

# An Experimental and Computational Study of 2-(3-Oxo-3H-benzo[f] chromen-1-ylmethoxy)-Benzoic Acid Methyl Ester

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Abstract The various methods for studying polarities are based on the use of probe molecules, whose molecular spectral profile is significantly affected by the polarity of the medium. The absorption and emission spectra and dipole moments ( $\mu_{\rm g}$  and  $\mu_{\rm e}$ ) of 2-(3-oxo-3H-benzo[f]chromen-1-ylmethoxy)-benzoic acid methyl ester (2BME) are studied in solvents of different polarities at room temperature. The determination of dipole moments by solvatochromic shift using various relations and the change in dipole moment  $(\Delta \mu)$  were determined using Stokes shift with the variation of the solvent polarity parameter  $(E_{\text{T}}^{\text{N}})$ . The value of  $\mu_e$  greater than  $\mu_g$  indicating that the probe is more polar in the higher state. DFT and TDDFT theoretical analysis of dipole moment in the vacuum and with solvent, solvent accessible surface (SAS) and molecular electrostatic potential (MEP) are also performed.

Keywords Solvatochromic method · Polarity function · Dipole moments · DFT calculations - HOMO–LUMO

## 1 Introduction

Coumarin and its derivatives [\[1\]](#page-18-0) have attracted considerable attention in recent years for their versatile properties in chemistry and pharmacology. The 3-carboxamide coumarins [[2\]](#page-18-0) have been described as small molecular weight FXIIa inhibitors. Synthetic coumarins

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[[3\]](#page-18-0) have been found to be active against the pancreatic cancers. These natural coumarins [[4\]](#page-18-0) isolated from Angelica decursiva and Artemisia capillaris have shown anti-Alzheimer's disease factors. Coumarins [[5](#page-18-0)] which have been in isolated from roots of Angelica dahurica cv. Hangbaizhi showed significant anti-inflammatory action; coumarins have some ability to regulate the diverse range of cellular pathways useful for some anticancer treatments [\[6\]](#page-18-0). Drug–metal interactions have been used to develop highly sensitive biosensors for bacteria [\[7](#page-18-0), [8](#page-18-0)].

The excited state dipole moment determination using the Lippert–Mataga equation is adopted widely [[9\]](#page-18-0); in this method, absorption and emission shifts, dielectric constant  $(\varepsilon)$ and the refractive index (n) are used. Ground ( $\mu_{\rm g}$ ) and excited ( $\mu_{\rm e}$ ) state dipole moments of a probe vary due to changes in the polarity or dielectric constant of the solvent, which leads to variations in  $\mu_{\rm g}$  and  $\mu_{\rm e}$ . Analysis of the effect of solvent gives more information and is very useful in studying  $\mu_e$  of the molecule. These studies are helpful in designing nonlinear optical devices [\[10\]](#page-18-0), probing the nature of photochemical transformations and the calculation of molecular polarizability [\[11\]](#page-18-0). The empirical results for  $\mu_e$  are also helpful for quantum chemical (QC) analysis.

We have studied the  $\mu_{\rm g}$  and  $\mu_{\rm e}$  of exalites [\[12](#page-18-0), [13\]](#page-18-0) and coumarin laser dyes [[14](#page-18-0), [15](#page-18-0)] using Guggenheim's and solvatochromic shift methods. These methods which were reported [[13](#page-18-0)–[15](#page-18-0)] a decade ago have been adopted by many researchers. Several research groups have made empirical and theoretical consideration on  $\mu_{\rm g}$  and  $\mu_{\rm e}$  using varoius methods on organic molecules like coumarin  $[16]$ , ketocyamine dye  $[17]$  and others [[18–24\]](#page-19-0).

It is known that the position of the bands of a probe in solutions varies due to specific and nonspecific interactions with the probe. The solvent induced spectral shift of a probe is difficult to understand because of the difficulty with the theoretical definition of the solvent polarity; for this reason different solvent polarity parameters and scales are used. Density functional theory (DFT) using the B3LYP functional and 6-311 basis set method was used to study the hydrogen bond by considering the difference between the structures with and without the bond; DFT predicts structural properties, lowest energy of the probe in gas phase and in the solvent.

In the present context, the photophysical properties of the 2BME molecule have been studied in different solvents of varying polarity and hydrogen bonding characteristics. The Stokes data of 2BME is used to determine  $\mu_e$  values of the singlet form, which have been verified using Lippert and Mataga bulk solvent functions [\[25,](#page-19-0) [26\]](#page-19-0) given by Kamlet et al. [[27–29\]](#page-19-0). The  $\mu_{\rm g}$  of 2BME was determined by the QC method and  $\mu_{\rm e}$  was measured using Lippert–Mataga (LME), Bakshiev's (BE) [\[30\]](#page-19-0), Kawski–Chamma–Viallet's (KCVE) and McRae (McRE) [\[31–36\]](#page-19-0) relations and an equation based on Reichardt's microscopic solvent polarity parameter  $E_{\text{T}}^{\text{N}}$  [[37](#page-19-0)]. The QC parameters have been adapted to analyze dipole moment, HOMO, LUMO, MEP and SAS of 2BME in vacuum and in ethanol as solvent [\[38\]](#page-19-0).

