



# Effect of vacuum annealing and heating–cooling cycle annealing on the soft magnetic properties at room- and high-temperatures for nanocrystalline FeCoAlSiBCuNb alloy

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## ABSTRACT

Effect of vacuum annealing and heating–cooling cycle annealing on the soft magnetic properties at room- and high-temperatures for  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{-Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy was investigated by the temperature dependence of initial permeability ( $\mu_i$ -T curves). The crystallization behavior and microstructural evolution of alloy with the increase of annealing temperature were studied by means of differential scanning calorimeter (DSC) and X-ray diffraction (XRD). After vacuum annealing at 617 °C for 0.5 h, the  $\mu_i$  can reach 30,300 and maintain more than 1000 until 640 °C because of larger crystalline volume fraction ( $V_{\text{cry}}$ ) of 81.9% and smaller intergranular amorphous layer ( $\Lambda$ ) of 0.930 nm, which is more excellent than that of other Fe–Co–Al–Si–B–Cu–Nb alloys annealed at their optimum temperatures. In addition, the crystallized interval temperature ( $\Delta T_x$ ) of 158 °C, superior to  $\text{Fe}_{73.5}\text{Si}_{13.5}\text{B}_9\text{Nb}_3\text{Cu}_1$  and  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{73.5}\text{Si}_{13.5}\text{B}_9\text{Nb}_3\text{Cu}_1$  alloys ever reported, is conducive to precipitating the bcc  $\alpha$ -(Fe,Co,Si) and  $\text{Fe}_3(\text{Si,Co,Al})$  soft magnetic crystalline phases and improving thermostability of alloy.

## 1 Introduction

Magnetism, one of the most basic physical properties, is ubiquitous among various substances. As indispensable functional materials in various fields, a variety of magnetic materials not only meet the development requirements of traditional industries,

but also play an increasingly important role in high-tech fields such as power, electronics and communications. The soft magnetic material is one of the most typical magnetic materials, having excellent comprehensive soft magnetic properties including high saturation magnetization ( $M_s$ ), high permeability ( $\mu_i$ ), low coercivity ( $H_c$ ) and low loss ( $W$ ) etc. [1], which

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meets the needs of the era of green, energy saving, environmental protection and high efficiency. What's more, typical Fe-based alloys are widely used in communication, computer and other high-tech fields as a result of the high initial permeability at room temperature. However, relatively low Curie temperature of amorphous phase ( $T_c^{\text{am}}$ ) and crystallized interval temperature ( $\Delta T_x$ ) limited the application of the alloys at high temperature [2, 3]. Some studies found that the adulteration of moderate Co increased both the  $\Delta T_x$  and the  $T_c^{\text{am}}$ , which was beneficial to form the single soft magnetic phase and improve the thermostability of the alloys [4–6]. And adding moderate Al to FeCo-based alloy obtained larger  $\Delta T_x$  above 210 °C [7, 8] and higher  $T_c^{\text{am}}$  than that of typical  $\text{Fe}_{73.5}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloys [3]. Therefore, the nanocrystalline  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy after vacuum annealing and heating–cooling cycle annealing was studied in this paper.

We mainly analyzed the changes of  $\mu_i$ -T curves by XRD patterns, Herzer's Model and magnetic coupling, and further explored the relationship between microstructure and soft magnetic properties in order to find optimal magnetic softness at room and high temperatures.

