Structure, piezoelectric, ferroelectric and dielectric properties of leadfree ceramics $0.67B$ **iFeO₃–0.33BaTiO₃–** xB **iGaO₃+0.0035MnO₂**

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Abstract

BF–BT–*x*BGa (shortened for 0.67BiFeO₃–0.33BaTiO₃–*xBiGaO₃+0.0035MnO₂)* lead-free piezoelectric were prepared by a conventional solid-state sintering method. The phase structure, miscostructure, piezoelectric, ferroelectric and dielectric properties were investigated. BiGaO₃ doping has little effect on the crystal structure of BF–BT–*x*BGa lead-free piezoelectric ceramics. All samples are perovskite structure. With the change of *x*, the structure of the sample is pseudo cubic structure. The ceramics (with $x = 0.02$ sintered at 990 °C) has $P_r = 25 \mu C/cm^2$, $E_c = 22.059 \text{ kV/cm}$ and high depolarization temperature $T_d = 422$ °C. It was found that the piezoelectric properties of the ceramic were improved by addition of BiGaO₃. The ceramic shows excellent electrical properties when *x* is 0.02: $d_{33} = 170$ pC/N and $T_c = 434$ °C. Proper amount of BiGaO₃ doping can improve the size of the crystal and promote the sintering of samples. These results show that the addition of $BiFeO₃–BaTiO₃$ based ceramics is promising lead-free piezoelectric ceramics for practical applications.

1 Introduction

Piezoelectric ceramic is a kind of polycrystalline material with mechanical and electrical energy conversion function. It has a wide range of applications in medicine, industry, aerospace and other fields $[1-4]$ $[1-4]$ $[1-4]$. Lead zirconate titanate (PZT) based ceramics have attracted much attention due to their good piezoelectric and electromechanical coupling properties [[5](#page-7-2)–[8\]](#page-7-3). However, lead toxicity poses a serious threat to human health and the environment, so it has led us to constantly search for other lead free alternatives [[9](#page-7-4)]. Hexagonal PbCr_xFe_{12−*x*}O₁₉ nanostructures were prepared by sol–gel self combustion method [[10](#page-7-5)]. La³⁺ doped ZnFe₂O₄ nanostructures were prepared by a 2-step sol–gel process [\[11](#page-7-6)]. Ziarati et al. have proved the preparation of ultrasound. The rare earth element based on ferrite and NIFE_{2−*x*}Eu_xO₄ is an effective catalyst for the synthesis of benzimidazole, benzoxazole and benzothiazole under ultrasonic irradiation

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[[12](#page-7-7)]. Rahimi-Nasrabadi et al. have been reported that CuFe_{2−*x*}Ce_{*x*}O₄ nanoparticles were successfully synthesized by simple self combustion sol–gel method [[13](#page-7-8)].

 $BiFeO₃$ has recently gained considerable interest both technologically and scientifically due to the existence of both ferroelectric and (anti)ferromagnetic ordering in the same phase in the material [\[14](#page-7-9)[–17](#page-7-10)]. Bi₂Fe₄O₉, Bi₄₆Fe₂O₇₂ and other heterozygous phases will be produced in the preparation of $BiFeO₃$ by the traditional solid state sinter-ing method [[18](#page-7-11)[–21](#page-7-12)]. The ionic radius of Ba^{2+} in BaTiO₃ is 1.61 Å (CN = 12) and the ionic radius of Ti^{4+} is 0.605 Å $(CN=6)$ [[22\]](#page-7-13). BaTiO₃ has a low tetragonal–cubic transition temperature. This tetragonal perovskite phase is stable over an extended temperature range, so it has good piezoelectric properties as a ferroelectric material $(d_{33} \sim 400 \text{ pC/N})$ [\[23](#page-7-14)]. However, its lower Curie temperature $T_c \sim 130$ °C hindered its application. In general, $BaTiO₃$ -based lead-free piezoelectric ceramics have the advantages of high piezoelectric constant, high dielectric constant, large electromechanical coupling coefficient, mechanical quality factor and small dielectric loss. However, it has the disadvantages of low Curie temperature, high sintering temperature and strong sensitivity of the piezoelectric properties to temperature changes.

In addition, many researchers have studied the crystal structure of the $(1-x)$ BiFeO₃–*x*BaTiO₃ ceramic system, but the results of the study are different. By changing the

