ORIGINAL PAPER



A variational principle for a thin film equation

Ji-Huan He¹ · Chang Sun²

Received: 29 June 2019 / Accepted: 26 August 2019 / Published online: 31 August 2019 © Springer Nature Switzerland AG 2019

Abstract

Thin film arises in various applications from electrochemistry to nano devices, many mathematical tools were adopted to study the problem, e.g. Lie symmetries and conservation laws, however, the variational approach is rare. This paper shows that the semi-inverse method is an effective approach to establishment of a variational formulation for the thin film equation. A detailed derivation process is given, a special skill for construction of a heuristic trial-functional is elucidated.

Keywords Variational theory \cdot Special function \cdot Trial-functional \cdot Nanoscale adhesion \cdot Coating \cdot Wetting

1 Introduction

Variational principle plays a key role in both numerical and analytical analyses of a practical problem, it suggests an energy conservation for the whole solution domain, and a variational-based numerical algorithm guarantees the energy conservation at each point, while a variational-based analytical solution is an optimal one for a given trial-solution and valid for the whole solution domain.

Recently Recio et al. [1] studied the following equation

$$u_t + (f(u)u_{xxx} + g(u)u_x)_x + h(u) = 0$$
(1)

where f, g and h are functions of u.

Equation (1) can describe a thin film problem, which can be found widely applications in electrochemistry [2], cell culture [3], fiber fabrication [4], nanoscale adhesion [5], coating [6], wetting [7] and micro/nano devices [8]. Many analytical methods and

[☑] Ji-Huan He hejihuan@suda.edu.cn

¹ School of Science, Xi'an University of Architecture and Technology, Xi'an 710055, China

² Qujing Normal University, Sanjiang Avenue, Economic Development Zone, Qujing, Yunnan, China

numerical methods were applied to study such problems [9-15]. This paper aims at establishing a variational formulation for Eq. (1) by the semi-inverse method [16-27].

2 The semi-inverse method and the variational formulation

The semi-inverse method [16-19] is to establish a variational principle directly from governing equations. In order to effectively use the method, we write Eq. (1) in a conservation form

$$u_t + [(F(u))_{xxx} + (G(u))_x + H(u)]_x = 0$$
⁽²⁾

where F, G and H satisfy the following relationships

$$(F(u))_{xxx} = f(u)u_{xxx} \tag{3}$$

$$(G(u))_x = g(u)u_x \tag{4}$$

$$(H(u))_u = h(u) \tag{5}$$

According to Eq. (2), we can introduce a special function φ , satisfying the following relations

$$\frac{\partial \varphi}{\partial x} = u \tag{6}$$

$$\frac{\partial \varphi}{\partial t} = -[(F(u))_{xxx} + (G(u))_x + H(u)]$$
⁽⁷⁾

The defined special function (φ) is potential-like function. By the semi-inverse method [16–19], we can establish a trial functional in the form

$$J(u,\varphi) = \iint \{u\varphi_t + [(F(u))_{xxx} + (G(u))_x + H(u)]\varphi_x + \sigma\}dxdt$$
(8)

where σ is an unknown function of u and its derivatives. The advantage of the above trial-functional is the stationary condition with respect to φ is Eq. (2), but we can not identify σ , so we modify Eq. (8) in the form

$$J(u,\varphi) = \iint \{au\varphi_t + b\varphi_x\varphi_t + [(F(u))_{xxx} + (G(u))_x + H(u)]\varphi_x + \sigma\}dxdt$$
(9)

where a and b are constants to be further determined.