## 2 Experimental procedures

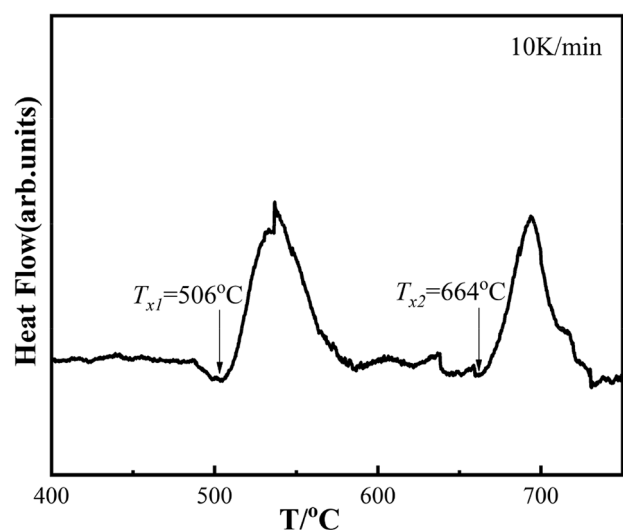
First, the smelted spindle was made into a series of amorphous ribbons with the width of 2 mm and thickness of 20  $\mu\text{m}$  by the single-roller melt-spinning technique, some of which were wound into toroidal cores with the outer diameter of 18 mm and inner diameter of 16 mm for cycle annealing and vacuum annealing, some of which were cut up for the thermal analysis of differential scanning calorimeter (DSC) using NETZSCH STA 409 PC/PG from room temperature to 1000 °C at a heating rate of 10 K/min. Next, the  $\mu_i$ -T curves of samples after different thermal treatment were measured in a furnace with Ar atmosphere protection in order to observe the temperature dependence of initial permeability by an HP4294A impedance analyzer at  $H = 0.4\text{A/m}$  and  $f = 10\text{ kHz}$ . Finally, the characterization of microstructure was obtained by X-ray diffraction (XRD) using *D/max-2500/PC* with  $\text{Cu-K}\alpha$  radiation ( $\lambda = 1.54056\text{ \AA}$ ) to analyze the relations between the microstructure and soft magnetic properties of the

alloys after vacuum annealing at 500–660 °C for 0.5 h (Fig. 1).

