

Effect of Ag layer thickness on optical and electrical properties of ion-beam-sputtered $TiO₂/Ag/TiO₂$ multilayer thin film

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ABSTRACT

 $TiO₂/Ag/TiO₂$ multilayer films have been deposited by Ion-beam sputtering technique with different Ag thickness. X-ray reflectivity technique with 3-layer fitting model has been used to calculate thickness, density, and roughness of the multilayer films. Gaussian fitting of XRD and Raman data is used to detect the crystalline growth of films while FTIR data confirm the presence of metal layer in the films. A phonon confinement effect originating from decrease in crystallite size is also observed with blue shift in some of the Raman active modes. Effect of Ag layer thickness on optical and electrical properties of multilayer films has been presented including a comparative study of our data with previously reported work on $TiO₂/Ag/TiO₂$ transparent conducting oxides. Transmittance of multilayer films in visible region (550 nm) has increased from 65 to 81% with increase in Ag thickness. The band gaps calculated from Tauc plot of these films follow Burstein–Moss effect. Sheet resistance has decreased from 15.41 to 5.92 Ω/\square while carrier concentration and mobility show an increasing trend from 5.94×10^{21} to 1.35×10^{22} and 6.83 to 13.72 cm²/Vs, respectively, with the increase in Ag mid-layer thickness. Haacke's Figure of Merit (FOM) is recorded to be 7.9 \times 10⁻³ Ω^{-1} .

1 Introduction

Transparent conducting oxides (TCOs) have attracted huge attention in thin film development area due to its wide applications including solar cells, optoelectronic devices, sensors, LCDs, and LEDs. A variety of metal oxides like TiO₂, ZnO, SnO₂, and $Nb₂O₅$ have been extensively exploited for developing efficient TCOs. In later phases, these are even doped with Al-, Nb-, Ga-like materials for enhanced electrical and optical properties. In order to replace ITO films for its toxic nature and cost ineffective reasons, metal oxides like $TiO₂$ and ZnO have been widely investigated and found promising results. TiO₂ have been studied for various applications due to their chemical stability, high photosensitivity, high photocatalytic activity,

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and low toxicity [[1\]](#page-10-0). Also, $TiO₂$ has wide band gap $($ > 3 eV) and high refractive index, which makes it highly photoactive in the UV region and an efficient solar cell material [[2,](#page-10-0) [3](#page-10-0)].

In recent development, multilayer thin films have been a great area of interest compared to single-layer TCOs due to its improved optical and electrical behavior. Multilayer TCO has a combination of thin metal layer sandwiched between different metal oxide layers (dielectrics). The thickness ratio of dielectric and metal layer plays an important role in electrical and optical properties of the film. $TiO₂$ is a desired dielectric material for multilayer films due to its wide band gap, high refractive index (2.7), high transparency in visible region, and better chemical stability. Ag is preferred as middle layer in oxide/ metal/oxide (O/M/O) because it results in low sheet resistance and high transmittance in visible region. The variable Ag thickness is important in finding the optimum electrical and optical balance of TCO for industrial application.

The precise thickness variation of mid-layer ranging within few nanometers is challenging and requires much controlled film growth. Although various sputtering methods, such as direct current (DC) magnetron sputtering and radio frequency (RF) magnetron sputtering, are used to develop $TiO_2/Ag/$ $TiO₂$ multilayer TCOs, ion-beam sputtering (IBS) has many advantages over others. It uses a high-energy ion-beam source over the sputtering target which produces a dense material on substrate resulting into high refractive index compared to other deposition techniques. The scattering loss is much lower for ionbeam-sputtered films than using other methods [[3\]](#page-10-0). It also provides better control over deposition parameters like current density and energy of source ion beam, for desired growth of films.

This work has demonstrated the development of $TiO₂/Ag/TiO₂$ multilayer thin films with variation in Ag layer thickness. A comparative analysis of electrical and optical properties of the developed TCOs with previous work has been presented.

