

Investigations of Molecular Interactions in the Binary Mixtures of 1‑Butyl‑3‑methylimidazolium bis(Trifuoromethanesulfonyl) Amide and 2‑Propoxyethanol from *T***=(298.15 to 323.15) K at Atmospheric Pressure**

G. Prasad^{1,5} · K. Muralidhar Reddy² · R. Padamasuvarna¹ · T. Madhu Mohan³ · **T. Vijaya Krishna³ · S. Govardhana Rao4**

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Abstract

Density and speed of sound values of pure and binary mixtures of 1-butyl-3-methylimidazolium bis(trifluoromethanesulfonyl) amide ($[Bmin][NTf₂]$) and 2-propoxyethanol (2-PR) were determined over the whole composition range as a function of temperature from 298.15 to 323.15 K) under atmospheric pressure. Heat capacity values of the pure compounds were also measured. The experimental values were used to calculate excess molar volume, excess isentropic compressibility, excess intermolecular free length and excess speed of sound values. The excess values were ftted with the Redlich–Kister polynomial equation to estimate the binary coefficients and standard deviation between the experimental and calculated values. The partial molar volumes at infnite dilutions were also calculated. The trends of variation of these properties were interpreted in light of the solute–solvent interactions occurring in the system. Further, the molecular interactions in the binary system were analyzed using experimental FT-IR spectra recorded at room temperature.

Keywords 1-Butyl-3-methylimidazolium bis(trifuoromethanesulfonyl) amide · 2-Propoxyethanol · Excess molar volume · Infnite dilution · FT-IR

1 Introduction

Ionic liquids gained popularity as "environmentally friendly solvents" and alternatives to "volatile organic solvents" [[1–](#page-24-0)[4\]](#page-24-1). In continuation of our previous work on ILs and their binary mixtures with industrially important organic solvents [\[5–](#page-24-2)[8](#page-24-3)], in the

 \boxtimes T. Madhu Mohan tmadhumohan@gmail.com

 \boxtimes T. Vijaya Krishna tvijayakrishna1980@gmail.com

Extended author information available on the last page of the article

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present study we have chosen the binary mixtures of the IL, 1-butyl-3-methylimidazolium bis(trifluoromethanesulfonyl) amide ($[Bmin][NTf₂]$) and 2-propoxyethanol (2-PR). The relatively low viscosity of imidazolium based ILs facilitates easy handling, adequate performance in extractive applications and thus these ILs are more environmentally friendly than the other ILs [[9\]](#page-25-0). [Bmim][NTf₂] is a preferred ionic liquid as electrolyte for capacitive energy storage and also in the process of nuclear fuel recycling [[10](#page-25-1)]. On the other hand, alkoxyethanols are amphiphilic organic solvents which can be used in many chemical processes. The mixtures containing alkoxyethanols are very important in view of their strong intermolecular efects produced by the presence of –O– and –OH groups in the same compound. 2-PR is used as a solvent in varnishes, dyes and pesticides.

It is well-known that the reaction medium plays an important role in determining reactivity, that is refected in thermodynamic, transport and spectral properties. The thermodynamic properties of solutions are very useful to obtain information on the intermolecular interactions and geometrical efects in the systems. Moreover, knowledge of the thermodynamic properties of binary mixtures is essential for the proper design of industrial processes. Thus the accurate knowledge of thermodynamic properties of the [Bmim] $[NTf₂] + 2-PR$ binary mixture can be of great mportance for their possible use in bioreactor technology, Li-ion batteries, synthesis of nano objects, separation technology, catalysis, etc. [[11](#page-25-2)].

2 Experimental Section

2.1 Chemicals Used

The ionic liquid, 1-butyl-3-methylimidazolium bis(trifuoromethanesulfonyl) amide was procured from Io-Li-Tec, Germany and 2-propoxyethanol was procured from Sigma Aldrich, USA. 2-PR was purified by distillation and $[\text{Bmin}][\text{NT}_2]$ was used without any further purifcation.

2.1.1 Analysis of the Water Contents of the Chemicals

The values of water content in $[Bmim][NTf₂]$ and 2-PR were determined using a Karl-Fischer titrator (Metrohm, 890 Titrando). This can detect water content from less than 10 ppm to 100%. All samples were dried for at least 72 h under atmospheric pressure and moderate temperature (beginning at room temperature and increasing gradually up to 333 K over a period of 6 h), before making the measurements. The provenance, CAS number, mass fraction purity and water content values of the pure compounds are given in Table [1](#page-2-0).

2.2 Apparatus and Procedure

2.2.1 Sample Preparation

All samples were prepared by mass and were stored in amber colored glass vials (8 mL) with screw caps having PFE septa, sealed with paraflm to prevent absorption of moisture from the atmosphere. The samples were prepared immediately prior to measurement, using an electronic balance (CPA-225D, Sartorius, Germany) with a readability of $\pm 1\times10^{-7}$ kg. The uncertainty in the mole fraction was estimated to be within $\pm 1 \times 10^{-4}$.