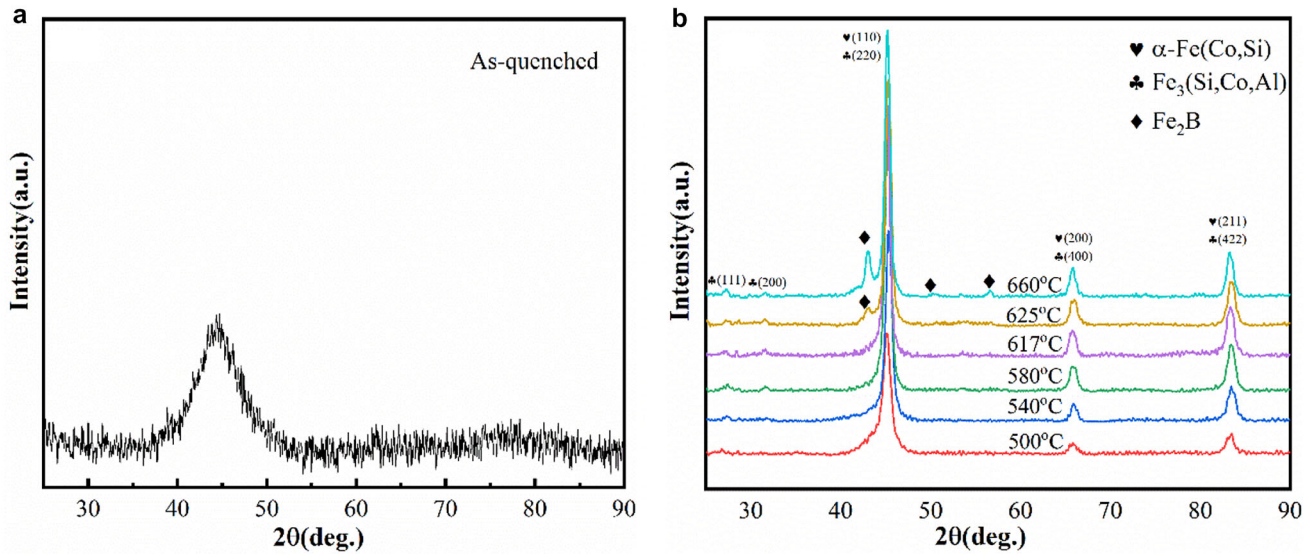
## 3 Results and discussion

The DSC curves of as-quenched  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloys show two separated exothermic peaks related to the enthalpy change and exothermic phase transformations [9, 10], the endothermic peak is corresponding to the melting of ribbons over the temperature range from 480 to 510 °C [11]. The onset primary crystallization temperature  $T_{x1}$  corresponds to the formation of crystalline  $\alpha\text{-Fe}(\text{Co},\text{Si})$  and  $\text{Fe}_3(\text{Si},\text{Co},\text{Al})$  phase shown in Fig. 2b, while the onset temperature of second crystallization  $T_{x2}$  relates to final crystallization of boride  $\text{Fe}_2\text{B}$  [12], the temperature interval  $\Delta T_x$  ( $\Delta T_x = T_{x2} - T_{x1}$ ) of 158 °C is superior to that of  $\text{Fe}_{73.5}\text{Si}_{13.5}\text{B}_9\text{Nb}_3\text{Cu}_1$  [2] and  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{73.5}\text{Si}_{13.5}\text{B}_9\text{Nb}_3\text{Cu}_1$  alloys [13], which is beneficial to obtain the single crystalline phase. It is well known that the size of particle relates to annealing temperature [14], higher annealing temperature leads to larger average grain size and crystalline volume fraction [15–18], which is substantiated in Table 1.

Figure 2 shows the XRD patterns of as-quenched and annealed  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy. It can be seen that the diffraction peak of as-quenched alloy is one broad scattering peak at around  $2\theta = 45^\circ$ ,



**Fig. 1** DSC curve of as-quenched  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloys at a heating rate of 10 K/min



**Fig. 2** XRD patterns of  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy as-quenched (a) and annealed at 500–660 °C for 0.5 h (b)

indicating that the as-quenched sample is complete amorphous state [19]. After vacuum annealing at 500–617 °C, the samples show the characteristic (110), (200) and (211) diffraction peaks of bcc  $\alpha$ -(Fe,Co,Si) crystalline phase and the characteristic (111), (200), (220), (400) and (422) diffraction peaks of  $\text{Fe}_3$ -(Si,Co,Al) crystalline phase, which suggests that the annealed samples are partially crystallized and form the desired two-phase nanocrystalline structure related to the crystallization phase and remaining amorphous phase.

The volume fraction of crystallization phase  $V_{\text{cry}}$  can be estimated by XRD spectra as follows [20]:

$$V_{\text{cry}} = \frac{I_{\text{cry}}}{I_{\text{cry}} + KI_{\text{am}}} \quad (1)$$

where  $I_{\text{cry}}$  and  $I_{\text{am}}$  denote the integral intensities of diffraction peaks of crystalline phase and amorphous phase respectively.  $K$  is a constant determined by experiment. With the increase of annealing

temperature ( $T_a$ ), the  $V_{\text{cry}}$  gradually increases from 77.3% for  $T_a=500$  °C to 81.9% for  $T_a = 617$  °C shown in Table 1. According to the full width at half maximum (FWHM) of the (110) and (220) diffraction peaks, the theoretical Scherrer formula [21–25] is used to calculate the average grain size ( $D$ ):

$$D = \frac{K\lambda}{\beta \cos \theta} \quad (2)$$

where ' $K$ ', ' $\lambda$ ', ' $\beta$ ' and ' $\theta$ ' are the Scherrer constant ( $K$  is nearly 1) [18, 26], wavelength of Cu-K $\alpha$  radiation ( $\lambda = 1.54056$  Å), the full width at half-maximum height (FWHM) and the Bragg diffraction angle, respectively. The results showed that the  $D$  gradually increases from 11.5 to 13.5 nm with the increase of  $T_a$  shown in Table 1. According to the  $V_{\text{cry}}$  and  $D$ , the thickness of amorphous layer ( $A$ ) is calculated as follows [27]:

$$A = D \left[ \left( \frac{1}{V_{\text{cry}}} \right)^{1/3} - 1 \right] \quad (3)$$

**Table 1** The average grain size  $D$ (nm), crystalline volume fraction ( $V_{\text{cry}}$ ) and intergranular amorphous layer ( $A$ ) of  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy annealed at 500–660 °C