The Euler-Lagrange equations of Eq. (9) are

$$-au_t - 2b\varphi_{xt} - [(F(u))_{xxx} + (G(u))_x + H(u)]_x = 0$$
(10)

$$a\varphi_t - \frac{\partial F}{\partial u}\varphi_{xxxx} - \frac{\partial G}{\partial u}\varphi_{xx} + \frac{\partial H}{\partial u}\varphi_x + \frac{\delta\sigma}{\delta u} = 0$$
(11)

where $\delta\sigma/\delta u$ is the variational derivative. In this paper it can be written in the form

$$\frac{\delta\sigma}{\delta u} = \frac{\partial\sigma}{\partial u} - \frac{\partial}{\partial x}(\frac{\partial\sigma}{\partial u_x}) - \frac{\partial^3}{\partial x^3}(\frac{\partial\sigma}{\partial u_{xxx}})$$
(12)

In view of Eq. (6), we can write Eq. (10) in the form

$$(a+2b)u_t + [(F(u))_{xxx} + (G(u))_x + H(u)]_x = 0$$
(13)

which should be Eq. (2), so we have

$$a + 2b = 1 \tag{14}$$

In view of Eqs. (6) and (7), we can write Eq. (11) in the form

$$-a[(F(u))_{xxx} + (G(u))_x + H(u)] - fu_{xxx} - gu_x + hu + \frac{\delta\sigma}{\delta u} = 0$$
(15)

or

$$-a[f(u)u_{xxx} + g(u)u_x + H(u)] - fu_{xxx} - gu_x + hu + \frac{\delta\sigma}{\delta u} = 0$$
(16)

In order to identify σ in Eq. (16), we set

$$\mathbf{a} = -1 \tag{17}$$

Equation (16) becomes

$$\frac{\delta\sigma}{\delta u} = -H(u) - h(u)u \tag{18}$$

From Eq. (18) we can determine σ easily, which satisfies the following relation

$$\frac{\partial\sigma}{\partial u} = -H(u) - h(u)u \tag{19}$$

We, therefore, obtain the following variational formulation

$$J(u,\varphi) = \iint \{-u\varphi_t + \varphi_x\varphi_t + [(F(u))_{xxx} + (G(u))_x + H(u)]\varphi_x + \sigma\}dxdt$$
(20)

where σ is defined in Eq. (19).

Remark if a = 1 as suggested in Eq. (8), we have difficulty in identifying σ from Eq. (16).

3 An example

We consider a special case of Eq. (1)

$$u_t + (u_{xxx} + uu_x)_x + u^2 = 0 (21)$$

Hereby f(u) = 1, g(u) = u, $h(u) = u^2$, F(u) = u, $G(u) = \frac{1}{2}u^2$, $H(u) = \frac{1}{3}u^3$, and

$$\frac{\partial \varphi}{\partial x} = u \tag{22}$$

$$\frac{\partial\varphi}{\partial t} = -(u_{xxx} + uu_x + \frac{1}{3}u^3) \tag{23}$$

$$\frac{\partial\sigma}{\partial u} = -H(u) - h(u)u = -\frac{1}{3}u^3 - u^3 = -\frac{4}{3}u^3$$
(24)

From Eq. (24) σ can be identified as

$$\sigma = -\frac{1}{3}u^4 \tag{25}$$

We, therefore, obtain the following variational principle for Eq. (21):

$$J(u,\varphi) = \iint \left\{ -u\varphi_t + \varphi_x\varphi_t + \left[u_{xxx} + \frac{1}{2}(u^2)_x + \frac{1}{3}u^3 \right] \varphi_x - \frac{1}{3}u^4 \right\} dxdt$$
(26)

which is subject to the constraint of Eq. (22).

Proof The stationary conditions of Eq. (26) with respect to φ and u are

$$u_t - 2\varphi_{xt} - \left[u_{xxx} + \frac{1}{2}(u^2)_x + \frac{1}{3}u^3\right]_x = 0$$
(27)

$$-\varphi_t - \varphi_{xxxx} - u\varphi_{xx} + u^2\varphi_x - \frac{4}{3}u^3 = 0$$
 (28)

In view of Eq. (22), we find Eqs. (27) and (28) are equivalent to Eq. (21) and Eq. (23), respectively.

4 Discussion and conclusion

A suitable construction of a trial functional is of great importance for the establishment of a variational principle. Equation (8) does not work, because if we set a = 1, Eq. (16) becomes

$$-[f(u)u_{xxx} + g(u)u_x + H(u)] - fu_{xxx} - gu_x + hu + \frac{\delta\sigma}{\delta u} = 0$$
(29)

or

$$-2f(u)u_{xxx} - 2g(u)u_x - H(u) + hu + \frac{\delta\sigma}{\delta u} = 0$$
(30)

In Eq. (30), we have difficulty in identification of σ due to the terms involving u_{xxx} and u_x .

