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Förster Resonance Energy Transfer Study of Cytochrome c—Lipid Interactions

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Abstract Specific interactions between a mitochondrial hemoprotein cytochrome c (cyt c) and cardiolipin, a lipid component of mitochondrial membrane, are crucial to electron shuttling and apoptotic activities of this protein. In the present study the Förster resonance energy transfer (FRET) between anthrylvinyl-labeled phosphatidylcholine as a donor and heme moiety of cyt c as an acceptor was employed to give a quantitative characterization of the protein binding to the model membranes from the mixtures of phosphatidylcholine (PC) with phosphatidylglycerol (PG), phosphatidylserine (PS) or cardiolipin (CL) in different molar ratios. The multiple arrays of the FRET data were globally analyzed in terms of the model of energy transfer in two-dimensional systems combined with the scaled particle adsorption model. The arguments in favor of the specificity of cyt c interactions with CL were obtained, including the higher adsorption potential and the deeper protein insertion in the lipid bilayer.

Keywords Cytochrome $c \cdot$ Cardiolipin \cdot Protein-lipid interactions \cdot Förster resonance energy transfer

Introduction

The Förster resonance energy transfer (FRET) has long been recognized as one of the most powerful fluorescence techniques capable of providing both qualitative and quantitative information on a diversity of processes occurring in biological membranes [1–4]. This technique proved to be particularly efficient in elucidating the nature of interactions between the protein and lipid membrane components [5]. FRET has been employed to gain insight into protein adsorption on a lipid bilayer [6], proximity relationships in the protein-lipid complexes [7], lipid domain formation [8], protein conformational changes in a membrane environment [9, 10], lipid-mediated protein aggregation and fibril formation [11, 12], etc. A great deal of the FRET membrane studies are carried out with the model protein-lipid systems composed of lipid vesicles of a given composition and isolated proteins. There are different formats of FRET measurements in which donors and acceptors may be intrinsic protein fluorophores, covalent protein or lipid labels, and fluorescent dyes non-covalently associating with proteins or lipids. One of these formats involves the case when the energy acceptors pertaining to the water-soluble protein distribute between the aqueous and lipid phases, while the donors are embedded in a lipid bilayer at a known depth through covalent labeling of lipid molecules in a certain position. In this case the efficiency of energy transfer depends both on transverse location of the protein molecule in a lipid bilayer and the acceptor surface concentration, which, in turn, is determined by the extent of the protein-membrane binding. Thus, in principle, it is possible to extract both structural (the depth of the protein insertion in the lipid bilayer) and binding (association constant and binding stoichiometry) parameters from the FRET measurements. As we demonstrated previously, the global analysis of two-dimensional array of FRET data acquired by varying both protein and lipid concentrations allows to circumvent the problem of parameter cross-correlation and determine both structural and binding characteristics of the protein-lipid complexes with high accuracy and statistical significance [13]. In the present work this approach was



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employed to gain further insights into lipid bilayer interactions of cytochrome c.

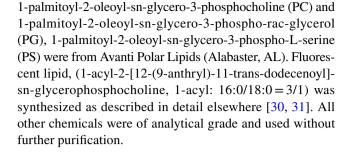
Cytochrome c (cyt c), a highly basic hemoprotein found in the intermembrane space of mitochondria, plays an essential role in the two extensively studied physiological processes – electron transfer in the respiratory chain of the inner mitochondrial membrane and programmed cell death, apoptosis [14–16]. Both these processes are thought to be mediated by cyt c interaction with cardiolipin (CL), an anionic mitochondrial glycerophospholipid with unique physicochemical properties [17]. The molecular mechanisms of cyt c binding to CL have long been a focus of considerable research efforts, shedding light on a variety of important aspects of this problem [18–20]. The current consensus points to a multistep character of cyt c – CL complexation that may be divided into at least three interdependent stages: (i) the protein adsorption at the lipid-water interface initially driven by electrostatic forces and then followed by hydrogen bonding and hydrophobic interactions [21–23]; (ii) CLinduced conformational changes of cyt c involving varying degrees of the protein unfolding and opening of the heme crevice [24–26]; (iii) structural transformations of the lipid bilayer associated with CL clustering and the formation of lipid phases with negative curvature [27–29].