2 Experimental details

 $TiO₂/Ag/TiO₂$ multilayer thin films were deposited on quartz and Si (100) substrates using Ion-beam sputtering system (IBS) consisting of hot-cathode Kaufman type 1.5 kV gridded ion source. Titanium target (99.99%) of size $100 \times 100 \times 1$ mm was used with similar dimension of silver target (99.99%). Oxygen and argon gases were used during the sputtering process. The base pressure inside the chamber was 2×10^{-5} Pa, while the working pressure was 2×10^{-3} Pa. Ti and Ag targets were bombarded with ion beam with beam current of 30 mA and 1000 eV ion-beam energy. Depending upon the deposition rate of $TiO₂$ and Ag, the precise sputtering time for both the targets was assigned alternately to obtain the multilayer film with desired thickness (Table 1). TiO₂ thickness of both the upper and lower layer was fixed at 30 nm while Ag (middle layer) thickness was varied at 2 nm, 4 nm, and 8 nm. The samples were named TAT A (single-layer $TiO₂$ — 60 nm), TAT B (TiO₂—30 nm/Ag—2 nm/TiO₂— 30 nm), TAT C $(TiO_2 \rightarrow 30 \text{ nm}/Ag \rightarrow 4 \text{ nm}/TiO_2 \rightarrow$ 30 nm), and TAT D $(TiO₂ - 30$ nm/Ag-8 nm/ $TiO₂$ —30 nm). After the deposition, TAT films were annealed at 500 $^{\circ}$ C in a high vacuum environment.

X-ray reflectivity (XRR) measurement was performed on multilayer films between 2 and 5° with step size of 0.01° using Bruker D8 Discover diffractometer with $CuK\alpha$ (0.154 nm) source. Parratt32 software was used for XRR data fitting and calculations of thickness, density, and roughness of the multilayer films. Crystallographic properties were analyzed using Bruker D8 Advanced X-ray diffractometer (XRD). Scanning electron microscope (SEM) and energy-dispersive X-ray (EDX) data were obtained by field emission scanning electron microscope (FESEM, Sigma 300, Carl Zeiss Microscope Ltd., Germany) operated at 20 kV. UV–Visible spectra of multilayer films were carried in transmission mode using Lambda 20, PerkinElmer UV–Visible spectrometer. Raman modes were identified using Renishaw Invia Raman spectrometer (514.6 nm Ar laser) across $100-1000$ cm^{-1} . Fourier transformed infrared (FTIR) analysis was done using PerkinElmer Frontier spectrometer across $400-4000$ cm^{-1} . Hall

Table 1 TiO₂/Ag/TiO₂ multilayer deposition parameters

| Sample | $O2$ (sccm) | Deposition time (min.) |
|--------------------------|-------------|------------------------|
| TAT A $(TiO2)$ | 2.5 | 42.52 |
| TAT B $(TiO_2/Ag/TiO_2)$ | 2.5 | 21.26/0.20/21.26 |
| TAT C $(TiO_2/Ag/TiO_2)$ | 2.5 | 21.26/0.39/21.26 |
| TAT D $(TiO2/Ag/TiO2)$ | 2.5 | 21.26/0.77/21.26 |

measurement was done by Van der Pauw's method using four-point probe technique. Keithley 2400 sourcemeter was used during the hall measurements with magnetic field of 1 Tesla.

3 Result and discussion

3.1 XRR

X-ray reflectivity is a powerful non-destructive technique to calculate the thickness, density, and roughness of thin films through electron density profile as a function of depth. Figure 1 represents the 3D schematic diagram of $TiO₂$ and $TiO₂/Ag/TiO₂$ multilayer films with different Ag thickness. Figure [2a](#page-3-0) depicts schematic diagram of interaction of X-rays with multilayer films in terms of reflectivity and momentum transfer. The reflected X-rays at each surface create interference fringes (Kiessig fringes) with numerous oscillations resulting into a reflectometry pattern. Thickness of the multilayer film is related to Bragg's law as

$$
n\lambda = 2(d_1\sin\alpha_{i1} + d_2\sin\alpha_{i1} + d_3\sin\alpha_{i3}),
$$
\n(1)

where, d_1 , d_2 , and d_3 are thickness of different layers of film. In XRR technique, a monochromatic beam of X-ray (α_i) is incident on thin film surface and reflected at α_f with wave vector k_i and k_f , respectively. The momentum transfer along the surface normal (z) is given by [[4,](#page-10-0) [5](#page-10-0)]

$$
Q_z = k_f - k_i = \frac{4\pi}{\lambda} \sin \alpha_i, \tag{2}
$$

where λ is the X-ray wavelength.

