

Thermal, structural and optical properties of TeO₂–Na₂O–TiO₂ glassy **system**

Jaqueline Valeski Gunha¹ · Anderson Gonçalves1,2 · Aloisi Somer1 [·](http://orcid.org/0000-0003-0610-6220) André Vitor Chaves de Andrade¹ · Daniele Toniolo Dias3 · Andressa Novatski[1](http://orcid.org/0000-0002-8327-6285)

Received: 27 February 2019 / Accepted: 8 May 2019 / Published online: 14 May 2019 © Springer Science+Business Media, LLC, part of Springer Nature 2019

Abstract

In this work, $80TeO_2-(20-x) Na_2O-xTiO_2 (x=0, 5, 10 and 20) mol% glasses by melting quenching technique were obtained.$ Density (ρ) and molar volume (V_m) of these glasses have been investigated. The characterization was performed by X-ray difraction (XRD), diferential scanning calorimetry (DSC), Fourier transform infrared (FTIR) and ultraviolet–visible (UV– Vis) spectroscopy. As expected, the V_m values decreased with x increasing, exhibiting a density opposite behavior. DSC, FTIR and UV–Vis results indicated that, with the increasing the $TiO₂$ content, the role of this oxide changes from network modifer to former. In addition, more covalent bonds are formed, reinforcing the network connectivity by increasing the Te–O–Ti inter-chain bounds. These features improved the glassy system thermal and optical properties.

1 Introduction

 $TeO₂$ -based glasses have increasing scientific and technological interest due to wide optical transmission window (0.350 to 5.0 μ m), consistent vitreous stability ($\Delta T > 100 \text{ cm}^{-1}$) and low phonon energy (between 600 to 850 cm⁻¹) [[1–](#page-5-0)[3\]](#page-5-1). Additionally, these materials behold density over 4 g/cm^{-3} and linear refractive index higher than 2 [\[4](#page-5-2)[–6](#page-6-0)]. These properties allows application for photonics and electronics, such as: optical fber, amplifers and sensors [[2,](#page-5-3) [7,](#page-6-1) [8\]](#page-6-2).

The $TeO₂$ -glassy is obtained only with ultra-rapid cooling rates [[9](#page-6-3)]. However, the addition of modifer oxides, such as alkali, alkaline earth or transition metal oxides drops the cooling rates to 1 kJ/s, making it possible to obtain these glasses by conventional melt-quenching technique [[10](#page-6-4)[–13](#page-6-5)].

The addition of $Na₂O$ in TeO₂-based glasses was first reported by Mochida et al. with a range of glass formation between $10 \le x \le 46.5$ mol% [[11\]](#page-6-6). The Na₂O oxide plays the

 \boxtimes Jaqueline Valeski Gunha jaquegunha@gmail.com

- ² Departamento de Física, Universidade Estadual de Maringá, Maringá, PR 87020-900, Brazil
- ³ Departamento Acadêmico de Física, Universidade Tecnológica Federal do Paraná, Ponta Grossa, PR 87030-900, Brazil

role of network modifer, causing the transformation of the basic structural unit TeO₄ to TeO₃₊₁, breaking the Te–O–Te linkages [[14\]](#page-6-7). Nishida et al. [\[15\]](#page-6-8) proposed that this transformation occurs due to the change of the glass matrix from a three- or two-dimensional network structure to a lower dimensional one. These structural changes gives rise to consequences on the thermal properties decreasing the vitreous transition temperature (T_g) [\[3](#page-5-1), [16](#page-6-9)].

On the other hand, $TiO₂$ is effective in reinforce TeO₂ glass network by the substitutional positioning of Te to Ti site [\[6,](#page-6-0) [17,](#page-6-10) [18](#page-6-11)]. In addition, the incorporation of this oxide increases third-order optical non-linearities, linear refractive index and thermal stability for systems like TeO₂–Li₂O–TiO₂. These features have been attributed to the change in coordination from TeO₃ to TeO₄ units by Te–O–Ti linkages [[4,](#page-5-2) [19–](#page-6-12)[22\]](#page-6-13).

In this sense, the correlation between structural with thermal and optical properties of $TeO_2-Na_2O-TiO_2$ glassy system were investigated, in order to understand the role of $TiO₂$ into the glass as Na₂O was exchanged for TiO₂.