We can easily determine σ for all even-order derivatives from the following equation

$$\frac{\delta\sigma}{\delta u} = m(u) + \alpha u_{xx} + \beta u_{xxxx} + \delta u_{xxxxxx}$$
(31)

where α , β , and δ are constants, m is a function of u. From Eq. (31) σ can be determined as

$$\sigma = M(u) - \frac{1}{2}\alpha(u_x)^2 + \frac{1}{2}\beta(u_{xx})^2 - \frac{1}{2}\delta(u_{xxx})^2$$
(32)

where M is defined as

$$\frac{\partial M}{\partial u} = m \tag{33}$$

In this paper, we find that the semi-inverse method provides an effective tool to finding a needed variational principle for a practical problem, the derivation process is explained step by step, so that it can be easily followed. Recently Wang et al. [27] successfully applied the semi-inverse method to fractal calculus [28, 29], and obtained a variational principle for wave traveling in a fractal space.

The variational principle is a foundation of the variational iteration method [30, 31], which is now widely applied in fractional calculus, and the present paper might give a hint for an effective identification of Lagrange multiplier in the fractional variational iteration method [30–36].

References

- 1. E. Recio, T.M. Garrido, R. de la Rosa et al., Conservation laws and Lie symmetries a (2+1)-dimensional thin film equation. J. Math. Chem. **57**(5), 1243–1251 (2019)
- X.-X. Li, D. Tian, C.-H. He, J.-H. He, A fractal modification of the surface coverage model for an electrochemical arsenic sensor. Electrochim. Acta 296, 491–493 (2019)
- J. Fan, Y.R. Zhang, Y. Liu et al., Explanation of the cell orientation in a nanofiber membrane by the geometric potential theory. Results Phys. 15, 102537 (2019)
- 4. Z.P. Yang, F. Dou, T. Yu et al., On the cross-section of shaped fibers in the dry spinning process: physical explanation by the geometric potential theory. Results Phys. 14, 102347 (2019)
- X.X. Li, J.H. He, Nanoscale adhesion and attachment oscillation under the geometric potential. Part 1: the formation mechanism of nanofiber membrane in the electrospinning. Results Phys. 12, 1405–1410 (2019)
- A. Saeed, Z. Shah, S. Islam et al., Three-dimensional casson nanofluid thin film flow over an inclined rotating disk with the impact of heat generation/consumption and thermal radiation. Coatings 9(4), 248 (2019)
- 7. C.-J. Zhou, D. Tian, J.-H. He, What factors affect lotus effect? Therm. Sci. 22, 1737–1743 (2018)

