

# **Synthesis of photoluminescent β-Ga<sub>2</sub>O<sub>3</sub> nanostructures using electrospinning method, and control of length‑diameter ratio by calcination heating rates**

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#### **Abstract**

Nanofiber precursors of PVP-Ga(NO<sub>3</sub>)<sub>3</sub> were synthesized through the electrospinning technique, and monoclinic  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> patterns were later obtained through the calcination method. During the annealing process, the nanofbers' pores decreased, and their lengths were uniform up to several micrometers due to the surface-to-core extension that comprises the crystallization through Ostwald ripening process. Synthesis on the structure and morphology of materials were investigated using scanning and transmission electron microscopy equipped with an energy dispersive spectrometer, X-ray difraction, Raman and Fourier-transform infrared (FTIR) spectroscopies. The  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> optical properties disclosed very broad and intense photoluminescence emission spectrum in the blue region of the wavelength, whose driving force was the presence of oxygen vacancies in the structures. Two types of  $Ga^{3+}$  ions  $(GaO_6)$  octahedral and  $GaO_4$  tetrahedral chains) were demonstrated to come from different vibrations of Ga–O bonds in the Raman and FTIR spectra. And  $Ga<sup>3+</sup>$ –CO adducts formed on coordinatively  $Ga^{3+}$  ion located at edges and corners of β-Ga<sub>2</sub>O<sub>3</sub> crystallites. Thus, successful results of this work included the control of length-diameter ratio by calcination heating rates, as well as the broad blue emission band, representing a strong potential of β-Ga<sub>2</sub>O<sub>3</sub> materials in optoelectronic applications.

## **1 Introduction**

Electrospinning is considered a method which is versatile for generating long fbers, ranging from dozens of nanometers to micrometers. This technique enables to manipulate the nanofber diameter, composition, and tune their high surface area to volume ratio and their porosity [\[1\]](#page-5-0). Through it, the materials growth has been described to be driven by aggregation and continuous growth of individual crystallites. Examples of these materials can be found as inorganic synthesis systems in solutions, with polymers as surface protection agents and the enhancement of aggregation and intern-particle interaction [[2](#page-5-1)]. Regarding the synthesis of one-dimensional (1 D) materials with diferent compositions, the electrospinning is simpler, convenient and costefective than other methods, like carbothermal reduction, metal-catalyzed vapor–liquid–solid, hydrothermal, and template-based methods [[1,](#page-5-0) [3](#page-5-2), [4\]](#page-5-3). Polymer, polymer/inorganic, composites, and inorganic nanofber are related to be the usual materials developed through the electrospinning technique [[5](#page-5-4)].

Concerning the light emitters, waveguides and light detectors, materials as nanorods, nanowires and nanofbers are often desired [\[4\]](#page-5-3). Due to unique electronic, optoelectronic and mechanical properties, their 1 D nanomaterials have attracted great attention. One example is the  $Ga<sub>2</sub>O<sub>3</sub>$  material, which has wide band gap, with a value of about 4.9 eV. This transparent oxide semiconductor, with a monoclinic phase in the visible spectrum to ultraviolet, have innumerous applications, specially within the optoelec-tronic devices, field effect transistors and gas sensors [\[6,](#page-5-5) [7](#page-5-6)]. The presence of both cation and anion vacancy sites are the responsible for the unique PL properties [\[8](#page-5-7)], as well as the 1 D, mesoporous and hierarquical structures with large surface area to volume ratio [[9,](#page-5-8) [10\]](#page-5-9).

The  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> monoclinic structure is the most thermodynamically stable phase, having other polymorphic phases (α-, γ-, δ- and ε-) ultimately transformed into β- through the calcination processes [[11](#page-5-10)]. Several β-Ga<sub>2</sub>O<sub>3</sub> physical

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properties are originated from the native point defects. The existence of shallow donor states is led by either oxygen vacancies or large number of dopants, providing the electrical conductivity source of the n-type β- $Ga_2O_3$ . Its large surface to volume ratio also provides interesting physical behaviors due to the large surface that interacts with the surrounding environment and promotes the desired surfacestates. For the purpose of obtaining superior physical properties than the ones from bulk materials arisen from the quantum effects, nanostructured  $Ga<sub>2</sub>O<sub>3</sub>$  are usually required [\[12\]](#page-6-0).