However, notwithstanding considerable advances in molecular-level understanding of cyt c-CL interactions, the precise mechanisms of this process still remain to be fully elucidated. One of the unresolved questions is how specific features of CL distinguishing it from other anionic phospholipids operate in determining the structural and functional peculiarities of cyt c-CL complexes. While addressing this issue in the present study, we employed the Förster resonance energy transfer to compare the lipid-associating characteristics of cyt c in the model protein-lipid systems containing as a lipid component lipid vesicles from phosphatidylcholine and varying proportions of one of the anionic phospholipids, cardiolipin, phosphatidylglycerol (PG) or phosphatidylserine (PS). To this end, the following objectives were pursued: (i) to explore the binding of cyt c to lipid bilayers by measuring the efficiency of energy transfer between anthrylvinyl-labeled PC (AV-PC) as a donor and heme group of the protein as an acceptor; (ii) to estimate the binding and structural parameters for cyt c-lipid complexes through the global analysis of the experimental data in terms of the combined adsorption-FRET model.

Materials and Methods

Chemicals

Bovine heart cardiolipin and horse heart cyt c (oxidized form) were purchased from Sigma (St. Louis, MO, USA).



Preparation of Lipid Vesicles

Large unilamellar vesicles were prepared by extrusion technique from PC mixture with PG (10, 20 and 40 mol %), PS (10, 20 and 40 mol %) or CL (5, 11 and 25 mol% CL). Appropriate amounts of lipid stock solutions were mixed in chloroform, evaporated to dryness under a gentle nitrogen stream, and then left under reduced pressure for 1.5 h to remove any residual solvent. The dry lipid residues were subsequently hydrated with 5 mM sodium-phosphate buffer, pH 7.4 at room temperature to yield lipid concentration of 1 mM. Thereafter, the sample was subjected to 15 passes through a 100 nm pore size polycarbonate filter (Millipore, Bedford, USA), yielding the liposomes of desired composition. AV-PC (0.3 mol % of total lipid) was added to the lipid mixture prior to the solvent evaporation. The concentration of fluorescent lipid was determined spectrophotometrically using the anthrylvinyl extinction coefficient $\varepsilon_{367} = 9 \times 10^3$ M⁻¹cm⁻¹(in ethanol) [31]. Hereafter, the liposomes containing 10, 20 or 40 mol% PG (PS) are referred to as PG10, PG20, PG40 (PS10, PS20, PS40), while the liposomes containing 5, 11 or 25 mol% CL are referred to as CL5, CL11, CL25, respectively.

Fluorescence Measurements

Fluorescence measurements were performed at 25 °C with LS-50 spectrofluorimeter (Perkin-Elmer Ltd., Beaconsfield, UK) using 10-mm path-length quartz cuvettes. AV-PC emission spectra were recorded with 367 nm excitation. Excitation and emission bandpasses were set at 5 nm. The FRET experiments were conducted with AV-PC as a donor and heme group of cyt c as an acceptor. Fluorescence intensity measured in the presence of cyt c at the maximum of AV emission (434 nm) was corrected for the reabsorption and inner filter effects using the following coefficients [32]:

$$k = \frac{\left(1 - 10^{-A_o^{ex}}\right) \left(A_o^{ex} + A_a^{ex}\right) \left(1 - 10^{-A_o^{em}}\right) \left(A_o^{em} + A_a^{em}\right)}{\left(1 - 10^{-(A_o^{ex} + A_a^{ex})}\right) A_o^{ex} \left(1 - 10^{-(A_o^{em} + A_a^{em})}\right) A_o^{em}}$$

where A_o^{ex} , A_o^{em} are the donor optical densities at the excitation and emission wavelengths in the absence of acceptor,



 A_a^{ex} , A_a^{em} are the acceptor optical densities at the excitation and emission wavelengths, respectively. The relative quantum yield of the donor (Q_r) was determined by measuring the decrease of AV fluorescence upon addition of cyt c:

$$Q_r = \frac{Q_{DA}}{Q_D} \tag{2}$$

where Q_D , Q_{DA} are the donor quantum yields in the absence and presence of acceptor, respectively. The steady-state fluorescence anisotropy of AV fluorophore was measured at the excitation and emission wavelengths of 367 and 434 nm, respectively, with excitation and emission bandpasses set at 10 nm.