## 2 Materials and Methods

### 2.1 Materials

The synthesis of 2-(3-oxo-3H-benzo[f]chromen-1-ylmethoxy)-benzoic acid methyl ester (2BME) was carried out by standard methods (Fig. [1](#page-2-0)) with a 97% yield [\[39\]](#page-19-0).

<span id="page-2-0"></span>Spectroscopic grade (SD Fine Chem. Ltd. India) solvents: acetone, ethanol, ethyl acetate, hexane, methanol, N,N-dimethylformamide (N,N-DMF), tetrahydrofuran, propanol, chloroform and butyl acetate, having 99% purity, were used without further purification. The concentrations of 2BME solutions were  $10^{-4}$ – $10^{-5}$  mol $\cdot$ dm<sup>-3</sup>. Dielectric constants  $\epsilon$ ) and refractive indices  $(n)$  of the solvents obtained from the literature  $[40-43]$  $[40-43]$  $[40-43]$  are presented in the Table [1](#page-3-0).

# 2.2 Experimental Methods

UV–vis absorption spectra (Shimadzu UV-1800 spectrophotometer) over a wavelength range of 200–800 nm and steady-state fluorescence spectra were obtained (Hitachi F-2700 Fluorescence Spectrophotometer) by selecting excitation wavelengths 319 and 347 nm. The measurements were carried out using  $1.0 \times 1.0$  cm quartz cell sample holder, all the experimental measurements were done at 27.00 °C.

## 2.3 Theoretical Background

From the literature, [[12](#page-18-0)–[15](#page-18-0)] the determination of the excited state dipole moments of different molecules by the solvatochromic shift method, using different assumptions, results in four independent equations for the estimation of the  $\mu_{e}$  for 2BME as follows:





<span id="page-3-0"></span>

Table 1 Values of solvent polarity parameters

<span id="page-4-0"></span>Lipper–Mataga equation [[44](#page-20-0)]

$$
\tilde{v}_a - \tilde{v}_f = m_1 F_1(\varepsilon, n) + \text{ constant} \tag{1}
$$

where  $\tilde{v}_a$  and  $\tilde{v}_f$  are the absorption and fluorescence maxima and

$$
F_1(\varepsilon, n) = \frac{\varepsilon - 1}{2\varepsilon + 1} \frac{n^2 - 1}{2n^2 + 1}.
$$

Bakshiev's equation [\[45\]](#page-20-0)

$$
\tilde{v}_a - \tilde{v}_f = m_2 F_2(\varepsilon, n) + constant \tag{2}
$$

where  $F_2(\varepsilon, n) = \frac{2n^2+1}{n^2+2} \left[ \frac{\varepsilon - 1}{\varepsilon + 2} - \frac{n^2 - 1}{n^2 + 2} \right]$ .

Kawski–Chamma–Viallet equation [[46](#page-20-0)]

$$
\frac{\tilde{v}_a + \tilde{v}_f}{2} = m_3 F_3(\varepsilon, n) + \text{ constant}
$$
\n(3)

where  $F_3(\varepsilon, n) = \left| \frac{2n^2 + 1}{2(n^2 + 2)} \left( \frac{\varepsilon - 1}{\varepsilon + 2} - \frac{n^2 - 1}{n^2 + 2} \right) \right|$  $\left(\frac{\varepsilon-1}{\varepsilon+2} - \frac{n^2-1}{n^2+2}\right) + \frac{3(n^4-1)}{2(n^2+2)^2}$  $2(n^2+2)^2$  $\begin{bmatrix} 1 & 1 & 1 \end{bmatrix}$ McRae equation [[47](#page-20-0)]

$$
\tilde{v}_a = -m_4 F_4(\varepsilon) + \text{ constant} \tag{4}
$$

where  $F_4(\varepsilon) = \left[\frac{2(\varepsilon - 1)}{\varepsilon + 2}\right]$ .