It is reported that the electrospinning technology provides excellent processability of polymeric materials [[13\]](#page-6-1). Within the photoluminescence issue, however, few researches describe the studies of the β- $Ga<sub>2</sub>O<sub>3</sub>$  nanofibers applying the electrospinning technique: strong emissions could be appeared with large band widths from 400 to 750 nm, depending on the defects created into the matrix during the temperature annealing processes [[3,](#page-5-2) [14\]](#page-6-2). Besides the β-Ga<sub>2</sub>O<sub>3</sub> nanofibers can exhibit smaller fiber diameter with decreasing solvent, e.g. ethanol, which is attributed to the sufficient time of electrostatic force to pull and stretch the polymer strand before reaching the fber collector. It is also known that the dielectric constant of the solution leads to higher Coulombic repulsion force, which decreases the diameter of formed fber via increased stretching forced exerted on the charged jet segment [\[7\]](#page-5-6). All that fndings could facilitate the exploration of the desired light-emitting materials. Therefore, the development of the related materials is the main goal of this work, aiming their intense PL for achieving targeted optical properties, once the nanoscale gallium oxide is related to have considered conduction and optical properties for being applied in optoelectronic devices.

#### **2 Materials and methods**

The nanofibers of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> materials were developed through the electrospinning followed by calcinations processes [[3](#page-5-2), [4](#page-5-3)]. As-prepared solutions were carried out by dissolving 15 wt% poly(vinyl pyrrolidone) (PVP, Sigma-Aldrich,  $M_w \approx 1,300,000$  in the mixed distilled water and ethanol with the weight ratio of 1:1, keeping the vigorous stirring for 3 h at room temperature. Into the solution, 0.1 g Ga( $NO_3$ )<sub>3</sub>·xH<sub>2</sub>0 (Sigma-Aldrich) was added, keeping the new solution in the vigorous stirring for 10 h, achieving the desired viscosity for the electrospinning technique.

The solution was inject into a plastic syringe equipped with stainless needle (18 inch), which was connected with a peristaltic pump to insure the fow rate of the solution at 0.035 mL min−1. The connected stainless needle was linked to a high voltage–power supply, and a plate of aluminum foil was placed with a distance of 15 cm from the needle tip. A high electric voltage of 20 kV was applied between the stainless needle and the collector, and the precursors were obtained on the aluminum foil as dense mats. The asprepared nanofbers were dried at 70 °C for 5 h and calcined at 900 °C in air for 1 h in order to obtain the β-Ga<sub>2</sub>O<sub>3</sub> fine nanofbers. The temperature heating rate was varied (1, 5 and 10 °C/min) aiming fne nanofbers.

The feld emission gun scanning electron microscope (FEG-SEM; JEOL, model JSM-7500F) equipped with an energy dispersive spectrometer (EDS) was used to analyze the morphology of the produced structures, and the transmission electron microscope (TEM; PHILIPS, model CM200) allowed us to correlate the morphology and structural characteristics of β-Ga<sub>2</sub>O<sub>3</sub> nanofibers.

The crystallinity and phase of the synthesized materials were analyzed by the X-ray difraction (XRD; RIGAKU RINT-2000) using Cu *Kα* radiation at a scan rate of 5 °C/ min, with  $2 \theta$  from  $7$  to  $90^{\circ}$  (D/teX Ultra; PSD).

The micro-Raman spectrum was measured using a 632.8 nm laser source (Horiba Jobin–Yvon, model Lab RAM HR) equipped with Raman microscopy (Horiba Jobin–Yvon´s confocal Raman microscope). The Fouriertransform infrared (FTIR) spectra were recorded in the range from 4000 to 400  $cm^{-1}$  (BRUCKER, model VERTEX 70).

Photoluminescence properties were studied using the excitation wavelength of 325 nm (Horiba Jobin–Yvon, model Fluorolog-3 FL3-122). Emission spectra were recorded in the range of 370–640 nm.

## **3 Results and discussion**

 $PVP/Ga(NO<sub>3</sub>)<sub>3</sub>$  nanofibers were obtained using the electrospinning technique. Burning organic components by annealing treatment in air, the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers could be formed. Due to the fact that the PVP is decomposed completely at approximately 450 °C, and a specifc phase of the fnal product was desired, the annealing temperature at 900 °C was applied. The FEG-SEM images (Fig. [1a](#page-2-0)–c) depicted diferent nanofbers' diameters, which used diferent heating rates, realizing that thinner diameter were formed at lower rate value [[3,](#page-5-2) [13\]](#page-6-1).