2 Materials and methods

The 80TeO₂–(20−x) Na₂O–xTiO₂ (x=0, 5, 10, 15 and 20 mol%) samples were prepared by means of melt-quenching technique in air atmosphere. The precursors reagents: TeO₂ (Alfa Aesar 99.99%), Na₂O (Sigma-Aldrich 99.5%)

¹ Departamento de Física, Universidade Estadual de Ponta Grossa, Ponta Grossa, PR 84016-210, Brazil

and $TiO₂$ (Synth 99.5%) were weighted and mixed with pistil and agate mortar during 40 min. The mixtures were melted in Pt crucible at a temperature of 840 °C for 30 min. The melt was quenched in a preheated brass mold at a temperature 20 °C lower than T_g and submitted to an annealing treatment for 240 min.

The density measurements were carried out by means the Archimedes's Principle with xylene ($\rho = 0.88$ g/cm³⁾ using a balance (AY220, Shimadzu) with 10^{-4} g precision. The difractograms were recorded on X-ray Automatic Difractometer (Ultima IV, Rigaku) at a scan rate of 2º/min (CuKα radiation $\lambda = 1.5418$ Å, 30 kV, 40 mA).

Differential Scanning Calorimetry (DSC) data were obtained with a 4 °C precision in calorimeter (Labsys Evo, Setaram Instrumentation). Such measures were realized under Argon atmosphere with 20 mL/min fux and heating rate of 10 °C/min, from ambient temperature to 600 °C using alumina crucible. In order to apply DRX and DSC, the samples were powdered and sieved (\leq 53 µm).

The structural analysis was carried out by means of Fourier transform infrared (FTIR) spectroscopy. For FTIR it was used a spectrometer, (IR-Prestige 21, Shimadzu). The glassy powdered samples were pulverized in an agate mortar with 100 mg of KBr in the 1:50 proportion. The resulting spectra were obtained in the absorbance model in the interval between 4000 and 400 cm⁻¹ using 64 scans with 4 cm⁻¹ resolution.

The optical absorption spectra were acquired using a spectrometer (Carry 50, Varian) in the wavelength range 300–700 nm. The samples thickness were 300 µm.

3 Results and discussion

3.1 Density, molar volume and oxygen density packing

The physical parameters density (ρ), molar volume (V_m) and oxygen density package (OPD) are presented in Table [1](#page-1-0). The ρ values observed in samples for x=0 and x=20 corroborate with the previously reported [[23](#page-6-14), [24](#page-6-15)]. The exchange of Na₂O by TiO₂ increases the ρ values, which is due to the substitution of a lower molecular weight substance $(Na₂O)$ 61.98 g/mol) with a higher molecular weight one (TiO₂) 79.87 g/mol), resulting in an increase in total molecular weight. Besides, the density increase is due to the structural modifcations occurring in the amorphous network. In sample with $x=0$, Na₂O plays the role of network modifier changing the Te atom coordination from TeO₄ \rightarrow TeO₃. With TiO₂ addition, more covalent bonds are formed, promoting a continuous glassy network by the Te–O–Ti linkages [\[17](#page-6-10)].

Other physical parameters which can contribute with valuable information about the changes in glass structure are molar volume and the oxygen packing density. V_m was calculated by using $V_m = \frac{M}{\rho}$, where M is the glass sample molecular weight. The OPD was determined using the equation: $OPD = 1000C_0 \times \frac{\rho}{M}$, where C₀ is the number of oxygen atoms per formula unit. As expected, the V_m values (Table [1\)](#page-1-0) decrease from 29.1 to 27.26 cm^3/mol , with increasing x, exhibiting density opposite behavior. The increase of OPD values (Table [1\)](#page-1-0) with x is an indicative of close packing glass network structure. This packing degree is dependent of cation feld intensity and the cation radius [\[25](#page-6-16)]. Since the ionic radius of Na⁺ and Ti⁴⁺ are 0.098 nm and 0.064 nm, respectively [[26\]](#page-6-17). Therefore, the average of the cation feld intensity increases with x, consequently increasing the OPD.