Figure [2b](#page-3-0) depicts the X-ray reflectivity profile of $TiO₂/Ag/TiO₂$ multilayer films along with $TiO₂$ monolayer. The reflectivity curve was fitted with theoretical model of independent $TiO₂$, Ag, and $H₂O$ adsorbate layers using parratt32 software [[6,](#page-10-0) [7](#page-10-0)]. A good fitting of reflectivity curve with appropriate theoretical model results in precise thickness determination of different layers. Critical angle (θ_c) is an important parameter for calculating electron density and mass density. In our experiment, critical angle ranged between 0.52 and 0.55° while the critical momentum transfer (Q_c) was observed to be 0.035–0.04 \AA^{-1} for TiO₂/Ag/TiO₂ multilayer films. Singh et al. (2019) have described the elaborated theoretical model to calculate the density and roughness of thin film [\[6](#page-10-0)]. Upon applying best suitable theoretical fitting model and calculations, the thickness, density, and roughness of our $TiO₂/Ag/$ $TiO₂$ multilayer films are presented in Table [2.](#page-3-0) Based on these data it was observed that, with the increase in Ag mid-layer thickness, density, and roughness of multilayer films also increased substantially.

3.2 XRD

XRD spectra of $TiO₂/Ag/TiO₂$ multilayer films with different Ag thickness (0–8 nm) are depicted in Fig. [3.](#page-3-0) Peaks were fitted with Gaussian profile to determine peak position, FWHM, and intensity. Table [3](#page-4-0) shows the crystallite size, micro-strain, and dislocation density of multilayer films. The crystallite size was determined using Scherrer formula and the microstrain was calculated using Williamson–Hall equation [[8,](#page-10-0) [9\]](#page-10-0). The inset image in Fig. [3](#page-3-0) represents the magnified version of TiO₂ and Ag peaks (35–58°) of TAT D multilayer spectra for better view and analysis. Upon deconvolution with Gaussian profile fitting, it was observed that all films were crystalline in nature. The diffracted peaks at 2 θ values 26.52 \degree , 37.98°, and 54.35° refer to $TiO₂$ anatase (JCPDS 84-1286) crystal planes (101), (004), and (211), respectively. Few low intensity yet prominent peaks of Ag were observed at 43.98° , 64.26° , and 77.18° in some of the films. These peaks correspond to (200), (220), and (311) planes, respectively (JCPDS 04-0783). TAT A contains only anatase (101) peaks, whereas TAT B, TAT C, and TAT D confirm the presence of both $TiO₂$ and Ag. With the increase in Ag thickness, the intensity of peaks increases along with improved

Fig. 2 (a) Schematic diagram presenting interaction of x-rays with different layers of film. (b) Electron density profile of TiO₂/Ag/TiO₂ multilayer with appropriate fitting model. Ag reflectivity curve is clearly visible (red circle) between the two $TiO₂$ curves

Fig. 3 XRD spectra of TiO₂/Ag/TiO₂ multilayer films. Crystalline $TiO₂$ with anatase phase is observed along with Ag peaks

crystallinity of multilayer films [\[10](#page-10-0), [11\]](#page-10-0). The preferential growth of crystalline $TiO₂$ structure occurs in the direction (101) plane and Ag grows in (200) plane.

The peak at 23.88° cannot be indexed as it neither matches with any of the $TiO₂$ phases nor with Ag nor

any of their mixed phases. Quartz do have a peak at a similar location. However, such a peak does not appear in the XRD profile of our quartz substrate (see Fig. 3). Quartz substrate is well known to have induced lattice mismatch against the deposited film giving rise to high lattice strain [[12\]](#page-10-0). This is an important factor in thin film crystalline growth as high lattice strain may lead to crystal defects and dislocations. Increase in film thickness reduces the strength of strain in the films [[13\]](#page-10-0).

