

Generalized Eigenvalue-Counting Estimates for the Anderson Model

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Abstract We generalize Minami’s estimate for the Anderson model and its extensions to n eigenvalues, allowing for n arbitrary intervals and arbitrary single-site probability measures with no atoms. As an application, we derive new results about the multiplicity of eigenvalues and Mott’s formula for the ac-conductivity when the single site probability distribution is Hölder continuous.

Keywords Anderson localization · Anderson model · Eigenvalue statistics · Poisson eigenvalue statistics · Spectral averaging · Minami estimate · Mott formula · Multiplicity of eigenvalues

1 Introduction

We consider the generalized Anderson model given by the random Hamiltonian

$$H_\omega = H_0 + V_\omega \quad \text{on } \ell^2(\mathbb{Z}^d), \quad (1.1)$$

where H_0 is a self-adjoint operator and V_ω is the random potential given by $V_\omega(j) = \omega_j$. Here $\omega = \{\omega_j\}_{j \in \mathbb{Z}^d}$ is a family of independent random variables; μ_j will denote the proba-

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bility distribution of the random variable ω_j . In this article we always assume that each μ_j has no atoms. We write \mathbb{E}_{ω_j} for the expectation with respect to the random variable ω_j , and write $\mathbb{E} = \mathbb{E}_\omega$ for the joint expectation. We also set $\omega_k^\perp = \{\omega_j\}_{j \in \mathbb{Z}^d \setminus \{k\}}$ and let $\mathbb{E}_{\omega_k^\perp}$ denote the corresponding expectation.

Restrictions of H_ω to finite volumes $\Lambda \subset \mathbb{Z}^d$ are denoted by $H_{\omega,\Lambda}$, a self-adjoint operator of the form

$$H_{\omega,\Lambda} = H_{0,\Lambda} + V_{\omega,\Lambda} \quad \text{on } \ell^2(\Lambda), \tag{1.2}$$

with $H_{0,\Lambda}$ a self-adjoint restriction of H_0 to the finite-dimensional Hilbert space $\ell^2(\Lambda)$ and $V_{\omega,\Lambda}(j) = \omega_j$ for $j \in \Lambda$. (The results discussed in this article are not sensitive to the choice of $H_{0,\Lambda}$.) Given a Borel set $J \subset \mathbb{R}$, we write $P_\omega^{(\Lambda)}(J) = P_{H_\omega}^{(\Lambda)}(J) = \chi_J(H_{\omega,\Lambda})$ for the associated spectral projection.

Minami [25] estimated the probability that $H_{\omega,\Lambda}$ has at least two eigenvalues in an interval I . Assuming that all μ_j have bounded densities ρ_j , Minami proved that

$$2\mathbb{P} \left\{ \text{tr } P_{H_\omega}^{(\Lambda)}(I) \geq 2 \right\} \leq \mathbb{E} \left\{ \left(\text{tr } P_{H_\omega}^{(\Lambda)}(I) \right)^2 - \text{tr } P_{H_\omega}^{(\Lambda)}(I) \right\} \leq (\pi \rho_\infty^{(\Lambda)} |I| |\Lambda|)^2, \tag{1.3}$$

where $\rho_\infty^{(\Lambda)} := \max_{j \in \Lambda} \|\rho_j\|_\infty$. Minami’s proof required H_0 to have real matrix elements, i.e., $\langle \delta_j, H_0 \delta_k \rangle \in \mathbb{R}$ for all j, k . This restriction was recently removed by Bellissard, Hislop and Stolz [5] and by Graf and Vaghi [18]. They also estimated the probability that $H_{\omega,\Lambda}$ has at least n eigenvalues in I for all $n \in \mathbb{N}$, assuming, as Minami, that all μ_j have bounded densities ρ_j .

Minami’s estimate has important consequences for the physical behavior of the Anderson model in the localized (insulator) regime. It is the crucial ingredient in Minami’s proof of the absence of eigenvalue repulsion, showing that the properly rescaled eigenvalues behave according to a Poisson process [25]. (See [20, 24, 26, 30, 31] for further developments.) It was shown to imply simplicity of eigenvalues by Klein and Molchanov [23]. It is an important ingredient in the derivation of a rigorous form of Mott’s formula for the ac-conductivity by Klein, Lenoble and Müller [22].

In [7] we introduced a new approach to Minami’s estimate. The crucial step in Minami’s proof, namely [25, Lemma 2], estimates the average of a determinant whose entries are matrix elements of the imaginary part of the resolvent; the proofs in [5, 18] have similar steps. In contrast, our method only averages spectral projections, which allowed us to finally prove a Minami estimate for the continuum Anderson Hamiltonian. As a consequence, we obtained Poisson eigenvalue statistics and simplicity of eigenvalues for the continuum Anderson Hamiltonian.

The new approach, in addition to providing a simple and transparent proof of Minami’s estimate for the Anderson model, also allows for arbitrary single-site probability measures with no atoms. Given a probability measure μ , we let $S_\mu(s) := \sup_{a \in \mathbb{R}} \mu([a, a + s])$, the concentration function of μ , and set

$$Q_\mu(s) := \begin{cases} \|\rho\|_\infty s & \text{if } \mu \text{ has a bounded density } \rho, \\ 8S_\mu(s) & \text{otherwise.} \end{cases}$$

(Note that the measure μ has no atoms if and only if $\lim_{s \downarrow 0} Q_\mu(s) = 0$.) For the generalized Anderson model H_ω as in (1.1), we let $Q_j = Q_{\mu_j}$ and set $Q_\Lambda(s) := \max_{j \in \Lambda} Q_j(s)$. In [7, Theorem 3.3] we obtained the following extension of (1.3):

$$\mathbb{E} \left\{ \left(\text{tr } P_\omega^{(\Lambda)}(I) \right) \left(\text{tr } P_\omega^{(\Lambda)}(I) - 1 \right) \right\} \leq (Q_\Lambda(|I|) |\Lambda|)^2. \tag{1.4}$$

(In [7] the proof of (1.4) is given for single-site probability measures with compact support, but (1.4) follows for arbitrary single-site probability measures by Lemma B.1. Note that the proof is valid for the generalized Anderson model.)

In this article we generalize Minami’s estimate and its extensions to n eigenvalues, allowing for n arbitrary intervals and arbitrary single-site probability measures with no atoms. We also give applications of (1.4), deriving new results about the multiplicity of eigenvalues and Mott’s formula for the ac-conductivity when the single site probability distribution is Hölder continuous.

The paper is organized as follows. In Sect. 2 we state our main results, namely our generalized eigenvalue-counting estimates. In Sect. 3, we consider the Anderson model with a Hölder continuous probability distribution, for which we extend previous results on the multiplicity of the spectrum and Mott’s formula for energies in the region of Anderson localization. In Sect. 4 we prove the results stated in Sect. 2. In Appendix A we provide proofs for the fundamental spectral averaging estimate (2.2). In Appendix B we prove an approximation lemma to go from probability measures with compact support to arbitrary probability measures.

2 Eigenvalue Counting Inequalities

In this section we state our main results. The proofs will be given in Sect. 4.