2-propoxyethanol Sigma Aldrich, USA 2807-30-9 0.994 Distillation 0.999 60 Gas liquid chromatography

0.994

Sigma Aldrich, USA 2807-30-9

 2 -propoxyethanol

Distillation

Gas liquid chromatography

 60

0.999

Table 1 List of chemicals with details of provenance, CAS number, mass fraction purity and water content $\ddot{}$ $\ddot{ }$ $\frac{1}{l}$ i, J ł

2.2.2 Measurement of Density and Speed of Sound

The density and speed of sound values of the pure compounds and their binary mixtures were measured using a digital Density and Sound Analyzer (DSA 5000 M, Anton Parr, Austria), which employs the well-known oscillating U-tube principle (for density measurement) and propagation time technique (for speed of sound measurement at a frequency of 3 MHz). The reproducibility of values of density and speed of sound were $\pm 1 \times 10^{-3}$ kg·m⁻³ and $\pm 1 \times 10^{-2}$ m·s⁻¹, respectively. The density meter was calibrated with double distilled degassed water (ρ =997.075 kg·m⁻³ at 298.15 K). After each measurement distiled water and anhydrous ethanol were used to clean the vibrating tube. The reported values are the average of three consecutive measurements carried out between the temperatures (298.15–323.15) K with intervals of 5 K, under atmospheric pressure. The standard uncertainty associated with the measurement of temperature is ± 0.01 K. The estimated uncertainties of density and speed of sound are within $\pm 0.05 \text{ kg} \cdot \text{m}^{-3}$ and $\pm 0.5 \text{ m} \cdot \text{s}^{-1}$, respectively. The comparison of density and speed of sound values with the NIST ThermoLit data at temperatures $T = (298.15 - 323.15)$ K is shown in Table [2](#page-4-0) $[12-21]$ $[12-21]$.

2.2.3 Measurement of Heat Capacity and Infrared Spectra

The heat capacity (C_p) values of pure [Bmim][NTf₂] and pure 2-PR were determined with a micro DSC 7 evo (SETARAM instrumentation, France) diferential scanning calorimeter. It has a double stage temperature controller which works in the temperature range (228.15–393.15) K, with a heating rate of $(0.001-2)$ K·min⁻¹. A "mixing batch cell" made up of Hastelloy C276 was used with a sample capacity of (10–50) μL. The instrument was calibrated by determining the enthalpy of fusion of naphthalene (147.78 J·g⁻¹) using the Joule effect method. The estimated uncertainty of heat capacity is $u(C_p)=(0.003\cdot C_p)$ J·K−1·mol−1. The infrared transmittance was measured by using a Shimadzu Fourier transform infrared (FT-IR) spectrometer equipped with attenuated total refectance (ATR) accessories. The spectral region was 500–3500 cm⁻¹ with resolution of 2 cm⁻¹ and 100 scans. At least fve repeated measurements were performed for each sample.

3 Results and Discussion

The density values obtained from the experiments were ftted well with the linear equation,

$$
\rho = a + bT \tag{1}
$$

where *a* and *b* represent the empirical parameters obtained by the method of least squares ftting. The ftting parameters and the average absolute relative deviation (*ARD*) defned by the Eq. [2](#page-3-0) are presented in supplemental Table S1.

$$
ARD = \left(\frac{1}{n}\sum \frac{\left|\rho_{\exp} - \rho\right|}{\rho_{\exp}}\right)100\tag{2}
$$

where *n* is the number of data points, $\rho_{\rm exp}$ is the experimental density and ρ is the calculated density.

The variation of density (ρ) values, as a function of [Bmim][NTf₂] concentration (x_1) and temperature, is shown in Fig. [1](#page-6-0). At all investigated temperatures, the density values

 \overline{a}

ł J

J

Fig. 1 Density ρ versus mole fraction of [Bmim][NTf₂] (x_1) in the binary mixture of [Bmim][NTf₂]+2-PR. 298.15 K (flled square), 303.15 K (red flled circle), 308.15 K (blue flled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green flled right sided triangle) (Color fgure online)

increase with increasing mole fraction of $[Bmin][NTf₂]$ in the mixture and decrease with increasing temperature. The variation of the speed of sound values (u) , as a function of $[Bmim][NTf₂]$ $[Bmim][NTf₂]$ $[Bmim][NTf₂]$ concentration and temperature, is shown in Fig. 2. The speeds of sound decrease with increasing mole fraction of $[Bmin][NTf₂]$ and as the temperature increases. The variations in the values of density and speed of sound with respect to mole fraction and temperature (Table [3\)](#page-8-0) demonstrate that the physical properties of $[Bmin][NTf₂]$ can be modifed in a controlled way by adding 2-PR or changing the temperature.

 V_m values (Table [4](#page-10-0)) of [Bmim][NTf₂] + 2-PR increase with [Bmim][NTf₂] concentration and with temperature. The increase in the molar volume of a system upon mixing of the components could be attributed to the dissociation of one component or both the components and formation of solute–solvent bonds.

The excess molar volume $(V_{\text{m}}^{\text{E}})$ was calculated from the equation,

$$
V_{\rm m}^{\rm E} = \left(\frac{x_1 M_1 + x_2 M_2}{\rho}\right) - \left(\frac{x_1 M_1}{\rho_1} + \frac{x_2 M_2}{\rho_2}\right) \tag{3}
$$

where $x_1, x_2, M_1, M_2, \rho_1$ and ρ_2 represent the mole fraction, molar mass and density of pure component 1 ([Bmim][NTf₂]) and pure component 2 (2-PR), respectively, and ρ is the density of the binary mixture.

Figure [3](#page-12-0) indicates that the V_{m}^{E} values exhibit a sigmoid variation, with negative values at low mole fractions of [Bmim][NTf2] and positive values at high mole fractions of [Bmim] [NTf2] in the binary mixture, over the investigated temperature range. The sign and magnitude of the volume changes that take place on mixing [Bmim][NTf2] and 2-PR are the result of several efects that operate in the same or in the opposite direction. In general, the physical contributions comprised of dispersion forces and non-specifc physical (weak) interactions lead to positive V_{m}^{E} values. The chemical contributions, like breaking up of the

Fig. 2 Speed of sound *u* versus mole fraction of $[\text{Bmin}][\text{NTf}_2]$ (x_1) in the binary mixture of $[\text{Bmin}]$ $[NTf₂] + 2-PR$. 298.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue filled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green filled right sided triangle) (Color figure online)

associates present in the pure liquids, result in positive V_{m}^{E} values. Whereas, the chemical contributions such as formation of H-bonding, charge transfer (donor–acceptor) complexes, strong dipole–dipole and ion–dipole interactions between the component species of the mixture result in negative V_{m}^{E} values [\[22\]](#page-25-13).