Apart from crystal dislocations, crystallite size plays major role in induced lattice strain in thin films due to grain boundary effect. When crystallite size decreases, the stress caused by internal surface tension induces lattice strain. However, when crystallite size increases, the lattice strain becomes weak due to lower stress field from grain boundaries [\[14](#page-10-0)]. The crystallite size of the multilayer films depicts similar behavior with the increase in Ag thickness. As the crystallite size increases from \sim 11 to 19 nm, a decrease in micro-strain from 1.27×10^{-2} to 1.03×10^{-2} is observed. Dislocation density, which indicates the amount of defects in the films, is seen to be decreasing with the increase in Ag thickness. This

Table 3 Average crystal size, micro-strain, and dislocation density of Ti $Ag/TiO₂$ multilayer films

shows that TAT D multilayer film with 8 nm Ag thickness has fewer defects compared to other films.

3.3 FESEM/EDX

FESEM images of $TiO₂/Ag/TiO₂$ multilayer films are shown in Fig. [4](#page-5-0). Microstructure of all the films are smooth, dense, and free from cracks. Pure $TiO₂$ and Ag (2 nm) film do not show much difference in surface morphology. Ag (4 nm) film looks more dense and shows slight agglomeration on the surface. Some dark spots are visible which appears due to Ag agglomeration over the $TiO₂$ surface, forming islandlike structure confirming the Volmer–Weber model of multilayer films [\[15](#page-10-0), [16\]](#page-10-0). These island agglomerate in Ag (8 nm) film, thereby, forming a continuous layer. Inset image of Fig. [4](#page-5-0)d represents the crystallite size of TAT D films using imageJ software. The average crystallite size is \sim 20 nm which is close to the values calculated by XRD spectra using Scherrer formula.

EDX compositional analysis of multilayer films is presented in Fig. [5.](#page-6-0) The elemental composition of films shows the presence of both Ti and Ag in multilayer films along with Si substrate due to high-energy EDX ion beam penetrating into the Si substrate. The wt% of Ag increases from 0.2 to 0.7% confirming the increase in Ag thickness in subsequent multilayers. The overall composition of Ag in $TiO₂/Ag/TiO₂$ Multilayer films ranges between 2.5 and 10.3 wt% which is close to the expected value of Ag composition in multilayer by volume from 3.2 to 11.7% [[17](#page-10-0)].

3.4 Raman

Figure [6](#page-6-0) presents micro-Raman analyses of $TiO₂/$ Ag/TiO2 multilayer films performed over 100–800 cm^{-1} to obtain the vibrational modes present in samples. Raman spectra were fitted with Gaussian profile to obtain peak positions, FWHM, intensity, and area under the peak. TiO₂ anatase has six active Raman modes: $1A_{1g}$, $2B_{1g}$, and $3E_g$. The Raman shifts observed at 144 cm⁻¹ (E_g), 197 cm⁻¹ (E_g), 395 cm⁻¹ (B_{1g}) , 515 cm⁻¹ (A_{1g}/B_{1g}) , and 636 cm⁻¹ (E_g) correspond to the characteristic peaks of tetragonal $TiO₂$ anatase phase [\[18](#page-10-0), [19](#page-10-0)]. The active mode E_g (144 cm⁻¹, 395 cm^{-1} , and 636 cm^{-1}) refers to O–Ti–O symmetric stretching vibrations, B_{1g} (395 cm⁻¹ and 515 cm⁻¹) relates to O–Ti–O symmetric bending vibrations, and A_{1g} (519 cm⁻¹) refers to anti-symmetric bending vibrations in anatase lattice [\[20](#page-10-0)]. B_{1g} (395 cm⁻¹) is not observed in TAT A monolayer but present in subsequent multilayers while B_{1g} (515 cm⁻¹) is slightly red shifted in TAT A. There is a blue shift observed in E_{ϱ} (multilayers) which may occur due to non-stoichiometry and phonon confinement effect [\[21](#page-10-0)]. This effect can be related to the decrease in crystallite size in multilayers as explained in XRD results. A red shift in B_{1g} mode is also observed in multilayer films suggesting distortion in $TiO₂$ lattice with the presence of inter-metallic layer [[22\]](#page-10-0). Increase in Raman peak intensity is also observed similar to XRD results confirming improved crystallinity with increase in Ag layer thickness. Also, the increase in total peak area, calculated upon deconvolution, implies the improvement in short-range atomic ordering and local microstructure upon the increase in Ag thickness [\[18](#page-10-0)].

