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# **Physical Properties of Sn1<sup>−</sup>***x***Fe***x***O2 Powders Using Solid State Reaction**

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Abstract Iron-doped SnO<sub>2</sub> diluted magnetic semiconducting powders (Sn1<sup>−</sup>*x*Fe*x*O2, *x* = 0*.*00, 0.03, 0.05, 0.07, 0.10, and 0.15) were synthesized by a simple solid state reaction followed by vacuum annealing and studied the effect of Fe dopant concentrations on structural, optical, and magnetic properties of the synthesized samples. From the X-ray diffraction, it was confirmed that the samples prepared at lower dopant concentrations were tetragonal in structure whereas the samples prepared at higher dopant concentration exhibited orthorhombic SnO and  $Fe<sub>2</sub>O<sub>3</sub>$  phases along with tetragonal  $SnO<sub>2</sub>$  structure. FT-IR spectrum has been used to confirm the formation of Sn–O bond. The optical band gap of the Sn1<sup>−</sup>*x*Fe*x*O2 powders was increased from 3.6 eV to 3.7 eV with increase of dopant concentration. Raman spectroscopy measurement revealed that the broadening of the most intense Raman peak observed at  $630 \text{ cm}^{-1}$ with Fe doping, conforming that the Fe ions are substituted at the Sn sites in the  $SnO<sub>2</sub>$  lattice. Vibrating sample magnetometer measurements confirmed that the Sn1<sup>−</sup>*x*Fe*x*O2 powders were ferromagnetic at room temperature.

**Keywords** Tin oxide · Solid state reaction · X-ray diffraction · Vacuum annealing

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#### **1 Introduction**

Since the discovery of room temperature ferromagnetism in Mn doped ZnO and GaN by Dietl et al. [[1\]](#page-6-0), research for the new kind of materials exhibiting ferromagnetism at/above room temperature are finding intense interest. A number of different semiconducting host materials have been investigated. Among which more attention was paid on (Ga, Mn)As [[2\]](#page-6-1) and (In, Mn)As [[3\]](#page-6-2) to make Curie temperature  $(T_c)$  above room temperature, but reported the highest Curie temperature  $(T_c)$  of 170 K for  $(Ga, Mn)As [4]$  $(Ga, Mn)As [4]$  $(Ga, Mn)As [4]$  and 35 K for  $(In, Mn)As [5]$  $(In, Mn)As [5]$  $(In, Mn)As [5]$ , which is not suitable for practical applications. Hence, intense interest is being put in developing a new type of DMS materials, which exhibit ferromagnetism at/above room temperature. Oxide-based diluted magnetic semiconductors (DMS) are ones which can be best suited for the applications such spintronics, nanoelectronics, magnetoelectronics, and microwave devices [[6,](#page-6-5) [7\]](#page-6-6). The advantages of oxide-based DMS are the wide band gap suits for applications of short wavelength light, capability to be grown on various different of substrates including plastic even at low temperatures, environment safety and durability, and low cost. Moreover, oxygen is expected to produce strong  $p - d$  exchange coupling between band carriers and localized spins [[8\]](#page-6-7). Wide band gap oxides materials, such as  $TiO_2$ ,  $SnO_2$ ,  $In_2O_3$ , and  $ZnO$  doped with different transition metals (TM) opened the doors for the above said electronic devices because of their high optical transparency, electrical conductivity along with ferromagnetism at room temperature  $[9-16]$  $[9-16]$ . The ferromagnetism had been observed in oxide-based DMS systems in powder, nanoparticles, and thin film forms [\[17](#page-6-10)[–19](#page-6-11)]. Among the different oxide materials,  $SnO<sub>2</sub>$  is the material, which finds many optoelectronic applications such as  $SnO<sub>2</sub>$  has been studied extensively because of its high optical transparency and high

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electrical conductivity with wide band gap (3.5 eV). Fitzgerald et al. [[20\]](#page-6-12) have reported the room temperature ferromagnetism in  $Sn_{1-x}M_xO_2$  (M = Mn, Fe, Co,  $x = 0.05$ ) ceramics and found that Sn1<sup>−</sup>*x*Co*x*O2 were paramagnetic whereas  $Sn_{1-x}Mo_xO_2$ ,  $Sn_{1-x}Fe_xO_2$  were ferromagnetic in nature. Sharma et al. [[21\]](#page-6-13) prepared Ni and Fe codoped  $SnO<sub>2</sub>$ nanoparticles by the simple wet chemical method and reported that the undoped  $SnO<sub>2</sub>$  as diamagnetic, whereas Ni and Fe doped  $SnO<sub>2</sub>$  nanoparticles were ferromagnetic in nature. The present study is an attempt to impart room temperature ferromagnetism in Fe doped  $SnO<sub>2</sub>$  and to study the origin of ferromagnetism. Different physical [\[22](#page-6-14)[–24](#page-6-15)] and chemical methods  $[25, 26]$  $[25, 26]$  $[25, 26]$  $[25, 26]$  have been applied for the synthesis of Fe doped  $SnO<sub>2</sub>$  powders. However, very few works have explored the solid state reaction method. A simple solid state reaction method is used here to prepare Fe doped  $SnO<sub>2</sub>$ powders and studied the influence of doping levels of Fe on structural, optical, and magnetic properties of  $(Sn<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>2</sub>$ , *x* = 0*.*00, 0.03, 0.05, 0.07, 0.10, and 0.15) powders.