$T_a/^\circ\text{C}$	$D/\text{nm}$	$V_{\text{cry}}/\%$	$A/\text{nm}$
500	11.5 ± 0.5	77.3 ± 0.5	1.032 ± 0.08
540	12.5 ± 0.5	78.2 ± 0.3	1.068 ± 0.06
580	13.5 ± 0.5	81.0 ± 0.3	0.983 ± 0.06
617	13.5 ± 0.5	81.9 ± 0.3	0.930 ± 0.06

The calculation results show that the  $A$  decreases from 1.032 nm for  $T_a = 500$  °C to 0.930 nm for  $T_a = 617$  °C, the higher  $V_{\text{cry}}$  and the thinner  $A$  can ensure the penetrating effect of strong exchange coupling between grains [13, 15, 17] and improve the soft magnetic properties of the alloy according to Herando's model [27].

Therefore, The 617 °C-annealed  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy with

larger  $V_{\text{cry}}$  of 81.9% and thinner  $\lambda$  of 0.930 nm exhibits superior soft magnetic properties, which is verified in Fig. 3b.

Figure 3a shows the  $\mu_i$ -T curve of as-quenched  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy. Amorphous structure with a sharp Hopkinson peak was observed near the Curie temperature of the amorphous phase  $T_c^{\text{am}}$  due to the faster decrease of magnetic anisotropy than that of the saturation magnetization when the temperature approaches  $T_c^{\text{am}}$  [28]. The  $\mu_i$  abruptly drops to zero just above  $T_c^{\text{am}}$  due to ferro-paramagnetic transition of the amorphous alloy [13]. The  $\mu_i$  rapidly rises from 0 to 2000 above 500 °C because of the precipitation of bcc  $\alpha$ -(Fe,Co,Si) and  $\text{Fe}_3(\text{Si,Co,Al})$  crystalline phases and it fleetly declines to zero at about 650 °C due to the precipitation of  $\text{Fe}_2\text{B}$  hard magnetic phase, which is consistent with the results of DSC curve.

The  $\mu_i$  at room temperature increases greatly with the increase of the annealing temperature from 8400 for  $T_a = 500$  °C to 30,300 for  $T_a = 617$  °C. When  $T_a$  is 500 °C and 540 °C, the  $\mu_i$  tends to decrease rapidly around  $T_c^{\text{am}}$ , which can be interpreted by Herzer's model [29]:

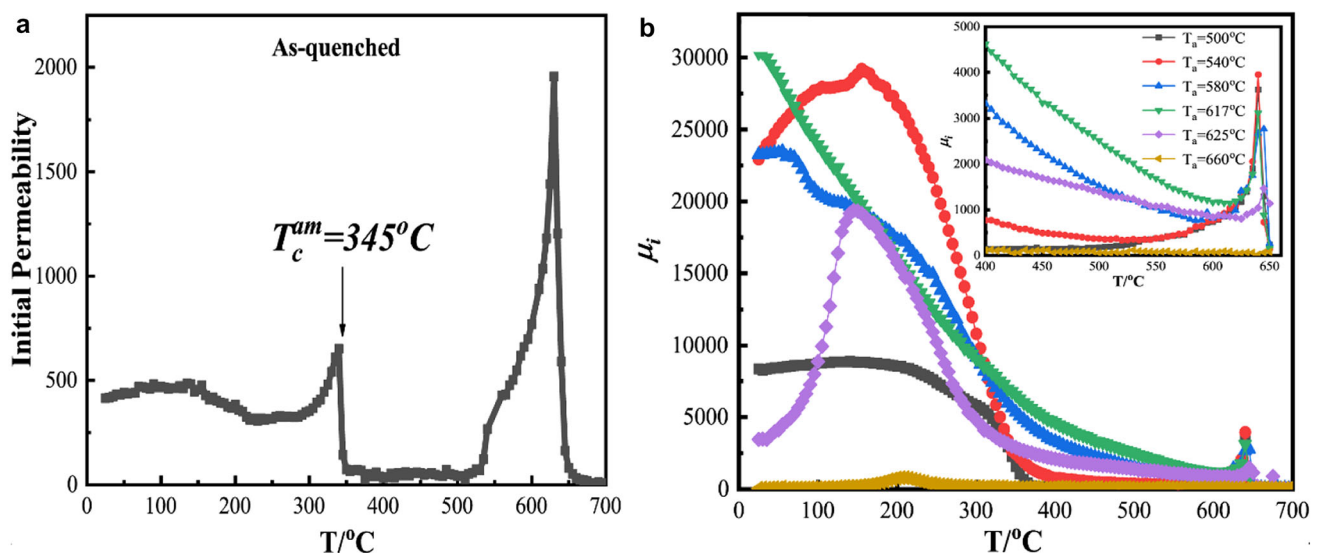
$$\langle K \rangle = \frac{1}{\phi^6} (1 - V_{\text{am}})^4 K_1^4 D^6 \left[ \frac{1}{\sqrt{A_{\text{cr}}}} + \frac{(1 - V_{\text{am}})^{-1/3} - 1}{\sqrt{A_{\text{am}}}} \right]^6 \quad (4)$$

