

Efects of heat treatment on the structural and electrical conductivity of Fe₂O₃-P₂O₅-PbO glasses

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Received: 30 June 2019 / Accepted: 23 September 2019 / Published online: 26 September 2019 © Springer Science+Business Media, LLC, part of Springer Nature 2019

Abstract

Glasses having compositions $xFe₂O₃-(80-x)P₂O₅ - 20PbO$ with $x = 20, 25, 30, 35$ mol% were prepared by traditional glass melt-quenching method. The efects of the heat treatment of glasses on their structural and electrical conductivity were studied by X-ray diffraction (XRD), differential scanning calorimeter, density (ρ) and dc conductivity (σ). The glass transition temperature (T_g) was observed to increase with Fe₂O₃ content. The density of the prepared glass and corresponding glass–ceramics increases with $Fe₂O₃$ content. XRD of as-quenched glass samples confirm the amorphous nature. Triclinic lead iron phosphate PbFe₂(P₂O₇)₂ is the main crystalline ceramic phase formed after heat treatment at 973 K, of 20 and 25 $Fe₂O₃$ mol% with average crystallite size of 74 nm. Hexagonal iron phosphate FePO₄ phases appeared with Fe₂O₃ exceeded 30 mol%. The high-temperature conductivity is well explained by polaronic hopping conduction model. It is observed that in this system non-adiabatic hopping conduction is present. The longitudinal optical phonon frequency (v_0) and density of state $N(E_F)$ of the prepared glass and corresponding glass–ceramics are reasonable for the localized states.

1 Introduction

Low softening temperature, low melt viscosity, high electrical conductivity and relatively high thermal expansion coefficients, makes phosphate glasses the object of deep scientifc investigation and excellent candidates for industrial applications $[1-4]$ $[1-4]$ $[1-4]$. The heavy lead inclusion in phosphate glass matrix can be used as nuclear waste storage media and a scintillators for gamma-ray and neutron detection [\[5\]](#page-6-2). Polyphosphate chains neighboring are linked together by cross bonding between the metal cations (Pb) and the nonbridging oxygen atoms of each PO_4 tetrahedron $[6–8]$ $[6–8]$ $[6–8]$ $[6–8]$. The P–O–P bond between PO_4 tetrahedra is generally much stronger than the cross bonding between chains via the metal cation. But due to their poor mechanical properties and chemical durability its commercial application is limited [\[3](#page-6-4)]. Recently, the elastic properties of lead phosphate glasses were studied using an ultrasonic technique [[9](#page-7-1)]. To improve

the stability of these glasses, it is necessity of incorporate cationic network modifers, such as alkali, alkaline-earth and transition metals to creation of nonbridging oxygen in the matrix [[10](#page-7-2)]. The addition of $Fe₂O₃$ enhance the poor chemical durability of phosphate and is more efective in stabilizing lead phosphate glasses more than other metal oxides [\[3](#page-6-4), [7](#page-7-3)].

Transition metal oxide glasses display semiconducting properties due to the presence of transition metal particles in more than one valence state. In these materials the electrical conduction is due to the transport of electrons. The iron ions in iron phosphate glasses and corresponding glass–ceramics coexists in two different valence states, $Fe³⁺$ has a 3d⁵ electronic configuration and Fe^{2+} has a 3d⁶ electronic configuration with $S = 5/2$ and 2, respectively $[11-13]$ $[11-13]$ $[11-13]$. Iron phosphate glasses are electronically conducting glasses with polaronic conduction mechanisms $[14–16]$ $[14–16]$. Conduction mechanism in alkali-free iron phosphate glasses is assumed to takes place by electrons migrating from low to high valence sites between Fe^{2+} and Fe^{3+} ions (hopping) usually termed small polaron hopping [\[17,](#page-7-8) [18\]](#page-7-9). Superior properties of glass–ceramics are useful for diferent technical wide range applications. Glass–ceramics transformation takes place by controlling the process of nucleation and crystallites growth [\[19,](#page-7-10) [20](#page-7-11)]. According to Musinu et al. [[21\]](#page-7-12), iron ion interacts with lead ions giving rise to a stable structural which could be responsible for