3.2 X‑ray difraction (XRD) results

Figure [1](#page-2-0) presents the XRD results of the studied samples. We can notice that the sample in the composition for $x=0$ is mostly amorphous, except for two sharp peaks at 13.3 and 26.5° superimposed on broad peaks due to the majority glassy phase. In fact, during the experimental studies, this sample was observed to be hygroscopic, such behavior is in accordance with literature for $(100-x)TeO₂–xNa₂O$ binary system. This feature can be attributed to the hydrated compound presence such as $Na₂Te₂O₅·2H₂O$ phase [\[23](#page-6-14)]. For $x=5$ we notice the predominant amorphous character with some sharp peaks superimposed at 13.3 and 26.5°, showing the same $Na₂Te₂O₅$ phase presence. The peaks at 26.1, 29.8 and 32.0° are attrib-uted to the TiTe₃O₈ phase [\[27](#page-6-18)]. The α TeO₂ phase presents peaks at 28.7 and 48.5°. For $x = 10$ and $x = 15$ the samples presented amorphous character. The sample with $x=20$ contained signifcant amount of crystalline phases and presented several

Table 1 Values of: molar volume (V_m) packing density (O glass transition (T_{φ}) crystallization (T_x) crystallization (T_c) stability $(ΔT)$; and Weinberg (H_w)

Fig. 1 Difractograms indicating the diferent phases formed in $x=20$, 15 and 0 samples

peaks attributed αTeO₂ (19.7, 28.7 and 48.5°) and TiTe₃O₈ (26.1, 29.8, 32.0, 34.6, 46.7 and 55.4°) [\[28](#page-6-19)].

Figure [2](#page-2-1) shows, the all sample photo that were characterized in Fig. [1.](#page-2-0) These samples change from yellowish to brownish with increasing x, being the glassy sample $(x=10)$ and $x=15$) and partially crystallized ($x=0$, 5 and 20).

3.3 Diferential scanning calorimetry (DSC) results

Thermal behavior of each sample was determined by DSC analysis. The obtained thermograms are presented in Fig. [3.](#page-3-0) From these data, the glass transition (T_g) , the onset crystallization (T_x) , the first exothermic peak at the crystallization (T_c) temperatures and the thermal stability were determined. The last one was determined using two criteria: $\Delta T = T_x - T_g$, which is frequently used for glass materials, and the Weinberg parameter $H_w = \frac{f_x - T_g}{T_g}$ [\[29,](#page-6-20) [30](#page-6-21)]. These results are presented in Table [1](#page-1-0). We can notice an increase in T_g from 245 to 367 °C

with the exchange of $Na₂O$ for TiO₂. This increase can be related to the reinforcement of the glass network due to the high coordination number of Ti^{4+} (4 or 6) [\[31\]](#page-6-22). Besides, the asymmetric crystallization peak for $x=0$, 5 and 20 suggests different distinct phase crystallizations. The shift in T_c from 370 to 340 °C with $x=5$ is an indicator of the increasing concentration of nuclei in the glass and consequently the probability of crystallization inside the glass matrix increases, this is a direct consequence of the αTeO2 presence, as observed in XRD results [\[32](#page-6-23)].

In binary glasses, excessive addition of a network modifer decreases vitrifcation behavior, due to the decreasing stability of TeO₃ units [\[23](#page-6-14)], increasing T_g values and decreasing glassforming ability. So, the addition of 5% mol of TiO₂ makes this oxide plays a role of modifer network. However, for 10 and 15% mol the $TiO₂$ assumes the role of former network, rearranging the breaking caused by Na_2O which creates TeO₃ and TeO₃₊₁ and does not allow the formation of the Na₂Te₂O₅ phase, facilitating the amorphization of the glass network. This directly impacts the shape of the peak T_c which is narrower for the samples $x = 10$ and 15 and is shifted to higher temperatures, indicating a lower concentration of nuclei that induce the crystallization.

In Table [1,](#page-1-0) ΔT decreases from $x=0$ to $x=5$ and slightly increases for $x=10$ and 15. This behavior is in accordance with XRD results, since the samples $x=10$ and 15 presented an amorphous behavior showing that the melting quenching technique was successfully performed for these two compositions. This result endorses the change of the role of $TiO₂$ oxide from network modifier to former. Besides, H_w (Table [1\)](#page-1-0) decreases with the exchange of $Na₂O$ for $TiO₂$, that suggest an increase in the crystallization tendency [[33\]](#page-6-24) which is an indicator of the substitution of Te–O–Te for Te–O–Ti bounds favoring the covalent character of the network [\[34](#page-6-25)]. These statements are reinforced by ρ , V_m and OPD results.

3.4 FTIR analysis

In order to obtain some insights about the structural change with composition, FTIR experiments were performed and the results are shown in Fig. [4](#page-3-1). According to the literature, in the region from 400 to 1000 cm−1, there are, at least, eight characteristic regions corresponding to the diferent structural units present in tellurite glasses with TiO₂: ~450, 580, 620, 660, 710, 750, 780, 800–830, 852–890. [[22,](#page-6-13) [23](#page-6-14), [27,](#page-6-18) [32](#page-6-23), [35–](#page-6-26)[38](#page-6-27)].