#### **2 Experimental**

Sn1<sup>−</sup>*x*Fe*x*O2 powder samples (*x* = 0*.*03, 0.05, 0.07, 0.10, and 0.15) were prepared by a solid state reaction followed by vacuum annealing. Commercially available  $SnO<sub>2</sub>$  and Fe<sub>2</sub>O<sub>3</sub> (M/S Sigma-Aldrich 99.99 % pure) were accurately weighed in required proportions and were mixed and ground thoroughly using an agate mortar and pestle to convert to very fine powders. The grinding of the mixtures was carried out for 16 hours for all the powder samples. The ground powder samples were loaded into a small one end closed quartz tube of diameter 10 mm and length of 10 cm, which was enclosed in a bigger quartz tube of diameter of 2.5 cm and length of 75 cm with provision to allow unwanted vapors to escape from the reaction chamber and evacuated at  $2 \times 10^{-3}$  mbar using a rotary pump was used for the synthesis of the present samples. The complete set up was placed in horizontal tubular microprocessor controlled furnace and fired for several hours at different temperatures. The firing temperature and firing periods were optimized at 900 °C and 10 hours. X-ray diffraction (X-ray diffractometer, D8 Advance, BRUKER) was used to establish structural aspects. Energy dispersive analysis spectroscopy (EDS) (OXFORD instrument inca penta FET X3) was used to carry out elemental analysis. The diffused reflectance spectra were recorded on UV-Vis-NIR Spectrophotometer (JASCO V-670). Fourier Transform Infrared (FT-IR) Spectroscopic analysis was carried using FT-IR Spectrophotometer (SHIMADZU). Magnetic measurements were carried out using Vibrating sample magnetometer (Lake Shore-7410, IIT Madras).

#### **3 Results and discussion**

#### 3.1 Structural properties

Figure  $1(a)$  $1(a)$  shows the X-ray diffraction patterns of the  $Sn_{1-x}Fe_xO_2$  at different Fe doping levels ( $x = 0.03, 0.05$ , 0.07, 0.10, and 0.15). All the diffraction peaks having higher and lower intensities were taken into the considerations. The diffraction peaks were found at diffraction angles of 26*.*62◦ (1 1 0), 33*.*89◦ (1 0 1), 37*.*98◦ (2 0 0), 39*.*09◦ (1 1 1), 42*.*65◦ (2 1 0), 51*.*80◦ (2 1 1), 54*.*79◦ (2 2 0), 57*.*85◦ (0 0 2), 61*.*92◦ (3 1 0), 64*.*75◦ (1 1 2), 65*.*98◦ (3 0 1), 71*.*39◦ (2 0 2), and 78*.*73◦ (3 2 1) were exactly coincided with tetragonal structure of  $SnO<sub>2</sub>$  [JCPDS No. 411445]. Among the above orientations, (1 1 0) is the predominant orientation. Same orientations were observed when the dopant concentration increased from 0.00 to 0.07. This clearly indicated that all  $Sn_{1-x}Fe_xO_2$  powder samples were polycrystalline in nature without having any secondary phase at lower Fe doping concentrations ( $x \le 0.07$ ). At a dopant concentration of 0.10, instead of single phase, other phases belonging to SnO orthorhombic were also identified. At higher Fe dopant concentration  $x = 0.15$ , in addition to orthorhombic SnO phases, phases belonging to  $Fe<sub>2</sub>O<sub>3</sub>$  were also observed. Further, as the Fe dopant concentration increased above 0.07, the intensity of the X-ray diffraction peaks decreased and FWHM increased, which show the degradation of crystallinity due to the additional impurities in the samples.