The above equations use quantum mechanical perturbation theory for  $\tilde{v}_a$ ,  $\tilde{v}_f$  and Stokes shifts in various solvents, where:

:

$$
m_1 = \frac{2(\mu_e - \mu_g)^2}{hca^3}
$$
 (5)

$$
m_2 = \frac{2(\mu_e - \mu_g)^2}{hca^3} \tag{6}
$$

$$
m_3 = \frac{2\left(\mu_{\rm e}^2 - \mu_{\rm g}^2\right)}{hca^3} \tag{7}
$$

$$
m_4 = \frac{\mu_{\rm g}(\mu_{\rm e} - \mu_{\rm g})}{hca^3} \tag{8}
$$

where c is the velocity of light in vacuum and h is Planck's constant. The  $m_1$ ,  $m_2$ ,  $m_3$  and  $m_4$  are the slopes of the plots of (1)  $(\tilde{v}_a - \tilde{v}_f)$  against  $F_1(\varepsilon, n)$ , (2)  $(\tilde{v}_a - \tilde{v}_f)$  against  $F_2(\varepsilon, n)$ , (3)  $(\frac{\tilde{v}_a - \tilde{v}_f}{2})$  against  $F_3(\varepsilon, n)$  and (4)  $\tilde{v}_a$  against  $F_4(\varepsilon)$ , respectively. The Onsager radius, a, of a solute molecule was estimated according to the Edward atomic increment method [[48](#page-20-0)]. If the  $\mu_{\rm g}$  and  $\mu_{\rm e}$  dipole moments are parallel, then:

$$
\mu_{\rm g} = \frac{m_3 - m_2}{2} \left[ \frac{hca^3}{2m_1} \right]^{\frac{1}{2}} \tag{9}
$$

$$
\mu_{\rm e} = \frac{m_3 + m_2}{2} \left[ \frac{hca^3}{2m_1} \right]^{\frac{1}{2}} \tag{10}
$$

$$
\mu_{\rm e} = \left[\frac{m_3 + m_2}{m_3 - m_2}\right] \mu_{\rm g}; \, m_3 > m_2. \tag{11}
$$

<span id="page-5-0"></span>If the  $\mu_{\rm e}$  and  $\mu_{\rm e}$  are non-parallel, forming an angle  $\varphi$ , then  $\varphi$  can be estimated using the formula:

$$
\cos \varphi = \frac{1}{2\mu_{g}\mu_{e}} \left[ \left( \mu_{e}^{2} + \mu_{g}^{2} \right) - \frac{m_{3}}{m_{2}} \left( \mu_{e}^{2} - \mu_{g}^{2} \right) \right]
$$
(12)

The slopes are  $m_1$ ,  $m_2$  and  $m_3$  and are obtained by plotting  $(\tilde{v}_a - \tilde{v}_f)$  and  $(\frac{\tilde{v}_a - \tilde{v}_f}{2})$  against solvent polarity functions.

# 2.3.1 Molecular Microscopic Solvent Polarity Parameter  $E_T^N$

Reichardt gave [[49](#page-20-0)] the polarity parameter  $E_{\text{T}}^{\text{N}}$  and using this the excited state dipole moment of a molecule can be calculated as:

$$
\tilde{v}_a - \tilde{v}_f = 11307.6 \left[ \left( \frac{\delta \mu}{\delta \mu_B} \right)^2 \left( \frac{a_B}{a} \right)^3 \right] E_T^N + \text{ constant}
$$
\n(13)

where  $\delta \mu_B$  is the variation in dipole moment and  $a_B$  the Onsager radius of the betaine molecule and  $\delta \mu$  and a the values for the solute molecule being studied. The value of  $\delta \mu$ can be determined as:

$$
\delta\mu = \mu_{\rm e} - \mu_{\rm g} = \sqrt{\frac{81m}{11307.6(6.2/a)^3}}
$$
(14)

where *m* is the slope of linear plot of Stoke shift against  $E_{\text{T}}^{\text{N}}$ .

#### 2.3.2 Kamlet–Taft Solvatochromic Parameters

The multiparametric approach was proposed by Kamlet and co-workers [\[50\]](#page-20-0). The magnitude of solute–solvent interactions are estimated using a solvatochromic method. The solvent effect on some measurable property A is assumed to be linearly dependent on three parameters of the medium: the polarizability parameter  $\pi^*$ , the hydrogen bond donor (HBD) parameter  $\alpha$ , and hydrogen bond acceptor parameter  $\beta$ . The effect of the solvent on the absorption, A, is expressed by the generalized equation:

$$
A = A_0 + s\pi^* + a\alpha + b\beta \tag{15}
$$

where  $A_0$  is the vapor phase wavenumber and the values of A are spectral band maxima in the different media. The Kamlet–Taft coefficients,  $s$ ,  $a$  and  $b$  give the relative sensitivities of A to the individual parameters.

#### 2.4 Computational Details

The theoretical simulations using the Gamess [\[51\]](#page-20-0) and optimization of geometry was done with DFT with the B3LYP functional and  $6-311$  (d,p) basis set. The vertical excitations, density graph and optimization of the  $\mu_{g}$  were analyzed with the help of TDDFT (timedependent DFT) at the B3LYP/6-311 (d,p) level. The HOMO and LUMO were calculated in vacuum and in ethanol as solvent.

## 3 Results and Discussion

#### 3.1 Effect Solvent on Absorption and Fluorescence Spectra

Absorption (Fig. 2) and fluorescence (Fig. [3\)](#page-7-0) studies were done for the 2BME molecule at room temperature. The absorption peak of 2BME varies from 231 to 348 nm in solvent medium which leads to  $\pi-\pi^*$  transition. The Stokes shifts values in different solvents indicate  $\mu_{\rm o}$  of the sample in protic solvents.