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the good chemical and mechanical properties of iron lead phosphate glasses. The average length of the polyphosphate chains decreases with the addition of iron oxide. The location of $Fe³⁺$ ions in the octahedral sites strength the cross-bonding between the polyphosphate chains which improve the durability of the glasses. The melting point of the glass former P_2O_5 is reduce by adding PbO while adding $Fe₂O₃$ improve the chemical resistance and to decrease the crystallization [\[22](#page-7-13), [23\]](#page-7-14). Also, iron phosphates usually have a low melting temperature [[24](#page-7-15)]. Iron ions act as both network former and modifer depending on the Fe²⁺/Fe³⁺ ratio, high content of Fe₂O₃ acts in the structure former [[18](#page-7-9), [25\]](#page-7-16).

The aim of the present work is to prepare iron lead phosphate glasses and corresponding glass–ceramics with diferent Fe₂O₃ composition by replacing of P_2O_5 with Fe₂O₃. The correlation between the structure and electrical properties have been examined. To investigate the structure and thermal stability of lead–iron phosphate glasses, X-ray difraction (XRD) and diferential scanning calorimetry (DSC) have been used. The dependence of the electrical conductivity upon temperature is discussed.

2 Experimental work

Four glass samples were prepared by melting reagent grade α -Fe₂O₃ (99.9%, Alfa Aesar), P₂O₅ and PbO, (99.9%, Sigma-Aldrich). These samples have the molar formula $xFe₂O₃-(80-x)P₂O₅ - 20PbO with x=20, 25, 30, 35 mol%.$ The glass systems were prepared by mixing specifed weights of nominal composition and heated in a silica crucible at 1200 °C for 15 min. The melt was then poured on a thick copper block and immediately quenched by pressing with another similar copper block to obtained bulk glass of about 2.0 mm in thickness.

XRD analysis was performed using a Siemens D5000 powder diffractometer equipped with Cu Ka radiation (wavelength=0.15406 nm) with a nickel flter at 40 kV and 30 mA. The difractometer was operated within the range of $5 < 20 < 60$ degree using a step time of one second and a step size of 0.05. Thermal analysis was performed in the temperature range from room temperature to 853 K using diferential scanning calorimetry NETZSCH DSC 204, with heating rate 10 K/min. The average density (ρ) of the glasses was measured at room temperature using Archimedes method with Toluene and Hexane as the immersion liquids. At least three samples of each glass were used to determine the density using the following relation;

$$
\rho = \rho_L \left(\frac{W_a}{W_a - W_L} \right) \tag{1}
$$

where W_a is the weight in air; W_L is the weight in liquid and ρ_L is the density of the liquid. The molar volume (V_m) was calculated from the formula;

$$
V_m = \sum \frac{X_i M_i}{\rho} \tag{2}
$$

where, X_i is the molar fraction of the oxide *i* and M_i is its relative molecular mass. Samples were grinded and polished with diferent grit sized emery paper to obtain parallel surfaces of thickness ranging between 1 and 1.3 mm. Silver paste electrodes deposited on both faces of the polished samples. A multimeter type Keithely 760 was used to collect the dc conductivity data. The dc conductivity of the present systems was measured at temperatures between room temperature and 423 K. The I–V characteristic between electrodes was confrmed.

3 Results and discussion

Figure [1](#page-1-0) shows the XRD patterns obtained from the asquenched glass samples with different $Fe₂O₃$ content. These patterns showed a characteristic broad hump without any sharp peaks which confrms the nature of present sample is pure amorphous and non-crystalline. Diferential scanning calorimetry (DSC) technique was used to confrm glassy state and to study the thermal behavior of prepared glass samples. Figure [2](#page-2-0) shows DSC curves of formed glasses with different Fe₂O₃ contents. Glass-transition temperature (T_o) at 748 K for 20% Fe₂O₃ sample followed by exothermic peak corresponds to the crystallization temperature (T_c) at 777 K. With increasing Fe₂O₃ content the T_g peak shifted to the higher temperature values up to 787 K for 35% Fe₂O₃ glass sample followed by exothermic peak (T_c) at 808 K. The