3.5 FTIR

Figure [7](#page-6-0) presents the FTIR analyses of $TiO₂/Ag/TiO₂$ multilayer films done along 400 cm^{-1} to 4000 cm^{-1} (stacked) to detect the trace of $TiO₂/Ag$ and functional bonding related to it. The peak at 483 cm^{-1} is related to bending vibration (Ti-O-Ti) bonds in TiO₂ lattice [[23,](#page-11-0) [24\]](#page-11-0). Weak band at 550 cm^{-1} relates to transverse optical phonon vibration of Ag–O bond [[25\]](#page-11-0). This Ag–O bond is present only in $TiO₂/Ag/$ $TiO₂$ multilayer films but not observed in $TiO₂$ monolayer film (TAT A) which is in accordance with the experiment. Although, no effect was observed in

Fig. 4 FESEM image of TiO₂/Ag/TiO₂ multilayer films with the increase in Ag thickness. (a) TiO₂ monolayer, (b) Ag 2 nm, (c) Ag 4 nm, and (d) Ag 8 nm. Inset image represents the crystallite size estimated by imageJ software

the absorption peak related to the increase in Ag thickness. Ananth et al. (2016) reported the presence of Ag–O metal oxide vibration between 650 and 800 cm^{-1} but it was not observed in our FTIR spectra. The peak at 775 cm^{-1} and 939 cm^{-1} is related to stretching vibrations of Ti–O–Ti and Ti–O bonding, respectively [\[26\]](#page-11-0). The inset of Fig. [7](#page-6-0) presents the same FTIR spectra (1200 and 4000 cm^{-1}) but in high scale for clarity. Peaks at 1433 cm⁻¹, 2860 cm⁻¹, 2960 cm⁻¹ are observed due to symmetric and anti-symmetric vibrations of surfactants present in the samples [[27\]](#page-11-0). Peak at 1640 cm^{-1} represents the bending vibrations of Ti–OH bonding due to chemically absorbed water molecule [[24,](#page-11-0) [28](#page-11-0)]. Gao et al. (2003) have reported the presence of OH group between 3100 and 3600 cm^{-1} while Latha et al. reported similar functional group between 3400 and 3600 cm^{-1} but very weak bonds were observed in our samples in those range [[23,](#page-11-0) [29\]](#page-11-0). The doublet peak around 2000 cm^{-1} (inset) may have originated due to instrumental grating error.

3.6 UV–Visible

Figure [8a](#page-7-0) presents UV–Visible spectra (%t) of TAT multilayer films with variation in Ag thickness. TAT multilayers show the transmittance between 65 and 80% at 550 nm wavelength with TAT D being the most optically transmitting film. But in higher wavelength region (1000–2500 nm), the transmittance changes as a parabolic function with TAT A showing the highest transmittance and TAT D lowest. This phenomenon happens due to the scattering of light at Ag and $TiO₂$ interface, while absorption of light is due to the electronic transition from 3d to 4s orbitals above Fermi level. But this behavior can only be observed in multilayers with mid-layer thickness

Fig. 5 EDX elemental composition of Ti and Ag (wt%) in multilayer films with the increase in Ag thickness. (a) TiO₂ monolayer, (b) Ag 2 nm, (c) Ag 4 nm, and (d) Ag 8 nm

Fig. 6 Raman spectra of $TiO₂/Ag/TiO₂$ multilayer films. Peaks Fig. 7 FTIR spectra of TiO₂/Ag/TiO₂ multilayer films scanned
were fitted with Gaussian profile $\frac{1}{2}$ and $\frac{1}{2}$

below 10 nm where Ag islands are discontinuous. These discontinuous islands limit the mean free path of conducting electrons resulting in the decrease in transmittance [[30\]](#page-11-0). The dielectric function depending upon the size of the Ag islands can be given by [\[31](#page-11-0)]

$$
\varepsilon_R = \varepsilon_m + \frac{C}{R},\tag{3}
$$

from 400 to 4000 cm^{-1} . Inset picture represents the scaled-up image of the spectra (1200 to 4000 cm^{-1}) for clarity

where ε_m is dielectric constant, R is particle radius, and C is constant. This leads to lower transmittance for smaller Ag islands below a critical value. It is observed that lower Ag mid-layer thickness leads to small discontinuous Ag islands forming below the

Fig. 8 (a) UV–Vis spectra of TiO₂/Ag/TiO₂ multilayer films in transmittance mode as function of mid-layer Ag thickness. (b) Tauc plot of multilayer films for optical band gap calculation

 $TiO₂$ layer. Due to the scattering of light from the $TiO₂/Ag$ interface, the average transmittance is low $(< 80\%)$. The multilayer thin films transmittance is strongly affected by the surface plasmonic effect due to high optical absorption. This results in decrease in average transmittance with small island size from lower Ag thickness. With the increase in Ag thickness above critical value, the continuous islands decrease the light scattering resulting in high transmittance [\[32](#page-11-0)].