Results and Discussion

Comparative Qualitative Analysis of FRET Between AV-PC and Heme Group of cyt c in PC/CL, PC/PS and PC/PG Model Membranes

The Förster resonance energy transfer from AV-PC to the heme group of cyt c was explored by monitoring the decrease of the donor fluorescence upon addition of the acceptor. The FRET experiments were designed to acquire the multiple data arrays required for the proper quantification of the obtained results. The relative quantum yield of AV-PC was measured as a function of the following variables: i) the total protein concentration (P), 12 points taken from the range 0.02–0.54 µM; ii) the total lipid concentration (L), 4 points, from the range 5–40 μ M; iii) the molar fraction of anionic phospholipid (0.1, 0.2, 0.4 for PG and PS; 0.05, 0.11, 0.25 for CL). In Figs. 1, 2, and 3 the relative quantum yield of AV-PC is plotted as a function of the protein and lipid concentration for different types of liposomes. One of the conspicuous features of the presented FRET profiles is that the most effective energy transfer occurs in the CL-containing systems. This finding cannot be explained by the higher surface electrostatic potential of PC/CL bilayers, because the molar proportion of CL was chosen to ensure the values of this parameter similar to those of PC/PG or PC/ PS bilayers. Although the degree of CL protonation has long been a matter of controversy, recent study of Malyshka et al. provided convincing evidence that CL is fully deprotonated at neutral pH, bearing the charge -2e [33]. Accordingly, the surface electrostatic potential calculated in terms of the Gouy-Chapman double layer theory for protein-free lipid vesicles possess the values $\{-90, -126, -162\}$ mV for {CL5, CL11, CL25} liposomes, and {-92, -126, -161} mV for {PG(PS)10, PG(PS)20, PG(PS)40} liposomes, respectively.

Also worthy of mention is the fact that the most pronounced differences between the FRET curves measured for the liposomes with similar surface charge are observed at the lowest content of anionic lipids (Fig. 1). This presumably arises from the distinctions between the different types of lipid vesicles in the extent of the protein binding and the heme position with respect to the lipid bilayer surface. The relative quantum yield of the donor (Q_r) is determined by the two main parameters: the amount of the membrane-bound acceptors that is equal to the concentration of bound protein (B) and the distance of the heme group from the bilayer midplane (d_c) . In turn, the concentration of the membrane-adsorbed protein is a function of the association constant (K_a) and the number of lipids per bound protein (n). Thus, the experimental $Q_r(P)$ profiles depend on d_c value and actually reflect the behavior of B(P) curves. Evidently, in the case of complete proteinmembrane binding d_c becomes the only parameter determining the differences between the obtained FRET profiles. To describe these differences in quantitative terms, we employed the theoretical approach combining the adsorption and FRET models purposefully developed for the protein-lipid systems.

Quantification of the Energy Transfer Profiles in Terms of the Combined Adsorption-FRET Model

The association of cyt c with the model lipid membranes was analyzed in terms of the scaled particle adsorption model which adequately describes the excluded area interactions between the adsorbing protein molecules [34, 35]. In terms of this model the activity coefficient of the adsorbed spherical ligand is given by:

$$\gamma = \frac{1}{1-\theta} \exp\left(\frac{3\theta}{1-\theta} + \left[\frac{\theta}{1-\theta}\right]^2\right); \ \theta = nB/L_{out}$$
 (3)

where θ is the fraction of surface area occupied by the adsorbed protein; B is the concentration of the bound species; n stands for the number of lipids per molecule of the adsorbed protein; L_{out} is the concentration of accessible lipids related to the total lipid concentration (L) as $L_{out} = 0.5L$; S_L is the mean area per lipid molecule:

$$S_L = f_{PC}S_{PC} + f_AS_A \tag{4}$$

here f_{PC} , f_A are the mole fractions of PC and one of the three anionic phospholipids; S_{PC} is the mean area per PC molecule (0.65 nm²); S_A is the mean area per molecule of anionic phospholipid taken as 0.65 nm² for PG or PS, and 1.2 nm² for CL. The adsorption isotherm is described by the following equation [34]:

$$K_a(P - B) = \theta \gamma(\theta) \tag{5}$$

To reproduce the experimental FRET curves, the above adsorption model was combined with the model of energy