The fluorescence spectra of 2BME is red-shifted between 285 and 427 nm from hexane to N,N-DMF. The effect of solvent on the fluorescence emission is much greater than on the absorption maximum. These results suggest that the ground state is less polar than the higher emission state [\[52–54](#page-20-0)]. The linearity between the  $E_T(30)$  and the Stokes shift was analyzed. The molecule 2BME show increases in Stokes shift from aprotic to protic medium due to increase in the solvent polarity and intramolecular charge transfer state (ICT).

The nature of 2BME molecule in solvents with various polarities may be explored using the differences in  $\mu_{\rm g}$  and  $\mu_{\rm e}$ . The charge separation on electronic excitation of 2BME is measured by analyzing the change dipole moment ( $\Delta \mu = \mu_e - \mu_g$ ). According to Lippert– Mataga (Eq. [1](#page-4-0)) [[55](#page-20-0)]:

$$
\Delta \tilde{v} = \tilde{v}_a - \tilde{v}_f = \left[ \frac{2(\mu_e - \mu_g)}{hca^3} \right] F_1(\varepsilon, n) + \text{ constant}
$$
 (16)

where  $F_1(\varepsilon, n) = \frac{\varepsilon - 1}{2\varepsilon + 1} \frac{n^2 - 1}{2n^2 + 1}$ , which shows the orientation polarizability and polarity parameter of the solvent; Fig. [4](#page-7-0) represents the Lippert–Mataga linear plot for various solvents.



Fig. 2 Absorption spectra of 2BME in different solvents

<span id="page-7-0"></span>

Fig. 3 Fluorescence spectra of 2BME in different solvents



**Fig. 4** Plot of Stokes shift versus solvent polarity function  $F_1$  ( $\varepsilon$ ,  $n$ ) using the Lippert equation for 2BME

The linear plots of the Bakshiev and Chamma–Viallet equations,  $E_T(30)$  against  $(\tilde{v}_a - \tilde{v}_f)$ ,  $E_T^N$  against  $(\tilde{v}_a - \tilde{v}_f)$  and the McRae equation are shown in Figs. [5,](#page-8-0) [6,](#page-8-0) [7](#page-9-0), [8](#page-9-0) and [9](#page-10-0).

The Kamlet and Taft [[56\]](#page-20-0) method was used to explore the effects of the various modes of solvation on the absorption and fluorescencent energies. The parameters  $\pi^*$ ,  $\alpha$  and  $\beta$  are used to characterize the polarizability and to describe the respective properties of a given solvent in Eq. [15](#page-5-0). Table [2](#page-10-0) lists the absorption and emission bands of 2BME. The solvent dipolar interaction  $(\pi^*)$  (Table [1](#page-3-0)) and hydrogen bond accepting property ( $\beta$ ) are the principal parameters which lead to the stabilization of the  $\mu_{g}$  and  $\mu_{e}$  of the 2BME molecule.

<span id="page-8-0"></span>

Fig. 5 Plot of Stokes shift versus solvent polarity function  $F_2$  (e, n) using the Bakshieve equation for 2BME



Fig. 6 Plot of Stokes shift versus solvent polarity function  $F_3(\varepsilon, n)$  using Kawski–Chamma–Villet's equation for 2BME

### 3.2 Estimation of Dipole Moments

The solvent parameters, solvent functions (Table [1\)](#page-3-0) and values of  $\mu_{\rm g}$  and  $\mu_{\rm e}$  of 2BME are determined using Eqs. [1–4](#page-4-0) as shown in Table [2](#page-10-0).

Plots of  $(\tilde{v}_a - \tilde{v}_f)$  against  $F_1(\varepsilon,n)$ ,  $(\tilde{v}_a - \tilde{v}_f)$  against  $F_2(\varepsilon, n)$ ,  $(\tilde{v}_a - \tilde{v}_f)/2$  against  $F_3(\varepsilon, n)$ ,  $(\tilde{v}_a - \tilde{v}_f)$  against  $E_T^N$ ,  $\tilde{v}_a$  against  $F_4(\varepsilon)$  $F_4(\varepsilon)$  $F_4(\varepsilon)$  and  $(\tilde{v}_a - \tilde{v}_f)$  against  $E_T(30)$  are shown in Figs. 4, 5, 6, [7,](#page-9-0) [8](#page-9-0) and [9](#page-10-0) and the intercepts, slopes and correlation coefficients are shown in Table [3](#page-10-0).

<span id="page-9-0"></span>

Fig. 7 Plot of Stokes shift versus solvent polarity function  $E_T^N$  using the Reichardt equation for 2BME



Fig. 8 Plot of Stokes shift versus solvent polarity function  $F_4(\varepsilon)$  using the McRae equation for 2BME

The plots are linear for eight solvents while other points varied from linearity due to specific solute–solvent interactions.