Spectral averaging is the basic ingredient for proving eigenvalue-counting inequalities for the generalized Anderson model. Consider the random self-adjoint operator

$$H_\omega = H_0 + \omega \Pi_\varphi \quad \text{on } \mathcal{H}, \tag{2.1}$$

where H_0 is a self-adjoint operator on the Hilbert space \mathcal{H} , $\varphi \in \mathcal{H}$ with $\|\varphi\| = 1$, and ω is a random variable with a non-degenerate probability distribution μ . By Π_φ we denote the orthogonal projection onto $\mathbb{C}\varphi$, the one-dimensional subspace spanned by φ . Let $P_\omega(J) = \chi_J(H_\omega)$ for a Borel set $J \subset \mathbb{R}$. There is a fundamental spectral averaging estimate: for all bounded intervals $I \subset \mathbb{R}$ we have

$$\mathbb{E}_\omega \{ \langle \varphi, P_\omega(I)\varphi \rangle \} := \int d\mu(\omega) \langle \varphi, P_\omega(I)\varphi \rangle \leq Q_\mu(|I|). \tag{2.2}$$

In full generality, i.e., μ arbitrary with $Q_\mu(s) = 8S_\mu(s)$, this is a recent result of Combes, Hislop and Klopp [10, (3.16)]. (We present a proof in Appendix A for completeness.) If μ has a bounded density ρ , (2.2) was known to hold with $Q_\mu(s) = \|\rho\|_\infty s$ (e.g. [6, 9, 14, 21, 32]; a simple proof is given in Appendix A). If μ is Hölder continuous, i.e., $S_\mu(s) \leq Cs^\alpha$ with $\alpha \in]0, 1[$, (2.2) was known with $Q_\mu(s) = C(1 - \alpha)^{-1}s^\alpha$ [6, Theorem 6.2]. We will thus always assume that Q_μ is as in (1.4), although all we will require of Q_μ is the validity of (2.2). (The estimate (2.2) is useful when the measure μ has no atoms, i.e., $\lim_{s \downarrow 0} Q_\mu(s) = 0$, which is always assumed in his paper.)

Now let H_ω be the generalized Anderson model. Note that we can rewrite the finite volume operator given in (1.2) as

$$H_{\omega,\Lambda} = H_{0,\Lambda} + \sum_{j \in \Lambda} \omega_j \Pi_j \quad \text{on } \ell^2(\Lambda), \quad \text{with } \Pi_j = \Pi_{\delta_j}. \tag{2.3}$$

The first eigenvalue-counting inequality for H_ω is the Wegner estimate [32], which measures the probability that $H_{\omega,\Lambda}$ has an eigenvalue in an interval I :

$$\mathbb{P} \left\{ \text{tr } P_{H_\omega}^{(\Lambda)}(I) \geq 1 \right\} \leq \mathbb{E} \left\{ \text{tr } P_{H_\omega}^{(\Lambda)}(I) \right\} \leq Q_\Lambda (|I|) |\Lambda|. \tag{2.4}$$

The Wegner estimate is an immediate consequence of (2.2):

$$\mathbb{E} \left\{ \text{tr } P_{H_\omega}^{(\Lambda)}(I) \right\} = \sum_{j \in \Lambda} \mathbb{E}_{\omega_j^\perp} \left\{ \mathbb{E}_{\omega_j} \left\{ \langle \delta_j, P_{H_\omega}^{(\Lambda)}(I) \delta_j \rangle \right\} \right\} \leq Q_\Lambda (|I|) |\Lambda|. \tag{2.5}$$

The second eigenvalue-counting inequality is the Minami estimate (1.4). It is generalized to two intervals in the following theorem.

Theorem 2.1 *Let H_ω be the generalized Anderson model, and fix a finite volume $\Lambda \subset \mathbb{Z}^d$. For any two bounded intervals I_1, I_2 we have*

$$\begin{aligned} & \mathbb{E} \left\{ (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2)) - \min \{ \text{tr } P_\omega^{(\Lambda)}(I_1), \text{tr } P_\omega^{(\Lambda)}(I_2) \} \right\} \\ & \leq 2Q_\Lambda (|I_1|) Q_\Lambda (|I_2|) |\Lambda|^2. \end{aligned} \tag{2.6}$$

If $I_1 \subset I_2$, we have

$$\mathbb{E} \left\{ (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2) - 1) \right\} \leq Q_\Lambda (|I_1|) Q_\Lambda (|I_2|) |\Lambda|^2. \tag{2.7}$$

Remark 2.2

- (i) The estimate (1.4), proved in [7, Theorem 3.3], is a particular case of (2.7).
- (ii) Note that

$$(\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2)) - \min \{ \text{tr } P_\omega^{(\Lambda)}(I_1), \text{tr } P_\omega^{(\Lambda)}(I_2) \} \geq 0. \tag{2.8}$$

- (iii) The intervals I_1 and I_2 in (2.6) may be disjoint. In this case the usual Minami’s estimate (1.3) would yield the bound $\pi^2(\rho_\infty^{(\Lambda)})^2|I|^2|\Lambda|^2$, with an interval $I \supset I_1 \cup I_2$, while under the same hypotheses the estimate (2.6) gives $2(\rho_\infty^{(\Lambda)})^2|I_1||I_2||\Lambda|^2$.

We now turn to the general case of n arbitrary intervals, extending the results of [5, 18]. Given $n \in \mathbb{N}$, we let \mathcal{S}_n denote the group of all permutations of $\{1, 2, \dots, n\}$, and recall that $|\mathcal{S}_n| = n!$. Given a finite volume $\Lambda \subset \mathbb{Z}^d$ and bounded intervals I_1, \dots, I_n (not necessarily distinct), we pick $\sigma_\omega = \sigma_\omega^{(\Lambda)}(I_1, \dots, I_n) \in \mathcal{S}_n$ such that

$$\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)}) \leq \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(2)}) \leq \dots \leq \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(n)}), \tag{2.9}$$

in which case we have

$$(\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)})) (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(2)} - 1)) \dots (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(n)} - (n - 1))) \geq 0. \tag{2.10}$$

To avoid ambiguity, we select σ_ω uniquely by requiring $\sigma_\omega(i) < \sigma_\omega(j)$ if $i < j$ and $\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(i)}) = \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(j)})$. (Note that the product in the left hand side of (2.10) is independent of the choice of $\sigma_\omega \in \mathcal{S}_n$ satisfying (2.9).) We let $\mathcal{S}_n(I_1, \dots, I_n)$ be the collection permutations $\sigma \in \mathcal{S}_n$ such that $\sigma = \sigma_\omega$ for some ω , and let $M(I_1, \dots, I_n)$ denote the cardinality of $\mathcal{S}_n(I_1, \dots, I_n)$. Note that $1 \leq M(I_1, \dots, I_n) \leq n!$, with $M(I_1, \dots, I_n) = 1$ if $I_1 \subset I_2 \subset \dots \subset I_n$.

Theorem 2.3 *Let H_ω be the generalized Anderson model, fix a finite volume $\Lambda \subset \mathbb{Z}^d$, let $n \in \mathbb{N}$, and consider n bounded intervals I_1, \dots, I_n (not necessarily distinct). Then, setting $\sigma_\omega = \sigma_\omega^{(\Lambda)}(I_1, \dots, I_n)$, we have*

$$\begin{aligned} &\mathbb{E} \left\{ \left(\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)}) \right) \left(\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(2)}) - 1 \right) \cdots \left(\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(n)}) - (n - 1) \right) \right\} \\ &\leq M(I_1, \dots, I_n) \left(\prod_{k=1}^n Q^{(\Lambda)}(|I_k|) \right) |\Lambda|^n. \end{aligned} \tag{2.11}$$

In the special case when $I_1 \subset I_2 \subset \dots \subset I_n$, we have

$$\begin{aligned} &\mathbb{E} \left\{ \left(\text{tr } P_\omega^{(\Lambda)}(I_1) \right) \left(\text{tr } P_\omega^{(\Lambda)}(I_2) - 1 \right) \cdots \left(\text{tr } P_\omega^{(\Lambda)}(I_n) - (n - 1) \right) \right\} \\ &\leq \left(\prod_{k=1}^n Q^{(\Lambda)}(|I_k|) \right) |\Lambda|^n. \end{aligned} \tag{2.12}$$

In particular, for any bounded interval I we have

$$\mathbb{E} \left\{ \left(\text{tr } P_\omega^{(\Lambda)}(I) \right) \left(\text{tr } P_\omega^{(\Lambda)}(I) - 1 \right) \cdots \left(\text{tr } P_\omega^{(\Lambda)}(I) - (n - 1) \right) \right\} \leq \left(Q^{(\Lambda)}(|I|) |\Lambda| \right)^n. \tag{2.13}$$

As a corollary, we get probabilistic estimates on the number of eigenvalues of $H_{\omega,\Lambda}$ in intervals.