A plausible qualitative interpretation of the variation of V_{m}^{E} with composition has been suggested. On mixing [Bmim][NTf2] and 2-PR, the disruption of the associations present in pure liquids takes place with subsequent formation of new interactions between unlike species. The negative V_{m}^{E} values at low mole fractions of ionic liquid ($x_1 < 0.4901$) suggest the formation of significant ion–dipole interactions between the ions $[Bmim]$ ⁺ and $[NTF_2]$ ⁻ formed by the ionic liquid and dipoles of 2-PR species. At higher mole fractions $(x_1 > 0.4901)$ the interactions between like species IL–IL and 2-PR–2-PR seem to dominate in the binary mixture.

Comparison of the present system with the literature can provide a clear picture about the nature of molecular interactions. In the 2-PR rich region, the present system ([Bmim] [NTf₂]+2-PR) was compared with the [Bmim][PF₆]+2-PR [\[18](#page-25-10)] and [Emim][EtSO₄]+2-PR [[21](#page-25-4)] systems. The V_{m}^{E} values of [Bmim][PF₆] + 2-PR and [Emim][EtSO₄] + 2-PR are negative over the complete composition range, at all temperatures. While, the $[Bmim][NTf₂] + 2-PR$ system shows a sigmoid variation. Further, more negative values of V_{m}^{E} are observed in the case of $[Bmin][NTf₂] + 2-PR$ compared with the remaining systems, indicating strong molecular interactions between the species. In the $[\text{Bmin}][\text{NTf}_2]$ rich region, the present system $([Bmim][NTf₂] + 2-PR)$ was compared with $[Bmim][NTf₂] + 2-metboxyethanol [23]$ $[Bmim][NTf₂] + 2-metboxyethanol [23]$ $[Bmim][NTf₂] + 2-metboxyethanol [23]$ and [Bmim][NTf₂]+2-ethoxyethanol [\[24\]](#page-25-15) systems. The V_{m}^{E} values of [Bmim][NTf₂]+2-methoxyethanol are negative over the whole composition range at all temperatures. While, [Bmim] $[NTF₂] + 2$ -ethoxyethanol and $[Bmim][NTf₂] + 2-PR$ systems show a sigmoid variation. Further, more positive values of V_{m}^{E} are observed in the case of $[\text{Bmim}][\text{NTf}_2]+2\text{-PR}$ compared with $[Bmim][NTf₂] + 2$ -ethoxyethanol system, indicating less packing.

Table 3 Density (ρ) and speed sound (*u*) with mole fraction (*x*1) of 1-butyl-3-methylimidazolium bis(trifuoromethanesulfonyl) amide in the binary mixture of 1-butyl-3-methylimidazolium bis(trifuoromethanesulfonyl) amide+ 2-propoxyethanol from *T*=(293.15 to 323.15) K at pressure *p*=0.1 MPa

Table 3 (continued)

Standard uncertainties u are $u(T) = \pm 0.01$ K, $u(x_1) = \pm 1 \times 10^{-4}$, $u(\rho)$ =±0.05 kg·m−3, u(*u*)=±0.5 m·s−1, u(*p*)=±0.002 MPa

Figure [3](#page-12-0) shows that V_{m}^{E} values become more negative as the temperature is increased in the 2-PR rich region. With increase in temperature the kinetic energy of the species increases, making more and more 2-PR free species available in the mixture due to the breaking of hydrogen bonds. These free alkanol species interact with the ions present in

Standard uncertainties u are $u(T) = \pm 0.01$ K, $u(x_1) = \pm 1 \times 10^{-4}$, $u(V_m) = \pm 0.05 \times 10^{-6}$ m³·mol⁻¹, $u(\kappa_s)$ $= \pm 0.05$ TPa⁻¹, u($\kappa_{s,m}$) = $\pm 0.05 \times 10^{-15}$ m³·Pa⁻¹· mol⁻¹, u(L_f) = $\pm 0.01 \times 10^{-11}$ m, u(*p*) = ± 0.002 MPa

Fig. 3 Excess molar volume V_{m}^{E} versus mole fraction of $[\text{Bmim}][\text{NTf}_2](x_1)$ in the binary mixture of $[\text{Bmim}][\text{NTf}_2](x_1)$ $[NTf₂] + 2-PR$ at 298.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue filled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green filled right sided triangle). The symbols represent experimental and the solid curves represent the values calculated from the Redlich–Kister equation (Eq. [11\)](#page-15-0) (Color figure online)

the dissociated ionic liquid at higher temperature. Therefore, at higher temperatures the ion–dipole interactions between unlike species (ions from ionic liquid and dipoles of 2-PR species) become more signifcant due to availability of a greater number of alkanol dipoles, leading to contraction in volume, thus decreasing the volume of the mixture [[18](#page-25-10)]. The reverse trend was observed in $[Bmim][NTf_2]$ rich region, i.e. that V_m^E values become more positive as the temperature increased. This trend can be attributed to the dispersion of forces and non-specifc interactions between the chemical species.