In order to further confirm that no Fe or  $Fe<sub>2</sub>O<sub>3</sub>$  phases were present at lower Fe doping concentrations, Rietveld refinement analysis was also carried out at higher and lower Fe doping concentrations. From Fig. [1\(](#page-2-0)b), it is clear that no other phases or impurities were observed when the Fe doping concentration was increased from  $x = 0.00$  to  $x = 0.07$ , whereas impurity phases were observed at higher Fe doping concentrations as shown in Fig.  $1(c)$  $1(c)$ . It clearly reflects that the miscibility of dopant impurity is limited in  $SnO<sub>2</sub>$  lattice. It is necessary to note that the small limiting solubility of various metal oxides is typical for  $SnO<sub>2</sub> [27]$  $SnO<sub>2</sub> [27]$ . Similar results were observed in Fe doped  $SnO<sub>2</sub>$  by Mathew et al. [\[28](#page-6-19)]. The diffraction peaks angle  $(2\theta)$  shifted slightly toward the lower angle with the Fe doping concentration. It may due to smaller ionic radius of Fe<sup>+3</sup> (0.63 Å) than that of ionic radius of  $\text{Sn}^{+4}$  (0.71 Å).

The crystallite size  $(G)$  was calculated by using the Debye–Scherrer formula [\[29](#page-6-20)],

$$
G = k\lambda/\beta\cos\theta\tag{1}
$$

where  $k$  is a constant,  $\lambda$  is the diffraction wavelength of CuK<sub>α</sub> ( $\lambda = 1.5406$  Å),  $\beta$  is the full width at half maximum



<span id="page-2-0"></span>**Fig. 1** (**a**) X-ray diffraction patterns of Sn1<sup>−</sup>*x*Fe*x*O2 powder samples at *x* = 0*.*00, 0.03, 0.05, 0.07, 0.10, and 0.15. (**b**) Rietveld refinement patterns of  $Sn_{1-x}Fe_xO_2$  powder samples at  $x = 0.00, 0.03, 0.05, 0.07$ . *Dots* represent the observed intensities, and *the solid line* is calculated ones. A difference (obs.-cal.) plot is shown beneath. *Vertical bars*

are the reflection position marks. (**c**) Rietveld refinement patterns of  $Sn_{1-x}Fe_xO_2$  powder samples at  $x = 0.10$  and 0.15. *Dots* represent the observed intensities, and *the solid line* is calculated ones. A difference (obs.-cal.) plot is shown beneath. *Vertical bars* are the reflection position marks

(FWHM), and  $\theta$  is the diffracted angle, respectively. The grain size of the powder samples increased from 45 nm– 92 nm with an increase of Fe doping concentration from 0.03 to 0.07. A decrease in lattice constant was observed from 4.7451 Å to 4.7358 Å when Fe doping concentration increased from 0.03 to 0.07. Further, the lattice volume decreased from 71.74  $\AA^3$  to 71.39  $\AA^3$  as the Fe doping concentration increased from  $x = 0.03$  to  $x = 0.15$ . A summary of Fe doping concentration, grain size, lattice parameters "*a*" and "*c*" and lattice volumes is given in Table [1.](#page-3-0)

In order to know the elemental compositions of tin, oxygen and iron, EDS spectra were recorded for all the Sn1<sup>−</sup>*x*Fe*x*O2 powder samples at *x* = 0*.*00, 0.03, 0.05, 0.07, 0.10, and 0.15. Figure [2](#page-3-1) show the EDS spectrum of Fe doped  $SnO<sub>2</sub>$  powder sample at  $x = 0.07$ . From the EDS spectra, it was found that the impurity was present in all the samples, <span id="page-3-1"></span>**Fig. 2** EDS spectrum of Sn1<sup>−</sup>*x*Fe*x*O2 powder sample at  $x = 0.07$ 



<span id="page-3-0"></span>**Table 1** A summary of doping concentration, 2*θ* (110), lattice constants *a* ( $\AA$ ), *c* ( $\AA$ ) and grain size *L* (nm), and lattice volume ( $\AA$ <sup>3</sup>)



which clearly represents the existence of Fe ions in doped sample and confirms the doping of Fe in SnO<sub>2</sub> host material. The elemental compositions of Fe, Sn, and O were found to be as 1.06 %, 16.33 %, and 82.61 %, respectively.

