



A Design Method of an Ultra-Wideband and Easy-to-Array Magic-T: A 6–14 GHz Scaled Model for a mm/submm Camera

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Abstract

We established a design method for a Magic-T with a single-layer dielectric/metal structure suitable for both wideband and multi-element applications for millimeter and submillimeter wave imaging observations. The design method was applied to a Magic-T with a coupled-line, stubs, and single-stage impedance transformers in a frequency-scaled model (6–14 GHz) that is relatively easy to demonstrate through manufacturing and evaluation. The major problem is that using the conventional perfect matching condition for a coupled-line alone produces an impractically large width coplanar coupled-line (CPCL) to satisfy the desired bandwidth ratio. In our study, by removing this constraint and optimizing impedances utilizing a circuit simulator with high computation speed, we found a solution with a $\sim 180 \mu\text{m}$ wide CPCL, which is approximately an order of magnitude smaller than the conventional analytical solution. Furthermore, considering the effect of transition discontinuities in the transmission lines, we optimized the line length and obtained a design solution with return loss < -20 dB, amplitude imbalance < 0.1 dB, and phase imbalance $< 0.5^\circ$ from 6.1 to 14.1 GHz.

Keywords Magic-T · Ratrace · Wideband · Single layer · Coplanar · Coupled-line

1 Introduction

In recent years, in order to perform wideband millimeter and submillimeter continuum wave imaging (typically cosmic microwave background observations), superconducting detector arrays have been developed in which planar antenna and read-out elements are integrated on a dielectric substrate and arranged on a focal plane (see review [1]).

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A Magic-T combines two signals with a phase difference of 0° or 180° at the sum (Σ) port or difference (Δ) port, respectively, and often used for dual polarization observation with an ortho-mode transducer (OMT), which separates each polarization with two pairs of probes. For example, it combines two signals of the fundamental mode (TE₁₁) at Δ port, which are in opposite phase and sent through the optics and horn antenna to an OMT on the detector substrate. The Magic-T also terminates undesirable higher-order even modes at the Σ port, which serves to shape the beam [2].

In previous studies, an ultra-wideband over $BWR \geq 2.3$ ($BWR :=$ bandwidth ratio, the highest/lowest frequency ratio in a band) has been realized using a multilayer structure or a crossover to a narrow stripline requiring high accuracy alignment [3, 4]. On the other hand, for detector array applications, it is desirable for the Magic-T to be manufactured as a single-layer dielectric/metal (hereafter referred to as *single layer*), which requires fewer manufacturing processes and is easier to fabricate.

In this paper, we present a design method for a Magic-T that is both ultra-wideband and single layer, without using a crossover to a stripline.

2 Wideband and Single-Layer Magic-T

A well-known basic structure of Magic-T is a ratrace coupler consisting of three $\lambda/4$ ($\lambda :=$ wavelength) and a $3\lambda/4$ transmission lines [5]. The bandwidth is narrow because the phase inversion depends on the difference in line length. Then, a wideband Magic-T was proposed to replace the $3\lambda/4$ line with a shorted coupled-line, which enables phase inversion by capacitive coupling [6]. Moreover, further broadbanding is demonstrated using stubs that compensate for the asymmetric short stubs that the coupled-line includes, and multi-stage impedance transformers (ITFs) that suppress impedance mismatch [3]. Theoretically, zero amplitude and phase imbalance (δA and $\delta\theta$; deviation from appropriate amplitude and phase difference at Σ or Δ port), perfect isolation (IL; signal leakage between two input ports. “perfect” means $-\infty$ dB), return loss (RL; signal loss due to reflection at the input port) ≤ -20 dB are obtained over $BWR \geq 2.3$.