Isentropic compressibility (κ_s) was calculated from the Newton–Laplace equation as,

$$
\kappa_s = \frac{1}{u^2 \rho} \tag{4}
$$

Molar isentropic compressibility $(\kappa_{\rm sm})$ is given by the equation,

$$
\kappa_{s,m} = V_m \kappa_s \tag{5}
$$

The non-linear variation of isentropic compressibility as a function of composition of liquid mixture is sufficient evidence for the existence of molecular interactions in solu-tions [[25](#page-25-16)]. In this work, a non-linear variation of κ_s was observed with the mole fraction of [Bmim][NTf₂] in the liquid mixture with the values of κ_s increasing with increasing tem-perature (Fig. [4\)](#page-13-0). This phenomenon of κ ^{*s*} indicates the formation of a complex in the liquid mixture. The increasing values of $\kappa_{s,m}$ (Table [4\)](#page-10-0), with the mole fraction of [Bmim][NTf₂] and temperature, show that the intermolecular interactions are leading to the formation of a complex in the mixture.

Excess isentropic compressibility (κ_s^E) was calculated from the following relation [[26](#page-25-17)[–29\]](#page-26-0),

Fig. 4 Isentropic compressibility κ _{*s*} versus mole fraction of [Bmim][NTf₂] (x ₁) in the binary mixture of [Bmim][NTf₂]+2-PR at 298.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue filled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green flled right sided triangle) (Color fgure online)

$$
\kappa_s^{\mathcal{E}} = \kappa_s - \kappa_s^{\text{id}} = \kappa_s - \sum_{i=1}^2 \phi_i \left[\kappa_{s,i} + \frac{T V_i(\alpha_i^2)}{C_{p,i}} \right] - \left\{ \frac{T \left(\sum_{i=1}^2 x_i V_i \right) \left(\sum_{i=1}^2 \phi_i \alpha_i \right)^2}{\sum_{i=1}^2 x_i C_{p,i}} \right\}
$$
(6)

where φ_i , $\kappa_{s,i}$, V_i , α_i , $C_{p,i}$, and x_i represent the volume fraction, isentropic compressibility, molar volume, coefficient of isobaric thermal expansion, heat capacity and mole fraction of pure components, respectively. From the density values, the isobaric coefficient of thermal expansion was calculated from the following equation,

$$
\alpha = -\frac{1}{\rho} \left(\frac{\partial \rho}{\partial T} \right)_p \tag{7}
$$

The variation of κ_s^E with volume fraction of [Bmim][NTf₂] (ϕ_1) in the binary mixture is shown in Fig. [5.](#page-14-0) It was observed that at all studied temperatures κ_s^E values of [Bmim] $[NTf_2] + 2-PR$ are negative over the complete composition range. The negative κ_s^E values indicate the strong attractive interactions due to the solvation of ions in the mixture. Further, these negative values indicate that the solvent species around the solute are less compressible than the solvent species in the bulk solutions. Increasing the concentration of IL decreases the negative values of κ_s^E at all temperatures. This might be due to the decrease in attraction of 2-PR and $[Bmim][NTf_2]$ species in the $[Bmim][NTf_2]$ rich concentration

Fig. 5 Excess isentropic compressibility κ_s^E versus volume fraction of $[Bmim][NTf_2](\phi_1)$ in the binary mixture of $[Bmin][NTf₂] + 2-PR$ at 298.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue filled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green flled right sided triangle). The symbols represent experimental and the solid curves represent the values calculated from Redlich–Kister equation (Eq. [11\)](#page-15-0) (Color figure online)

region. These discrepancies may vary from IL to IL, solvent to solvent and also depend on the nature as well as structural arrangement of IL and solvent.

To understand the nature of intermolecular interactions, Jacobson suggested an empirical relation for calculating the intermolecular free length (L_f) [[30](#page-26-1)],

$$
L_f = K_j \sqrt{\kappa_s} \tag{8}
$$

where K_j is the Jacobson's constant, which depends on temperature. Its value is $(93.875 + 0.375T) \times 10^{-8}$

The intermolecular free length values of $[Bmim][NTf₂] + 2-PR$, investigated at various temperatures, are shown in Table [4.](#page-10-0) It is observed that, in the studied temperature range, the L_f values of 2-PR are greater than the L_f values of [Bmim][NTf₂]. But, as the mole fraction of $[Bmim][NTf₂]$ is increasing in the mixture the intermolecular free length values are decreasing. This behavior shows the presence of molecular interactions between the compounds of the liquid mixture. Further, as the temperature increases the L_f values increase. This increase in the intermolecular free length increases the distance between the surfaces of the two species, which in turn leads to the decrease in the speed of sound.

Excess intermolecular free length (L_f^E) was calculated from the following equation [\[31\]](#page-26-2),

$$
L_f^{\rm E} = L_f - K_j \sqrt{\kappa_s^{\rm id}} \tag{9}
$$

The variations of L_f^E can be attributed to various types of molecular interactions. The physical effects like dispersion forces or weak dipole-dipole interactions give positive $L_f^{\rm E}$ values and the chemical efects, like the formation of charge transfer complexes or other complex forming interactions, including strong dipole–dipole interactions between component species,

give negative L_f^E values [\[7\]](#page-24-4). In this work, the negative values of L_f^E (Fig. [6](#page-15-1)) designate specific interactions between different species in the binary mixture. The large negative values of L_f^E indicate structural readjustments in the mixture, which may result in a less compressible liquid phase and closer packing.