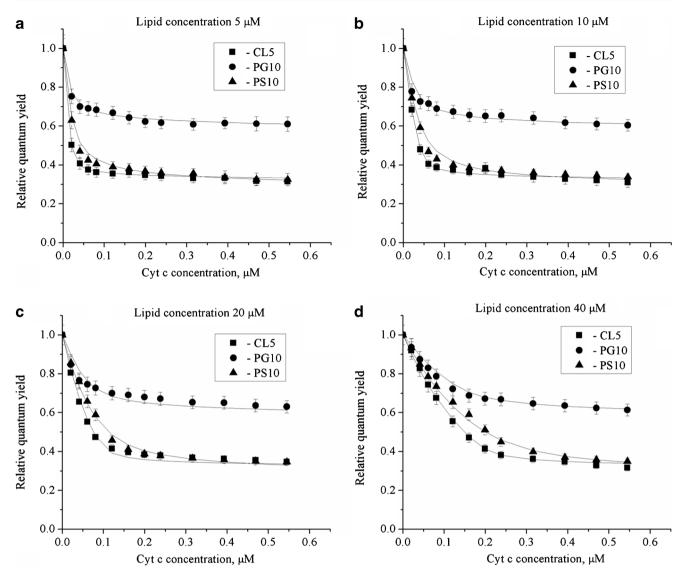


Fig. 1 FRET profiles representing the relative quantum yield of AV-PC as a function of cytochrome c concentration for CL5, PS10 and PG10 liposomes. Lipid concentration was 5 μ M (a), 10 μ M (b),

 $20~\mu M$ (c), $40~\mu M$ (d). (b) and (c). Solid lines show theoretical curves calculated by numerical solution of Eqs. (3)–(11) with the model parameters presented in Table 1

transfer in two-dimensions, initially developed by Fung and Stryer [36] and extended in our previous studies to allow for the distance dependence of the orientation factor in the membrane systems and the possible heterogeneity of the donor or acceptor populations [37]. The examined cyt c – lipid systems were treated as containing one acceptor plane located at the distance d_c from the bilayer center, and two donor planes separated by a distance d_t . In this case, the relative quantum yield of the donor is given by:

$$S_1(\lambda) = \int_{d_c - 0.5d_r}^{\infty} \left[1 - \exp\left(-\lambda \kappa_1^2(R) \left(\frac{R_o^r}{R}\right)^6\right) \right] 2\pi R dR \quad (7)$$

$$S_2(\lambda) = \int_{d_c + 0.5d_t}^{\infty} \left[1 - \exp\left(-\lambda \kappa_2^2(R) \left(\frac{R_o^r}{R}\right)^6\right) \right] 2\pi R dR \quad (8)$$

$$Q_r = 0.5 \times \left(\int_0^\infty \exp(-\lambda) \exp\left[-\frac{B}{S_L L_{out}} S_1(\lambda) \right] d\lambda + \int_0^\infty \exp(-\lambda) \exp\left[-\frac{B}{S_L L_{out}} S_2(\lambda) \right] d\lambda \right)$$
 (6)