Fig. 1 XRD patterns of $xFe₂O₃-(80-x)P₂O₅-20PbO$ glass samples for different $Fe₂O₃$ content

Fig. 2 DSC curve of $xFe₂O₃-(80-x)P₂O₅-20PbO$ glass samples for different $Fe₂O₃$ content

Fig. 3 XRD patterns of $xFe₂O₃-(80-x)P₂O₅-20PbO$ glass samples for different Fe₂O₃ content after heat treatment at 873 K

larger the difference between the T_c and T_g is, the higher is the thermal stability of the glasses. As shown in Fig. [2](#page-2-0), with increasing $Fe₂O₃$ content the thermal stability decreases.

Lead metaphosphate glass can be completely crystallized by heating the glass in air at 300 $\rm{^{\circ}C}$ (573 K) for several hours, the addition of iron oxide to pure lead phosphate sup-pressed the evidence of crystallization [\[7](#page-7-3)]. $43Fe₂O₃ - 57P₂O₅$ and $10ZnO-30Fe₂O₃ - 60P₂O₅$ glasses have glass-transition temperature at \sim 790 K followed by the two exothermic peaks at \sim 910 and \sim 1070 K that are ascribed to the crystallization temperatures [[11,](#page-7-4) [19](#page-7-10)].

After the heat treatment, the glasses were examined using X-ray diffraction powder diffractometer, Fig. [3.](#page-2-1) For all glasses, there was no evidence of crystallization or phase separation as a result of the heat treatment from 673 to 773 K for 2 h. Following the same heat treatment conditions at 873 K, X-ray analysis indicated that slight crystallization and phase separation had occurred on the surface of the higher $Fe₂O₃$ content samples and phase/s crystallinity increases with $Fe₂O₃$ content. According to Reis and Martinelli $[22]$ $[22]$, Fe₂O₃ addition decrease the lead phosphate glasses crystallization. From Fig. [3,](#page-2-1) 20% Fe₂O₃ glass sample was the most stable sample with temperature, higher percentages show good tendency to glass crystallization. This XRD results confrms thermal stability by DSC. Lead–iron–phosphorus–oxygen network becomes stronger in compositions with an O/P ratio \sim 3.9 and (Fe + Pb)/P ratio \sim 0.90 [\[26\]](#page-7-17).

Figure [4](#page-3-0)a shows XRD patterns of heat treated $25Fe₂O₃ - 55P₂O₅ - 20PbO$ glass samples at different temperatures for 2 h up to 973 K. Crystalline triclinic (or anorthic) lead iron phosphate single phase $PbFe₂(P₂O₇)₂$ (ICDD-JCPDS, No.: 01-087-1904) is formed after heat treatment. By using Rietveld refnement, the calculated lattice parameters of the formed triclinic phase were 0.4801 nm, 0.7088 nm, 0.7877 nm, 89.68°, 87.26° and 73.75°. As shown in Fig. [4](#page-3-0)a, the crystallinity increases with increasing heat treatment temperature. The lattice strain (ε) and crystallite sizes (D) can be estimated from the X-ray line broadening by using Williamson–Hall equation [\[27](#page-7-18)];

$$
\beta \cos \theta = \frac{k\lambda}{D} + 4\varepsilon \sin \theta \tag{3}
$$

where β is the effective broadening, k is the shape factor (0.9), and λ is the wavelength of Cu-k_α radiation (0.15405 nm). Average crystallite size of \sim 74 nm and 0.4 lattice strain % were estimated from the plot of *β*cos*θ* with respect to 4sin*θ*. Figure [4](#page-3-0)b, shows XRD patterns of $35Fe₂O₃ - 45P₂O₅ - 20PbO$ glass samples at the same heat treatment conditions where crystalline monoclinic lead iron phosphate $Pb_3Fe_2(PO_4)_4$ (ICDD-JCPDS No.: 00-053-0175) and hexagonal iron phosphate $FePO₄$ (ICDD-JCPDS No.: 00-050-1635) are formed. The average crystallite size and strain % were 76 nm and 0.2, respectively.