Excess speed of sound (u^E) was estimated from the following equation [[32\]](#page-26-3),

$$
u^{\mathcal{E}} = u - u^{\text{id}} = u - (\rho^{\text{id}} \kappa_s^{\text{id}})^{-1/2}
$$
 (10)

where $\rho^{\text{id}} = \sum_{i=1}^{2} \phi_i \rho_i$

Figure [7](#page-16-0) shows the variations of excess speed of sound with mole fraction and temperature. The sign and magnitude of u^E plays an important role in describing the molecular rearrangements occurring as a result of specifc interactions in the liquid mixture. Generally, the positive values of *u*^E indicate the presence of strong interactions and the negative values of \mathbf{u}^{E} indicate the presence of weak interactions between unlike species in the mixture [\[33](#page-26-4)]. In the present work, positive u^E values were observed for the entire composition range at all the studied temperatures, indicating the presence of strong ion–dipole interactions between the components of the binary mixture.

All the excess values $(V_{\text{m}}^{\text{E}}, \kappa_{s}^{\text{E}}, L_{f}^{\text{E}}$ and $u^{\text{E}})$ were fitted by the method of non-linear least squares to a Redlich–Kister polynomial equation of the type [\[34](#page-26-5)],

$$
Y^{\rm E} = x_1 x_2 \sum_{i=0}^{n} A_i (1 - 2x_1)^i \tag{11}
$$

Fig. 6 Excess intermolecular free length L_f^E versus mole fraction of [Bmim][NTf₂] (x_1) in the binary mixture of $[Bmim][NTf₂] + 2-PR$ at 298.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue filled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green flled right sided triangle). The symbols represent experimental and the solid curves represent the values calculated from Redlich–Kister equation (Eq. [11\)](#page-15-0) (Color figure online)

Fig. 7 Excess speed of sound u^E versus mole fraction of [Bmim][NTf₂] (x_1) in the binary mixture of [Bmim][NTf₂]+2-PR at 298.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue filled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green flled right sided triangle). The symbols represent experimental and the solid curves represent the values calculated from Redlich–Kister equation (Eq. [11](#page-15-0)) (Color figure online)

where x_1 is the mole fraction of [Bmim][NTf₂] and x_2 is the mole fraction of 2-PR. A_0 , A_1 , A_2 and A_3 are adjustable binary coefficients. The coefficients A_i were estimated using multi-parametric regression analysis based on the least-squares method. The number *n* of *Ai* parameters was optimized using the *F*-test. In each case, the optimum number of coefficients A_i was determined from an examination of the variation of standard deviation (σ) . The coefficients and standard deviations obtained from Redlich–Kister polynomial equation are provided in Table [5.](#page-17-0) The standard deviation values show satisfactory results for the experimental values.

The partial molar volumes, $\overline{V}_{m,1}^{\circ}$ of component 1 ([Bmim][NTf₂]) and $\overline{V}_{m,2}^{\circ}$ of component 2 (2-PR) in the mixture over the entire composition range were calculated using the following relations [[35\]](#page-26-6),

$$
\overline{V}_{m,1}^{\text{o}} = V_{m}^{\text{E}} + V_{m,1}^{*} + x_{2} \left(\frac{\partial V_{m}^{\text{E}}}{\partial x_{1}} \right)_{T,p}
$$
(12)

$$
\overline{V}_{\text{m},2}^{\text{o}} = V_{\text{m}}^{\text{E}} + V_{\text{m},2}^{*} - x_{1} \left(\frac{\partial V_{\text{m}}^{\text{E}}}{\partial x_{1}} \right)_{T,p} \tag{13}
$$

where $\overline{V}_{m,1}^*$ and $\overline{V}_{m,2}^*$ are the molar volumes of the pure components [Bmim][NTf₂] and 2-PR, respectively. The derivative $\left(\frac{\partial V_m^E}{\partial x_1}\right)$ in Eqs. 12 and 13 is obtained by differentiating) in Eqs. [12](#page-16-1) and [13](#page-16-2) is obtained by differentiating Eq. [11,](#page-15-0) which gives the following equations for $\overline{V}_{m,1}^{\circ}$ and $\overline{V}_{m,2}^{\circ}$.

Table 5 Coefficients (A_0, A_1, A_2, A_3) of the Redlich–Kister equation (Eq. [11\)](#page-15-0) and corresponding standard deviations (*𝜎*) for 1-butyl-3-methylimidazolium bis(trifuoromethanesulfonyl) amide+2-propoxyethanol from $T = (298.15 \text{ to } 323.15) \text{ K}$