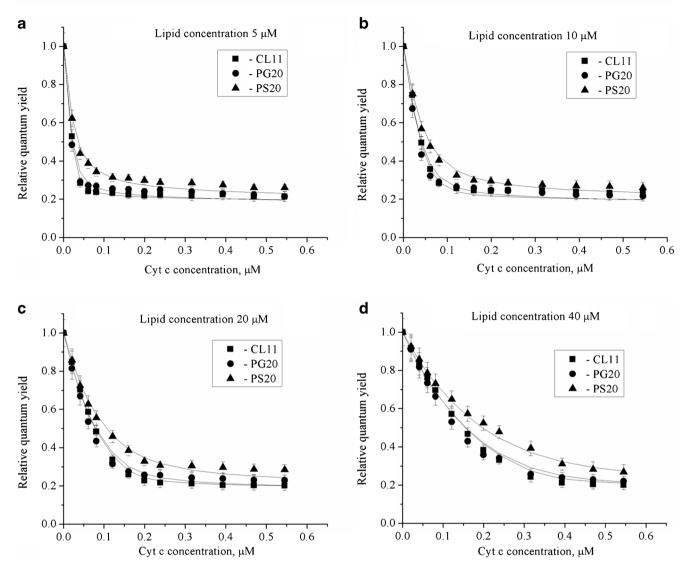


Fig. 2 FRET profiles representing the relative quantum yield of AV-PC as a function of cytochrome c concentration for CL11, PS20, PG20 liposomes. Lipid concentration was 5 μ M (a), 10 μ M (b),

 $20~\mu M$ (c), $40~\mu M$ (d). (b) and (c). Solid lines show theoretical curves calculated by numerical solution of Eqs. (3)–(11) with the model parameters presented in Table 1

where R is the donor–acceptor separation, $\lambda = t/\tau_d$; τ_d is the lifetime of excited donor in the absence of acceptor. By representing the Förster radius as $R_o = \left[\kappa^2(R)\right]^{1/6} \cdot R_o^r$, it follows that

$$R_o^r = 979 \left(n_r^{-4} Q_D J \right)^{1/6} J = \frac{\int\limits_0^\infty F_D(\lambda) \epsilon_A(\lambda) \lambda^4 d\lambda}{\int\limits_0^\infty F_D(\lambda) d\lambda}$$
(9)

here n_r is the refractive index of the medium $(n_r=1.37)$, Q_D is the donor quantum yield (0.8), J is the overlap between the donor emission $(F_D(\lambda))$ and acceptor absorption $(\varepsilon_A(\lambda))$ spectra [38]. Assuming that the donor emission and acceptor absorption transition moments are symmetrically distributed within the cones about certain axes D_x and A_x , the distance-dependent orientation factor is given by:

$$\kappa_{1,2}^{2}(R) = d_{D}d_{A}\left(3\left(\frac{d_{c} \mp 0.5d_{t}}{R}\right)^{2} - 1\right) + \frac{1 - d_{D}}{3} + \frac{1 - d_{A}}{3} + \left(\frac{d_{c} \mp 0.5d_{t}}{R}\right)^{2}\left(d_{D} - 2d_{D}d_{A} + d_{A}\right)$$

$$\tag{10}$$



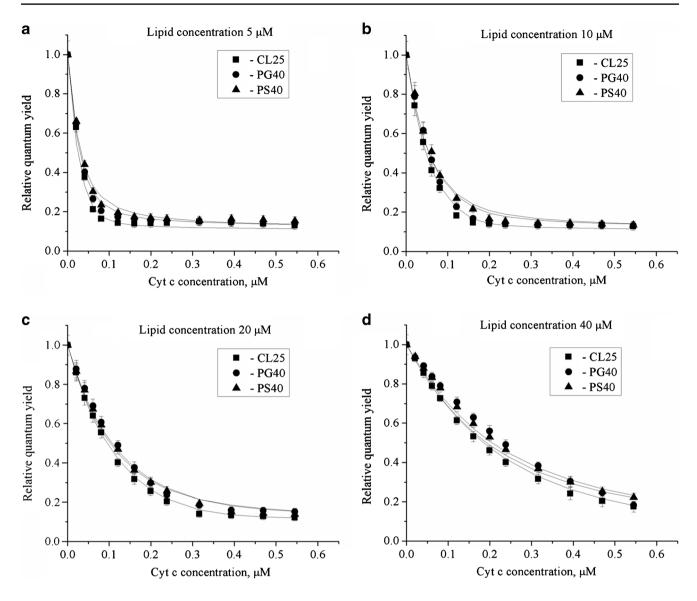


Fig. 3 FRET profiles representing the relative quantum yield of AV-PC as a function of cytochrome c concentration for CL25, PS40, PG40 liposomes. Lipid concentration was 5 μ M (a), 10 μ M (b),