Corollary 2.4 *Let H_ω be the generalized Anderson model, and fix a finite volume $\Lambda \subset \mathbb{Z}^d$. For all $n \in \mathbb{N}$ and I a bounded interval, we have*

$$\mathbb{P} \left\{ \text{tr } P_\omega^{(\Lambda)}(I) \geq n \right\} \leq \frac{1}{n!} \left(Q^{(\Lambda)}(|I|) |\Lambda| \right)^n. \tag{2.14}$$

Furthermore, for all bounded intervals I_1, \dots, I_n we get

$$\begin{aligned} &\mathbb{P} \left\{ \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)}) \geq 1, \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(2)}) \geq 2, \dots, \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(n)}) \geq n \right\} \\ &\leq M(I_1, \dots, I_n) \left(\prod_{k=1}^n Q^{(\Lambda)}(|I_k|) \right) |\Lambda|^n, \end{aligned} \tag{2.15}$$

and, in the special case when $I_1 \subset I_2 \subset \dots \subset I_n$, we have

$$\mathbb{P} \left\{ \text{tr } P_\omega^{(\Lambda)}(I_1) \geq 1, \text{tr } P_\omega^{(\Lambda)}(I_2) \geq 2, \dots, \text{tr } P_\omega^{(\Lambda)}(I_n) \geq n \right\} \leq \left(\prod_{k=1}^n Q^{(\Lambda)}(|I_k|) \right) |\Lambda|^n. \tag{2.16}$$

Remark 2.5 Given bounded intervals I_1 and I_2 , let $d(I_1, I_2)$ denote the distance between the two intervals. It follows from (2.16) that

$$\begin{aligned} &\mathbb{P} \left\{ \text{tr } P_\omega^{(\Lambda)}(I_1) \geq 1 \text{ and } \text{tr } P_\omega^{(\Lambda)}(I_2) \geq 1 \right\} \\ &\leq \left(\min \left\{ Q^{(\Lambda)}(|I_1|), Q^{(\Lambda)}(|I_2|) \right\} Q^{(\Lambda)}(d(I_1, I_2) + |I_1| + |I_2|) \right) |\Lambda|^2. \end{aligned} \tag{2.17}$$

Note that (2.17) does not generally hold if the right hand side is replaced by the more desirable $C Q^{(\Lambda)}(|I_1|) Q^{(\Lambda)}(|I_2|) |\Lambda|^2$, see [4].

3 Applications to Hölder Continuous Distributions

The (standard) Anderson model is given by H_ω as in (1.1), with $H_0 = -\Delta$, the centered discrete Laplacian, and $\omega = \{\omega_j\}_{j \in \mathbb{Z}^d}$ a family of independent identically distributed random variables with joint probability distribution μ , which we assume to have no atoms. Localization for the Anderson model has been well studied, mostly for μ with a bounded density ρ , cf. [1, 2, 11–14, 29] and many others, as well as for probability distributions μ that are Hölder continuous [3, 6, 12, 15, 19], i.e., $Q_\mu(s) \leq Us^\alpha$ for s small, for some constants U and $\alpha \in]0, 1[$. If the probability distribution μ has a bounded density, Minami’s estimate (1.3) was a crucial ingredient in Klein and Molchanov’s proof of simplicity of eigenvalues [23] and in Klein, Lenoble and Müller derivation of a rigorous form of Mott’s formula for the ac-conductivity [22]. In this section we show that with (1.4) these proofs extend to the case when μ is only Hölder continuous.

3.1 Multiplicity of the Spectrum

Let H_ω be a generalized Anderson model as in (1.1), let $\alpha \in]0, 1[$, and assume that the probability distributions μ_j are uniformly α -Hölder continuous, i.e., there is a constant U and $s_0 > 0$ such that

$$\sup_{j \in \mathbb{Z}^d} Q_j(s) \leq Us^\alpha \quad \text{for all } s \in [0, s_0]. \tag{3.1}$$

In this case we say that H_ω is an α -Hölder continuous generalized Anderson model.

We say that the generalized Anderson model H_ω exhibits Anderson localization in some interval I if, with probability one, H_ω has pure point spectrum in I and the corresponding eigenfunctions decay exponentially. Given $x \geq 0$, we let $[x]$ denote the integer part of x . Following Klein and Molchanov [23], we prove the following result.

Theorem 3.1 *Let H_ω be an α -Hölder continuous generalized Anderson model. Suppose H_ω exhibits Anderson localization in some interval I . Then, with probability one, every eigenvalue of H_ω in I has multiplicity $\leq [\alpha^{-1}]$. In particular, if $\alpha > \frac{1}{2}$, with probability one every eigenvalue of H_ω in I is simple.*

Remark 3.2 For the standard Anderson model, where $\mu_j = \mu$ for all $j \in \mathbb{Z}^d$, this theorem was originally proved by Simon [28] when μ has a bounded density (i.e., $\alpha = 1$). Our proof is based on a simple proof later provided by Klein and Molchanov [23], based on Minami’s estimate (1.3). For singular measures (i.e., $\alpha < 1$), the best previously known result for the standard Anderson model is the finite multiplicity of the eigenvalues [8, 17]; uniform boundedness of the multiplicity was not previously known. Thus Theorem 3.1 improves on both results.

Proof of Theorem 3.1 We proceed as in [23]. We call $\varphi \in \ell^2(\mathbb{Z}^d)$ α -fast decaying if it has β -decay, that is, $|\varphi(x)| \leq C_\varphi (1 + |x|)^{-\beta}$ for some $C_\varphi < \infty$, with

$$\beta > \left(\frac{1}{2} + \left(\alpha - ([\alpha^{-1}] + 1)^{-1} \right)^{-1} \right) d. \tag{3.2}$$

To prove the theorem, we will show that, with probability one, an α -Hölder continuous generalized Anderson model H_ω cannot have an eigenvalue with $[\alpha^{-1}] + 1$ linearly independent α -fast decaying eigenfunctions.

We set $N = [\alpha^{-1}] + 1$, so that $N\alpha > 1$. For a given open interval I we pick

$$q > \frac{Nd}{N\alpha - 1} = \left(\alpha - ([\alpha^{-1}] + 1)^{-1}\right)^{-1} d. \tag{3.3}$$

Given a scale $L > 0$, we let Λ_L denote the cube of side L centered at 0, and cover I by $2([\frac{L^q}{2}|I|] + 1) \leq L^q|I| + 2$ intervals of length $2L^{-q}$, in such a way that any subinterval $J \subset I$ with length $|J| \leq L^{-q}$ will be contained in one of these intervals. We consider the event $\mathcal{B}_{L,I,q}$, which occurs if there exists an interval $J \subset I$ with $|J| \leq L^{-q}$ such that $\text{tr } P_\omega^{(\Lambda_L)}(J) \geq N$. Its probability can be estimated, using (2.14) and (3.1), by

$$\mathbb{P}\{\mathcal{B}_{L,I,q}\} \leq \frac{1}{N!} (L^q|I| + 2) (U (2L^{-q})^\alpha L^d)^N \leq (|I| + 1) \frac{(2^\alpha U)^N}{N!} L^{-(N\alpha - 1)q + Nd}. \tag{3.4}$$

In view of (3.3), taking scales $L_k = 2^k$, it follows from the Borel-Cantelli Lemma that, with probability one, the event $\mathcal{B}_{L_k,I,q}$ eventually does not occur.