The region at ~450 cm⁻¹ can be related to Te–O–Te and/ or Te–O–Ti stretching vibrations. Besides, for $x = 0$ we notice a shoulder at 451 cm⁻¹ and for $x=20$ at 440 cm⁻¹.

Fig. 2 Samples photograph in bulk after melting quenching

Fig. 3 DSC scans of samples

Fig. 4 FTIR spectra of samples

This shoulder has been related to symmetric stretching of O–Te–O in α-TeO₂ structure. The 467 cm⁻¹ peak for $x=0$ is related to Te–O–Te stretching vibrations. For $x = 20$, Te–O–Ti stretching vibrations appear at 486 cm⁻¹. For $x = 5$, 10 and 15 these modes were superimposed and presented a broad band in this region [[27,](#page-6-18) [32,](#page-6-23) [35\]](#page-6-26).

The region at 2900 - 3500 cm^{-1} is attributed to the vibration of the elongation of the hydroxyl group and hydrogen. This high intensity region decreases with increasing x. In this case the high $Na₂O$ content confirm the sample tendency to be hygroscopic [\[36](#page-6-28)].

FTIR results in the region from 500 to 950 cm−1 were detailed into Gaussian ftting as presented in Fig. [5](#page-3-2). The band centered at 580 cm−1, which is attributed to Te–O stretching vibrations of TeO₄ and TeO₃₊₁ units, remains at the same position but decreases in area from $x=10$ to $x=20$. According to Dimitriev [\[38](#page-6-27)], the addition of modifier network stimulates TeO_{3+1} units formation. Our results showed that the $TiO₂$ acts as network modifier for the sample with $x=5$. On the other hand, for $x=10$, 15 and 20 its role changed from network modifer to former reinforcing the DSC results as shown in Fig. [3](#page-3-0).

The band area at 620 cm^{-1} increases with x content. This mode, which is attributed to Te–O bonds of deformed $TeO₄$ units, increases from $x=0$ to $x=5$, remains constant from $x=5$ to $x=15$ and increases to $x=20$. This behavior is an indicative that $TiO₂$ inhibits the Na2O action in creates more TeO₃₊₁ and TeO₃ from TeO₄ units [[3,](#page-5-1) [23\]](#page-6-14). The 660 cm⁻¹ band is associated to asymmetric stretching vibration of Te–O bonds in TeO₃ pyramidal units [[3\]](#page-5-1). The intensity of this band remains approximately constant up to $x=15$ and

Fig. 5 Deconvoluted FTIR spectra of samples

decreases for $x=20$. The 710 cm⁻¹ band increases up to $x=10$ and decreases with increasing x, this band can be attributed to TeO₃ units [\[38](#page-6-27)]. The band at 750 cm⁻¹ is attributed to TeO_{3+1} unit and increases with x content. The band at 773 cm⁻¹ is well resolved for $x=0$ and $x=5$ samples and it is attributed to αTeO_2 [\[36\]](#page-6-28). For the other samples this mode vibration is superimposed with the 800 cm−1 which it is attributed to TeO₃ units $[23]$ $[23]$. This band decrease in intensity increasing x content in the samples. In Fig. [5,](#page-3-2) shifts from 800 to 852 cm⁻¹ and from 831 to 889 cm⁻¹, with x increasing from 5 to 20 can be perceived. These bands are related to stretching vibration of Ti–O–Ti bridges in $TiO₄$ units $[36]$ $[36]$. FTIR analyses indicate that when Ti is incorporated in Te–O–Te interchain, Te–O–Ti bonds are formed, and for relative high $TiO₂$ composition and after a saturation of Te–O–Ti interchains, occurs the TiO₄ units formation $[22, 1]$ $[22, 1]$ $[22, 1]$ [23](#page-6-14), [27](#page-6-18), [35](#page-6-26), [36](#page-6-28)].

3.5 Optical absorption

The analysis of the optical absorption edge is useful to obtain some insights about the optical transition and optical bandgap of materials. According to the literature [[39](#page-6-29)], in glasses, the UV absorption edge shifts to longer wavelengths by increasing the molecular weight of their precursor oxides. The yellow color of the glasses, as observed in Fig. [2,](#page-2-1) is due to a variable shift of the ultraviolet absorption edge towards the visible range, which depends on the glass composition. For the analyzed samples, it was not possible to analyze the $x=20$ sample UV edge, since it was completely opaque even for 300 µm thickness. For the other samples, the optical band gap and Urbach´s energy were determined.