As it is shown in Fig. [1](#page-2-0)a–c, pores appeared along the nanofbers. The pores decreased at slower heating rate, and the length was uniform up to several micrometers due to the surface-to-core extension [\[10](#page-5-9)], which is so-called reversed crystal growth mechanism that comprises the crystallization from the surface to core through Ostwald ripening process. The surface-to-core extension is related to participate of a large number of inorganic compounds crystallization, such as zeolites, perovskites, metal, and metal oxides [\[2](#page-5-1)]. A scheme of the reversed crystal growth is illustrated in Fig. [1d](#page-2-0).



<span id="page-2-0"></span>**Fig. 1** FEG-SEM images of  $β$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers synthesized at 1 °C (a), 5  $\degree$ C (b), 10  $\degree$ C (c) per minute, and later calcined at 900  $\degree$ C; reversed crystal growth (d); FEG-SEM image of the precursor

annealed at 70 °C (**e**); EDS (**f**); TEM (**g**); SAED (**h**); and HRTEM that illustrates the atomic structure in (−202) crystal plane (**i**) of the gallium oxide produced at 1 °C



<span id="page-3-0"></span>**Fig. 2** Standard (JCPS card no. 87-1901) and experimental XRDs of  $β$ -Ga<sub>2</sub>O<sub>3</sub>, as well as the as-prepared nanofibers XRD

Comparing the nanofibers before and after being annealed at the higher temperature, it is observed smaller length-diameter ratio for the nanofibers formed at  $70^{\circ}$ C (Fig. [2e](#page-3-0)), with the PVP covering the  $Ga(NO<sub>3</sub>)<sub>3</sub>$ .xH<sub>2</sub>O material. At higher annealing temperature as 900 °C, the fber diameters decreased, changes took place on the surfaces due to the evaporation of ethanol, and the degradation of PVP and decomposition of  $Ga(NO<sub>3</sub>)<sub>3</sub>$  occurred, as previously described in the literature [[15\]](#page-6-3). FEG-SEM images of β-Ga<sub>2</sub>O<sub>3</sub> calcined at 1 °C, 5 °C and 10 °C per min are depicted in Fig. [1](#page-2-0)a–c, respectively, showing that they morphologically appeared as nanofbers, with large length-diameter ratio.

The material composition was characterized using the EDS technique, exhibiting the Ga and O peaks from the materials in analysis (Fig. [1f](#page-2-0)). One dimensionally β- $Ga<sub>2</sub>O<sub>3</sub>$ nanofbers were formed, as demonstrated. Both low-magnifcation TEM (Fig. [1g](#page-2-0)) and FEG-SEM images demonstrated diameters of approximately 60 nm for thinner nanofbers, which were synthesized at slower heating rate. The nanofiber nature of β-Ga<sub>2</sub>O<sub>3</sub> was investigated using the SAED image (Fig. [1h](#page-2-0)). High-resolution image (HRTEM, Fig. [1i](#page-2-0)) demonstrated well-resolved lattice fringes, with a distance between each other as 2.8 Å, corresponding to the interplanar spacing value of (−202) atomic plane of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> [\[16\]](#page-6-4), in agreement with the literature (JCPS no. 87-1901) and the experimental results, appearing as one of the stronger XRD peak intensities.

After the thermal treatment, the crystallinity and phase of the materials synthesized were analyzed using the XRD equipment, which demonstrated that the nanofbers were monoclinic structures [[17](#page-6-5)]. The difraction peaks of the  $β$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers were obtained and fully indexed in the monoclinic structure (JCPS no. 87-1901), as depicted in Fig. [2,](#page-3-0) showing that the cell parameters corresponded to a = 12.2140 Å, b = 3.0371 Å, c = 5.79810 Å, and  $\beta$  = 103.83 Å, and the space group corresponded to C2/m. The Miller indices (hkl) of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> were shown in the XRD plots. Other impurity phases were not detected [[14](#page-6-2)]. Before the thermal treatment at high temperature, the precursor appeared in the XRD as an amorphous phase. According to L. Wang (2009), one of the difraction peaks appeared with broad width at around  $2\theta = 22^{\circ}$  corresponded to the semicrytalline PVP [[1](#page-5-0)], which is depicted in Fig. [2](#page-3-0) of this work.