In visible region, higher transmittance was observed at 700 nm wavelength due to coupling between incident light and surface plasmon of the Ag islands. The decrease in optical transmittance in IR region is observed due to the reflection depending upon surface plasmon absorption. The plasmon frequency (ω_p) can be described from Drude–Lorentz free electron model as

$$
\omega_p = \sqrt{\frac{ne^2}{me_0}},\tag{4}
$$

where *n* is the carrier density and ε_0 is the permittivity in free space [\[33](#page-11-0)].

Figure 8b presents the band gap of TAT multilayer films calculated by tauc plot method ranging between 3.47 and 3.54 eV. The shift in band gap of multilayer films is due to the increase in free charge carriers with Ag layer thickness which leads to Burstein–Moss effect [\[34](#page-11-0)]. In this effect, free charge carriers fill the states near the bottom of conduction band leading to shift in Fermi level resulting in the widening of band gap. The shift in Fermi level is given by

$$
\Delta E = \frac{h^2}{8m_{VC}^*} \left(\frac{3N}{\pi}\right)^{2/3} \tag{5}
$$

So, widening of band gap is calculated by

$$
E_0 = E_G + \Delta E, \tag{6}
$$

where m_{VC}^* is the effective mass of conduction and valence band, N is free electron density, E_G is the minimum energy separation of thin film band [[35\]](#page-11-0).

3.7 Electrical properties

Electrical properties of TAT multilayer largely depend upon the highly conductive metallic midlayer. Sheet resistance, mobility, and carrier concentration are calculated as a factor of Ag layer thickness (Fig. [9\)](#page-8-0). Results from few other authors are important in the discussion for a comparative study. Loka et al. (2017) have reported the decrease in sheet resistance as a factor of -3.07 (slope) with the increase in percentage Ag thickness in $TiO₂/Ag/TiO₂$ multilayer (total thickness) from 5.9 to 15.8% [\[36](#page-11-0)]. Similarly, the carrier concentration and mobility is reported to be increasing as a factor of 1.12 and 0.48, respectively. Dhar et al. (2013) have reported the decrease in sheet resistance to factor of -9.95 with Ag content from 8.3 to 21.7% [[31\]](#page-11-0). The author has also reported the increase in carrier concentration and mobility as a factor of 1.01 and 1.14, respectively. Kim et al. (2015) have reported similar trend for sheet

Fig. 9 (a) Sheet resistance, (b) carrier concentration, and (c) Mobility of TiO₂/Ag/TiO₂ multilayer films compared with references against the increase in Ag layer thickness

resistance, carrier concentration, and mobility as a factor of -0.38 , 0.46, and 0.88, respectively, along 15.8–23.8% Ag content [[10\]](#page-10-0). Jia et al. (2004) have reported similar trend in sheet resistance with a factor of -0.78 with Ag content from 23.1 to 45.9% [[37\]](#page-11-0). The high mid-layer Ag thickness (6–17 nm) between 10-nm-thick $TiO₂$ does not show a sharp drop in sheet resistance rather a very low decreasing rate. Zhu et al. (2016) have reported similar decline in sheet resistance with factor of -0.732 along 11.8–25% Ag thickness [[38\]](#page-11-0). This work reports the decrease in sheet resistance from 15.4 to 5.92 Ω/\square as a factor of -1.47 from percentage Ag thickness of 6.3–14.3%. Carrier concentration varies from 5.94×10^{21} to 1.35×10^{22} cm⁻³, whereas mobility varies from 6.8 to 13.7 cm² V⁻¹ S⁻¹ with a factor of 1.23 and 1.11, respectively.