 $20~\mu M$ (c), $40~\mu M$ (d). (b) and (c). Solid lines show theoretical curves calculated by numerical solution of Eqs. (3)–(11) with the model parameters presented in Table 1

$$d_{D,A} = \left\langle d_{D,A}^{x} \right\rangle \left(\frac{3}{2}\cos^{2}\alpha_{D,A} - \frac{1}{2}\right) \left\langle d_{D,A}^{x} \right\rangle = \left(\frac{3}{2}\cos^{2}\psi_{D,A} - \frac{1}{2}\right)$$
(11)

where $\psi_{D,A}$ are the cone half-angles, $\alpha_{D,A}$ are the angles made by D_x and A_x with the bilayer normal N. The axial depolarization factors $\langle d_D^x \rangle$ and $\langle d_A^x \rangle$ are related to the experimentally measurable steady-state (r) and fundamental (r_0) anisotropies of the donor and acceptor [39]:

$$d_{D,A}^{x} = \pm \left(r_{D,A}/r_{0D,A}\right)^{1/2} \tag{12}$$

The donor depolarization factor d_D was calculated from Eq. 12 using the results of AV-PC anisotropy measurements. The r_D values appeared to fall in the range from 0.05 to 0.07.

The fundamental anisotropy of anthrylvinyl fluorophore was taken as 0.08 [30]. According to the linear dichroism data the transition moment of cyt c lies within the porphyrin plane [40], i.e. $\psi_A = \pi/2$. If cyt c adopts a certain specific orientation relative to the lipid-water interface, the porphyrin plane makes an angle α_A with the bilayer surface. The results of FRET measurements presented here were analyzed within the two limiting assumptions: i) $\alpha_A = 0$ (the porphyrin plane is parallel to the bilayer surface) and ii) $\alpha_A = \pi/2$ (the porphyrin plane is perpendicular to the bilayer surface), thus allowing to estimate the upper and lower bounds of the d_c value.

Altogether, the set of the above equations was used to perform a global analysis of the FRET data arrays obtained upon varying the protein and lipid concentrations for a given



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System	Binding stoichiom- etry (n)	Adsorption potential J_A , in units of RT per anionic lipid	Adsorption potential J_S , is units of RT per nm ²	Heme distance from bilayer center d_c , nm	f
CL ₅	48 ^{±2}	$7.6^{\pm0.3}$	$0.6^{\pm0.03}$	2.4 ^{±0.1}	7.4×10^{-4}
CL_{11}	$28^{\pm 2}$	$5.9^{\pm0.3}$	$0.9^{\pm0.05}$	$2.4^{\pm0.1}$	1.7×10^{-3}
CL ₂₅	$14^{\pm 3}$	$5.2^{\pm0.3}$	$1.6^{\pm0.09}$	$2.7^{\pm0.2}$	1.9×10^{-3}
PS ₁₀	$28^{\pm 2}$	$5.1^{\pm0.2}$	$0.8^{\pm0.03}$	$3.1^{\pm0.1}$	1.2×10^{-3}
PS ₂₀	18 ^{±2}	$4.0^{\pm0.2}$	$1.2^{\pm0.06}$	$3.2^{\pm0.1}$	1.3×10^{-3}
PS ₄₀	10 ^{±2}	$3.6^{\pm0.3}$	$2.2^{\pm0.18}$	$3.2^{\pm0.3}$	2.8×10^{-3}
PG10	$50^{\pm 6}$	$2.9^{\pm0.4}$	$0.4^{\pm0.06}$	$4.0^{\pm0.4}$	3.5×10^{-3}
PG20	24 ^{±4}	$3.4^{\pm0.5}$	$1.0^{\pm 0.15}$	$2.7^{\pm0.4}$	3.9×10^{-3}
PG40	10 ^{±2}	$3.9^{\pm0.5}$	$2.4^{\pm0.31}$	$3.4^{\pm0.3}$	3.7×10^{-3}

Table 1 The binding and structural parameters derived for cyt c – lipid systems

cyt c – lipid system. The numerical solution of these equations yielded the set of parameters $\{K_a, n, d_c\}$ providing the best agreement between experiment and theory (Table 1). The data fitting procedure was based on the minimization of the function:

$$f = \frac{1}{N} \sum_{i=1}^{N} \left(Q_{ri}^{e} - Q_{ri}^{t} \right)^{2} \tag{13}$$