stoichiometry of BiFeO₃ and BaTiO₃ in (1–*x*)BiFeO₃–*x*Ba- $TiO₃$ ceramic system, Kumar et al. found that the crystal structure and lattice parameters of $(1-x)BiFeO₃–xBaTiO₃$ ceramic system were closely related to the stoichiometry of BiFeO₃ and BaTiO₃ [[24](#page-7-15)]. In 2009, Leontsev and Eitel reported that Mn doped $(1-x)BiFeO₃ - xBaTiO₃$, at $x = 0.25$, the piezoelectric coefficient d_{33} reached 116 pC/N, and its ferroelectric transition temperature was above 450 $^{\circ}$ C [\[25](#page-7-16)]. In 2015, Lee et al. studied $(1-x)[0.67Bi_{1.05}FeO₃-0.33B]$ aTiO₃]–*x*Bi_{1.05}(Zn_{0.5}Ti_{0.5})O₃ ($x = 0.03$, BF33BT–3BZT) and $0.67Bi_{1.05}(Fe_{1-x}Ga_x)O_3-0.33BaTiO_3$ ($x = 0.03$, BF33BT–3BG) piezoelectric ceramics. The piezoelectric constant d_{33} and the Curie temperature T_c are respectively BZT (BF33BT–3BZT, $d_{33} = 324$ pC/N, $T_c = 466$ °C) and BG (BF33BT–3BG, $d_{33} = 402$ pC/N, $T_c = 454$ °C) [[26](#page-7-17)]. $BiFeO₃$ –BaTiO₃ has been shown to have excellent piezoelectric properties and high temperature resistance. Zhou et al. has been studied that Ga doped $BiFeO₃ – BaTiO₃$ lead-free piezoelectric ceramics. It was found that a small amount of Ga could improve the grain size and piezoelectric properties of the sample ceramics [[27](#page-7-18)]. In this work, the ceramics of $BiGaO₃$ doped BF–BT–*x*BGa solid solution were prepared by conventional sintering technique. The microstructure, the piezoelectric properties, ferroelectric properties, dielectric properties and thermal stability of the ceramics were investigated.

2 Experimental procedures

BF–BT–*x*BGa ceramics were prepared by a conventional ceramic fabrication technique. Reagent-grade oxide of $Bi₂O₃$ (99.99%), Fe₂O₃ (99.2%), TiO₂ (99.99%), BaCO₃ (99.99%), and $MnO₂$ (99.99%) were used as the starting materials. The mixture was ball milled with addition of alcohol for 24 h. The samples were dried at a temperature of 800 °C for 4 h after drying, drying and screening after adding PVA binder. The samples were heated to 600° C by sintering 4 h for the glue, then sintered for 2 h at 970–1010 \degree C. The crystal structure of the sample was detected by XRD (D8 Advance; Bruker Inc., Karlsruhe, Germany). The surface morphology of the sample was observed by scanning electron microscope (SEM, JSM-5610LV, JEOL, Japan). The piezoelectric properties and dielectric properties were measured by the quasi static d_{33} tester (ZJ-3A, CAS, Shanghai, China) and the precision impedance analyzer (4294A, Agilent). The polarization–electric field (*P–E*) loops at room temperature was characterized by a ferroelectric tester (Premier II, Radiant Technologies Inc., Albuquerque, NM, USA). The temperature spectrum was measured by the intelligent temperature control system and the precision impedance analyzer (4294A, Agilent). The depolarization temperature of the sample was measured by *ex situ* depolarization method.

3 Results and discussion

Figure [1](#page-1-0)a shows the X-ray diffraction patterns of BF–BT–*x*BGa (*x*=0, 0.01, 0.02, 0.03, 0.04) system piezoelectric ceramics at 10°–80°. All the samples were characterized by XRD. The characteristic peaks were (100), (110), (111), (200), (210), (211), (220), (221), (310), and some small clutter peaks were observed in the XRD pattern, indicating that impure phase was formed. These small clutter peaks represent $Bi_2Fe_4O_9$. All sample ceramics are perovskite structure. According to the 20-0169 JCPDS card number $[28]$ $[28]$, we know that pure BiFeO₃ is a three party structure. According to the 5-626 JCPDS card number [\[29](#page-7-20)], we know that pure $BaTiO₃$ is a tetragonal structure. In this paper, the pseudo cubic structure of samples shows that the mixture of $BiFeO₃$ and $BaTiO₃$ can form pseudo cubic struc-ture. Figure [1](#page-1-0)b was shown the BF–BT– $xBGa$ ($x=0$, 0.01, 0.02, 0.03, 0.04) piezoelectric ceramic at 36°–50° XRD diffraction pattern. It can be seen from the figure that the (111) single diffraction peak at about 39° and the (200) single diffraction peak at 46° indicate that the crystal structure of the ceramic sample does not change significantly with

Fig. 1 XRD pattern of the BF–BT–*x*BGa ceramics

the increase of *x*. It is a pseudo-cubic structure. With the increase of *x*, the height of the peaks at 39° and 46° increases gradually and moves to a high angle. It indicates that the spacing of crystal structure decreases gradually. The lattice parameters of BF–BT–*x*BGa are shown in Table [1.](#page-2-0) With the increase of x , the lattice parameters (a, b, c) gradually decrease, and the intersection angle (α = 90, β = 90, γ = 120) remain unchanged. It is obvious that the doping of BGa has little effect on the lattice parameters of the sample ceramics.