From Table [3,](#page-10-0) the r values vary from 0.8162 to 0.9933 depending on the solvents used, which represents the linearity for these correlations. The linear behavior of the Stokes shifts versus solvent polarity function indicates general solvent effects as a function of dielectric constant and refractive index. The Onsager cavity radii of 2BME was measured

<span id="page-10-0"></span>

Fig. 9 Plot of Stokes shift versus solvent polarity function  $E_T(30)$  for 2BME

Solvent	$\lambda_{\rm a}$ (nm)	$\lambda_e$ (nm)	$\tilde{v}_a$ (cm <sup>-1</sup> )	$\tilde{v}_f$ (cm <sup>-1</sup> )	$\tilde{v}_a - \tilde{v}_f$ (cm <sup>-1</sup> )
Hexane	262	285.0	38138.8	35087.7	3051.1
Chloroform	304	417.5	32954.4	23952.1	9002.3
Ethyl acetate	347	421.5	28843.4	23724.8	5118.6
Butyl acetate	270	419.7	37064.1	23821.3	13243.0
Tetrahydrofuran	348	427.0	28727.4	23419.2	5308.2
Acetone	347	422.5	28835.1	23668.6	5166.4
Propanol	231	416.0	43384.0	24038.5	19345
Ethanol	268	411.0	37369.2	24330.9	13038
Methanol	319	356.0	31348.0	28089.9	3258.1
$N.N$ DMF	269	328.5	37174.7	30441.4	6733.3

Table 2 Experimental absorption and fluorescence properties of 2BME in different solutions

Table 3 Statistical treatment of the correlations of the spectral shift of 2BME

Equations of the form equation	Slope $(m)$	Intercept	Correlation coefficient (r)	No. of data
Lippert-Mataga correlation	$m_1 = 11637.97$	3122.72	0.8162	10
Bakhshieve correlation	$m_2 = 3248.76$	3846.29	0.8370	8
Kawski–Chamma–Viallet	$m_3 = 19311.13$	39140.0	0.8447	11
$E_{\rm T}^{\rm N}$ correlation	$m_A = 2673.746$	4650.57	0.8562	8
McRae	$m_5 = 2.20695$	$-0.27577$	0.9933	12
$E_{\rm T(30)}$	$m_6 = 843.25$	4650.57	0.863	8

<span id="page-11-0"></span>by the atomic increment method and found to be 4.16  $\AA$ . The ground state dipole moment is determined from the  $m_2$  and  $m_3$  slopes of the BE and KCVE methods, using Eqs. [6](#page-4-0) and [7](#page-4-0) and also the  $\mu_e$  values are obtained from  $m_1$ ,  $m_2$ ,  $m_3$  and  $m_4$  of LME, BE, KCVE, MRE and  $E_{\rm T}^{\rm N}$  by applying Eqs. [9](#page-4-0), [10](#page-4-0) and [14](#page-5-0) and the results are presented in Table 4.

Table 4 indicates that  $\mu_e$  (2.553 D) is found to be 1.13 times larger than  $\mu_g$  (2.266 D).  $\mu_{g}$ ; differences are in the range of 2.411–7.0871 D, which shows that the molecule is more polar in the excited state than in the ground state. The differences between  $\mu_{\rm g}$  and  $\mu_{\rm e}$  can give the properties of the emitting state, charge transfer and change occurring under excitation. The  $\mu_e$  value of 7.0878 D obtained from BE is higher than the values from the other equations. The angle between the  $\mu_{g}$  and  $\mu_{e}$  vectors calculating using Eq. [12](#page-5-0) is found to be 9.236°, its direction depends on the center of charges. The small change in angle results in the higher charge distribution across the molecule in the excited state.

The linear correlation values are greater than 0.8 (Table [3\)](#page-10-0), which gives linearity of the Stokes shift; the deviations from linearity could be from solute–solvent interactions. For the 2BME molecule the interaction with nonpolar solvents is due to the dipole–induced dipole interactions, the solute–solvent interaction depend on the higher dipole–dipole interactions (also pointed out by others [[57](#page-20-0), [58](#page-20-0)]) to use  $E_T$  (30) parameter, which is the experimental solvent parameter [\[43\]](#page-19-0) for the polarization dependent spectral behavior. However,  $E_T(30)$  data have dimensions of kJ·mol<sup>-1</sup> [\[40\]](#page-19-0); hence, the normalized  $E_T^N$  data has been considered. The  $\mu_{g}$  analysis by the B3LYP/6-311 (d,p), PCM and SCF techniques