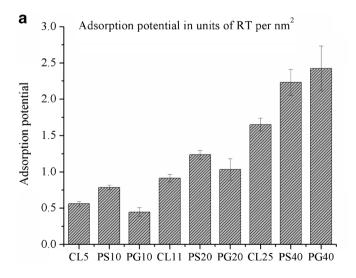
where N is the total number of the experimental points involved in the global analysis (N=64), Q_r^e is the experimental Q_r value, Q_r^t is the relative quantum yield calculated by the numerical integration of Eqs. (3)–(8).

To compare the binding characteristics of cyt c for the systems differing in the chemical nature of the anionic phospholipid headgroup, the recovered association constants were transformed into the adsorption potentials reflecting the variations in both quantities, K_a and n. These quantities

were expressed in the units of RT per nm² (J_S) or per anionic lipid (J_A):

$$J_S = \frac{\ln K_a}{nS_L}; J_A = \frac{\ln K_a}{nf_A} \tag{14}$$

As seen in Fig. 4, the adsorption potential calculated per unit area increases with increasing the surface electrostatic potential for a given type of anionic lipid, pointing to the prevailing role of electrostatic interactions in the cyt c - lipid binding. At the same time, the J_A values decrease in the rows: {CL5>PS10>PG10}, {CL11>PS20>PG20}, {CL25>PG40≥PS40}, clearly showing the preference of cyt c for CL. It is worth noting at this point that our findings are in harmony with those of earlier studies [41–43]. More specifically, the affinity of cyt c for anionic phospholipids assessed by the surface plasmon resonance was found to decrease in the row: CL>PS>PC [43]. The solid-state 31 P NMR studies of cyt c interactions with lipid CL, PS and PG bilayers showed that regardless the main driving force of



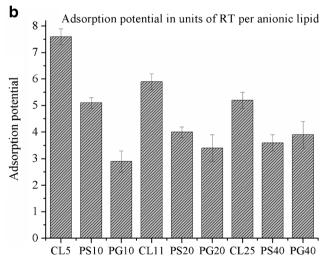


Fig. 4 Adsorption potentials of cytochrome c binding to different types of liposomes expressed in units of RT per nm² (a) and in units of RT per anionic lipid (b)



all these interactions is electrostatic in nature, the structural alterations of the bound protein are lipid dependent [18, 41]. PS and CL were found to produce significant destabilization of cyt c structure, while in complexes with PG the protein retains a native-like conformation [41]. According to our FRET results, PG displays a binding behavior distinct from that of CL or PS, as most clearly seen in Fig. 1, showing the FRET curves measured at the lowest content of anionic lipid when the extent of protein binding was far from being complete. Interestingly, the resemblance between CL and PS in their effect on the structural state of cyt c was established in the early Raman [44, 45] and ³¹P NMR [18] studies. The two different conformational states of cyt c adsorbed on the negatively charged surfaces have been identified – the native-like state I and state II, with the opened heme crevice. The equilibrium between these two states was supposed to be governed by the electrostatic interactions between lysine residues located around the heme crevice and the lipid phosphate groups. It was found that CL and PS promote the conformational equilibrium over the same temperature interval, while PG turned out to be less efficient in this respect, suggesting that cyt c conformation in complexes with PG is close to native [18]. A thorough recent work of Pandiscia & Schweitzer-Stenner provides evidence for an equilibrium between partially unfolded and native-like conformers of cyt c adsorbed on CL-containing bilayers, one of which forms electrostatic contacts with CL headgroups while the other one associates with lipids either via hydrogen bonding or hydrophobic interactions [26].