$$\mu_i = P_\mu \frac{J_s^2}{\mu_0 \langle K \rangle} \quad (5)$$

where  $\phi$ ,  $A_{\text{cr}}$ ,  $A_{\text{am}}$ ,  $K_1$ ,  $V_{\text{am}}$ ,  $\langle K \rangle$ ,  $p_\mu$ ,  $\mu_0$  and  $J_s$  represent the spin rotation angle over the exchange-correlation length, exchange stiffness of crystalline phases, exchange stiffness of amorphous phases, magnetocrystalline anisotropy, amorphous volume fraction, effective magnetic anisotropy, a dimensionless prefactor, vacuum magnetic permeability and average saturation magnetic polarization respectively. In the dual-phase nanostructure,  $A_{\text{am}}$  drops to a much smaller value than  $A_{\text{cr}}$  when the temperature approaches to the  $T_c^{\text{am}}$  [30], which makes  $\langle K \rangle$  increase and  $\mu_i$  decrease rapidly. While the increasing rate of  $\langle K \rangle$  is much slower at higher  $V_{\text{cry}}$ , showing the smoother decreasing rate of  $\mu_i$  at higher  $T_a$  [31] such as annealing at 580 °C and 617 °C.

With further increasing annealing temperature from 625 to 660 °C, the initial permeability at room temperature decreases significantly from 3500 down to nearly zero due to the successive precipitating  $\text{Fe}_2\text{B}$  hard magnetic phase from amorphous matrix and the higher magnetocrystalline anisotropy of Fe-borides at room temperature [16]. But the  $\mu_i$  of 625 °C-annealed sample gradually increases to 19,300 around 150 °C, and then decreases significantly with temperature further increasing. The  $\mu_i$  peak may correspond to the minimum value of magnetocrystalline anisotropy for  $\text{Fe}_2\text{B}$  hard magnetic phase [16].

In conclusion, although 625 °C-annealed sample has a small amount of hard magnetic phase, there are few effects on the high-temperature magnetic