Compound	2BME
Onsager radius Å	4.160
$\mu_{\rm g}^{\rm a}$ (D)	2.266
$\mu_{\rm e}^{\rm b}$ (D)	2.553
$\mu_{\rm e}^{\rm c}$ (D)	7.087
$\mu_{\rm e}^{\rm d}$ (D)	2.496
$\mu_{\rm e}^{\rm e}$ (D)	4.397
$\mu_{\rm e}^{\rm f}$ (D)	2.411
$\Delta \mu_{\rm e}^{\rm g}$ (D)	0.286
$\frac{\mu_{\rm e} {\rm h}}{\mu_{\rm g}}\left({\rm D}\right)$	0.887
$\Delta\mu_{\rm e}^{\rm i}$ (D)	2.405
$\varphi^j$ (rad)	0.161

Table 4 Ground state and singlet excited state dipole moments of 2BME in Debye (D)

<sup>a</sup> Ground state dipole moment from quantum chemical calculation

<sup>b</sup> Excited state dipole moment calculated from the Kawski–Chamma–Viallet equation

<sup>c</sup> Excited state dipole moment calculated from the Bakhshieve equation

<sup>d</sup> Excited state dipole moment calculated from the Lippert–Mataga equation

 $e^e$  Excited state dipole moment calculated from the  $E_T^N$  equation

<sup>f</sup> Excited state dipole moment calculated from the McRae equation

<sup>g</sup> Change in dipole moments from  $\mu_e$  and  $\mu_g$ 

h The ratio of excited state and ground dipole moments

<sup>i</sup> The difference in ground state and excited dipole using the  $E_{\text{T}}^{\text{N}}$  method

<sup>j</sup> The orientation of the molecule in excited state relative to the ground state

are obtained as 2.203 D (vacuum), and 2.269 D (ethanol), respectively. The empirical and the theoretical results vary because theoretical analysis is for the isolated gaseous molecule and the empirical values are in the solid-state.

# 3.2.1 Molecular–Microscopic Solvent Polarity Parameter ( $E_T^N$ )

The graph of Stoke's shift of 2BME against  $E_T^N$  values for various media is shown in Fig. [7](#page-9-0). The  $\mu_e$  value calculated using  $E_T^N$  according to Eq. [13](#page-5-0) is given in Table [4](#page-11-0). The  $\mu_e$  value 4.397 D by the  $E_{\rm T}^{\rm N}$  method is 46.85% smaller than that from the Bakhshieve method. This is because BE does not consider specific solute solvent interaction while these details are included in the technique based on  $E_{\text{T}}^{\text{N}}$  [[37](#page-19-0)]. The increasing in solvent polarity in UV–vis and fluorescence bands give a bathochromic shift. This shows the ICT absorption of the less dipolar  $\mu_{g}$  molecule with dominant mesomeric structure results in greater dipolarity in  $\mu_e$  of the molecule with prominent structure. Hence, the  $\mu_e$  for 2BME is more polar than the  $\mu_{\rm g}$  state due to ICT.

### 3.2.2 Kamlet-Taft Solvatochromic Parameters

For explaining the solvent polarity dependence and hydrogen bonding ability of 2BME, the analysis used the Kamlet–Taft equation, which takes account of the polarity and hydrogen bond ability of the solvents. To HBD and HBA abilities of  $\tilde{v}_a$ ,  $\tilde{v}_f$  and  $\Delta \tilde{v} = (\tilde{v}_a - \tilde{v}_f)$  are related to  $\alpha$ ,  $\beta$  and  $\pi^*$  by multiple regression. This analysis results along with (r) arepresented in the below relations:

$$
\tilde{v}_a (cm^{-1}) = 38046.29 - 3426.01 \pi^* + 7851.22 \alpha + 7263.5 \beta (r = 0.957)
$$
  

$$
\tilde{v}_f (cm^{-1}) = 24267.68 + 1152.4 \pi^* + 712.315 \alpha + 640.714 \beta (r = 0.819)
$$
  

$$
\Delta \tilde{v} (cm^{-1}) = 3163.67 + 7865.0 \pi^* - 1657.10 \alpha + 6422.16 \beta (r = 0.813).
$$

The absorption and emission spectra ( $v_{\text{max}}$ ) values vary and the polarizability parameter (s) increases from  $\mu_{\rm g}$  to  $\mu_{\rm e}$ , which gives the stability of the molecule. It can be concluded from the above relations that the non-specific dielectric interaction  $(\pi^*)$  strongly affects the solute. Also, the equation indicates that HBD ( $\alpha$ ) has a greater effect than HBA ( $\beta$ ) for  $\tilde{v}_a$ and  $\tilde{v}_f$ , whereas, in  $\Delta \tilde{v}$ , the HBD ( $\alpha$ ) influence is less than that of HBA ( $\beta$ ).