Davis and Mott [[40](#page-6-30)] presented an expression for the absorption coefficient (α) as a function of photon energy $(h\nu)$, as follow:

$$
\alpha(hv) \propto \left(hv - E_G\right)^m \tag{1}
$$

where α is the absorption coefficient, (hv) is the energy of the photon, E_g is the bandgap energy and *m* indicates the type of electronic transition involved in the absorption process. The values for *m* depend on the bandgap characteristic: $m=1/2$ for direct allowed transition; $m=3/2$ for direct forbidden transition, m=2 for indirect allowed transition and m=3 for indirect forbidden transition. The absorption spectra presented in Fig. [6a](#page-4-0) were used to determine the optical bandgap (E_{\circ}) of samples using Eq. [1](#page-4-1). Figure [6](#page-4-0)b presents the curves of $(\alpha.h\nu)^{1/m}$ versus photon energy, with the E_g values being estimated by extrapolation of the linear portion of the curves, where $(\alpha.h\nu)^{1/m}$ = 0. The best fit was obtained for $m=2$, indicating an indirect allowed transition [\[41\]](#page-6-31). Another method to determine the bandgap is the intersection point between lines that extrapolate absorption coefficient values

Fig. 6 a Optical absorption spectra of the studied glasses. **b** $(\alpha \cdot h \nu)^{1/2}$ versus photon energy showing the linear regression for E_g determination

Table 2 Bandgap values (E_g and E'_g), Urbach's energy (ΔE), and electronic ion polarizability of the samples

Sample	$E_{\rm g}$ (eV) (\pm 0.003)	E_g (eV)	ΔE (eV) (± 0.004)	$\gamma_O^2(E_g)$ \AA^3
$x=0$	3.394	3.46	0.488	3.02 ± 0.03
$x = 5$	2.915	2.99	0.171	3.04 ± 0.01
$x = 10$	2.857	2.93	0.185	2.922 ± 0.007
$x = 15$	2.845	2.89	0.198	2.831 ± 0.006

in the small hv range and at the linear absorption edge (E'_g) λ [[42\]](#page-6-32). Table [2](#page-4-2) presents the values of E_g and E'_g in eV. The E_g values changes from 3.394 to 2.845 eV with increasing *x* content. In principle, this behavior would be related to the increase in defects of glassy network, however from FTIR and DSC results we noticed that the exchange of $Na₂O$ for $TiO₂$ is reinforcing the network. In fact, this behavior can be explained by means of ion oxide polarizability and it will be discussed later.

The optical absorption coefficient near the fundamental absorption band has an exponential dependency, which is given by the Urbach´s rule [\[41\]](#page-6-31):

$$
\alpha = \alpha_0 \exp\left(\frac{h\nu}{\Delta E}\right) \tag{2}
$$

in which α_0 is a constant and ΔE is the Urbach's energy. ΔE corresponds to the optical transitions between localized tail states adjacent to the valence band and extended states in the conduction band above the mobility edge [\[43](#page-6-33)]. The absorption coefficient values (in log scale) were plotted against the photon energies. The ΔE values were obtained by linear ftting from those plots and are presented in Table [2](#page-4-2). It can be notice that the ΔE values decrease and then increase from $x=5$ to $x=15$. According to Mott and Davis [\[40](#page-6-30)] model, this energy depends on the degree of the both disorder and defects of the amorphous structure. If the glass has a more polymerized network, Urbach´s energy is lower.

According to the literature [[43\]](#page-6-33), the behavior of the polarizability of the electronic ion are calculated by the following Eq. [3](#page-5-4):

$$
\gamma_{O^2-}(E_g) = \left[\left(\frac{V_m}{2,52} \right) \left(1 - \sqrt{\frac{E_g}{20}} \right) - \sum p \gamma_i \right] (q)^{-1} \tag{3}
$$

were γ is a polarizability of the cation, V_m is the molar volume, p denotes the number of cations and q oxide ions respectively, in de chemical formula $A_p O_q$. Figure [7](#page-5-5) presents the bandgap values (E_g) , Urbach's energy (ΔE), and the polarizability (γ_O^2) versus TiO₂ content.