Finally, from the XRD patterns of 25 mol% Fe₂O₃, the only detected crystalline phase is the triclinic lead iron phosphate $Fe₂Pb(P₂O₇)₂$, with iron valence +3. The Fe²⁺ is present in residual amorphous phase [[28](#page-7-19)]. This triclinic structure consists of $(Fe₃O₁₂)^{16–}$ clusters, with $Fe²⁺$ ions in trigonal prism coordination are sandwiched between two $Fe³⁺$ ions in octahedral coordination which are connected by $(P_2O_7)^{4-}$ groups [\[26\]](#page-7-17). With higher Fe₂O₃ content, FePO₄ structure corresponds to ferric orthophosphate is formed. These results indicate that the origin of electrical conductivity of the present samples is due to small polaron hopping (SPH) between iron atoms of diferent valance states, i.e., a step-by-step electron hopping from Fe^{2+} to Fe^{3+} . According to [[3,](#page-6-4) [18,](#page-7-9) [25\]](#page-7-16), phosphate glass structure based on cornersharing PO4 tetrahedra. When $Fe₂O₃$ is added to a phosphate glass, the P–O–P bonds are replaced by P–O–Fe²⁺ and/or

Fig. 4 XRD patterns of **a** $25Fe_2O_3-55P_2O_5-20PbO$ and **b** $35Fe_2O_3-45P_2O_5-20PbO$ glass samples after heat treatment at different temperatures

Fig. 5 Dependence of molar volume and density with $Fe₂O₃$ content of glasses and corresponding glass–ceramics

P–O–Fe³⁺ bonds where two valence state exists, as Fe^{2+} and Fe³⁺ and the Fe²⁺ \Leftrightarrow Fe³⁺ equilibrium depends upon the melting conditions.

Figure [5](#page-3-1) investigates the density and molar volume dependence on the $Fe₂O₃$ content in the glasses and corresponding glass ceramics. In glass and glass–ceramics,

density gradually increases with increasing $Fe₂O₃$ content while molar volume shows opposite behavior. The molar volume of glass samples decreases from 46.88 to 41.68 cm³/ mol with increasing Fe₂O₃. This could be due to the substitution of P by Fe atoms which makes the glass network structure denser. Glass–ceramics samples show the same trend, where the molar volume decreases from 33.71 to 30.84 cm³/mol with increasing Fe₂O₃. PbO up to 50 mol% enter the structure preferentially as network modifer form-ing non-bridging oxygen (NBO) ions (P=O–Pb²⁺O⁻) [\[29](#page-7-20)]. With increasing $Fe₂O₃$ content, the glasses and corresponding glass–ceramic structure becomes less open, allowing for the probable formation of decreasing number of nonbridging oxygen ions (NBOs). This may be due to the largest ionic radii and bond strength of $Fe₂O₃$ as compared to glass former P_2O_5 . As the Fe₂O₃ content increased the glass matrix becomes more compaction.

The logarithmic dc conductivity, σ , as function of reciprocal temperature of the present glasses and corresponding glass ceramics are shown in Fig. [6a](#page-4-0), b. It is observed from the fgure a linear temperature dependence up to a certain temperature ($\theta_{\rm D}/2$, $\theta_{\rm D}$ being the Debye temperature). Such behavior arises from the hopping of electrons or polarons between mixed valance states [\[4](#page-6-1), [8,](#page-7-0) [18](#page-7-9)] as confrmed from XRD results. So the experimental conductivity data above

Fig. 6 Conductivity as a function of temperature for different Fe₂O₃ content of **a** glasses, **b** heat treated glasses (glass ceramics) at 600 °C/2 h