The above analytical solution uniquely determines the impedance solution by imposing the perfect matching condition on the coupled-line alone [3]. However, when using a coplanar coupled-line (CPCL), which is a single layer and does not require a crossover to a stripline, it is difficult to realize the analytical solution because it requires a wide full width ($:= g + 2w + 2s$, see Sect. 3, Fig. 1a,b) close to the ratrace scale. However, the perfect matching condition can be removed [7], and even in that case perfect isolation and zero imbalance are preserved.

With these in mind, we propose a method to obtain a solution that narrows the full width of a CPCL by removing the perfect matching condition expanding the range of impedance solutions, and exhaustively searching by utilizing a fast circuit simulator.

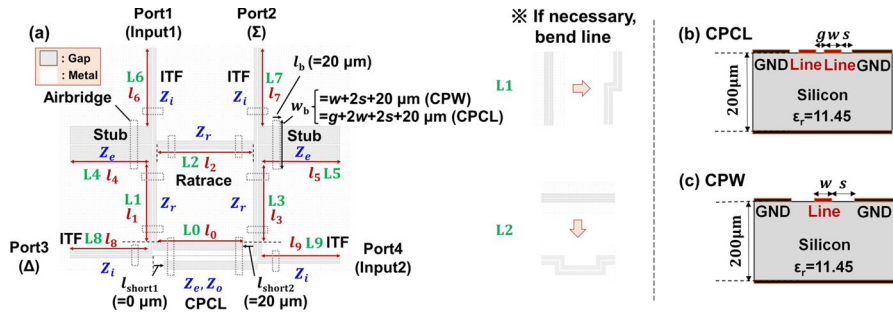


Fig. 1 **a** The basic structure of the Magic-T used in this study (top view; not to scale) and definitions of impedance (blue), number of transmission lines (green), and line lengths (brown). The height of the airbridge h_b is fixed at 10 μm . Note that L1 and L2 are bent according to the optimization results. **b**, **c** Definition of line width and gap space of the CPCL and CPW (cross-sectional view). ϵ_r means relative permittivity of silicon

3 Requirements and Basic Structure

Here, we summarize the design conditions of our Magic-T. The bandwidth ratio criterion was given to be $BWR = 2.3$. This value corresponds to the limited instantaneous bandwidth of horn-coupled OMTs [8]. For return loss and phase imbalance, the values obtained in [3] are used as the standard, and amplitude imbalance is set to be equivalent to the amplitude deviation due to phase imbalance. From the above, it was determined that the requirements for the Magic-T are $RL \leq -20$ dB, $\delta\theta \leq 1.5^\circ$, and $\delta A \leq 0.2$ dB be satisfied over $BWR \geq 2.3$. In this study, to verify the design method through relatively simple manufacturing and performance evaluation in the future, the center frequency was set to 10 GHz and the minimum line width of the transmission line was set to 3 μm , which can be fabricated reliably using stepper lithography. Furthermore, the input impedance of the Magic-T is 100 Ω to match the impedance of CPW feed lines connected to the OMT probe via impedance transformation.

In order to meet the required $BWR \geq 2.3$, the basic structure of our Magic-T adopted the circuit model proposed in [3], which includes an coupled-line, stubs, and 1-stage ITFs. A schematic diagram of the structure is shown in Fig. 1a. Here, the even and odd mode impedances of the coupled-line are Z_e and Z_o , respectively, the impedance of the ratrace section is Z_r , and the impedance of the ITF is Z_i . Due to the matching condition of the ratrace section, $Z_r = 2Z_e Z_o / (Z_e - Z_o)$ is fixed. In addition, the impedance of the stub is fixed to Z_e so that theoretically zero imbalance is achieved. The line numbers are defined as shown in Fig. 1a; green characters. As transmission lines, a CPCL and coplanar waveguides (CPWs) are used and the line width and the gap space are defined as shown in Fig. 1b,c. Since the full width increases as g of the CPCL increases, we fix $g = 3$ μm in this study. The substrate dielectric is assumed to be a silicon wafer (permittivity = 11.45 [9], height = 200 μm), and dielectric losses are assumed to be negligible. Niobium was assumed as a superconductor, and the kinetic inductance is about one order of magnitude smaller than the CPW geometric inductance (calculated from supplementary of [10]). To save computation time, we neglected the sheet impedance and film thickness.