Functions	A_0	A_1	A ₂	A_3	σ
$T = 298.15 K$					
$V_{\text{m}}^{\text{E}} \times 10^{6}$ (m ³ ·mol ⁻¹)	0.183	-10.569	3.636	6.381	0.0518
$\kappa_c^E \times 10^{12}$ (Pa ⁻¹)	-23.877	-54.221	-37.690		0.0170
$L_{\rm f}^{\rm E} \times 10^{13}$ (m)	-14.115	-23.518	-20.224	-12.208	0.0668
u^{E} (m·s ⁻¹)	37.726	27.923	49.356	48.226	0.1756
$T = 303.15$ K					
$V_{\text{m}}^{\text{E}} \times 10^6 \text{ (m}^3 \cdot \text{mol}^{-1})$	0.1227	-11.298	3.541	6.680	0.0324
$\kappa_s^E \times 10^{12}$ (Pa ⁻¹)	-34.855	-51.299	-34.145		0.0942
$L_f^{\rm E} \times 10^{13}$ (m)	-19.629	-25.564	-21.833	-5.224	0.0912
u^{E} (m·s ⁻¹)	50.206	28.526	51.384	32.112	0.2412
$T = 308.15$ K					
$V_{m}^{E} \times 10^{6}$ (m ³ ·mol ⁻¹)	0.218	-12.643	4.148	8.126	0.0514
$\kappa_s^E \times 10^{12}$ (Pa ⁻¹)	-44.704	-53.040	-25.160		0.0321
$L_f^{\rm E} \times 10^{13}$ (m)	-24.650	-27.448	-17.721	-1.839	0.0621
u^{E} (m·s ⁻¹)	61.232	28.679	40.524	21.612	0.1532
$T = 313.15 K$					
$V_{\text{m}}^{\text{E}} \times 10^6 \text{ (m}^3 \cdot \text{mol}^{-1})$	0.221	-13.612	4.384	9.034	0.0615
$\kappa_c^E \times 10^{12}$ (Pa ⁻¹)	-54.270	-58.940	-23.290		0.0182
$L_f^{\rm E} \times 10^{13}$ (m)	-29.613	-31.452	-18.137	0.302	0.1421
u^{E} (m·s ⁻¹)	72.143	32.108	41.364	22.423	0.3138
$T = 318.15 K$					
$V_{\text{m}}^{\text{E}} \times 10^6 \text{ (m}^3 \cdot \text{mol}^{-1})$	0.217	-15.071	4.914	10.123	0.0913
$\kappa_c^E \times 10^{12}$ (Pa ⁻¹)	-62.888	-65.019	-25.643		0.0722
$L_f^{\rm E} \times 10^{13}$ (m)	-35.672	-37.134	-18.792	7.122	0.1316
u^{E} (m·s ⁻¹)	81.789	40.143	38.561	9.541	0.2296
$T = 323.15$ K					
$V_{m}^{E} \times 10^{6}$ (m ³ ·mol ⁻¹)	0.198	-15.621	4.912	10.152	0.1023
$\kappa_c^E \times 10^{12}$ (Pa ⁻¹)	-72.737	-69.765	-27.840		0.0675
$L_{\rm f}^{\rm E} \times 10^{13}$ (m)	-38.504	-38.024	-27.324	-0.172	0.0926
u^{E} (m·s ⁻¹)	88.561	33.128	51.763	22.712	0.2352

$$
\overline{V}_{m,1}^{o} = V_{m,1}^{*} + x_2^{2} \sum_{i=0}^{n} A_i (1 - 2x_1)^{i} - 2x_1 x_2^{2} \sum_{i=1}^{n} i A_i (1 - 2x_1)^{i-1}
$$
(14)

$$
\overline{V}_{m,2}^{o} = V_{m,2}^{*} + x_{1}^{2} \sum_{i=0}^{n} A_{i} (1 - 2x_{1})^{i} + 2x_{1}^{2} x_{2} \sum_{i=1}^{n} i A_{i} (1 - 2x_{1})^{i-1}
$$
(15)

The excess partial molar volumes, $\overline{V}_{m,1}^{\text{oE}}$ and $\overline{V}_{m,2}^{\text{oE}}$, over the whole composition range were calculated by using the following relations,

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$$
\overline{V}_{m,1}^{\text{oE}} = \overline{V}_{m,1}^{\text{o}} - V_{m,1}^* \tag{16}
$$

$$
\overline{V}_{m,2}^{\text{oE}} = \overline{V}_{m,2}^{\text{o}} - V_{m,2}^* \tag{17}
$$

Another important property describing the behavior of the solution is the partial molar volume of solution at infinite dilution. By setting $x_2 = 1$ in Eq. [14](#page-17-1), we get the partial molar volume of $[Bmim][NTf₂]$ at infinite dilution in 2-PR.

$$
\overline{V}_{m,1}^{\text{o},\infty} = V_{m,1}^* + \sum_{i=0}^n A_i (1)^i
$$
 (18)

Similarly, by putting $x_2 = 0$ in Eq. [15,](#page-17-2) we get the partial molar volume of 2-PR at infinite dilution in $[Bmin][NTf₂]$.

$$
\overline{V}_{m,2}^{0,\infty} = V_{m,2}^* + \sum_{i=0}^n A_i (-1)^i
$$
 (19)

The excess partial molar volumes, $\overline{V}_{m,1}^{E,\infty}$ of [Bmim][NTf₂] and $\overline{V}_{m,2}^{E,\infty}$ of 2-PR at infinite dilution were determined using

$$
\overline{V}_{m,1}^{\text{oE},\infty} = \overline{V}_{m,1}^{\text{o},\infty} - V_{m,1}^*
$$
\n(20)

$$
\overline{V}_{m,2}^{oE,\infty} = \overline{V}_{m,2}^{o,\infty} - V_{m,2}^{*}
$$
 (21)