 $\theta_{D}/2$ were fitted with the SPH model proposed by Austin and Mott [[30,](#page-7-21) [31](#page-7-22)]. The high temperature activation energy was computed from the slope of each curve in the highest range of the temperature measured. The dc conductivity (σ) was calculated from;

$$
\sigma = \frac{Id}{VA} \tag{4}
$$

where I is the series current at temperature T, V is the applied dc voltage, d is the specimen thickness and A is the electrode area. The experimental conductivity data in such a situation is well described by activation energy for conduction given by Mott formula [[31\]](#page-7-22);

$$
\sigma = \sigma_o \exp\left(-\frac{W}{kT}\right) \tag{5}
$$

and

$$
\sigma_o = \begin{cases} v_o Ne^2R^2C(1-C) \exp(-2\alpha R) & \text{for non-adiabatic} \\ v_oNe^2R^2C(1-C) & \text{for adiabatic} \end{cases}
$$
(6)

where σ_0 is the pre-exponent, W is the activation energy for the dc conductivity, k is the Boltzmann constant, T is the absolute temperature, R is interionic distance α , is the tunneling factor (the ratio of wave function decay), c the fraction of reduced transition metal ion.

The electrical conductivity of both glass and heat-treated glasses (glass–ceramics) increases with the increase of $Fe₂O₃$ content (Fig. [6](#page-4-0)a, b). A pronounced decrease in the electrical conductivity of the formed glasses is obtained after heat treatment at 873 K for 2 h, Fig. [6b](#page-4-0). The activation energies were determined from the linear part of the high-temperature region of the plots, Fig. [6](#page-4-0)a, b, that follows the Mott equation. The activation energy values for glass samples between 0.575 and 0.782 eV are very closed to that of glass–ceramics (0.535 and 0.777 eV). Also, Fig. [6](#page-4-0) shows that the natural logarithm of conductivity (ln σ) at 373 K increases with $Fe₂O₃$ content while the activation energy (*W*) shows a reversed behavior in all glass and glass–ceramics samples. The addition of iron oxide to the phosphate network leads to breakdown of the P=O bonds and the P–O–P bonds are more chemically durable $P-O-Fe^{2+}$ and/or $P-O-Fe^{3+}$ bonds [\[3](#page-6-4), [15](#page-7-23)]. The conductivity of glasses and corresponding glass–ceramics could be due the electron hopping from iron ions Fe^{+2} and Fe^{+3} .

Figure [7](#page-4-1) presents conductivity variation with $Fe₂O₃$ content at a fixed temperature of glasses and corresponding glass–ceramics. It is interesting to note that glasses show high conductivity compared to glass ceramics. This can be

Fig. 7 Activation energy and conductivity dependence with $Fe₂O₃$ content of glasses and corresponding glass–ceramics

attributed to the improved crystallinity leading to reduction in the traps or size of the crystallites in the samples. The electrical conductivity is afected by grain boundary scattering. This grain boundaries act as traps which capture electrons and lead to the formation of a potential barrier. Building up of a negative charge and a potential barrier at the grain boundaries also causes scattering of conduction electrons, which decreases the conductivity. Based on the grain size efects of electrical conductivity, when the size of the crystalline grains is small, the grain boundary scattering increases signifcantly and can increase the conductivity. The electronic conductivity in lead iron phosphate glasses depends not only on the Fe^{2+}/Fe_{tot} ratio, but also on the formation of the conduction pathways along and/or through the crystalline grains [\[32–](#page-7-24)[34\]](#page-7-25).

It is reported [\[33](#page-7-26)] that dc conductivity in iron phosphate glasses is electronic and strongly depends on the interionic distance (R) according to the relation;

$$
R = \left(\frac{4\pi N}{3}\right)^{-\frac{1}{3}}\tag{7}
$$

According to summarized values of R in Tables [1](#page-5-0) and [2,](#page-5-1) R decreases with increasing $Fe₂O₃$ content. As R decreases it is reasonable to fnd an increase in dc conductivity.

The values of the polaron radii r_p calculated from [[35\]](#page-7-27);

$$
r_p = \left(\frac{\pi}{6}\right)^{\frac{1}{3}} \frac{R}{2} \tag{8}
$$

The density of states for thermally activated electron hopping near the Fermi level from basic principles is given as;

$$
N(E_F) = \frac{3}{4\pi R^3 W} \tag{9}
$$

The dc conductivity of these glasses is discussed on the basis of a polaron hopping conduction mechanism where the activation energy *W* is given by [\[30](#page-7-21), [35](#page-7-27)].