It can be noted in Table [4,](#page-11-0) that the dipole moment of the excited state was greater than ground state dipole moment with all methods. The Kawski–Chamma–Viallet equation gives the determination of ground state and excited state dipole moments. Using the Lippert–Mataga method the excited state dipole moment was estimated in the 10 solvents, the values are nearly equal to those from the Kawski–Chamma–Viallet method, this indicates significant electronic rearrangement of the probe by using the  $E_{\text{T}}^{\text{N}}$  method. We obtained dipole difference value of (2.405) using 8 solvents [[59](#page-20-0)]. This indicates that the Lippert–Mataga method may not account for the specific interactions between solute and solvent molecules. It was found that solvent polarity and polarizability played important roles in the solvatochromic studies of 2BME.

Atom	DFT/B3LYP (in vacuum)		TDDFT/B3LYP (in ethanol)		
	Electron density	Atomic charge	Electron density	Atomic charge	
1 <sup>C</sup>	6.175301	$-0.175301$	6.190029	$-0.190029$	
2C	6.096015	$-0.096015$	6.108274	$-0.108274$	
3C	6.210216	$-0.210216$	6.212495	$-0.212495$	
4 C	5.731672	0.268328	5.732683	0.267317	
5 C	6.043765	$-0.043765$	6.060744	$-0.060744$	
6 C	6.195517	$-0.195517$	6.193690	$-0.193690$	
7C	5.483920	0.516080	5.471095	0.528905	
8 O	8.484246	$-0.484246$	8.486323	$-0.486323$	
9 C	6.274832	$-0.274832$	6.290689	$-0.290689$	
10 O	8.339091	$-0.339091$	8.387616	$-0.387616$	
11 O	8.480159	$-0.480159$	8.499185	$-0.499185$	
12 C	6.253597	$-0.253597$	6.269003	$-0.269003$	
13 C	5.799339	0.200661	5.777201	0.222799	
14 C	6.016320	$-0.016320$	6.013977	$-0.013977$	
15 C	5.868805	0.131195	5.859370	0.140630	
16 O	8.493942	$-0.493942$	8.513951	$-0.513951$	
17 C	6.308176	$-0.308176$	6.320663	$-0.320663$	
18 C	5.491001	0.508999	5.473997	0.526003	
19 C	5.868186	0.131814	5.865773	0.134227	
20 C	6.109450	$-0.109450$	6.118691	$-0.118691$	
21 C	6.129503	$-0.129503$	6.133109	$-0.133109$	
22 C	6.147192	$-0.147192$	6.163080	$-0.163080$	
23 C	6.135687	$-0.135687$	6.146753	$-0.146753$	
24 C	6.154623	$-0.154623$	6.166039	$-0.166039$	
25 C	6.140157	$-0.140157$	6.150608	$-0.150608$	
26 C	6.197266	$-0.197266$	6.200703	$-0.200703$	
27 O	8.352046	$-0.352046$	8.427377	$-0.427377$	
28 H	0.810394	0.189606	0.792024	0.207976	
29 H	0.802332	0.197668	0.796404	0.203596	
30 H	0.802365	0.197635	0.797466	0.202534	
31 H	0.815452	0.184548	0.807508	0.192492	
32 H	0.844338	0.155662	0.822609	0.177391	
33 H	0.845078	0.154922	0.817285	0.182715	
34 H	0.795256	0.204744	0.772876	0.227124	
35 H	0.849901	0.150099	0.823089	0.176911	
36 H	0.845237	0.154763	0.823680	0.176320	
37 H	0.835540	0.164460	0.812267	0.187733	
38 H	0.788923	0.211077	0.779369	0.220631	
39 H	0.834610	0.165390	0.807752	0.192248	
40 H	0.813058	0.186942	0.801213	0.198787	
41 H	0.764291	0.235709	0.786417	0.213583	

<span id="page-13-0"></span>Table 5 Mulliken atomic charges (a.u) and electron density (a.u) of the 2BME molecule in vacuum and ethanol calculated using DFT/B3LYP/6-311 and TDDFT/B3LYP/6-311 methods

Atom	DFT/B3LYP (in vacuum)		TDDFT/B3LYP (in ethanol)	
	Electron density	Atomic charge	Electron density	Atomic charge
42 H	0.786578	0.213422	0.763677	0.236323
43 H	0.786620	0.213380	0.763248	0.236752

Table 5 continued

### 3.3 Quantum Chemical Analysis

The DFT and TDDFT [\[51\]](#page-20-0) analysis explore the electronic geometry to explain the empirical results. The solvent effectiveness in ethanol was analyzed by PCM, which is a SCR technique [\[60,](#page-20-0) [61](#page-20-0)] using Gamess. The atomic charges and electron densities of 2BME give an idea of the dipole moment and distribution of charges. Electron density and Mulliken atomic charges are listed in Table [5](#page-13-0) and plotted in Fig. 10. The atomic charges in vacuum were calculated in the previous study [\[62\]](#page-20-0). In vacuum  $C_1$ ,  $C_2$ ,  $C_3$ ,  $C_4$ ,  $C_6$ ,  $C_9$  and all oxygen atoms have negative charges (donor atoms) whereas all hydrogens are positive, as are  $C_{12}$ ,  $C_{19}$ ,  $C_{15}$ ,  $C_{13}$ ,  $C_7$  (acceptor atoms). In the ethanol solvent, an increase in the atomic charges of all the atoms is observed; the charges are presented in Table [5.](#page-13-0) The charge distribution represents the existence of strong polar nature, directing towards the  $O_{11}$  ( $Q = -0.480$ ) atom. Such a charge transfer is the characteristic that gives the dipole moment difference between the  $\mu_{\rm g}$  and  $\mu_{\rm e}$  states. It is noted that TDDFT/B3LYP indicates