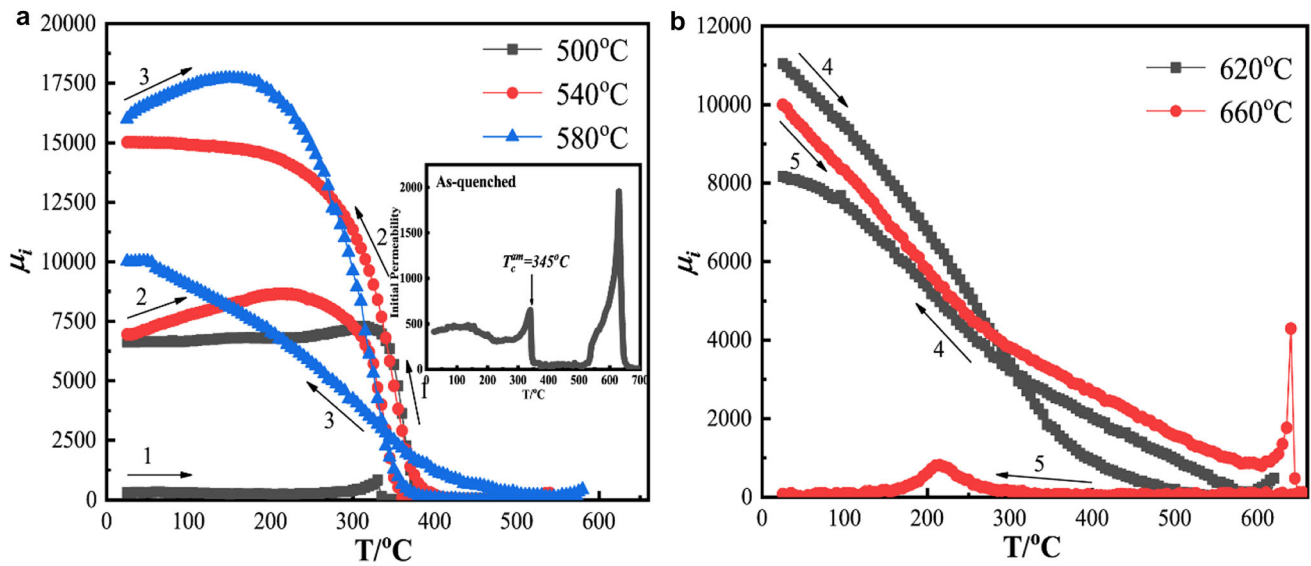


**Fig. 3**  $\mu_i$ -T curves of  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy as-quenched (a) and annealed at 500–660 °C (b)



**Table 2** The initial permeability at room and high temperatures for Fe–Co–Al–Si–B–Cu–Nb alloys annealed at their optimal temperatures and conditions

Composition	Conditions	$\mu_{i25}$	$\mu_{i500}$	$\mu_{i550}$	$\mu_{i600}$
(Fe <sub>0.5</sub> Co <sub>0.5</sub> ) <sub>73.5</sub> Al <sub>0.6</sub> Si <sub>13.5</sub> B <sub>8.4</sub> Cu <sub>1</sub> Nb <sub>3</sub> [8]	Vacuum annealing at 600 °C	1150	1100	1100	1100
(Fe <sub>0.5</sub> Co <sub>0.5</sub> ) <sub>72.7</sub> Al <sub>0.8</sub> Si <sub>17.5</sub> B <sub>5</sub> Cu <sub>1</sub> Nb <sub>3</sub> [17]	Heating–cooling cycle annealing at 550 °C	1600	1500	1500	1500
(Fe <sub>0.8</sub> Co <sub>0.2</sub> ) <sub>72.7</sub> Al <sub>0.8</sub> Si <sub>17.5</sub> B <sub>5</sub> Cu <sub>1</sub> Nb <sub>3</sub> [32]	Vacuum annealing at 640 °C	8100	2000	1500	500
(Fe <sub>0.9</sub> Co <sub>0.1</sub> ) <sub>72.7</sub> Al <sub>0.8</sub> Si <sub>13.5</sub> B <sub>9</sub> Cu <sub>1</sub> Nb <sub>3</sub>	Vacuum annealing at 617 °C	30,300	2500	1700	1200



**Fig. 4**  $\mu_i$ - $T$  curves of (Fe<sub>0.9</sub>Co<sub>0.1</sub>)<sub>72.7</sub>Al<sub>0.8</sub>Si<sub>13.5</sub>B<sub>9</sub>Cu<sub>1</sub>Nb<sub>3</sub> alloy during continuous heating–cooling cycle annealing at 500–660 °C for 15 min (the top temperature of each cycle is 500 °C, 540 °C, 580 °C, 620 °C and 660 °C respectively)

softness shown in partially enlarged inset of Fig. 3. The reason should be attributed to stronger magnetic coupling between grains for nanocrystalline alloys at higher  $T_a$ . What is more, the (Fe<sub>0.9</sub>Co<sub>0.1</sub>)<sub>72.7</sub>Al<sub>0.8</sub>Si<sub>13.5</sub>B<sub>9</sub>Cu<sub>1</sub>Nb<sub>3</sub> alloy has the highest  $\mu_i$  of 30,300 and higher  $\mu_i$  above 1000 until to 640 °C after optimal vacuum annealing at 617 °C, which is superior to other Fe–Co–Al–Si–B–Cu–Nb alloys in Table 2.