- J.H. He, From micro to nano and from science to technology: nano age makes the impossible possible. Micro Nanosyst. 12(1), 1–2 (2010)
- J. Manafian, C.T. Sindi, An optimal homotopy asymptotic method applied to the nonlinear thin film flow problems. Int. J. Numer. Methods Heat Fluid Flow 28(12), 2816–2841 (2018)
- N. Faraz, Y. Khan, Thin film flow of an unsteady Maxwell fluid over a shrinking/stretching sheet with variable fluid properties. Int. J. Numer. Methods Heat Fluid Flow 28(7), 1596–1612 (2018)
- F. Ghani, T. Gul, S. Islam et al., Unsteady magnetohydrodynamics thin film flow of a third grade fluid over an oscillating inclined belt embedded in a porous medium. Therm. Sci. 21(2), 875–887 (2017)
- 12. Q.T. Ain, J.H. He, On two-scale dimension and its applications. Therm. Sci. 23(3B), 1707–1712 (2019)
- J.H. He, F.Y. Ji, Two-scale mathematics and fractional calculus for thermodynamics. Therm. Sci. 57(8), 1932–1934 (2019)
- J.H. He, F.Y. Ji, Taylor series solution for Lane–Emden equation. J. Math. Chem. (2019). https://doi. org/10.1007/s10910-019-01048-7
- 15. J.H. He, The simplest approach to nonlinear oscillators. Results Phys. 15, 102546 (2019)
- J.H. He, Variational principles for some nonlinear partial differential equations with variable coefficients. Chaos Solitons Fractals 19(4), 847–851 (2004)
- J.H. He, J. Zhang, Semi-inverse method for establishment of variational theory for incremental thermoelasticity with voids, in *Variational and Extremum Principles in Macroscopic Systems*, ed. by S. Sieniutycz, H. Farkas (Elsevier, Amsterdam, 2005), pp. 75–95
- J.H. He, A modified Li–He's variational principle for plasma. Int. J. Numer. Methods Heat Fluid Flow (2019). https://doi.org/10.1108/HFF-06-2019-0523
- J.H. He, Lagrange crisis and generalized variational principle for 3D unsteady flow. Int. J. Numer. Methods Heat Fluid Flow (2019). https://doi.org/10.1108/HFF-07-2019-0577
- Y. Wu, J.H. He, A remark on Samuelson's variational principle in economics. Appl. Math. Lett. 84, 143–147 (2018)
- 21. J.H. He, Hamilton's principle for dynamical elasticity. Appl. Math. Lett. 72, 65-69 (2017)
- K. Libarir, A. Zerarka, A semi-inversevariational method for generating the bound state energy eigenvalues in a quantum system: the Dirac Coulomb type-equation. J. Mod. Opt. 65(8), 987–993 (2018)
- J. Manafian, P. Bolghar, A. Mohammadalian, Abundant soliton solutions of the resonant nonlinear Schrodinger equation with time-dependent coefficients by ITEM and He's semi-inverse method. Opt. Quant. Electron. 49(10), 322 (2017)
- O.H. El-Kalaawy, New variational principle-exact solutions and conservation laws for modified ionacoustic shock waves and double layers with electron degenerate in plasma. Phys. Plasmas 24(3), 032308 (2017)
- A. Biswas, Q. Zhou, S.P. Moshokoa et al., Resonant 1-soliton solution in anti-cubic nonlinear medium with perturbations. Optik 145, 14–17 (2017)
- Y. Li, C.H. He, A short remark on Kalaawy's variational principle for plasma. Int. J. Numer. Methods Heat Fluid Flow 27(10), 2203–2206 (2017)
- Y. Wang, J.Y. An, X.Q. Wang, A variational formulation for anisotropic wave traveling in a porous medium. Fractals 27(4), 1950047 (2019)
- J.H. He, A tutorial review on fractal spacetime and fractional calculus. Int. J. Theor. Phys. 53(11), 3698–3718 (2014)
- 29. J.H. He, Fractal calculus and its geometrical explanation. Results Phys. 10, 272–276 (2018)
- N. Anjum, J.H. He, Laplace transform: making the variational iteration method easier. Appl. Math. Lett. 92, 134–138 (2019)
- J.H. He, Some asymptotic methods for strongly nonlinear equations. Int. J. Mod. Phys. B 20, 1141–1199 (2006)
- D. Baleanu, H.K. Jassim, H. Khan, A modified fractional variational iteration method for solving nonlinear gas dynamic and coupled KdV equations involving local fractional operator. Therm. Sci. 22, S165–S175 (2018)
- D. Dogan Durgun, A. Konuralp, Fractional variational iteration method for time-fractional nonlinear functional partial differential equation having proportional delays. Therm. Sci. 22, S33–S46 (2018)
- M. Inc, H. Khan, D. Baleanu et al., Modified variational iteration method for straight fins with temperature dependent thermal conductivity. Therm. Sci. 22, S229–S236 (2018)
- H. Jafari, H.K. Jassim, J. Vahidi, Reduced differential transform and variational iteration methods for 3-D diffusion model in fractal heat transfer within local fractional operators. Therm. Sci. 22, S301–S307 (2018)

36. Y. Wang, Y.F. Zhang, Z.J. Liu, An explanation of local fractional variational iteration method and its application to local fractional modified Kortewed-de Vries equation. Therm. Sci. **22**, 23–27 (2018)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.