4 Search for Coplanar Coupled-Line Solutions

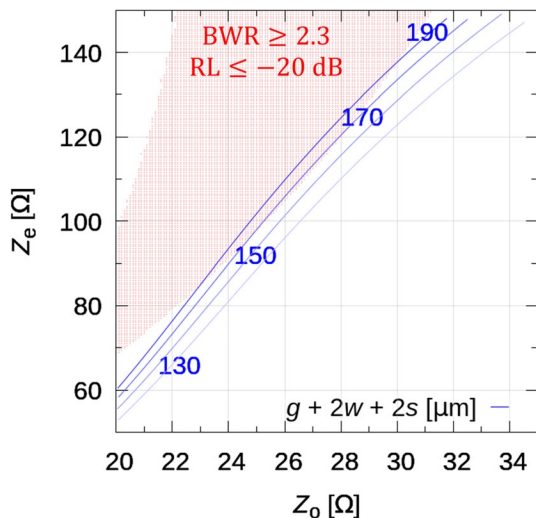
The full width range of the CPCL that can achieve $BWR \geq 2.3$ was estimated by removing the perfect matching condition ($Z_e Z_o = Z_r^2$) for the coupled-line alone. Transmission line models of Quite Universal Circuit Simulator (QUCS) were used in calculations of the Magic-T performance. The relationship between impedance and full width of a CPCL was calculated using analytical solution [11]. The impedance search range in the calculation of bandwidth ratio was set to $Z_e = 50\text{--}150\text{--}0.5 \Omega$, $Z_o = 15\text{--}45\text{--}0.1 \Omega$ and $pZ_i \equiv Z_i/Z_r = 0.60\text{--}1.30\text{--}0.002$ (min.-max.-step). As for the bandwidth ratio, the widest one among the pZ_i search ranges was selected. Note that pZ_i was taken to be close to 1, and it was confirmed that optimal solution was not obtained at the edge of the range of pZ_i . Here, the effective even mode and odd mode permittivity of the coupled-line (ϵ_{re} and ϵ_{ro}) were assumed to be both 1, and the electrical length was $\lambda/4$, which is an ideal case. Figure 2 shows the relationship between the full width and the region where $BWR \geq 2.3$ for each Z_e and Z_o in the above search range.

From Fig. 2, we confirmed that the lower limit of the CPCL full width is about $170 \mu\text{m}$. This is about one order of magnitude lower than the analytical solutions $Z_e = 163.2 \Omega$ and $Z_o = 28.0 \Omega$ given by [3]. Based on the results, we will attempt to optimize the impedance and line length of the Magic-T in the next section.

5 Optimization

In this section, we first determine the cross-sectional CPCL and CPWs geometries by optimizing impedances based on bandwidth and full width criteria. Then we determine $l_1, l_2, l_3, l_4, l_6, l_9$ by minimizing the imbalance. Lastly we determine l_5, l_7 and l_8 to minimize both return loss and insertion loss.

Fig. 2 Full width of a CPCL and bandwidth ratio of the Magic-T versus impedance Z_e and Z_o of the coupled-line. (Blue line): Full width calculated from the analytical solution for $g = 3 \mu\text{m}$ [11]. (Red region): Region of impedance above the bandwidth ratio of 2.3 calculated by the equivalent circuit model of QUCS. Contours with full width $\geq 170 \mu\text{m}$ overlap the region where $BWR \geq 2.3$