The PL spectrum of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> (Fig. [3\)](#page-3-1) showed an emission band located at a maximum value of approximately 420 nm under an excitation wavelength of 325 nm at room temperature. This blue band emission has been described to come from the recombination of an electron on a donor formed by the oxygen vacancies  $(V_0^x)$  and a hole on an acceptor formed by the gallium vacancies  $(V_{Ga}^{''''})$  or by galliumoxygen vacancy pairs  $(V_o, V_{Ga}''')'$ . The holes trapped on the acceptors recombined radiatively, emitting photons around the blue region of the wavelength. Usually, high temperatures of calcinations originate emissions in the blue range of the wavelength due to the presence of oxygen vacancies, which are easily formed because of higher surface to volume ratio. The asymmetry of the emission band can also be explained by the electron–phonon interactions [\[5](#page-5-4), [18](#page-6-6)].

Unlike it was not found here in this work, another sort of emission could be present in matrixes of  $Ga<sub>2</sub>O<sub>3</sub>$ : recent researches have described a red light emission that came from the nitrogen absorbed into the host lattice. In this case,  $Ga_2O_3$  particles synthesized in air were doped with N and post-annealed in ammonia atmosphere with diferent nitrogen doping concentrations, whose photoluminescence spectra showed an emission band in the red range of the wavelength. The particles prepared without nitrogen



<span id="page-3-1"></span>**Fig. 3** PL emission spectra of  $β$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers at room temperature using an excitation wavelength of 325 nm

doping presented no red light emission [[18](#page-6-6)]. The red light has also appeared after  $Ga<sub>2</sub>O<sub>3</sub>$  samples were annealed in oxygen ambient environment at high temperatures [[19\]](#page-6-7). N–N bond lengths and N–N stretching frequencies of N-complexes higher activation is also found through the enzymatic reduction of ammonia, another possible process that applies enzyme nitrogenase to fxes dinitrogen into the matrix using specialized catalysts, high pressure and high temperature  $[20, 21]$  $[20, 21]$  $[20, 21]$  $[20, 21]$ . It can be inferred that intrinsic emission band comes from the recombination of an electron trapped on a donor due to the oxygen vacancy and a hole trapped on an acceptor due to the nitrogen doping. The electron trapped on the donor, induced by the nitrogen doping, recombine radiatively, emitting photons around the red region of the wavelength [[18,](#page-6-6) [19\]](#page-6-7). A representative scheme of the PL mechanism for blue, UV, green and red emissions is depicted in Fig. [4](#page-4-0). Through the energy band diagram, the origin of possible emissions is shown.

Generated by the recombination of an electron trapped on a donor due to  $V_0^x$  and a hole trapped on an acceptor due to N impurity from the atmosphere, a red emission is observed. The N impurity enables a creation of a deep acceptor level inside the band gap. This level lies near the middle of the band gap [\[22](#page-6-10)]. Therefore, the main recombination channel occurs via defect level inside the band gap. Thus, created electrons, composed by holes, are frstly relaxed to the band edges, being trapped by the defect levels inside the band gap, and the carriers are recombined, producing a photon with the energy  $h\nu_1$  that is inferred to the red emission [[23](#page-6-11)]. In case of blue emission, this band can be originated by  $V_{Ga}^{\prime\prime\prime}$ ,  $V_{\rm o}^{\rm x}$ , and  $(V_0, V_{Ga}''')'$  in  $Ga_2O_3$ . The blue emission is described to come from the recombination of a donor–acceptor pair (DAP) through a tunneling process. And the UV emission is reported to be originated by the recombination of selftrapped excitons, which are created when an electron at the donor level (formed by  $V_0^x$ ) is captured by a hole at the acceptor level [formed by  $V_{Ga}^{\prime\prime\prime}$  or  $(V_o, V_{Ga}^{\prime\prime\prime})'$ ] [[22](#page-6-10), [24](#page-6-12)].



<span id="page-4-0"></span>**Fig. 4** Scheme of PL for  $β$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers: blue, UV, green and red emissions (Color figure online)

Besides the blue and the green emissions generation can mainly be associated with  $V_{Ga}^{\prime\prime\prime}$ ,  $(V_o, V_{Ga}^{\prime\prime\prime})'$ , and oxygen interstitials. The green PL emission can come from the neutral oxygen interstitial defects, and from isolated gallium vacancies (octahedral site with −2 charge state or tetrahedral site with  $-1$  charge state) [[25\]](#page-6-13).