# 3.2 Optical properties

Figure [3](#page-3-2)(a) shows the diffused reflectance spectra of the Sn<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>2</sub> powder samples at  $x = 0.00, 0.03, 0.05, 0.07$ . In addition to reflectance, absorbance spectra was also plotted (not shown here) for calculating the band gap of the Fe doped SnO2. The absorption coefficient was calculated using Kubelka–Munk equation [\[30](#page-6-21)],

$$
\alpha = (1 - R)^2 / 2R \tag{2}
$$

where  $\alpha$  is the absorption coefficient,  $R$  is reflectance of the samples. As secondary phases belonging to SnO and  $Fe<sub>2</sub>O<sub>3</sub>$ were observed at Fe dopant concentration of  $x = 0.10$  and  $x = 0.15$ , the diffused reflectance spectra and band gap values were not taken into considerations. The optical band gap (*E*g) of the powders were determined by Tauc's plots using equation [[31\]](#page-6-22),

$$
\alpha h v = A (E_g - h v)^n \tag{3}
$$



<span id="page-3-2"></span>**Fig. 3** (**a**) Diffused reflectance spectra of Sn1<sup>−</sup>*x*Fe*x*O2 powder samples at  $x = 0.00, 0.03, 0.05,$  and 0.07. (**b**) Plots of  $(\alpha h \nu)^2$  versus  $h \nu$ of Sn1<sup>−</sup>*x*Fe*x*O2 powder samples at *x* = 0*.*00, 0.03, 0.05, and 0.07

<span id="page-4-0"></span>**Fig. 4** FT-IR spectra of  $Fe<sub>2</sub>O<sub>3</sub>$ and Sn1<sup>−</sup>*x*Fe*x*O2 powder sample at *x* = 0*.*00, 0.07, and 0.15



where *n* depend on the kind of optical transition that prevails. Here,  $n = 1/2$ , as SnO<sub>2</sub> is directly allowed n-type degenerate semiconductor. The optical band gap is obtained by plotting  $(\alpha h v)^2$  versus the photon energy  $(hv)$  and by extrapolating of the linear region of the plots to zero absorption ( $\alpha = 0$ ). The optical band gap  $E<sub>g</sub>$  is obtained by plotting  $(\alpha h \nu)^2$  versus the photon energy  $(h \nu)$  and by extrapolating the linear region of the plots to zero absorption ( $\alpha = 0$ ). The optical band gap of the powder samples increased from 3.6 eV to 3.7 eV when the Fe doping concentration increased from  $x = 0.03$  to  $x = 0.07$ . At higher doping concentrations, SnO and Fe<sub>2</sub>O<sub>3</sub> phases were present, hence the optical band gap at higher doping concentrations is not considered here. A similar increase in the optical band gap was also observed in Mn doped  $SnO<sub>2</sub>$  nanoparticles and the observed band gap is lesser when compared with Mn doped  $SnO<sub>2</sub>$  nanoparticles [\[32](#page-6-23)].

## 3.3 FT-IR studies

Fourier transform infrared spectra were recorded to pure Fe<sub>2</sub>O<sub>3</sub> and Sn<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>2</sub> powder samples at  $x = 0.00, 0.07$ , and 0.15 as shown in Fig. [4.](#page-4-0) The bands observed in the wavenumber range of 430 cm<sup>-1</sup> to 620 cm<sup>-1</sup> related to the vibration of antisymmetric O–Sn–O bridge functional groups of  $SnO<sub>2</sub>$ . The change in peak position, shape, and size indicate that incorporation of Fe took place in  $SnO<sub>2</sub>$ host material.

## 3.4 Raman analysis

Raman scattering is one of the most powerful tool to study the crystallinity, phase composition, and defect structures associated with the material. Figure [5](#page-5-0) shows the Raman spectra of  $\text{Sn}_{1-x}\text{Fe}_x\text{O}_2$  powders at  $x = 0.00, 0.03, 0.05$ , and 0.07. Raman spectra of Sn<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>2</sub> powders at higher Fe doping concentrations are not shown here as secondary phases were observed at higher doping levels. The representation of the normal vibration modes at the center of the Brillouin zone were explained using Hong et al. [\[33](#page-6-24)]. The pure  $SnO<sub>2</sub>$  spectrum shows the classic cassirite  $SnO<sub>2</sub>$  vibration at 476, 630, and 776 cm<sup>-1</sup>. The peaks at 621 cm<sup>-1</sup> can be assigned to the A<sub>1g</sub> mode, which is shifted to 608 cm<sup>-1</sup>, 606 cm<sup>-1</sup>, and 620 cm<sup>-1</sup> for Sn<sub>1-*x*</sub>Fe<sub>*x*</sub>O<sub>2</sub> powders at *x* = 0*.*00, 0.03, 0.05, and 0.07, respectively. The Raman bands at 466 cm<sup>-1</sup> and 762 cm<sup>-1</sup> are assigned to be the vibration modes of  $E<sub>g</sub>$  and  $B<sub>2g</sub>$ , respectively.