The polarizability behavior should be inverse to the bandgap; however, this behavior does not occur to our system. The substitution of Na₂O by TiO₂ generated a decrease in the bandgap, but there is a polarizability atypical behavior, which does not accompany inversely the bandgap.

Fig. 7 a Bandgap and ΔE energy values and **b** polarizability values $γ_{02}$ versus TiO₂ content

Therefore, this same behavior which decreases the polarizability is related to the increase of mean simple bonding [[43,](#page-6-33) [44](#page-6-34)]. This increase in average simple bonding is observed when a certain amount of $TiO₂$ added, suggesting that there is an *onset* of composition which allow the formation of strong chemical bonds in the glass network. This statement is reinforced by the DSC results, which show a decrease in thermal stability for the samples. Finally, the feature of changing the role of $TiO₂$ from network modifier to former drastically infuences the thermal and optical properties in TeO₂–Na₂O–TiO₂ glassy system.

4 Conclusion

We conclude that the structural of the $TeO_2-Na_2O-TiO_2$ system investigated by FTIR analysis, with the equimolar substitution of Na_2O by TiO₂ showed a structural transformation in the TeO2 network. The $Na₂O$ performs the role of network modifer, changing the coordination of the atom of TeO₄ \rightarrow TeO₃. With the addition of TiO₂, more covalent bonds are formed, promoting a continuous glass network through the Te–O–Ti bond. The XRD results corroborate with FTIR, and the presence of the hydrated compound, such as the $Na₂Te₂O₅·2H₂O$ phase and the TiTe₃O₈ phase.

In addition, the decrease in polarization is related to the increase in the mean simple bonding. This strong bond increase, which is related to decreases of the bandgap and polarizability values, indicates the increase of Te–O–Ti bonds (band at 450 cm⁻¹). This increase also agrees with the DSC results, which show a decrease in the thermal stability for the samples. Finally, we conclude that $TiO₂$ changes its role from network modifer to former after 5 mol%, which drastically infuences the thermal and optical properties through increasing the network connectivity.

Acknowledgments The authors thank to FINEP, CAPES, CNPQ and Fundação Araucária for the partial fnancial support. To LABMU-UEPG for technical support.