It is observed that conduction through the dielectric dominates below critical Ag thickness forming the discontinuous islands. With the increase in island size and Ag thickness, decrease in resistivity is observed due to quantum tunneling between the island gaps. The conductivity (σ) in lower Ag island condition is expressed by the equation

$$
\sigma \propto \exp\biggl(-2\beta L - \frac{W}{k_{\beta}T}\biggr),\tag{7}
$$

where W denotes the charging energy of Ag island, L is island separation, β is the tunneling exponent of electron wave functions in the insulator, k is the Boltzmann constant, and T is the temperature. With the increase in Ag thickness, the island forms a uniform layer. In this case, the conductivity in multilayer film is derived by the combination of carrier concentration and mobility of conducting electrons due to surface boundary scattering. The conductivity in this case is given by

$$
\frac{\sigma}{\sigma_0} \propto \frac{3}{4} (1 - p) k \log \frac{1}{k},\tag{8}
$$

where p is the electron distribution function at surface, σ_0 is conductivity of bulk material, $k = t/\lambda_0$ [\[31](#page-11-0), [39\]](#page-11-0). Haacke's figure of merit (FOM) was calculated by the formula:

$$
\varphi_{TC} = \frac{T^{10}}{R_S},\tag{9}
$$

where T is Transmittance and R_S is sheet resistance. TAT D multilayer with Ag thickness 8 nm gives the best figure of merit with $7.9 \times 10^{-3} \Omega^{-1}$ [[40\]](#page-11-0).

4 Conclusion

 $TiO₂/Ag/TiO₂$ multilayer films were deposited using ion-beam sputtering technique with mid-layer Ag thickness as 2 nm, 4 nm, 8 nm, while $TiO₂$ thickness is 60 nm. XRR analysis showed the increase in both roughness and densities of multilayer with the increase in Ag thickness. XRD analyses confirmed the presence of both $TiO₂$ (anatase) and Ag in multilayer films with preferential growth in [101] and [200] planes, respectively. Crystallite size is in multilayers compared to $TiO₂$ monolayer. The crystallite size of both $TiO₂$ and Ag increased while micro-strain and dislocation density decreased with the increase in Ag thickness. A lattice mismatch between quartz substrate and multilayer film is also observed giving rise to sharp peak at lower 2θ value. FESEM images confirm dense, smooth, and crack-free films. Ag islands were visible as dark spots over surface in 4 and 8 nm Ag films. EDX analyses confirmed the presence of $TiO₂$ and Ag in our films. The composition (wt%) of Ag ranged between 2.5 and 10.3% by volume. Raman spectra confirmed $TiO₂$ anatase phase with 6 Raman active modes detected in the films, and their intensities increased with Ag thickness confirming crystalline nature of multilayer films. Blue shift of E_g band in multilayers confirmed phonon confinement effect due to the increase in Ag midlayer thickness. FTIR detected the presence of Ag–O bond at 550 cm^{-1} confirming the presence of Ag in multilayer films. UV–Visible results showed transmittance ranging between 65 and 81% in visible

region (550 nm) and band gap was calculated to be 3.47–3.54 eV for TAT multilayer films. Sheet resistance of TAT multilayers decreased with Ag thickness while carrier concentration and mobility increased with Ag thickness. TAT D with 8 nm Ag thickness showed lowest sheet resistance of 5.92 Ω/\square with carrier concentration 1.35×10^{22} cm⁻³ and mobility 13.72 cm²/Vs. Haacke's Figure of Merit (FOM) was recorded highest for TAT D to be $7.9 \times 10^{-3} \Omega^{-1}$ at 8 nm Ag thickness. A well-matched comparative plot of electrical properties with other previous reported data is presented in the manuscript. Optical and electrical properties have further room of improvement by optimal selection of $TiO₂$ and Ag thickness including other deposition parameters for an optimal TCO. These findings may help in the growth of high-performance TCOs for many future applications.

Author contributions

All authors contributed to the study conception and design. Material preparation, data collection, and analysis were performed by RS. The first draft of the manuscript was written by SKM and all authors commented on previous versions of the manuscript. All authors read and approved the final manuscript.

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Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author on reasonable request.

Declarations

Conflict of interest We wish to confirm that there are no known conflicts of interest associated with this publication and there has been no significant financial support for this work that could have influenced its outcome.

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