# 3.5 Magnetic properties

Figure [6\(](#page-5-1)a)–(c) shows the M–H curves for  $Sn_{1-x}Fe_xO_2$ powder sample with different Fe concentrations. Magnetization measurements were carried out for all the samples including pure  $SnO<sub>2</sub>$ . The intrinsic  $SnO<sub>2</sub>$  exhibited diamagnetism. By introducing the Fe atoms into  $SnO<sub>2</sub>$ , the magnetic behavior of semiconductor changed from diamagnetic to ferromagnetic. It may be due to substitution of Fe atoms in Sn site in  $SnO<sub>2</sub>$  lattice. The saturation magnetization increased with an increase of Fe doping concentration and



<span id="page-5-0"></span>**Fig. 5** Raman spectra of  $\text{Sn}_{1-x}\text{Fe}_x\text{O}_2$  powder sample at  $x = 0.00$ , 0.03, 0.05, and 0.07

became high enough for the sample at a doping concentration of  $x = 0.07$ . The samples exhibited a magnetization of 95.99 memu/g at Fe doping concentration of  $x = 0.03$  and a highest of magnetization 151.41 memu/g at a Fe dopant concentration of  $x = 0.07$  was observed. After that, it again decreased even though it had Fe<sub>2</sub>O<sub>3</sub> phases at higher dopant concentrations. When the Fe doping concentration further increased, the saturation magnetic moment decreased. From the X-ray diffraction, it is clear that when Fe doping concentration increased to  $x = 0.15$ , Fe<sub>2</sub>O<sub>3</sub> along with SnO phase was also observed. But the saturation magnetic moment did not increase even  $Fe<sub>2</sub>O<sub>3</sub>$  phase was present in the sample. It reflects that the ferromagnetism is not due to  $Fe<sub>2</sub>O<sub>3</sub>$  clusters where as it is an intrinsic property rather than impurities. This confirms that the observed ferromagnetism in Sn<sub>1−*x*</sub>Fe<sub>*x*</sub>O<sub>2</sub> powder samples at *x* = 0.03, 0.05, 0.07 are purely intrinsic in nature. It may be due to ferromagnetic interactions developed in the samples at lower dopant concentrations. Ferromagnetism was also observed in noncrystalline Fe doped  $SnO<sub>2</sub>$  due to presence of secondary phases [\[34](#page-6-25)]. The decrease in magnetic moment at higher Fe doping concentration was observed, which may be due to antiferromagnetic super exchange interaction among the nearest neighbor in Fe atoms. The observed magnetic moments are higher when compared to Fe doped  $SnO<sub>2</sub>$  at a doping concentration of  $x = 0.10$ . From Fig. [6](#page-5-1)(c), it can be seen that the M–H curves are along the origin as the observed magnetizations are very small. The observed magnetizations are equal to that of pure  $SnO<sub>2</sub>$ .

# **4 Conclusions**

Iron-doped SnO2 diluted magnetic semiconducting powders (Sn1<sup>−</sup>*x*Fe*x*O2, *x* = 0*.*00, 0.03, 0.05, 0.07, 0.10, and 0.15)



<span id="page-5-1"></span>**Fig. 6** (a) M–H curve of  $\text{Sn}_{1-x}\text{Fe}_x\text{O}_2$  powder sample at  $x = 0.07$ . (**b**) M–H curve  $\text{Sn}_{1-x}\text{Fe}_x\text{O}_2$  powder sample at  $x = 0.15$ . (**c**) M–H curves of  $Sn_{1-x}Fe_xO_2$  powder samples at  $x = 0.00, 0.03, 0.05, 0.07$ , and 0.15

were synthesized by a simple solid state reaction followed by vacuum annealing and studied the effect of Fe doping concentration on structural, optical, and ferromagnetic properties of the synthesized powder samples. The powder samples exhibited tetragonal structure at lower Fe dopant concentrations. At higher dopant concentrations, in addition to tetragonal  $SnO<sub>2</sub>$ ,  $SnO$  and  $Fe<sub>2</sub>O<sub>3</sub>$  phases were also observed. The samples at lower and higher Fe dopant concentrations exhibited ferromagnetism at room temperature and a decrease in magnetic moment was observed at higher Fe doping concentrations.

<span id="page-6-1"></span><span id="page-6-0"></span>**Acknowledgements** Authors are highly thankful to VIT-SIF for providing XRD, DRS facilities to carry out the present work. The authors also thank the Sophisticated Advanced Instruments Facility (SAIF), IIT Madras, Tamilnadu, India, for providing vibrating sample magnetometer facility.

## <span id="page-6-5"></span><span id="page-6-4"></span><span id="page-6-3"></span><span id="page-6-2"></span>**References**

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