Moreover, researches have related diferent emissions of the  $Ga<sub>2</sub>O<sub>3</sub>$  nanostructures due to different regions observed in the materials: a crystalline  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> core and an amorphous shell. The blue and yellow emissions are originated from the crystalline  $β$ -Ga<sub>2</sub>O<sub>3</sub> core. In regard to the  $V_0^x$ , the UV and blue emission intensities are dependent on them: the intensity of UV-blue emission decreases as the number of  $V_0^x$ decreases. Reaction time also affects the  $V_0^x$ , consequently, the emission intensity: as long as the reaction time increases, there is an improvement of the crystallinity, which facilitates the reduction of the  $V_0^x$  concentration, decreasing the emis-sion intensity [\[22](#page-6-10), [24](#page-6-12)].

The detailed structural characterizations were studied through the analysis of micro-Raman spectrum, as depicted in Fig. [5.](#page-4-1) The Raman peaks were in a good agreement with the literature. According to the researches, a large number of defects can be detected through red-shifted peaks. The monoclinic  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is related to belong to the space group of  $C_{2h}$  where 15 Raman modes are expected to be observed in the material vibration spectrum. The Raman peaks which appeared in this work were indexed in the calculated Raman active modes as reported in the literature (Local Density approximation)  $[26]$  $[26]$ . In this work, the experimental peaks at 113, 145, 169, and 200  $\text{cm}^{-1}$  corresponded to low frequency modes, which come from the libration and translation of tetrahedral-octahedral chains. Mid-frequency modes that are located in the range of approximately 500–300 cm−1 correspond to the deformation of  $GaO<sub>6</sub>$  octahedral modes (here, the peaks were located at 320, 347, 416, and 475  $cm^{-1}$ ).



<span id="page-4-1"></span>**Fig. 5** Micro-Raman spectrum of β- $Ga_2O_3$  nanofibers



<span id="page-5-11"></span>**Fig. 6** FTIR spectra of  $β$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers

And high-frequency modes that are located in the range of approximately 770–500  $cm^{-1}$  correspond to the stretching and bending  $GaO<sub>4</sub>$  tetrahedral modes [\[27](#page-6-15), [28](#page-6-16)]. In this work, they were located at 629, 653, and 766  $cm^{-1}$ .

Confirming the formation of crystalline  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers and identifying any adsorbed elements on their surface, FTIR was measured from 4000 to 400 cm−1, as depicted in Fig. [6.](#page-5-11) Intensities at 440 cm<sup>-1</sup> and 670 cm<sup>-1</sup> are inferred to be regions correspondent to Ga–O stretching vibrations from  $GaO<sub>6</sub>$  octahedra and O–Ga–O bending vibrations from GaO<sub>4</sub> tetrahedra, respectively. The region at 773 cm<sup>-1</sup> corresponded to the stretching and bending vibrations of GaO<sub>4</sub> tetrahedra [\[29](#page-6-17)[–31](#page-6-18)]. The peak at high-energy of a maximum value as  $2361 \text{ cm}^{-1}$  can be attributed to the adsorption of CO that came from the air atmosphere. According to M.R. Delgado (2003), there is a formation of  $Ga^{3+}$ –CO adducts on coordinatively  $Ga^{3+}$  ion located at edges and corners of  $β$ -Ga<sub>2</sub>O<sub>3</sub> crystallites. Lateral interactions between adsorbed CO molecules occur due to the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> has a crystal structure where the  $Ga^{3+}$  ions occupy both distorted tetrahedral and octahedral sites [[32,](#page-6-19) [33\]](#page-6-20). Another weak and broad band, with a maximum at 3728  $cm^{-1}$  is ascribed to –OH bending vibrations of H<sub>2</sub>O molecules absorbed on the β-Ga<sub>2</sub>O<sub>3</sub> surface, suggesting that the Ga–OH radicals could be incorporated into the gallium oxide. Both weaker regions showed hydrogen stretch modes which came from the  $H_2O$  adsorbed by the β-Ga<sub>2</sub>O<sub>3</sub> nanofibers investigated at room temperature in air atmosphere [[19,](#page-6-7) [34](#page-6-21)].

 $β$ -Ga<sub>2</sub>O<sub>3</sub> nanofibers could be obtained as the final product in this proposed work, with Ga–O stretching vibrations and O–Ga–O bending vibrations appearing at 440 and 670 cm−1, respectively, as depicted in Fig. [6](#page-5-11).