Now, suppose that for some ω there exists $E \in I$ which is an eigenvalue of H_ω with N linearly independent α -fast decaying eigenfunctions, so they all have β -decay for some β as in (3.2). It follows, as in [23, Lemma 1], that for L large enough the finite volume operator H_{ω,Λ_L} has at least N eigenvalues in the interval $J_{E,L} = [E - \varepsilon_L, E + \varepsilon_L]$, where $\varepsilon_L = CL^{-\beta + \frac{d}{2}}$ for an appropriate constant C independent of L . In view of (3.2), we can pick q , satisfying (3.3), such that $\beta - \frac{d}{2} > q$, and hence $\varepsilon_L < L^{-q}$ for all large L . But with probability one this is impossible since the event $\mathcal{B}_{L_k,I,q}$ does not occur for large L_k . \square

3.2 Generalized Mott’s Formula

Let $\alpha \in]0, 1[$, and consider the Anderson model H_ω with a single-site probability distribution μ of compact support and uniformly α -Hölder continuous:

$$Q_\mu(s) \leq Us^\alpha \quad \text{for all } s \in [0, s_0], \tag{3.5}$$

where U and $s_0 > 0$ are constants. The fractional moment method can be applied to such measures, leading to exponential decay of the expectation of some fractional power of the Green’s function [3, 19]. We may then define the region of complete localization Ξ^{CL} , introduced in [16, 17], as in [22, Definition 2.1]. However, [22, (4.1), (4.3) and (4.4)] have not been derived from the fractional moment method for μ with compact support satisfying only the condition (3.5). ([3, 19, Appendix A] assumes that μ has a bounded density in the derivation of such estimates.) But in this region of complete localization we can always perform a multiscale analysis as in [15] with only hypothesis (3.5), and get the estimates [22, (4.1), (4.3)] with sub-exponential decay [15, 17], and hence conclude that [22, Assumption 3.1 and (4.4)] are satisfied. Thus, given a Fermi energy $E_F \in \Xi^{\text{CL}}$, the analysis in [22] applies and we may define the *average in phase conductivity* $\bar{\sigma}_{E_F}^{\text{in}}(\nu)$ as in [22, (2.17)]. We have the following extension of [22, Theorem 2.3]. Note that we get $\bar{\sigma}_{E_F}^{\text{in}}(\nu) \leq C\nu^{2\alpha}(\log \frac{1}{\nu})^{d+2}$ for small ν , consistent with $C\nu^2(\log \frac{1}{\nu})^{d+2}$ for $\alpha = 1$ as in [22].

Theorem 3.3 *Given $\alpha \in]0, 1[$, let H_ω be an Anderson model with a single-site probability distribution μ of compact support and uniformly α -Hölder continuous as in (3.5). Consider a Fermi energy in its region of complete localization: $E_F \in \Xi^{\text{CL}}$. Then*

$$\limsup_{\nu \downarrow 0} \frac{\bar{\sigma}_{E_F}^{\text{in}}(\nu)}{\nu^{2\alpha}(\log \frac{1}{\nu})^{d+2}} \leq B\ell_{E_F}^{d+2}, \tag{3.6}$$

where ℓ_{E_F} is given in [22, (2.3)], and the constant B depends only on d, U and α .

Proof The proof of [22, Theorem 2.3] applies, with modifications due to the use of (2.4) and (1.4) with $Q(s) = Q_\mu(s)$ as in (3.5). The modifications are as follows (we use the notation of [22]):

- (1) We systematically use (2.2) instead of [22, (4.5)]. In particular, [22, (4.10)] becomes

$$\langle \langle Y_{E_F}, \chi_B(\mathcal{H}_L) Y_{E_F} \rangle \rangle \leq W_\beta Q(|B|)^\beta \quad \text{for all Borel sets } B \subset \mathbb{R}. \tag{3.7}$$

To derive this estimate, we use the sub-exponential decay of the Fermi projection given in [17, Theorem 3], i.e., we use [22, (4.1)] but with sub-exponential decay. This can be done because we only use summability of this decay. As a consequence, [22, (4.6)] becomes

$$\Psi_{E_F}(B_+ \times B_-) \leq W_\beta (\min \{Q(|B_+|), Q(|B_-|)\})^\beta. \tag{3.8}$$

- (2) Using (1.4), [22, (4.51)] becomes

$$\mathbb{E} \{ \langle \delta_0, F_{-,L} X_1 F_{+,L} X_1 F_{-,L} \delta_0 \rangle \} \leq \frac{1}{4} Q(|J|)^2 L^{d+2}. \tag{3.9}$$

- (3) In [22, Lemma 4.6], we cannot use the estimate [22, (4.28)]. But proceeding as in the proof of [22, (4.30)], we can replace it by

$$\begin{aligned} \mathbb{E} \{ |\langle \delta_x, F_\pm \delta_y \rangle|^p \} &\leq C_I \{ \{f_\pm\} \}_2 e^{-\frac{1}{\ell}|x-y|} \\ &\text{for all } p \in [1, \infty[\text{ and } x, y \in \mathbb{Z}^d. \end{aligned} \tag{3.10}$$

As a consequence, the right hand side of [22, (4.27)] becomes

$$C \left((\{f_+\}_2 \{f-\}_3 \{f-\}_4)^{\frac{1}{3}} + (\{f-\}_3 \{f_+\}_4)^{\frac{1}{3}} \right) L^{\frac{4}{3}d} e^{-\frac{1}{12\ell}L}. \tag{3.11}$$

- (4) Using (1.4) instead of [22, (4.47)], and taking into account the above modifications, [22, (4.61)] becomes

$$\begin{aligned} \Psi_{E_F}(I_+ \times I_-) &\leq \frac{1}{4} Q(2v)^2 L^{d+2} + C' v^{-15} L^{\frac{4}{3}d} e^{-\frac{1}{12\ell}L} + 4W_{\frac{1}{2}} Q(v^4)^{\frac{1}{2}} \\ &\leq 4^{\alpha-1} U^2 v^{2\alpha} L^{d+2} + C' v^{-15} L^{\frac{4}{3}d} e^{-\frac{1}{12\ell}L} + 4U^{\frac{1}{2}} W_{\frac{1}{2}} v^{2\alpha}, \end{aligned} \tag{3.12}$$

where we used (3.5).

- (5) As in [22, (4.62)], we choose $L = A\ell \log \frac{1}{v}$, where A is some suitable constant, depending on d, α and U , such that, similarly to [22, (4.63)], we get

$$\Psi_{E_F}(I_+ \times I_-) \leq B \ell^{d+2} v^{2\alpha} \left(\log \frac{1}{v} \right)^{d+2} + C'' v^{2\alpha}, \tag{3.13}$$

where B and C'' are constants, with B depending only on d, α and U , from which (3.6) follows. □

4 Proofs of Eigenvalue Counting Inequalities

In this section we prove Theorems 2.1 and 2.3 and Corollary 2.4. Since we always have $\text{tr } P_\omega(I) \leq |\Lambda|$, it follows from Lemma B.1 that it suffices to prove the theorems when all the probability measures μ_j have compact support, which is assumed in the proofs of Theorems 2.1 and 2.3.

Our proofs are based on the fundamental spectral averaging estimate (2.2) and [7, Lemma 3.2], which we now state.

Lemma 4.1 [7] *Consider the self-adjoint operator $H_s = H_0 + s\Pi_\varphi$ on the Hilbert space \mathcal{H} , where H_0 is a self-adjoint operator on \mathcal{H} , $\varphi \in \mathcal{H}$ with $\|\varphi\| = 1$, and $s \in \mathbb{R}$. Let $P_s(J) = \chi_J(H_s)$ for an interval J , and suppose $\text{tr } P_0([-\infty, c]) < \infty$ for all $c \in \mathbb{R}$. Then, given $a, b \in \mathbb{R}$ with $a < b$, we have*

$$\text{tr } P_s([a, b]) \leq 1 + \text{tr } P_t([a, b]) \quad \text{for all } 0 \leq s \leq t. \tag{4.1}$$

Let $\Lambda \subset \mathbb{Z}^d$ finite. Given $\omega \in \mathbb{R}^{\mathbb{Z}^d}$, we set $P_\omega^{(\Lambda)}(I) = \chi_I(H_{\omega, \Lambda})$. Given $j \in \Lambda$, we write $\omega = (\omega_j^\perp, \omega_j)$ and $P_{\omega_j=s}^{(\Lambda)}(I) = P_{(\omega_j^\perp, s)}^{(\Lambda)}(I)$ when we want to make explicit the value of ω_j .

