2 - \tau - t|^{1+2\alpha}} d\tau.$$

It remains to note that $2 - \tau - t \geq |t - \tau|$ for any $0 < t, \tau < 1$. □

Lemma 4.5. *Assume that $a < b$, $1/2 < \alpha < 1$, $w(t) \in W_2^\alpha(a, b)$, $w(a) = 0$, and*

$$z(t) = \begin{cases} 0, & t < a, \\ w(t) - w(a), & t \in (a, b). \end{cases}$$

Then $z(t) \in W_2^\alpha(2a - b, b)$.

Proof. Without loss of generality we can assume, $a = 0$, $b = 1$. Then

$$\int_{-1}^1 dt \int_{-1}^1 \frac{(z(t) - z(\tau))^2}{|t - \tau|^{1+2\alpha}} d\tau \leq \int_{-1}^1 dt \int_0^1 \frac{(w(t) - w(\tau))^2}{|t - \tau|^{1+2\alpha}} d\tau + C \int_0^1 \frac{(w(t) - w(a))^2}{t^{2\alpha}} dt.$$

Note, that $w(t) - w(a) = J_a^\alpha \partial_a^\alpha w(t)$. Thus the required statement follows from the inequality (3.17) in [3])

$$\int_0^1 \frac{(w(t) - w(a))^2}{t^{2\alpha}} dt \leq C \|w(t)\|_{W_2^\alpha(a,b)}^2.$$

The lemma is proved. □

Theorem 4.2. *Let the condition (3.9) holds. The problem (3.1)-(3.5) has at most one solution satisfying (3.10) provided that the functions $k(0, x)$ and $k(T, x)$ do not change sign.*

Proof. Let $f(t, x) \equiv 0$, and $u(t, x)$ a solution to problem (3.1)-(3.5), satisfying (3.10). We have four slightly different cases to consider.

Case 1. Assume that $k(0, x) \geq 0$ and $k(T, x) \geq 0$. The required statement follows from Corollary 4.1.

Case 2. Assume that $k(0, x) \leq 0$ and $k(T, x) \geq 0$. Assume also that $0 < \delta < T/2$, $v(t, x) = u_\delta(t, x)$, and $F(\delta, t, x) = k(t, x)v_t(t, x) - \rho_\delta(t) * (k(t, x)u_t(t, x))$. Then for $(t, x) \in Q_{\delta, T-\delta}$

$$k(t, x)v_t(t, x) - J^\nu \Delta v(t, x) + \gamma_1 J^\nu v_t(t, x) + \gamma_2 v(t, x) = \chi_0 + F(\delta, t, x),$$

$$(k(t, x)v_t(t, x))_t - D^\mu \Delta v(t, x) + \gamma_1 D^\mu v_t(t, x) + \gamma_2 v_t(t, x) = F_t(\delta, t, x).$$

As above, multiplying by $2v_t(t, x)$ and integrating over $Q_{\delta, T-\delta}$, we get

$$\begin{aligned} & \int_{\Omega} (k(T, x)v_t^2(T - \delta, x) - k(0, x)v_t^2(\delta, x)) d\Omega + \gamma_0 \int_{Q_{0, T-\delta}} v_t^2 dQ \\ & \leq \int_{\Omega} (k(T, x) - k(T - \delta, x))v_t^2(T - \delta, x) d\Omega - \int_{\Omega} (k(0, x) - k(\delta, x))v_t^2(\delta, x) d\Omega \\ & + 2 \int_{Q_{0, \delta}} v_t(-D^\mu \Delta v(t, x) + \gamma_1 D^\mu v_t(t, x) + \gamma_2 v_t(t, x)) dQ \end{aligned}$$

$$+ \int_{Q_{\delta, T-\delta}} |F_t(\delta, t, x)v_t(t, x)| dQ = J_1 + J_2 + J_3.$$

By Lemmas 4.1 and 4.2, we have $|J_1| = O(\delta^\beta)$ and $|J_3| = o(1)$ as $\delta \rightarrow 0$. By the inequality (2.3),

$$|J_2| \leq C(\|v_t\|_{W_2^{\mu/2}(0, \delta; L_2(\Omega))}^2 + \|v\|_{W_2^{(1+\mu)/2}(0, \delta; \dot{W}_2^1(\Omega))}^2).$$

Passing to the limit as $\delta \rightarrow 0$, we get $u(t, x) \equiv 0$.

Case 3. Assume that $k(0, x) \leq 0$ and $k(T, x) \leq 0$. First of all, we extend the equality (3.7) inside $Q_{0, 2T}$. By Lemma 4.4, $u(t, x) \in W_2^{\frac{1+\mu}{2}}(0, 2T; \dot{W}_2^1(\Omega))$ and (see [3, Theorem 11.6]) $u_t(t, x) \in W_2^{\mu/2}(0, 2T; L_2(\Omega))$ since $\mu < 1$. We set

$$\begin{aligned} H(t, x) &= -J^\nu \Delta u(t, x) + \gamma_1 J^\nu u_t(t, x) + \gamma_2 u(t, x), \\ \tilde{H}(t, x) &= H(t, x) + H(2T - t, x). \end{aligned}$$

By (3.6) and (3.8), $\chi_1(x) \equiv 0$ and $H(T, x) = \chi_0$. Next, let $T < t < 2T$. By definition,

$$k(t, x)u_t(t, x) = -k(2T - t, x)u_t(2T - t, x) + \tilde{k}(t, x)u_t(2T - t, x),$$

where $\tilde{k}(t, x) = k(2T - t, x) - k(t, x)$. So, for $T < t < 2T$

$$k(t, x)u_t(t, x) + H(t, x) = \tilde{k}(t, x)u_t(2T - t, x) + \tilde{H}(t, x) - \chi_0(x).$$