5.1 Impedance Optimization

The impedances of the Magic-T were optimized using QUCS to determine all line widths. Electrical lengths of CPWs were fixed to $\lambda/4$. For CPCL, the even and odd mode permittivities are slightly different. Therefore, the physical length was adjusted so that for the ideal model with the $\lambda/4$ electrical lengths for both modes, the sum of squared residuals of the S-matrix is minimized. Under these conditions, in order to prevent not only a CPCL but also stubs from becoming wider, we searched for a pair of impedances that minimizes the full width of a CPCL and stubs that satisfy the requirements. Referring also to the results obtained in section 4, independent variables and their ranges were $Z_e = 80\text{--}110\text{--}0.1 \ \Omega$, $Z_o = 21\text{--}28\text{--}0.1 \ \Omega$, $pZ_i = 0.8\text{--}1.2\text{--}0.005$ (min.-max.-step). Note that the return loss criterion here is set to about -20.4 dB ($= 0.095 \text{ mag}$) to provide margin for later optimization.

The optimized impedance and line widths are shown in Table 1. The obtained optimal CPCL and stub linewidths are $(g, w, s) = (3.0, 41.9, 45.8) \ \mu\text{m}$ and $(w, s) = (3.0, 16.9) \ \mu\text{m}$, respectively. Hence, the full width of the CPCL was less than $180 \ \mu\text{m}$. Note that the line widths w and s of the CPCL were adjusted within $+1.0 \ \mu\text{m}$ to match the impedance calculated by an electromagnetic simulator (Sonnet), respectively. Furthermore, it was confirmed that the optimal solution was not obtained at the edge of the impedance range. From Fig. 3a, the solution satisfied the requirements, and the power leakage to Σ port in odd mode input and to Δ port in even mode are $\leq -40 \text{ dB}$.

5.2 Line Length Optimization

As a preparation for the optimization of the line length ($l_0\text{--}l_9$), an initial estimate of the line length is determined by compensating for the discontinuity at the connection of lines (transition). S-matrices of the transitions between cross junctions in the ratrace section, between stubs and the GNDs, and between ports and the ITFs were calculated by Sonnet, respectively. The dimensions of the $100 \ \Omega$ CPW extending from the port were $(w, s) = (3, 22.7) \ \mu\text{m}$. In addition, the line lengths of the equivalent circuit model of each of the transitions were fitted so that the residual sum of squares of the S-matrices between the transition and the equivalent circuit model

Table 1 Optimized parameters in Sect. 5. In Sect. 5.1, the impedance Z and the dimensions w , s , and g of each transmission line were determined, and in Sect. 5.2, the physical line lengths l_p were determined. Note that the lengths are rounded to the nearest $10 \ \mu\text{m}$, which is less than the step in the optimization

Param	L0	L1	L2	L3	L4	L5	L6	L7	L8	L9
$Z \ [\Omega]$	$(Z_e, Z_o) =$ (93.5, 24.3)	65.7	65.7	65.7	93.5	93.5	65.0	65.0	65.0	65.0
$w \ [\mu\text{m}]$	41.9	3.0	3.0	3.0	3.0	3.0	3.0	3.0	3.0	3.0
$s \ [\mu\text{m}]$	45.8	4.3	4.3	4.3	16.9	16.9	4.2	4.2	4.2	4.2
$g \ [\mu\text{m}]$	3.0	–	–	–	–	–	–	–	–	–
$l_p \ [\mu\text{m}]$	2900	2830	2950	2820	3020	3020	2980	2920	2950	2940