We also write $P_{\omega_j \rightarrow s}^{(\Lambda)}(I)$ to denote that ω_j was replaced by s .

Since we assumed that the measures μ_j have no atoms, it follows from (2.4) that $\mathbb{E}_\omega\{\text{tr } P_\omega^{(\Lambda)}(\{c\})\} = 0$ for any $c \in \mathbb{R}$. Thus it does not matter if the intervals are open or closed at the endpoints, so in the proofs we may take all intervals to be of the form $]a, b]$, which allows the use of Lemma 4.1.

Theorem 2.1 is a particular case of Theorem 2.3, but in order to illustrate the simplicity of our approach we first give a proof of Theorem 2.1 and then prove the general case.

Proof of Theorem 2.1 Fix a finite volume $\Lambda \subset \mathbb{Z}^d$ and let I_1, I_2 be bounded intervals. Using Lemma 4.1, for $\tau_j \geq \omega_j$ we always have

$$\begin{aligned} & (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2) - 1) \\ &= \sum_{j \in \Lambda} \{ \langle \delta_j, P_\omega^{(\Lambda)}(I_1) \delta_j \rangle (\text{tr } P_\omega^{(\Lambda)}(I_2) - 1) \} \\ &\leq \sum_{j \in \Lambda} \left\{ \left\langle \delta_j, P_{(\omega_j^\perp, \omega_j)}^{(\Lambda)}(I_1) \delta_j \right\rangle \left(\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I_2) \right) \right\}. \end{aligned} \tag{4.2}$$

We now take $\tau_j \geq \max \text{supp } \mu_j$ for all $j \in \Lambda$, and average over the random variables $\omega = \{\omega_j\}_{j \in \mathbb{Z}^d}$, where each ω_j has the probability distribution μ_j . Using (2.2), we get

$$\begin{aligned} & \mathbb{E}_\omega \{ (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2) - 1) \} \\ &\leq \sum_{j \in \Lambda} \mathbb{E}_{\omega_j^\perp} \left\{ \left(\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I_2) \right) \left(\mathbb{E}_{\omega_j} \left\{ \left\langle \delta_j, P_{(\omega_j^\perp, \omega_j)}^{(\Lambda)}(I_1) \delta_j \right\rangle \right\} \right) \right\} \\ &\leq Q_\Lambda (|I_1|) \sum_{j \in \Lambda} \mathbb{E}_{\omega_j^\perp} \left\{ \text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I_2) \right\}. \end{aligned} \tag{4.3}$$

This holds for all $\tau_j \geq \max \text{supp } \mu_j$, $j \in \Lambda$, so we now take $\tau_j = \max \text{supp } \mu_j + \tilde{\omega}_j$, where $\tilde{\omega} = \{\tilde{\omega}_j\}_{j \in \mathbb{Z}^d}$ and $\omega = \{\omega_j\}_{j \in \mathbb{Z}^d}$ are two independent, identically distributed collections of

random variables, and average over these random variables. We get

$$\begin{aligned} & \mathbb{E}_\omega \left\{ (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2) - 1) \right\} \\ &= \mathbb{E}_\omega \left\{ \mathbb{E}_\omega \left\{ (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2) - 1) \right\} \right\} \\ &\leq Q_\Lambda (|I_1|) \sum_{j \in \Lambda} \mathbb{E}_{(\omega_j^\perp, \tilde{\omega}_j)} (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I_2)) \leq Q_\Lambda (|I_1|) Q_\Lambda (|I_2|) |\Lambda|^2, \end{aligned} \tag{4.4}$$

where we used the Wegner estimate (2.4).

The estimates (2.7) and (1.4) follow immediately from (4.4). To get (2.6), we use (4.4) and the obvious estimate

$$\begin{aligned} & (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2)) - \min \{ \text{tr } P_\omega^{(\Lambda)}(I_1), \text{tr } P_\omega^{(\Lambda)}(I_2) \} \\ &\leq (\text{tr } P_\omega^{(\Lambda)}(I_1)) (\text{tr } P_\omega^{(\Lambda)}(I_2) - 1) + (\text{tr } P_\omega^{(\Lambda)}(I_2)) (\text{tr } P_\omega^{(\Lambda)}(I_1) - 1). \end{aligned} \tag{4.5}$$

□

We now turn to the general case.

Proof of Theorem 2.3 Fix a finite volume $\Lambda \subset \mathbb{Z}^d$. We first prove (2.13), a particular case of (2.11), since it has a simpler proof. We fix the bounded interval I and proceed by induction on n . The case $n = 1$, is just Wegner’s inequality (2.4). Let us assume that (2.13) holds for n , for all possible probability distributions μ_j with compact support. Then, given $j \in \Lambda$ and $\tau_j \geq \max \text{supp } \mu_j$, we have, using (4.1), that for all $k = 1, 2, \dots, n$,

$$\text{tr } P_\omega^{(\Lambda)}(I) - k \leq 1 + \text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - k = \text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - (k - 1). \tag{4.6}$$

Note that $(\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I))(\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - 1) \dots (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - (n - 1)) \geq 0$. Since either $\text{tr } P_\omega^{(\Lambda)}(I)(\text{tr } P_\omega^{(\Lambda)}(I) - 1) \dots (\text{tr } P_\omega^{(\Lambda)}(I) - n) = 0$ or $\text{tr } P_\omega^{(\Lambda)}(I) - k > 0$ for $k = 0, 1, \dots, n$, it follows that we always have

$$\begin{aligned} & (\text{tr } P_\omega^{(\Lambda)}(I)) (\text{tr } P_\omega^{(\Lambda)}(I) - 1) \dots (\text{tr } P_\omega^{(\Lambda)}(I) - n) \\ &\leq \sum_{j \in \Lambda} \langle \delta_j, P_\omega^{(\Lambda)}(I) \delta_j \rangle (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I)) (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - 1) \dots \\ &\quad \times (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - (n - 1)). \end{aligned} \tag{4.7}$$

Using (2.2), we have

$$\begin{aligned} & \mathbb{E} \left\{ \langle \delta_j, P_\omega^{(\Lambda)}(I) \delta_j \rangle (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I)) (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - 1) \dots (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - (n - 1)) \right\} \\ &= \mathbb{E}_{\omega_j^\perp} \left\{ \mathbb{E}_{\omega_j} \left\{ \langle \delta_j, P_\omega^{(\Lambda)}(I) \delta_j \rangle (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I)) (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - 1) \dots \right. \right. \\ &\quad \left. \left. \times (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - (n - 1)) \right\} \right\} \\ &\leq Q_\Lambda (|I|) \mathbb{E}_{\omega_j^\perp} \left\{ (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I)) (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - 1) \dots (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - (n - 1)) \right\}. \end{aligned} \tag{4.8}$$