Finally, we have

$$k(t, x)u_t(t, x) - J^\nu \Delta u(t, x) + \gamma_1 J^\nu u_t(t, x) + \gamma_2 u(t, x) - \chi_0(x) = G(t, x),$$

where

$$G(t, x) = \begin{cases} 0, & t < T, \\ \tilde{k}(t, x)u_t(2T - t, x) + \tilde{H}(t, x) - 2H(T, x), & t > T. \end{cases}$$

Let $0 < \delta < T/2$, $v(t, x) = u_\delta(t, x)$. Then for $(t, x) \in Q_{\delta, T}$

$$(k(t, x)v_t(t, x))_t - D^\mu \Delta v(t, x) + \gamma_1 D^\mu v_t(t, x) + \gamma_2 v_t(t, x) = F_t(\delta, t, x) + G_{\delta t}(t, x).$$

As above, we multiply this equation by $2v_t(t, x)$ and integrate over $Q_{\delta, T}$. Since $v_t(T, x) = 0$, we get

$$\begin{aligned} & \int_{\Omega} -k(0, x)v_t^2(\delta, x) d\Omega + \gamma_0 \int_{Q_{0, T}} v_t^2 dQ \leq \int_{\Omega} (k(\delta, x) - k(0, x))v_t^2(\delta, x) d\Omega \\ & + 2 \int_{Q_{0, \delta}} v_t(-D^\mu \Delta v(t, x) + \gamma_1 D^\mu v_t(t, x) + \gamma_2 v_t(t, x)) dQ \\ & + \int_{Q_{\delta, T}} |F_t(\delta, t, x)v_t(t, x)| dQ + \int_{Q_{T-\delta, T}} |G_{\delta t}(t, x)v_t(t, x)| dQ. \end{aligned}$$

Compared to Case 2, we need to estimate the term

$$J = \int_{Q_{T-\delta,T}} |G_{\delta t}(t, x)v_t(t, x)| dQ.$$

Let $H_1(t, x) = J^\nu \Delta u(t, x) - J^\nu \Delta u(T, x)$, $H_2(t, x) = J^\nu u_t(t, x) - J^\nu u_t(T, x)$, $H_3(t, x) = u_t(t, x) - u_t(T, x)$,

$$G_0(t, x) = \begin{cases} 0, & t < T, \\ \tilde{k}(t, x)u_t(2T - t, x), & t > T, \end{cases}$$

and for $k = 1, 2, 3$

$$G_k(t, x) = \begin{cases} 0, & t < T, \\ H_k(t, x) + H_k(2T - t, x), & t > T. \end{cases}$$

Then we can write

$$|J| \leq \sum_{k=0}^3 \int_{Q_{T-\delta,T}} |G_{k\delta t}(t, x)v_t(t, x)| dQ = J_0 + J_1 + J_2 + J_3.$$

It is obvious that

$$|J_3| \leq C \|u_t\|_{L_2(Q_{T-2\delta,T})}^2.$$

By Lemma 4.3,

$$|J_0| \leq C \|u_t\|_{L_2(Q_{T-\delta,T})} \|u_t\|_{L_2(Q_{T-2\delta,T})}.$$

We consider the term $G_2(t, x)$. By Lemmas 4.4 and 4.5, $G_2(t, x) \in W_2^{\nu+\mu/2}(T-\delta, T; L_2(\Omega))$, and, by (2.3),

$$|J_2| \leq C \|v_t\|_{W_2^{\mu/2}(T-\delta, T; L_2(\Omega))} \|J_{T-\delta}^{\mu/2} G_{2\delta t}(t, x)\|_{L_2(Q_{T-2\delta,T})} \leq C \|u_t\|_{W_2^{\mu/2}(T-\delta, T; L_2(\Omega))}.$$

Finally, we consider the term $G_1(t, x)$. By Lemmas 4.4 and 4.5, $G_{1\delta t}(t, x) \in W_2^{\nu/2}(0, T; W_2^{-1}(\Omega))$ and $G_{1\delta t}(t, x) = 0$ if $t < T - \delta$. Using (2.3), we get

$$|J_1| \leq C \|J_T^{\nu/2} v\|_{L_2(T-\delta, T; \dot{W}_2^1(\Omega))} \|G_{1\delta t}(t, x)\|_{W_2^{\nu/2}(0, T; W_2^{-1}(\Omega))} \leq C \|v\|_{W_2^{(1+\mu/2)}(T-\delta, T; \dot{W}_2^1(\Omega))}.$$

Thus, we have established all the required estimates and can pass to the limit as $\delta \rightarrow 0$.

Case 4. Assume that $k(0, x) \geq 0$ and $k(T, x) \leq 0$. According to (3.6), $\chi_0 \equiv 0$ and

$$k(t, x)u_t(t, x) - J_{-T/2}^\nu \Delta u(t, x) + \gamma_1 J_{-T/2}^\nu u_t(t, x) + \gamma_2 u(t, x) = G(t, x).$$

Now, we multiply this equation by $2v_t(t, x)$ and integrate over the cylinder $Q_{-T/2, T}$. The rest of the proof is the same as in Case 3. \square

Declarations

Data availability This manuscript has no associated data.

Ethical Conduct Not applicable.

Conflicts of interest The author declare that there is no conflict of interest.

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