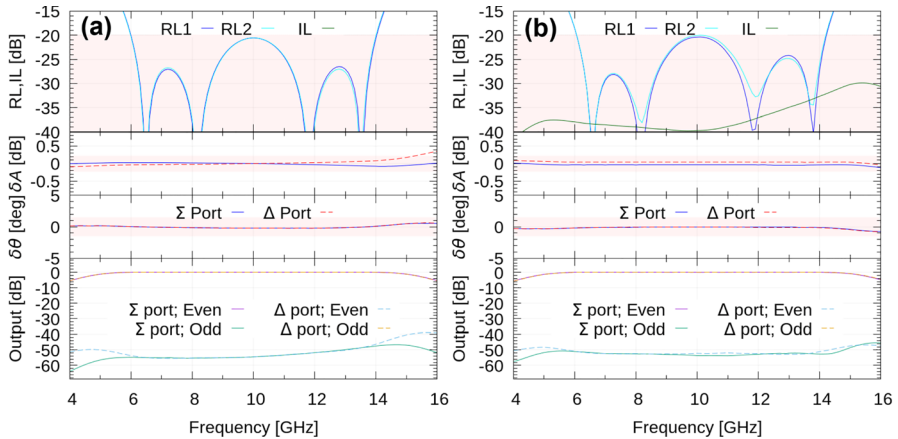


Fig. 3 **a** Calculated performance of the Magic-T after impedance optimization with idealized transmission lines in QUCS. The effective permittivity of the coupled-line was calculated by equations from [11]. **b** Calculated performance after optimizing the transmission line length considering the transition. The S-matrices of all transmission lines and transitions are computed in Sonnet and cascaded in QUCS. (top): return loss (RL) in input 1,2 and isolation (IL). (second): amplitude imbalance (δA), (third): phase imbalance ($\delta\theta$). (bottom): transmitted power in even and odd modes at the Σ and Δ ports. Thin red boxes indicate the range where the required performance is satisfied. Note that the IL in (a) is not shown since it is less than -40 dB at all frequencies in the axis

is minimized, and the initial values of the line lengths of the Magic-T were set to the values compensated for the deviations. Then, the line length was exhaustively searched as follows. In the first step, l_1, l_2, l_3, l_4, l_6 and l_9 , which affect the imbalance at the output, were determined to be the values that minimize the imbalance. Next, l_5, l_7 and l_8 , which have little effect on the imbalance at the output, were determined to minimize the in-band average of return loss and isolation power. Note that l_0 was fixed in the same way as in Sect. 5.1. The search range of line length was $0.93\text{--}1.07\text{--}0.01$ times (min.-max.-step) the initial value, and it was confirmed that the optimal solution was not obtained at the edge of the range. Finally, the idealized transmission lines in QUCS were replaced with CPWs and a CPCL S-matrices computed by Sonnet. The optimized line lengths are listed in Table 1, and the performance is shown in Fig. 3b.

Figure 3b shows that even considering the transition discontinuity, $RL < -20$ dB, $IL < -30$ dB, $\delta A < 0.1$ dB, and $\delta\theta < 0.5^\circ$ from 6.1 to 14.1 GHz (BWR = 2.3), which satisfy the required performance. Moreover, the power leakage to the Σ port in the odd mode and to the Δ port in the even mode can be negligible (~ -50 dB).

6 Conclusion

We established a design method of a dielectric/metal single-layer Magic-T suitable for wideband and detector arraying and applied this method to a scaled model Magic-T (~ 10 GHz). The conditions of conventional analytical solutions were partially removed, and fast circuit simulations were utilized to search for a solution with

a narrower CPCL full width. Furthermore, we showed that a design solution that satisfies the requirements can be obtained by optimizing the impedance and line length of the Magic-T, taking into account the discontinuity of transitions.

In this study, the sheet impedance was neglected to save the computation time. However, even if it is considered, almost the same method can be used by calculating the CPCL impedance in advance from an electromagnetic simulator.

Our next step is to fabricate and measure the design solution. In the future, the design method will be applied to millimeter and submillimeter waves. A potential challenge is that the shorter length of the ratrace requires a narrower CPCL full width (i.e., narrower center gap g). Due to manufacturing issues, it may be difficult to scale $g = 3 \mu\text{m}$ in the 10 GHz model directly to millimeter or submillimeter waves. In this case, a trade-off must be considered between a wider g (i.e., easier fabrication) and larger transition discontinuities due to a wider CPCL full width. However, the performance degradation could be minimized using the established method, which can minimize the CPCL full width and compensate for the transition discontinuities.

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