We now take $\tau = \{\tau_j = a_j + \tilde{\omega}_j\}_{j \in \Lambda}$, where $\tilde{\omega} = \{\tilde{\omega}_j\}_{j \in \Lambda}$ are independent random variables, independent of ω , such that $\tilde{\omega}_j$ has μ_j for probability distribution, and $a_j = \max \text{supp } \mu_j$. Using (4.7) and (4.8), plus the induction hypothesis, we get

$$\begin{aligned} & \mathbb{E}_\omega \left\{ \text{tr } P_\omega^{(\Lambda)}(I) (\text{tr } P_\omega^{(\Lambda)}(I) - 1) \cdots (\text{tr } P_\omega^{(\Lambda)}(I) - n) \right\} \\ &= \mathbb{E}_{(\omega, \tau)} \left\{ \text{tr } P_\omega^{(\Lambda)}(I) (\text{tr } P_\omega^{(\Lambda)}(I) - 1) \cdots (\text{tr } P_\omega^{(\Lambda)}(I) - n) \right\} \\ &\leq Q_\Lambda(|I|) \sum_{j \in \Lambda} \mathbb{E}_{(\omega_j^\perp, \tau_j)} \left\{ (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I)) (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - 1) \cdots \right. \\ &\quad \left. \times (\text{tr } P_{(\omega_j^\perp, \tau_j)}^{(\Lambda)}(I) - (n - 1)) \right\} \\ &\leq Q_\Lambda(|I|) \sum_j (Q_\Lambda(|I|) |\Lambda|)^n = (Q_\Lambda(|I|) |\Lambda|)^{n+1}. \end{aligned} \tag{4.9}$$

We now turn to the proof of (2.11). The case $n = 1$ is just (2.4), and $n = 2$ is (2.6), so we assume $n \geq 3$. Let I_1, I_2, \dots, I_n be bounded intervals. For a fixed ω , we have (2.9) and (2.10). Let us suppose

$$\Xi_\omega := (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)})) (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(2)} - 1) \cdots (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(n)} - (n - 1)) > 0, \tag{4.10}$$

and note that in this case we must have

$$\begin{aligned} & \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(k)} - k + 1 \geq 1, \quad \text{i.e.,} \quad \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(k)}) \geq k, \\ & \text{for all } k = 1, \dots, n. \end{aligned} \tag{4.11}$$

Then, using Lemma 4.1 repeatedly, we get

$$\begin{aligned} \Xi_\omega &\leq \sum_{j_1 \in \Lambda} \left\{ \left\langle \delta_{j_1}, P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)}) \delta_{j_1} \right\rangle \left(\text{tr } P_{\omega^{(j_1)}}^{(\Lambda)}(I_{\sigma_\omega(2)}) \right) \cdots \left(\text{tr } P_{\omega^{(j_1)}}^{(\Lambda)}(I_{\sigma_\omega(n)} - n) \right) \right\} \\ &\leq \cdots \\ &\leq \sum_{j_1, j_2, \dots, j_{n-1} \in \Lambda} \left\{ \left\langle \delta_{j_1}, P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)}) \delta_{j_1} \right\rangle \left\langle \delta_{j_2}, P_{\omega^{(j_1)}}^{(\Lambda)}(I_{\sigma_\omega(2)}) \delta_{j_2} \right\rangle \cdots \right. \\ &\quad \left. \times \left\langle \delta_{j_{n-1}}, P_{\omega^{(j_1, j_2, \dots, j_{n-2})}}^{(\Lambda)}(I_{\sigma_\omega(n-1)}) \delta_{j_{n-1}} \right\rangle \left(\text{tr } P_{\omega^{(j_1, j_2, \dots, j_{n-1})}}^{(\Lambda)}(I_{\sigma_\omega(n)}) \right) \right\}, \end{aligned} \tag{4.12}$$

where $\omega^{(j_1)}$ is ω with $\omega_{j_1} \rightarrow \tau_{j_1}^{(1)}$, $\omega^{(j_1, j_2)}$ is $\omega^{(j_1)}$ with $\omega_{j_2}^{(j_1)} \rightarrow \tau_{j_2}^{(2)}$, ..., $\omega^{(j_1, j_2, \dots, j_{n-1})}$ is $\omega^{(j_1, j_2, \dots, j_{n-2})}$ with $\omega_{j_{n-1}}^{(j_1, j_2, \dots, j_{n-2})} \rightarrow \tau_{j_{n-1}}^{(n-1)}$. To be able to apply Lemma 4.1 we must have

$$\begin{aligned} & \omega_{j_1} \leq \tau_{j_1}^{(1)}, \quad \omega_{j_2}^{(j_1)} \leq \tau_{j_2}^{(2)}, \quad \dots, \quad \omega_{j_{n-1}}^{(j_1, j_2, \dots, j_{n-2})} \leq \tau_{j_{n-1}}^{(n-1)}, \\ & j_k \in \Lambda, \quad k = 1, 2, \dots, n - 1. \end{aligned} \tag{4.13}$$

We then take

$$\tau = \left\{ \tau_{j_k}^{(k)} = a_{j_k}^{(k)} + \omega_{j_k}^{(k)}; j_k \in \Lambda, k = 1, 2, \dots, n - 1 \right\},$$

where $\widehat{\omega} = \{\omega_{j_k}^{(k)}; j_k \in \Lambda, k = 1, 2, \dots, n - 1\}$ are independent random variables, independent of ω , such that $\omega_{j_k}^{(k)}$ has μ_{j_k} for probability distribution, and the real numbers $a_{j_k}^{(k)}$ are chosen such that (4.13) holds $\mathbb{P}_{(\omega, \widehat{\omega})}$ -almost surely.

Since the last expression in (4.12) is obviously nonnegative, it follows that (4.12) holds also when $\Xi_\omega = 0$, an hence it holds $\mathbb{P}_{(\omega, \widehat{\omega})}$ -almost surely.

Given $\sigma \in \mathcal{S}_n$, let

$$\begin{aligned} \Xi_{\omega, \tau, \sigma} := & \sum_{j_1, j_2, \dots, j_{n-1} \in \Lambda} \left\{ \left\langle \delta_{j_1}, P_\omega^{(\Lambda)}(I_{\sigma(1)}) \delta_{j_1} \right\rangle \left\langle \delta_{j_2}, P_{\omega(j_1)}^{(\Lambda)}(I_{\sigma(2)}) \delta_{j_2} \right\rangle \cdots \right. \\ & \left. \times \left\langle \delta_{j_{n-1}}, P_{\omega(j_1, j_2, \dots, j_{n-2})}^{(\Lambda)}(I_{\sigma(n-1)}) \delta_{j_{n-1}} \right\rangle \left(\text{tr } P_{\omega(j_1, j_2, \dots, j_{n-1})}^{(\Lambda)}(I_{\sigma(n)}) \right) \right\}. \end{aligned} \tag{4.14}$$

It follows that $\mathbb{P}_{(\omega, \widehat{\omega})}$ -almost surely we have

$$\Xi_\omega \leq \sum_{\sigma \in \mathcal{S}_n(I_1, \dots, I_n)} \Xi_{\omega, \tau, \sigma}, \tag{4.15}$$

and hence

$$\mathbb{E}_\omega \{ \Xi_\omega \} = \mathbb{E}_{(\omega, \widehat{\omega})} \{ \Xi_\omega \} \leq \sum_{\sigma \in \mathcal{S}_n(I_1, \dots, I_n)} \mathbb{E}_{(\omega, \widehat{\omega})} \{ \Xi_{\omega, \tau, \sigma} \}. \tag{4.16}$$

By performing the integrations in the right order, using (2.2) $n - 1$ times, and then using the Wegner estimate (2.4), we get

$$\mathbb{E}_{(\omega, \widehat{\omega})} \{ \Xi_{\omega, \tau, \sigma} \} \leq \left(\prod_{k=1}^n Q^{(\Lambda)}(|I_k|) \right) |\Lambda|^n. \tag{4.17}$$

Since $M(I_1, \dots, I_n) = |\mathcal{S}_n(I_1, \dots, I_n)|$, the estimate (2.11) follows. □