The values of $\overline{V}_{m,1}^{\circ}$ and $\overline{V}_{m,2}^{\circ}$ as functions of mole fraction of [Bmim][NTf₂] and temper-ature are given in Table [6](#page-19-0). We can observe that the values of $\overline{V}_{m,1}^{\circ}$ and $\overline{V}_{m,2}^{\circ}$, in the mixture, are diferent from their respective individual values in the pure states. The lower values indicate the contraction of volume and the higher values indicate the expansion of volume, upon mixing of [Bmim][NTf₂] with 2-PR. The variations of $\overline{V}_{m,1}^{\circ}$ and $\overline{V}_{m,2}^{\circ}$ with mole fraction of $[Bmim][NTf₂]$ and temperature are shown in Figs. [8](#page-21-0) and [9](#page-21-1) respectively. In general, the negative $\overline{V}_{m,1}^{\text{OE}}$ and $\overline{V}_{m,2}^{\text{OE}}$ values indicate the presence of significant solute–solvent interactions between unlike species and the positive $\overline{V}_{m,1}^{\text{OE}}$ and $\overline{V}_{m,2}^{\text{OE}}$ values indicate the presence of signifcant solute–solute or solvent–solvent interactions between like species in the mix-ture [\[36\]](#page-26-7). We can observe from Figs. [8](#page-21-0) and [9](#page-21-1) that the values of $\overline{V}_{m,1}^{\text{oe}}$ and $\overline{V}_{m,2}^{\text{oe}}$ are negative for the lower concentrations of $[Bmim][NTf_2]$ and positive for the higher concentrations of [Bmim][NTf₂] in the mixture. The negative $\overline{V}_{m,1}^{oE}$ and $\overline{V}_{m,2}^{oE}$ values indicate that [Bmim] [NTf₂]–2-PR interactions are stronger than [Bmim][NTf₂]–Bmim][NTf₂] or 2-PR–2-PR interactions and the positive $\overline{V}_{m,1}^{\text{oE}}$ and $\overline{V}_{m,2}^{\text{oE}}$ values indicate that [Bmim][NTf₂]–2-PR interactions are weaker than $[\text{Bmin}][\text{NTf}_2] - [\text{Bmin}][\text{NTf}_2]$ or 2-PR–2-PR interactions.

The values of $\overline{V}_{m,1}^*$, $\overline{V}_{m,1}^{o,\infty}$, $\overline{V}_{m,2}^*$, $\overline{V}_{m,2}^{o,\infty}$ and $\overline{V}_{m,2}^{o,\infty}$ for the pure compounds [Bmim] $[NTf₂]$ and 2-PR, at the investigated temperatures, are listed in Table [7](#page-22-0). The negative

Standard uncertainties u are u(*T*)=±0.01 K, u(*x*₁) =±1×10⁻⁴, u($\overline{V}_{m,1}^{\circ}$)
=±0.05×10⁻⁶ m³·mol⁻¹, u($\overline{V}_{m,2}^{\circ}$) =±0.05×10⁻⁶ m³·mol⁻¹, u(*p*) $=\pm 0.002$ MPa

Table 6 (continued)

Fig. 8 Excess partial molar volume $\overline{V}_{m,1}^{OE}$ of [Bmim][NTf₂] versus mole fraction of [Bmim][NTf₂] (*x*₁) in the binary mixture of $[Bmim][NIT_2]+2-PR$ at 298.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue flled triangle), 313.15 K (green flled inverted triangle), 318.15 K (pink flled left sided triangle) and 323.15 K (light green flled right sided triangle) (Color fgure online)

Fig. 9 Excess partial molar volume $\overline{V}_{m,2}^{oE}$ of 2-PR versus mole fraction of [Bmim][NTf₂] (x_1) in the binary mixture of $[Bmin][NTf₂] + 2-PR$ at $2\overline{98}$.15 K (filled square), 303.15 K (red filled circle), 308.15 K (blue filled triangle), 313.15 K (green filled inverted triangle), 318.15 K (pink filled left sided triangle) and 323.15 K (light green flled right sided triangle) (Color fgure online)

values of $\overline{V}_{m,1}^{\text{oE},\infty}$ indicate that the molar volume of [Bmim][NTf₂] in the mixture is smaller than its molar volume in the pure state, i.e., there is a contraction in the volume on mixing $[Bmin][NTf_2]$ with 2-PR. The increase in the negative values of $\overline{V}_{m,1}^{\text{DE},\infty}$ with the increase in temperature is attributed to the expansion in volume of the mixture, which allows more

Table 7 Molar volumes of 1-butyl-3-methylimidazolium bis(trifluoromethanesulfonyl) amide $(V_{m,1}^*)$, 2-propoxyethanol ($V_{m,2}^*$), partial molar older at infinite dilugion for 1-butyl-3-methylimidazolium
bis(trifluoromethanesulfonyl) amide ($\overline{V}_{m,1}^*$), 2-propoxyethanol ($\overline{V}_{m,2}^*$) and excess partial molar volu infinite dilution for 1-butyl-3-methylimidazolium bis(trifluoromethanesulfonyl) amide $(\overline{V}_{m,1}^{\text{ob},\infty})$, 2-propoxyethanol $(\overline{V}_{m,2}^{oE,\infty})$ from *T*=(293.15 to 323.15) K at pressure *p*=0.1 MPa

T/K	$V_{\text{m,d}}^* \times 10^6$ (m ³ ·mol ⁻¹)	$\overline{V}_{m,1}^{0,\infty} \times 10^6$ (m ³ ·mol ⁻¹⁾	$\overline{V}^{oE,\infty}$ $\times 10^6$ $\binom{m}{m}$ mol ⁻¹)	$V_{\text{m2}}^* \times 10^6$ (m ³ ·mol ⁻¹)	$\frac{\overline{V}_{m,2}^{0,\infty} \times 10^6}{(m^3 \cdot mol^{-1})}$	$\overline{L^2}^{\text{OE},\infty}$ $\times 10^{6}$ $\binom{m}{m^3 \cdot mol^{-1}}$
298.15	291.81	291.34	-0.47	114.80	122.87	8.06
303.15	292.75	292.27	-0.48	115.36	123.53	8.17
308.15	293.79	293.26	-0.53	115.92	124.72	8.80
313.15	294.74	294.17	-0.56	116.49	125.64	9.16
318.15	295.77	295.07	-0.69	117.06	126.41	9.35
323.15	296.67	295.92	-0.75	117.64	127.94	10.30

Standard uncertainties u are $u(T) = \pm 0.01$ K, $u(x_1) = \pm 1 \times 10^{-4}$, $u(V_{m,1}^*, V_{m,2}^*) = \pm 0.05 \times 10^{-6}$ m³·mol⁻¹, $u(\overline{V}_{m,1}^{0,\infty}, \overline{V}_{m,2}^{0,\infty}) = \pm 0.05 \times 10^{-6} \text{ m}^3 \cdot \text{mol}^{-1}, u(\overline{V}_{m,1}^{0,E,\infty}, \overline{V}_{m,2}^{0,E,\infty}) = \pm 0.05 \times 10^{-6} \text{ m}^3 \cdot \text{mol}^{-1}, u(p) = \pm 0.002 \text{ MPa}$

favorable fitting of smaller 2-PR species into the voids created by larger $[\text{Bmin}][\text{NTf}_2]$ species at higher temperature.