$$
W = \begin{cases} W_H + \frac{W_D}{2} & \text{for } T > \theta_D/2\\ W_D & \text{for } T < \theta_D/4 \end{cases}
$$
 (10)

where W_H is the polaron hopping energy and is equal to $W_p/2$, W_p is the polaron binding energy and W_d is the disorder energy arising due to the energy diference of the neighboring sites at low temperature. $\theta_{\rm D}$ is the Debye temperature and is defined by $h\nu_0 = k\theta_D$, ν_0 is the optical phonon frequency, *h* being Planck's constant and k is the Boltzmann constant. To determine v_0 for all different compositions, the Debye temperature θ_D was estimated by T > $\theta_D/2$ (Eq. [10](#page-5-2)). v_o of the present glasses and corresponding glass ceramics where obtained to be in the range of 666–680 K and 703–711 K, respectively. The measured and calculated values of physical properties like content of iron ions (N), iron interionic distance (R), activation energy (*W*), polaron radii (r_n) , Debye temperature (θ_n) , longitudinal optical phonon frequency (ν_0) and density of state N(E_F) of the prepared glass and glass–ceramics are summarized in Tables [1](#page-5-0) and [2.](#page-5-1)

To explore the of hopping conduction nature (adiabatic or non-adiabatic), a plot of ln σ against activation energy W at a fxed experimental temperature (373 K) for glasses and corresponding glass–ceramics of diferent compositions has to be plotted, Fig. [8.](#page-6-5) This plot shows a straight-line nature, which indicate that the following dc conductivity equation is valid.

$$
\sigma = c(1 - c)\frac{e^2}{RkT} \exp(-2\alpha R) \exp\left(-\frac{W}{kT}\right) = \sigma_o \exp\left(-\frac{W}{kT}\right)
$$
\n(11)

From the slope, 1/kT, if the value of the temperature is nearly equal to the observed temperature, then the hopping conduction is adiabatic in nature and it is mainly controlled by the activation energy. If the estimated and observed

Table 1 Physical par different $Fe₂O₃$ conte samples

ceramics samples

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Fig. 8 Conductivity variation with activation energy at fxed temperature (373 K) of glasses and corresponding glass–ceramics

Fig. 9 Pre-exponential factor dependence with the $Fe₂O₃$ content of glasses and corresponding glass–ceramics

temperature values are very diferent; then this will suggest that the nature of the hopping conduction is non-adiabatic [\[36](#page-7-28)]. Figure [8](#page-6-5) shows the relation between ln *σ* and activation energy W for glass and corresponding glass ceramics at fxed temperature (373 K). The estimated temperature calculated from the slope of the plot is difer than the experimental temperature (501 and 667 K) confrming that, the conduction mechanism in the present samples is due to non-adiabatic SPH of electrons [\[9](#page-7-1)]. Figure [9](#page-6-6) presents the effect of Fe₂O₃ content on the pre-exponential factors (σ_o) obtained from the least squares straight line fits of the data σ_{ρ} . The figure indicates a decrease in σ ^{*o*} with Fe₂O₃ content. From the above results, we conclude that again the conduction mechanism in the present samples is due to non-adiabatic hopping of the polarons.

From Tables 1 and 2 , it can be observed that with increasing $Fe₂O₃$ content, the average interionic distance (R) decreases. Consequently, the addition of $Fe₂O₃$ leads to compaction of glass network which confrms density and molar volume results. The activation energy decreases with addition of $Fe₂O₃$. Also, as the average interionic distance (R) decrease, the activation energy decreases.

4 Conclusion

- From the XRD and DSC results all the samples were fully amorphous in nature, no crystalline phases were detected.
- The activation energy increases as the distance between iron ions, R, increases.
- In glasses and corresponding glass–ceramics, dc conductivity decreases with heat-treatment temperature.
- In glasses and corresponding glass–ceramics, the dc conductivity was found to increase as the iron content increases while the activation energy decreases.
- The glass transition temperatures were found to increase with the $Fe₂O₃$ content.
- The conduction is attributed to non-adiabatic hopping of small polaron (SPH).
- The parameters such as polaron radius and average separation between iron ions obtained from the fts of the experimental data to this model appear reasonable and are consistent with glasses and glass ceramics.

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