Proof of Corollary 2.4 The estimate (2.14) follows from (2.13) and the inequality

$$\begin{aligned} & \mathbb{P} \{ \text{tr } P_\omega^{(\Lambda)}(I) \geq n \} \\ & \leq \mathbb{P} \{ (\text{tr } P_\omega^{(\Lambda)}(I)) (\text{tr } P_\omega^{(\Lambda)}(I) - 1) \cdots (\text{tr } P_\omega^{(\Lambda)}(I) - (n - 1)) \geq n! \} \\ & \leq \frac{1}{n!} \mathbb{E} \{ (\text{tr } P_\omega^{(\Lambda)}(I)) (\text{tr } P_\omega^{(\Lambda)}(I) - 1) \cdots (\text{tr } P_\omega^{(\Lambda)}(I) - (n - 1)) \}. \end{aligned} \tag{4.18}$$

To obtain (2.15), we use (2.11) with

$$\begin{aligned} & \mathbb{P} \{ \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)}) \geq 1, \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(2)}) \geq 2, \dots, \text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(n)}) \geq n \} \\ & \leq \mathbb{E} \{ (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(1)})) (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(2)}) - 1) \cdots (\text{tr } P_\omega^{(\Lambda)}(I_{\sigma_\omega(n)}) - (n - 1)) \}. \end{aligned} \tag{4.19}$$

Similarly, (2.16) follows from (2.12). □

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Appendix A: The Fundamental Spectral Averaging Estimate

For the reader convenience we present a proof of the fundamental spectral averaging result (2.2). Consider the random self-adjoint operator on a Hilbert space \mathcal{H} given in (2.1): $H_\omega = H_0 + \omega \Pi_\varphi$, where H_0 is a self-adjoint operator on the Hilbert space \mathcal{H} , $\varphi \in \mathcal{H}$ with $\|\varphi\| = 1$, Π_φ is the orthogonal projection onto the one-dimensional subspace spanned by φ , and ω is a random variable with a non-degenerate probability distribution μ . Given z with $\Im z > 0$, we have, as in [6, Proof of Lemma 6.1], that

$$\langle \varphi, (H_\omega - z)^{-1} \varphi \rangle = (\langle \varphi, (H_0 - z)^{-1} \varphi \rangle^{-1} + \omega)^{-1}. \tag{A.1}$$

A.1 The Probability Distribution μ Has a Bounded Density ρ

In this case, we use

$$\int_{\mathbb{R}} d\omega \Im \langle \varphi, (H_\omega - z)^{-1} \varphi \rangle = \pi, \tag{A.2}$$

a consequence of (A.1) (cf. [6, Proof of Lemma 6.1]). It then follows from Stone’s formula (cf. [27, Theorem VII.13]) that

$$\int_{\mathbb{R}} d\omega \left\langle \varphi, \frac{1}{2} \{ P_\omega([a, b]) + P_\omega(]a, b[) \} \varphi \right\rangle \leq (b - a). \tag{A.3}$$

In particular, $\int_{\mathbb{R}} d\omega \langle \varphi, P_\omega(\{c\}) \varphi \rangle = 0$, and hence for any bounded interval I we get

$$\int_{\mathbb{R}} d\omega \langle \varphi, P_\omega(I) \varphi \rangle \leq |I|. \tag{A.4}$$

Since μ has a bounded density ρ , we get

$$\int_{\mathbb{R}} d\mu(\omega) \langle \varphi, P_\omega(I) \varphi \rangle = \int_{\mathbb{R}} d\omega \rho(\omega) \langle \varphi, P_\omega(I) \varphi \rangle \leq \|\rho\|_\infty |I|. \tag{A.5}$$

Remark A.1 The reader may notice that [6] has an extra factor of π in the right hand side of (A.5); the difference comes from using Stone’s formula instead of the simple estimate (A.10). Since in Theorem 2.1 we obtain (1.4) as a consequence of (A.5), and (1.3) is a particular case of (1.4), we do not have the factor of π^2 in the right hand side (1.3): the estimate is just $(\rho_\infty^{(\Lambda)} |I| |\Lambda|)^2$.

A.2 Arbitrary Probability Distribution μ

We consider an interval $I = [E - \varepsilon, E + \varepsilon]$, $\varepsilon > 0$, and set $z = E + i\varepsilon$ and $R_0(z) = (H_0 - z)^{-1}$. Given $\kappa > 0$, we define real numbers a and b by

$$a - ib = \frac{\kappa}{2\varepsilon} \langle \varphi, R_0(z) \varphi \rangle^{-1}, \tag{A.6}$$

and note that we always have

$$\frac{2}{\kappa} b = \frac{\frac{1}{\varepsilon} \Im \langle \varphi, R_0(z) \varphi \rangle}{|\langle \varphi, R_0(z) \varphi \rangle|^2} = \frac{\|R_0(z) \varphi\|^2}{|\langle \varphi, R_0(z) \varphi \rangle|^2} \geq 1. \tag{A.7}$$

From (A.6) and (A.1) we get,

$$\varepsilon \Im \langle \varphi, (H_\omega - z)^{-1} \varphi \rangle = \frac{\kappa}{2} \frac{b}{(a + \frac{\kappa}{2\varepsilon} \omega)^2 + b^2}. \tag{A.8}$$

Proceeding as in [10], and using (A.8), [10, Lemma 3.1] and (A.7), we get

$$\begin{aligned} & \varepsilon \int d\mu(\omega) \Im \langle \varphi, (H_\omega - z)^{-1} \varphi \rangle \\ &= \frac{\kappa}{2} \sum_{n \in \mathbb{Z}} \int_{[n \frac{2\varepsilon}{\kappa}, (n+1) \frac{2\varepsilon}{\kappa}]} d\mu(\omega) \frac{b}{(a + \frac{\kappa}{2\varepsilon} \omega)^2 + b^2} \\ &\leq \frac{\kappa}{2} S_\mu \left(\frac{2\varepsilon}{\kappa} \right) \sum_{n \in \mathbb{Z}} \sup_{y \in [0, 1]} \frac{b}{(a + n + y)^2 + b^2} \leq \frac{\kappa}{2} \pi \left(1 + \frac{1}{b} \right) S_\mu \left(\frac{2\varepsilon}{\kappa} \right) \\ &\leq \pi \left(1 + \frac{\kappa}{2} \right) S_\mu \left(\frac{2\varepsilon}{\kappa} \right). \end{aligned} \tag{A.9}$$

We may now get (2.2) in two ways. Using the simple inequality

$$P_\omega(I) \leq 2\varepsilon \Im (H_\omega - z)^{-1}, \tag{A.10}$$

(2.2) follows immediately from (A.9) with

$$Q_\mu(|I|) = \inf_{\kappa > 0} \left(\pi(2 + \kappa) S_\mu \left(\frac{1}{\kappa} |I| \right) \right) \leq 3\pi S_\mu(|I|). \tag{A.11}$$

We can improve the constant slightly by using the more sophisticated inequality given in [10, (3.1)], that is,

$$P_\omega(I) \leq \frac{4}{\pi} \int_I d\lambda \Im (H_\omega - \lambda - i|I|)^{-1}, \tag{A.12}$$

together with (A.9), getting (2.2) with

$$Q_\mu(|I|) = \inf_{\kappa > 0} \left(4(1 + \kappa) S_\mu \left(\frac{1}{\kappa} |I| \right) \right) \leq 8S_\mu(|I|). \tag{A.13}$$

Appendix B: An Approximation Lemma

Lemma B.1 *Let F be a bounded, nonnegative Borel measurable function on \mathbb{R}^N , $s_1, s_2, \dots, s_q > 0$, and $\omega = \{\omega_j\}_{j=1,2,\dots,N}$ a family of independent random variables, μ_j denoting the probability distribution of the random variable ω_j . We write $\mu = \{\mu_j\}_{j=1,2,\dots,N}$, and denote the corresponding expectation by \mathbb{E}_μ . Let $Q_\mu(s) := \max_{j=1,2,\dots,N} Q_{\mu_j}(s)$. Suppose there exists a constant $K > 0$ such that when μ_j has compact support for all $j = 1, 2, \dots, N$ we have*

$$\mathbb{E}_\mu \{F(\omega)\} \leq K \prod_{i=1}^q Q_\mu(s_i). \tag{B.1}$$

Then (B.1) holds for arbitrary probability distributions $\mu = \{\mu_j\}_{j=1,2,\dots,N}$.