References

- 1. H. Bürger, W. Vogel, V. Kozhukharov, IR transmission and properties of glasses in the TeO₂–RnOm, RnXm, Rn $(SO₄)$ m, Rn (PO_3) mand B_2O_3] systems. Infrared Phys. 25, 395–409 (1985)
- 2. J.S. Wang, E.M. Vogel, E. Snitzer, Tellurite glass: a new candidate for fber devices. Opt. Mater. (Amst) **3**, 187–203 (1994)
- 3. J. Heo, D. Lam, G.H. Sigel Jr., E.A. Mendoza, D.A. Hensley, Spectroscopic analysis of the structure and properties of alkali tellurite glasses. J. Am. Ceram. Soc. **75**, 277–281 (1992)
- 4. F. Chen, T. Xu, S. Dai, Q. Nie, X. Shen, J. Zhang, X. Wang, Linear and non-linear characteristics of tellurite glasses within $TeO₂$ Bi2O3–TiO2 ternary system. Opt. Mater. (Amst) **32**, 868–872 (2010)
- 5. H. Takebe, S. Pujino, K. Morinaga, Refractive-index dispersion of tellurite glasses in the region from 0.40 to 1.71 μm. J. Am. Ceram. Soc. **77**, 2455–2457 (1994)
- 6. M. Udovic, P. Thomas, A. Mirgorodsky, O. Durand, M. Soulis, O. Masson, T. Merle-Méjean, J.-C. Champarnaud-Mesjard, Thermal characteristics, Raman spectra and structural properties of new tellurite glasses within the Bi_2O_3 –TiO₂–TeO₂ system. J. Solid State Chem. **179**, 3252–3259 (2006)
- 7. N. Boubata, A. Roula, I. Moussaoui, Thermodynamic and relative approach to compute glass-forming ability of oxides. Bull. Mater. Sci. **36**, 457–460 (2013)
- 8. S. Manning, H. Ebendorff-Heidepriem, T.M. Monro, Ternary tellurite glasses for the fabrication of nonlinear optical fbres. Opt. Mater. Express. **2**, 140–152 (2012)
- 9. P.T. Sarjeant, R. Roy, New glassy and polymorphic oxide phases using rapid quenching techniques. J. Am. Ceram. Soc. **50**, 500–503 (1967)
- 10. R. El-Mallawany, The optical properties of tellurite glasses. J. Appl. Phys. **72**, 1774–1777 (1992)
- 11. N. Mochida, K. Takahashi, K. Nakata, S. Shibusawa, Properties and structure of the binary tellurite glasses containing mono-and di-valent cations. J. Ceram. Assoc. Jpn. **86**, 317–326 (1978)
- 12. H. Bürger, K. Kneipp, H. Hobert, W. Vogel, V. Kozhukharov, S. Neov, Glass formation, properties and structure of glasses in the TeO2–ZnO system. J. Non. Cryst. Solids. **151**, 134–142 (1992)
- 13. D. Souri, Y. Shahmoradi, Calorimetric analysis of non-crystalline TeO₂–V₂O₅–Sb₂O₃. J. Therm. Anal. Calorim. **129**, 601–607 (2017)
- 14. T. Sekiya, N. Mochida, A. Ohtsuka, M. Tonokawa, Raman spectra of $MO_1/2TeO_2$ (M=Li, Na, K, Rb, Cs and Tl) glasses. J. Non. Cryst. Solids. **144**, 128–144 (1992)
- 15. T. Nishida, S. Saruwatari, Y. Takashima, Structural study of Na₂O– TeO₂ glasses by Mössbauer spectroscopy and differential thermal analysis. Bull. Chem. Soc. Jpn **61**, 4093–4097 (1988)
- 16. Y. Himei, A. Osaka, T. Nanba, Y. Miura, Coordination change of Te atoms in binary tellurite glasses. J. Non. Cryst. Solids. **177**, 164–169 (1994)
- 17. T. Hayakawa, M. Koduka, M. Nogami, J.R. Duclère, A.P. Mirgorodsky, P. Thomas, Metal oxide doping efects on Raman spectra and third-order nonlinear susceptibilities of thallium–tellurite glasses. Scr. Mater. **62**, 806–809 (2010)
- 18. M. Soulis, A.P. Mirgorodsky, T. Merle-Méjean, O. Masson, P. Thomas, M. Udovic, The role of modifer's cation valence in structural properties of TeO₂-based glasses. J. Non. Cryst. Solids. 354, 143–149 (2008)
- 19. S.M. Lima, W.F. Falco, E.S. Bannwart, L.H. Andrade, R.C. Oliveira, J.C.S. Moraes, K. Yukimitu, E.B. Araújo, E.A. Falcão, A. Steimacher et al., Thermo-optical characterization of tellurite glasses by thermal lens, thermal relaxation calorimetry and interferometric methods. J. Non. Cryst. Solids. **352**, 3603–3607 (2006)
- 20. H. Nasu, O. Matsushita, K. Kamiya, H. Kobayashi, K. Kubodera, Third harmonic generation from Li_2O –TiO₂–TeO₂ glasses. J. Non. Cryst. Solids. **124**, 275–277 (1990)
- 21. W.A. Capanema, K. Yukimitu, J.C.S. Moraes, F.A. Santos, M.S. Figueiredo, S.M. Sidel, V.C.S. Reynoso, O.A. Sakai, A.N. Medina, The structure and optical dispersion of the refractive index of tellurite glass. Opt. Mater. (Amst) **33**, 1569–1572 (2011)
- 22. W. Stambouli, H. Elhouichet, M. Ferid, Study of thermal, structural and optical properties of tellurite glass with different $TiO₂$ composition. J. Mol. Struct. **1028**, 39–43 (2012)