Table 2 shows the initial permeability of different Fe–Co–Al–Si–B–Cu–Nb alloys at 25 °C, 500 °C, 550 °C and 600 °C named as  $\mu_{i25}$ ,  $\mu_{i500}$ ,  $\mu_{i550}$  and  $\mu_{i600}$  to compare their magnetic softness at room- and high-temperatures. The results point out that the room-temperature initial permeability of (Fe<sub>0.9</sub>Co<sub>0.1</sub>)<sub>72.7</sub>Al<sub>0.8</sub>Si<sub>13.5</sub>B<sub>9</sub>Cu<sub>1</sub>Nb<sub>3</sub> alloy is distinctly higher than that of other alloys, and the high-temperature initial permeability at 500–600 °C is clearly superior to others except (Fe<sub>0.5</sub>Co<sub>0.5</sub>)<sub>72.7</sub>Al<sub>0.8</sub>Si<sub>17.5</sub>B<sub>5</sub>Cu<sub>1</sub>Nb<sub>3</sub>

alloy at 600 °C. In summary, the (Fe<sub>0.9</sub>Co<sub>0.1</sub>)<sub>72.7</sub>Al<sub>0.8</sub>Si<sub>13.5</sub>B<sub>9</sub>Cu<sub>1</sub>Nb<sub>3</sub> alloy after vacuum annealing at 617 °C has more excellent magnetic softness at room and high temperatures than other Fe–Co–Al–Si–B–Cu–Nb alloys annealed at their optimal temperatures and conditions.

The  $\mu_i$ - $T$  curves of (Fe<sub>0.9</sub>Co<sub>0.1</sub>)<sub>72.7</sub>Al<sub>0.8</sub>Si<sub>13.5</sub>B<sub>9</sub>Cu<sub>1</sub>Nb<sub>3</sub> alloy during continuous heating–cooling cycle annealing at 500–660 °C are showed in Fig. 4. The shape of  $\mu_i$ - $T$  curves is related to phase structure of the alloys. A sharp Hopkinson peak appears at  $T_c^{am}$  for as-quenched alloy shown in enlarged inset of Fig. 4a. The cooling curve of 580 °C and vacuum annealing at same temperature have the same trend, the Hopkinson peak disappears and a higher  $\mu_i$  is observed due to the formation of two-phase nanocrystalline structure [33]. In the fifth cycle, the alloy exhibited improved high-temperature magnetic softness during heating curve, the higher  $\mu_i$

above 1000 can keep up to 600 °C as a result of stronger magnetic coupling, and a sharp Hopkinson peak appears near 640 °C relating to the Curie temperature of crystalline phase. However, the  $\mu_i$  is very low nearly zero shown in cooling curve of 660 °C as a result of precipitating the hard magnetic phase.

## 4 Conclusion

Effect of vacuum annealing and heating–cooling cycle annealing on microstructure and soft magnetic properties of nanocrystalline  $(\text{Fe}_{0.9}\text{Co}_{0.1})_{72.7}\text{-Al}_{0.8}\text{Si}_{13.5}\text{B}_9\text{Cu}_1\text{Nb}_3$  alloy was investigated. With the increase of vacuum annealing temperature from 500 to 617 °C, the  $V_{\text{cry}}$  gradually increases from 77.3 to 81.9% and the  $\lambda$  decreases from 1.032 to 0.930 nm, which is conducive to enhancing the exchange coupling between magnetic particles and improving the soft magnetic properties. So 617 °C-annealed sample shows more excellent magnetic softness at room- and high-temperatures than that of other Fe–Co–Al–Si–B–Cu–Nb alloys ever reported. Furthermore, the variational trends of  $\mu_i$ -T curves after cyclic annealing also exhibited that the excellent high-temperature magnetic softness was obtained by annealing around 620 °C.

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