Proof Given a probability distribution μ and $M \in \mathbb{N}$, we set $\chi_M = \chi_{[-M, M]}$ and

$$\mu^{(M)} = c_{\mu^{(M)}} \chi_M \mu, \quad \text{with } c_{\mu^{(M)}} = (\mu \{[-M, M]\})^{-1}. \tag{B.2}$$

Note that $\mu^{(M)}$ is a probability measure with compact support for all $M \in \mathbb{N}$, and

$$\lim_{M \rightarrow \infty} c_{\mu^{(M)}} = 1. \tag{B.3}$$

Moreover, we have

$$Q_{\mu^{(M)}}(s) \leq c_{\mu^{(M)}} Q_{\mu}(s) \quad \text{for all } s > 0. \tag{B.4}$$

Now, given probability distributions $\mu = \{\mu_j\}_{j=1,2,\dots,N}$, and $M \in \mathbb{N}$, we consider the probability distributions $\mu^{(M)} = \{\mu_j^{(M)}\}_{j=1,2,\dots,N}$, and set $c_{\mu^{(M)}} = \prod_{j=1}^N c_{\mu_j^{(M)}}$. We have

$$\mathbb{E}_{\mu^{(M)}} \{F(\omega)\} = c_{\mu^{(M)}} \mathbb{E}_{\mu} \left\{ \left(\prod_{j=1}^N \chi_M(\omega_j) \right) F(\omega) \right\}, \tag{B.5}$$

and hence it follows from the bounded convergence theorem and (B.3) that

$$\lim_{M \rightarrow \infty} \mathbb{E}_{\mu^{(M)}} \{F(\omega)\} = \mathbb{E}_{\mu} \{F(\omega)\}. \tag{B.6}$$

Since (B.1) holds for $\mu^{(M)}$, the lemma follows from (B.3), (B.4), and (B.6). □

References

1. Aizenman, M.: Localization at weak disorder: some elementary bounds. *Rev. Math. Phys.* **6**, 1163–1182 (1994)
2. Aizenman, M., Molchanov, S.: Localization at large disorder and extreme energies: an elementary derivation. *Commun. Math. Phys.* **157**, 245–278 (1993)
3. Aizenman, M., Schenker, J., Friedrich, R., Hundertmark, D.: Finite volume fractional-moment criteria for Anderson localization. *Commun. Math. Phys.* **224**, 219–253 (2001)
4. Aizenman, M., Warzel, S.: On the joint distribution of energy levels of random Schrödinger operators. Preprint
5. Bellissard, J., Hislop, P., Stolz, G.: Correlation estimates in the Anderson model. *J. Stat. Phys.* **129**, 649–662 (2007)
6. Carmona, R., Klein, A., Martinelli, F.: Anderson localization for Bernoulli and other singular potentials. *Commun. Math. Phys.* **108**, 41–66 (1987)
7. Combes, J.M., Germinet, F., Klein, A.: Poisson statistics for eigenvalues of continuum random Schrödinger operators. [arXiv:0807.0455v3](https://arxiv.org/abs/0807.0455v3) [math-ph]
8. Combes, J.M., Hislop, P.D.: Localization for some continuous, random Hamiltonians in d -dimension. *J. Funct. Anal.* **124**, 149–180 (1994)
9. Combes, J.-M., Hislop, P.D., Klopp, F.: Regularity properties for the density of states of random Schrödinger operators. In: *Waves in Periodic and Random Media*, South Hadley, MA, 2002. *Contemp. Math.*, vol. 339, pp. 15–24 (2003)
10. Combes, J.M., Hislop, P.D., Klopp, F.: Optimal Wegner estimate and its application to the global continuity of the integrated density of states for random Schrödinger operators. *Duke Math. J.* **140**, 469–498 (2007)
11. Delyon, F., Lévy, Y., Souillard, B.: Anderson localization for multidimensional systems at large disorder or large energy. *Commun. Math. Phys.* **100**(4), 463–470 (1985)
12. von Dreifus, H., Klein, A.: A new proof of localization in the Anderson tight binding model. *Commun. Math. Phys.* **124**, 285–299 (1989)
13. Fröhlich, J., Martinelli, F., Scoppola, E., Spencer, T.: Constructive proof of localization in the Anderson tight binding model. *Commun. Math. Phys.* **101**, 21–46 (1985)

14. Fröhlich, J., Spencer, T.: Absence of diffusion with Anderson tight binding model for large disorder or low energy. *Commun. Math. Phys.* **88**, 151–184 (1983)
15. Germinet, F., Klein, A.: Bootstrap multiscale analysis and localization in random media. *Commun. Math. Phys.* **222**, 415–448 (2001)
16. Germinet, F., Klein, A.: A characterization of the Anderson metal-insulator transport transition. *Duke Math. J.* **124**, 309–351 (2004)
17. Germinet, F., Klein, A.: New characterizations of the region of complete localization for random Schrödinger operators. *J. Stat. Phys.* **122**, 73–94 (2006)
18. Graf, G.-M., Vaghi, A.: A remark on an estimate by Minami. *Lett. Math. Phys.* **79**, 17–22 (2007)
19. Hundertmark, D.: On the time-dependent approach to Anderson localization. *Math. Nachr.* **214**, 25–38 (2000)
20. Killip, R., Nakano, F.: Eigenfunction statistics in the localized Anderson model. *Ann. H. Poincaré* **8**, 27–36 (2007)
21. Kirsch, W.: An invitation to random Schrödinger operators. In: *Panorama et Synthèses* (2008)
22. Klein, A., Lenoble, O., Müller, P.: On Mott’s formula for the ac-conductivity in the Anderson model. *Ann. Math.* **166**, 549–577 (2007)
23. Klein, A., Molchanov, S.: Simplicity of eigenvalues in the Anderson model. *J. Stat. Phys.* **122**, 95–99 (2006)
24. Kritchevski, E.: Poisson statistics of eigenvalues in the hierarchical Anderson model. *Ann. H. Poincaré* **9**, 685–709 (2008)
25. Minami, N.: Local fluctuation of the spectrum of a multidimensional Anderson tight binding model. *Commun. Math. Phys.* **177**, 709–725 (1996)
26. Nakano, F.: The repulsion between localization centers in the Anderson model. *J. Stat. Phys.* **123**, 803–810 (2006)
27. Reed, M., Simon, B.: *Methods of Modern Mathematical Physics I: Functional Analysis*. Academic Press, New York (1980), revised and enlarged edition
28. Simon, B.: Cyclic vectors in the Anderson model. Special issue dedicated to Elliott H. Lieb. *Rev. Math. Phys.* **6**, 1183–1185 (1994)
29. Simon, B., Wolff, T.: Singular continuum spectrum under rank one perturbations and localization for random Hamiltonians. *Commun. Pure Appl. Math.* **39**, 75–90 (1986)
30. Stoiciu, M.: The statistical distribution of the zeros of random paraorthogonal polynomials on the unit circle. *J. Approx. Theory* **139**, 29–64 (2006)
31. Stoiciu, M.: Poisson statistics for eigenvalues: from random Schrödinger operators to random CMV matrices. In: *Probability and Mathematical Physics*. CRM Proc. Lecture Notes, vol. 42, pp. 465–475. Amer. Math. Soc., Providence (2007)
32. Wegner, F.: Bounds on the density of states in disordered systems. *Z. Phys. B* **44**, 9–15 (1981)