Along with the binding characteristics, the quantification of the presented experimental data in terms of the combined adsorption-FRET model yielded the structural parameter of the examined systems - the heme distance from the membrane midplane. The upper meaningful limit of this parameter, corresponding to a superficial location of cyt c, can be estimated as $d_c^{\text{max}} = 0.5d_m + r_{cyt} + r_t \sim 4.7 \text{ nm}$, where d_m is the lipid bilayer width (4.6 nm), r_t is the displacement of heme group off the molecule center (ca. 0.7 nm). As seen from Table 1, in all CL-containing membranes the heme group is located closer to the membrane midplane, at the distance ~ 2.4–2.7 nm, while in the other liposomal membranes the heme resides at the greater distance from the bilayer center. Thus, cyt c - CL complexation is distinguished by the deeper membrane penetration of the protein compared to PS- and PG-containing systems. To interpret this particular behavior of CL we must consider some specific properties of this phospholipid, one of which consists in the formation of non-bilayer structures such as hexagonal (H_{II}) phase [28]. This process involves gathering of CL molecules in the regions with negative curvature, i.e. hexagonal cylinders around cyt c and displacement of PC molecules from the protein vicinity. Vladimirov et al. have defined the hexagonal state of cyt c – CL complex as a spherical nanoparticle with a diameter of ~ 11 nm, where the protein is completely covered by the headgroups of 35–40 CL molecules [46]. In any case, the minimum separation between the heme moiety and AV group attached to methyl terminus of PC acyl chain would be greater than 4 nm, that is hard to reconcile with our d_c estimates.

Another specific property of CL is associated with its ability to adopt an extended conformation in which one [22, 23] or two [21] acyl chains protrude out of a membrane entering the hydrophobic cavity of cyt c. This type of protein-lipid interactions is specific for CL and is governed by its tendency to minimize the bending stress created by a high negative curvature. Notably, the ability of CL to adopt both the above configurations, extended anchorage and H_{Π} , is supposed to originate from its conical shape inducing a negative curvature strain in the lipid bilayer [47]. A mechanistic model for the extended lipid anchorage proposed by Kalanxhi and Wallace [22] assumes that ionic contacts between phosphate group and Lys₇₂, Lys₇₃ anchor CL in cyt c hydrophobic channel lined with the residues Tyr₆₇ – Pro₇₁ and Phe₈₂ – Ile₈₅. Later, Sinibaldi et al. put forward the idea that this channel, together with the nonpolar pocket close to Asn₅₂ can simultaneously accommodate two of four hydrocarbon chains of CL [21]. Regardless of the exact mode of cyt c – CL complexation, the d_c values ~ 2.4–2.7 nm imply that the center of heme moiety is located in the aqueous phase close to the surface of PC/ CL bilayers. On the other hand, in the native fold of cyt c the heme is separated from the surface of the protein molecule on average by 1.7 nm. If one assumes that CL-bound cyt c possesses native-like structure, the interfacial position of the heme group is consistent with the membrane embedment of the protein up to the level of terminal groups of hydrocarbon chains. However, this seems unlikely given the strongly basic nature of cyt c. Alternatively, our findings provide indirect arguments in favor of substantial reorganization of the protein structure accompanied by the exposure of nonpolar residues and the changes in the heme environment, that are in line with increasing evidence for CL-induced unfolding of cyt c [24–26]. In comparison with CL, PS and PG seem to provoke less pronounced destabilization of the protein molecule, resulting in the shallower location of cyt c relative to the lipid-water interface.

Conclusions

Cumulatively, the measurements of the Förster resonance energy transfer between anthrylvinyl-labeled phosphatidyl-choline as a donor and heme group as an acceptor revealed the differences in cytochrome c interactions with PC/CL, PC/PS and PC/PG lipid bilayers. The global analysis of the FRET data obtained at varying the molar fraction of anionic phospholipid, protein and lipid concentrations was



performed in terms of the SPT adsorption model combined with the model of energy transfer in two-dimensional systems. The main advantages of the employed theoretical approach are: (i) allowance for the area-excluding interactions between the adsorbed protein molecules and (ii) taking into account the distance dependence of the orientation factor in the membrane systems. The differences between CL on the one hand, and PS/PG on the other hand were quantitatively expressed in terms of the following parameters: the number of lipids per bound protein at saturation; adsorption potential and position of heme group relative to the lipid-water interface. The specificity of cyt c for CL is displayed in the higher adsorption potential per anionic lipid and deeper penetration of the protein in the lipid bilayer. The revealed specific features of cyt c – CL interactions may have important biological implications as a means of modulating the electron-transfer and apoptotic functions of this protein.

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