Fig. 10 Calculated atomic charge and electron densities of the 2BME molecule by using DFT/B3LYP/6- 311 (d,p) and TDDFT/B3LYP//6-311 (d,p) methods in vacuum and ethanol

<span id="page-15-0"></span>an increase in atomic charges because of the inductive effect of the solvent. The change in hydrogen atom gives increased positive values showing charge transfer from hydrogen to the other atoms.

An atomic charge represents the amount of electron cloud situated on an atom and the values of the electron density (8 a.u.) on oxygen atoms are in the order  $O(16)$  $O(8)$  >  $O(11)$  >  $O(4)$  >  $O(5)$ . The solvent medium effect tends to decrease the electron cloud of  $O(16)$  atom due to p-bond of  $O(15)$ , whereas there is an increase in the electron cloud of the other oxygen atoms and  $C(17)$  has the highest electron density (6.308 Å). HOMO and LUMO analysis give the facts about charge transfer within the molecule. Usually, molecular chemical activity can be understood from the HOMO and LUMO energy gap. The smaller the energy gap and the more easily HOMO electrons can be excited to the LUMO levels. The HOMO and LUMO energies are calculated to be  $-0.2280$  and  $-0.0800$  eV in vacuum and -0.2290 and -0.0814 eV in ethanol, respectively. The difference in energy of LUMO– HOMO in vacuum is slightly more than in ethanol solvent with the value of 0.0004 eV, this



Fig. 11 3D plots of HOMO and LUMO calculated by DFT and TDDFT in vacuum and ethanol for the 2BME molecule

represents the stabilization energy of ethanol. In the LUMO and the HOMO 3D plot, the positive (red) and the negative (green) phases are presented in Fig. [11,](#page-15-0) which are not corrected for the methyl moiety. In the HOMO topology, density is found to be localized across the double bonds in the aromatic rings. In the LUMO the charge density is not localized across double bonds but is prominently across the single bonds in the aromatic ring; even in the LUMO charge density is found to be absent in the methyl moiety.

The SAS relative to atomic colors and electron densities are shown in Fig. 12a and b. As can be seen from Fig. 12a, SAS is mapped by MEP from charges in the light green regions indicating the interaction of hydrogen atoms in the solvent, the highest region is in the greenish mat yellow zone showing the interaction of hydrogen atoms and medium. It is observed that the interaction of 2BME and medium is controlled predominately by the hydrogen atoms and electron density surface presented in Fig. 12 matches this data.



Fig. 12 Solvent accessible surfaces according to atomic charges of the 2BME molecule a in vacuum and b in ethanol solvent



Fig. 13 3D plot of molecular electrostatic potential of 2BME molecule **a** in vacuum and **b** in ethanol solvent

The MEP (Fig. 13) displays the constant electron density surface and helps in obtaining the position for electrophilic, hydrogen bonding interaction and nucleophilic reactions in solvent [[63](#page-20-0), [64\]](#page-20-0). Electrostatic potentials at the surface are shown in various colors. The negative charge (red) area of MEP corresponds to electrophilic reactivity and the positive (blue) area represents nucleophilic reactivity (Fig. 13a). The MEP area, the negative cloud which represents surrounding oxygen atoms is not observed and positive regions of potential are observed above the atoms of hydrogen.

## 4 Conclusions

We obtained the  $\mu_{\rm g}$  state by using the Bakshive, Kawski-Chamma method and  $\mu_{\rm e}$  state by adopting Kawski–Chamma–Viallet, Bakshieve's, Lippert–Mataga,  $E_{\text{T}}^{\text{N}}$  and McRae solvatochromic shift techniques. It is found that  $\mu_e$  is greater than  $\mu_g$ . This indicates that the <span id="page-18-0"></span>molecule is more polar in the excited state than in the ground state, which shows the molecule is sensitive to solvent effects. The absorption and emission data have been studied using solvent parameters. Density functional theory B3LYP/6-311(dp) calculations were carried out. The calculated result shows that the optimized geometry can well reproduce the molecular structure. The MEP map shows that the negative and positive potential positions are on electronegative atoms which are around the atoms. These positions provide the data about the regions from where the molecule can have intermolecular interactions. DFT-TDDFT calculations represents that the 2BME molecule is reactive in interacting with the ethanol medium.

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