- 23. K.B. Kavaklioğlu, S. Aydin, M. Çelikbilek, A.E. Ersundu, The TeO₂–Na₂O system: thermal behavior, structural properties, and phase equilibria. Int. J. Appl. Glas. Sci. **6**, 406–418 (2015). [https://](https://doi.org/10.1111/ijag.12103) doi.org/10.1111/ijag.12103
- 24. G. Jayasinghe, D. Coppo, P. Bandaranayake, J.L. Souquet, Electrical properties of TeO₂ glasses with $Na₂O$ as network modifier. Solid State Ionics **76**, 297–300 (1995)
- 25. A. Bachvarova-Nedelcheva, R. Iordanova, S. Ganev, Y. Dimitriev, Glass formation and structural studies of glasses in the $TeO₂–ZnO-$ Bi2O3–Nb2O5 system. J. Non. Cryst. Solids. **503**, 224–231 (2019)
- 26. V. Dimitrov, T. Komatsu, An interpretation of optical properties of oxides and oxide glasses in terms of the electronic ion polarizability and average single bond strength. J. Univ. Chem. Technol. Met. **45**, 219–250 (2010)
- 27. I. Shaltout, Crystallization kinetics and structure of $(TeO₂-TiO₂)$ Fe2O3) glasses. J. Mater. Sci. **35**, 323–329 (2000)
- 28. N. Elkhoshkhany, R. El-Mallawany, E. Syala, Mechanical and thermal properties of TeO₂–Bi₂O₃–V₂O₅–Na₂O–TiO₂ glass system. Ceram. Int. **42**, 19218–19224 (2016)
- 29. A.F. Kozmidis-Petrović, Theoretical analysis of relative changes of the Hruby, Weinberg, and Lu–Liu glass stability parameters with application on some oxide and chalcogenide glasses. Thermochim. Acta **499**, 54–60 (2010)
- 30. A.F. Kozmidis-Petrović, Sensitivity of the Hruby, Lu–Liu, Fan, Yuan, and Long glass stability parameters to the change of the ratios of characteristic temperatures Tx/Tg and Tm/Tg. Thermochim. Acta **510**, 137–143 (2010)
- 31. F. Chen, Q. Yu, B. Qiao, S. Dai, Q. Zhang, Influence of $TiO₂$ on thermal stability and crystallization kinetics of tellurite glasses within $TeO_2-Bi_2O_3-Nb_2O_5$ pseudo-ternary system. J. Non. Cryst. Solids. **404**, 32–36 (2014)
- 32. E. Idalgo, E.B. Araújo, K. Yukimitu, J.C.S. Moraes, V.C.S. Reynoso, C.L. Carvalho, Efects of the particle size and nucleation temperature on tellurite $20Li_2O-80TeO_2$ glass crystallization. Mater. Sci. Eng., A **434**, 13–18 (2006)
- 33. M.C. Weinberg, Assessment of glass stability criteria. Phys. Chem. Glas. **35**, 119–123 (1994)
- 34. J.-C. Sabadel, P. Armand, P.-E. Lippens, D. Cachau-Herreillat, E. Philippot, Mössbauer and XANES of TeO₂–BaO–TiO₂ glasses. J. Non Cryst. Solids. **244**, 143–150 (1999)
- 35. M. Çelikbilek, A. Erçin Ersundu, S. Aydin, Glass formation and characterization studies in the TeO₂–WO₃–Na₂O System. J. Am. Ceram. Soc. **96**, 1470–1476 (2013)
- 36. I.N. Sapian, M.I.M. Yusof, A.K. Yahya, Elastic and structural properties of (95–x) TeO₂-5La₂O₃-xTiO₂ lanthanum tellurite glass system. Chalcogenide Lett. **11**, 10 (2014)
- 37. S.Q. Mawlud, M.M. Ameena, M.R. Sahar, Z.A.S. Mahraz, K.F. Ahmed, Thermal stability and Judd-Ofelt analysis of optical properties of Sm3+-doped sodium tellurite glasses, in: AIP Conf. Proc., 2017: p. 20032
- 38. Y. Dimitriev, V. Dimitrov, M. Arnaudov, IR spectra and structures of tellurite glasses. J. Mater. Sci. **18**, 1353–1358 (1983)
- 39. M.A. Villegas, J.M.F. Navarro, Physical and structural properties of glasses in the TeO_2 –Ti O_2 –Nb₂ O_5 system. J. Eur. Ceram. Soc. 27, 2715–2723 (2007)
- 40. E.A. Davis, N.F. Mott, Conduction in non-crystalline systems V. Conductivity, optical absorption and photoconductivity in amorphous semiconductors. Philos. Mag. **22**, 903–922 (1970)
- 41. F. Urbach, The long-wavelength edge of photographic sensitivity and of the electronic absorption of solids. Phys. Rev. **92**, 1324 (1953)
- 42. M. Nowak, B. Kauch, P. Szperlich, Determination of energy band gap of nanocrystalline SbSI using difuse refectance spectroscopy. Rev. Sci. Instrum. **80**, 46107 (2009)
- 43. V. Dimitrov, S. Sakka, Electronic oxide polarizability and optical basicity of simple oxides. J. Appl. Phys. **79**, 1736–1740 (1996)
- 44. V. Dimitrov, T. Komatsu, Classifcation of oxide glasses: a polarizability approach. J. Solid State Chem. **178**, 831–846 (2005)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.