In summary, broadness, intensity, and emission bands in the visible spectrum show that the materials are promising candidates for optoelectronic applications, as they were demonstrated in the literature [\[35\]](#page-6-22).

## **4 Conclusions**

Electrospinning was applied with successful, synthesizing nanofibers of β-Ga<sub>2</sub>O<sub>3</sub> in a monoclinic structure and beta phase after thermal treatment at 900 °C for 1 h in the oven. The diameter and morphology of electrospun nanofbers were controlled by applying the temperature heating rate as 1, 5 and 10 °C/min, aiming fne nanofbers. Aging temperature and time of synthesis control the formation of oxygen vacancies, which were the driving forces of the existence of the strong blue emission. The nanofbers' pores decreased at slower heating rate, and the length was uniform up to several micrometers due to the surface-to-core extension, also called reversed crystal growth mechanism that occurred through the Ostwald ripening process.  $GaO<sub>6</sub>$  octahedral and  $GaO<sub>4</sub>$ tetrahedral chains were the two types of  $Ga^{3+}$  ions that came from diferent vibrations of Ga–O bonds, demonstrated in the Raman and FTIR spectra. There was also a formation of  $Ga^{3+}$ –CO adducts on coordinatively  $Ga^{3+}$  ion located at edges and corners of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> crystallites, showed in FTIR spectrum. Thus, successful results of this work included the synthesis and the broad, intense, and visible emission in the blue range of the wavelength, representing strong potential of β-Ga<sub>2</sub>O<sub>3</sub> materials for applications as light-emitting materials in optoelectronic devices.