The imidazolium salts contain a delocalized 3-center-4-electron confguration across the N₁–C₂–N₃ moiety, a double bond between C₄ and C₅ at the opposite side of the ring and a weak delocalization in the central region [[37](#page-26-8)]. The hydrogen atoms C_2-H , C_4-H and C_5 –H carry almost the same charge, but the carbon C_2 is positively charged owing to the electron deficit in the C=N bond, whereas C_4 and C_5 are practically neutral. The resulting acidity of the hydrogen atoms is the key to understand the properties of these ionic liquids (Scheme [1\)](#page-22-1).

FT-IR spectra of pure 2-PR, pure $[Bmim][NTf₂]$ and also for different mole fractions of [Bmim][NTf₂] $(x_1=0.2980, 0.4901$ and 0.8003) in the binary mixture were recorded from 500 to 3500 cm−1 (Fig. [10\)](#page-23-0). In order to study the efects of molecular interactions on the [Bmim] cation, the CH stretching region between 2800 and 3200 cm⁻¹ was analyzed. The signals between 2800 and 3000 cm^{-1} result from aliphatic CH groups in the butyl moiety [[38](#page-26-9)[–42\]](#page-26-10). The signals between 3000 and 3200 cm⁻¹ can be assigned to CH modes originating from the aromatic imidazolium ring and also from C_2 –H to $C_{4,5}$ –H stretching vibrations [[41](#page-26-11)]. The addition of 2-PR to the ionic liquid leads to a blue shift in imidazolium C–H

Scheme 1 Schematic structures for **a** [Bmim][NTf₂] and **b** 2PE

Fig. 10 FT-IR spectra of **a** pure 2-PR, **b** $x_1 = 0.2980$, **c** $x_1 = 0.4901$, **d** $x_1 = 0.8003$ and **e** pure [Bmim][NTf₂]

stretching vibrations. Similar molecular interactions were observed in the binary mixtures of imidazolium ionic liquids with other molecular solvents [\[6](#page-24-5), [43](#page-26-12)[–45\]](#page-26-13). In the case of the [NTf₂] anion, FT-IR spectra in the range from 630 to 680 cm⁻¹ were analyzed. The peak at 653 cm⁻¹ in the pure [Bmim][NTf₂] can be assigned to the S–N–S bending mode. The transmittance decreased with increasing mole fractions of $[Bmim][NTf₂]$, indicating the presence of interactions between $[NTf₂]$ anions and dipoles of 2PE.

The interaction mechanism of 2PE molecules with the cation and anion is diferent in nature. The C–H groups in the [Bmim] cation and nitrogen in $[NTF_2]$ anion are characteristic groups interacting with 2-PR molecules. By studying the ATR-FTIR, we can predict that the addition

of $[Bmim][NT_2]$ to 2-PR influenced (i) C-H vibrations of $[Bmim]$ ⁺ cation; (ii) S-N-S bending mode of [NTf₂] anion; (iii) O–H stretching of 2-PR. Thus the FT-IR spectroscopy study indicates ion–dipole interactions in the binary mixture of $[Bmim][NTf₂]$ and 2-PR.

4 Conclusions

Densities and speeds of sound were measured for the pure and binary mixture of [Bmim] [NTf₂]+2-PR for various mole fractions at temperatures from $T = (298.15 \text{ to } 323.15) \text{ K}$ under atmospheric pressure. From these experimental values, various acoustical and thermodynamic parameters were calculated. The intermolecular interactions between the binary mixture of $[Bmin][NT₂]$ and 2-PR were analyzed by studying the variations in the parameters determined. The negative excess molar volume indicate that more efficient packing are taking place when $[Bmin][NTf_2]$ was added with 2-PR. The negative values of excess intermolecular free length designated the structural readjustments in the mixture leading to a less compressible liquid phase and closer packing of species. The positive excess speed of sound values designated the presence of strong ion–dipole interactions between the components of the binary mixture. Satisfactory correlations for the excess parameters were obtained by ftting them with the Redlich–Kister polynomial equation. The FT-IR spectra analysis indicated the formation of hydrogen bonding between $[Bmin][NTf₂]$ and 2-PR.

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Afliations

G. Prasad^{1,5} · K. Muralidhar Reddy² · R. Padamasuvarna¹ · T. Madhu Mohan³ · **T. Vijaya Krishna³ · S. Govardhana Rao4**

- ¹ Department of Physics, JNTUA-College of Engineering, Ananthapuramu, Andhra Pradesh, India
- ² Department of EIE, R.G.M. College of Engineering & Technology, Nandyal, Andhra Pradesh, India
- ³ Department of Physics, Vasireddy Venkatadri Institute of Technology, Nambur, Andhra Pradesh, India
- ⁴ Department of Physics, St. Mary's Group of Institutions Guntur, Guntur, Andhra Pradesh, India
- ⁵ Department of Physics, SASI Institute of Technology & Engineering, Tadepalligudem, Andhra Pradesh, India