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#### **References**

- <span id="page-5-0"></span>1. L. Wang, Z. Hou, Z. Quan, H. Lian, P. Yang, J. Lian, Mater. Res. Bull. **44**, 1978–1983 (2009)
- <span id="page-5-1"></span>2. H.F. Greer, F.J. Yu, W.Z. Zhou, Sci. China **54**, 1867–1876 (2011)
- <span id="page-5-2"></span>3. T. Zhou, P. Chen, S. Hu, Y. Yan, W. Pan, H. Li, Ceram. Int. **42**, 6467–6474 (2016)
- <span id="page-5-3"></span>4. J. Zhao, W. Zhang, E. Xie, Z. Ma, A. Zhao, Z. Liu, Appl. Surf. Sci. **257**, 4968–4972 (2011)
- <span id="page-5-4"></span>5. J.G. Zhao, Z.X. Zhang, Z.W. Ma, H.G. Duan, X.S. Guo, E.Q. Xie, Chin. Phys. Lett. **25**, 3787–3789 (2008)
- <span id="page-5-5"></span>6. C. Han, W. Mao, K. Bao, H. Xie, Z. Jia, L. Ye, Int. J. Hydrog. Energy **42**, 19913–19919 (2017)
- <span id="page-5-6"></span>7. Y. Yoon, K.I. Han, B.H. Kim, I.G. Lee, Y. Kim, J.P. Kim, W.S. Hwang, Thin Solid Films **645**, 358–362 (2018)
- <span id="page-5-7"></span>8. B. Fernandes, M. Hegde, P.C. Stanish, Z.L. Mišković, P.V. Radovanovic, Chem. Phys. Lett. **684**, 135–140 (2017)
- <span id="page-5-8"></span>9. Y. Zhang, J. Yang, Q. Li, X. Cao, J. Cryst. Growth **308**, 180–184 (2007)
- <span id="page-5-9"></span>10. H. Chen, G.D. Li, M. Fan, Q. Gao, J. Hu, S. Ao, C. Wei, X. Zou, Sens. Actuators, B **240**, 689–696 (2017)
- <span id="page-5-10"></span>11. L. Cui, H. Wang, B. Xin, G. Mao, Appl. Phys. A **123**, 634 (2017)
- <span id="page-6-0"></span>12. D. Calestani, A.B. Alabi, N. Coppedè, M. Villani, L. Lazarinni, F. Fabbri, G. Salviati, A. Zappettini, J. Cryst. Growth **457**, 255–261 (2017)
- <span id="page-6-1"></span>13. S.J. Park, N.A.M. Barakat, K.U. Jeong, H.Y. Kim, Polym. Int. **60**, 322–326 (2011)
- <span id="page-6-2"></span>14. C. Sun, J. Deng, L. Kong, L. Chen, Z. Chen, Y. Cao, H. Zhang, X. Wang, in *5th Annual International Conference on Materials Science and Engineering*, vol. 275 (2017), p. 012046
- <span id="page-6-3"></span>15. J. Zhao, W. Zhang, E. Xie, Z. Liu, J. Feng, Z. Liu, Mater. Sci. Eng., B **176**, 932–936 (2011)
- <span id="page-6-4"></span>16. G. Sinha, S. Chaudhuri, Mater. Chem. Phys. **114**, 644–649 (2009)
- <span id="page-6-5"></span>17. H.M. Lam, M.H. Hong, C.L. Gan, T.C. Chong, in *Fifth International Symposium on Laser Precision Microfabrication*, vol. 5662 (2004), pp. 62–66
- <span id="page-6-6"></span>18. T. Zhang, J. Lin, X. Zhang, Y. Huang, X. Xu, Y. Xue, J. Zou, C. Tang, J. Lumin. **140**, 30–37 (2013)
- <span id="page-6-7"></span>19. Y. Cheng, J. Chen, K. Yang, Y. Wang, Y. Yin, H. Liang, G. Du, J. Vac. Sci. Technol., B **32**, 03D119 (2014)
- <span id="page-6-8"></span>20. S. Hinrichsen, H. Broda, C. Gradert, L. Soencksen, F. Tuczek, Annu. Rep. Prog. Chem. Sec. A **108**, 17–47 (2012)
- <span id="page-6-9"></span>21. B.M. Flöser, F. Tuczek, Coord. Chem. Rev. **345**, 263–280 (2017)
- <span id="page-6-10"></span>22. S. Kumar, V. Kumar, T. Singh, A. Hähnel, R. Singh, J. Nanopart. Res. **16**, 2189 (2014)
- <span id="page-6-11"></span>23. G. Pozina, M. Forsberg, M.A. Kaliteevski, C. Hemmingsson, Sci. Rep. **7**, 42132 (2017)
- <span id="page-6-12"></span>24. J.M. Jeong, Y.J. Kwon, H.Y. Cho, H.G. Na, H.W. Kim, J. Ceram. Process. Res. **15**, 428–432 (2014)
- <span id="page-6-13"></span>25. C. Liu, Y. Berencén, J. Yang, Y. Wei, M. Wang, Y. Yuan, C. Xu, Y. Xie, X. Li, S. Zhou, Semicond. Sci. Technol. **33**, 095022 (2018)
- <span id="page-6-14"></span>26. A. Khan, S.N. Khan, W.M. Jadwisienczak, M.E. Kordesch, Sci. Adv. Mater. **1**, 236–240 (2009)
- <span id="page-6-15"></span>27. I. Nowak, M. Misiewicz, M. Ziolek, A. Kubacka, V.C. Coberán, B. Sulikowski, Appl. Catal. A **325**, 328–335 (2007)
- <span id="page-6-16"></span>28. S. Kumar, C. Tessarek, S. Christiansen, R. Singh, J. Alloys Compd. **587**, 812–818 (2014)
- <span id="page-6-17"></span>29. Y. Quan, D. Fang, X. Zhang, S. Liu, K. Huang, Mater. Chem. Phys. **121**, 142–146 (2010)
- 30. K. Girija, S. Thirumalairajan, G.S. Avadhani, D. Mangalaraj, N. Ponpandian, C. Viswanathan, Mater. Res. Bull. **48**, 2296–2303 (2013)
- <span id="page-6-18"></span>31. H.J. Bae, T.H. Yoo, Y. Yoon, I.G. Lee, J.P. Kim, B.J. Cho, W.S. Hwang, Nanomaterials **8**, 594 (2018)
- <span id="page-6-19"></span>32. M.R. Delgado, C.O. Areán, Mater. Lett. **57**, 2292 (2003)
- <span id="page-6-20"></span>33. C.O. Areán, A.L. Bellan, M.P. Mentruit, M.R. Delgado, G.T. Palomino, Microporous Mesoporous Mater. **40**, 35 (2000)
- <span id="page-6-21"></span>34. H. Xiao, H. Pei, W. Hu, B. Jiang, Y. Qiu, Mater. Lett. **64**, 2399– 2402 (2010)
- <span id="page-6-22"></span>35. A.V. Rodrigues, M.O. Orlandi, Ceram. Int. **45**, 5023–5029 (2019)

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