Figure [2](#page-2-1) is the micromorphology of $BF-BT-xBGa$ piezoelectric ceramics at different Ga doping $(x=0, 0.01,$ 0.02, 0.03, 0.04). It can be seen from the diagram that all the surfaces of the samples become denser, and no obvious holes are formed. The grains are regular polygons, some of them are diamond shaped, and the grain boundaries are clearer. When x is 0, the average grain size is about 3.6 μ m, indicating that the doping of Ga increases the grain size of the ceramic. when the doping amount *x* is 0.01 or 0.02, the grains grow gradually, the average grain size $(x=0.02)$ reach about 7.2 µm. It is probable that with the increase of Ga content, the intergranular mobility increases during the sintering process, which promotes the movement of the

grain boundaries and increases the grain size. This indicates that Ga doping in a certain range is conducive to the growth of ceramic sample grains. When *x* exceeds 0.02, the grain size of the sample ceramics decreases. This may be because excessive Ga accumulates around the grain boundaries, hindering the grain growth.

Figure [3](#page-3-0) is a EDS diagram of BF–BT– $xBGa$ ($x=0.02$) piezoelectric ceramics. From the graph, we can see that when $x = 0.02$, the weight percentage of Ga is 0.84%, and the atomic percentage is 0.70%. We can conclude that the percentages of Ga elements in the samples are relatively low.

Figure [4a](#page-4-0) shows the d_{33} variation of Ga-doped BF–BT–*x*BGa under different compositions of different temperatures. From the figure we can see that the d_{33} value of the sample first increases and then decreases with the increase of temperature and reaches the maximum at 990 °C. When the composition point is $x=0.02$ and the sintering temperature is 990 °C, the d_{33} reaches 170 pC/N. This is mainly due to the increase of sintering temperature, which is beneficial to increase the density of ceramics and grain size. The increase of grain size is conducive to the turning of domain in polarization, which makes the anisotropy produced after polarization decrease, and the polarization effect increase, resulting in the increase of extrinsic piezoelectric properties. However, the increasing temperature easily leads to the volatilization of Bi, resulting in oxygen vacancies. At the same time, the density is reduced, the intrinsic piezoelectric properties are reduced and the total piezoelectric properties are decreased. Figure [4b](#page-4-0), c shows the relationship between the piezoelectric constant d_{33} , the electromechanical coupling coefficient k_p and the mechanical quality factor Q_m of the BF–BT–*x*BGa piezoelectric ceramic at 990 °C with the content of the doping element Ga. From the diagram, we can see that d_{33} , k_p and Q_m increase first and then decrease with the increase of *x*, and reach the maximum $d_{33} = 170$ pC/N, $k_p = 0.306$ and $Q_m = 42.877$ at $x = 0.02$.

Figure [5](#page-4-1) shows the dielectric constant and dielectric loss of Ga-doped BF–BT–*x*BGa piezoelectric ceramics at room temperature in relation to *x*. As is shown in Fig. [5](#page-4-1)a, the dielectric constant decreases firstly with the amount of BGa doping, reaches the minimum (ε_r =620.79 at *x* = 0.02), and then increases further with the increase of *x*. Dielectric loss also decreases with the increase of *x*, reaches the minimum $(tan\delta = 0.053$ at $x = 0.02$), and then increases gradually with the increase of *x*.

Figure [6](#page-5-0) shows the hysteresis loop of BF–BT–*x*BGa lead-free piezoelectric ceramics under AC electric field of 30–60 kV/cm. From the figure we can see that when the polarization electric field is 30 kV/cm, the hysteresis loop is not saturated and becomes a flat ellipse. With the increase of the polarization voltage, the hysteresis loop becomes saturated gradually, and the shape evolves into a rectangle gradually, the remanent polarization of the sample increases gradually, indicating that as the polarization electric field gradually increases to 60 kV/cm, the ferroelectric properties of the sample increase gradually. On the other hand, under the same electric field, with the increase of *x*, the ferroelectric properties of the samples increase firstly and then decrease gradually. Figure [6](#page-5-0)f is the change of P_r and *E*c of BF–BT–*x*BGa lead-free piezoelectric ceramics with the increase of x at 60 kV/cm. As is shown in the figure, the coercive field Ec first decreases with the increase of *x*, reaches the minimum $(E_c=22.059 \text{ kV/cm at } x=0.02)$. As *x* increases, E_c increases to $E_c = 24.03$ kV/cm at $x = 0.04$. The remanent polarization P_r increases with the increase of x and reaches the maximum of $(P_r = 25 \mu C/cm^2$ at $x = 0.02$), then decreases with the increase of *x*. It is indicated that the proper doping of Ga can improve the remanent polarization of BF–BT–*x*BGa sample ceramics and improve the ferroelectric properties.

Figure [7a](#page-6-0)–e shows the relative permittivity ε_r and dielectric loss tanδ of BF–BT–*x*BGa system lead-free piezoelectric ceramics at 1, 10 and 100 kHz with temperature.

Fig. 3 EDS analysis of BF–BT–*x*BGa ceramics

Fig. 4 Piezoelectric constant d_{33} , the planar electromechanical coupling factor k_p and mechanical quality factor Q_m for the BF–BT– *x*BGa ceramics

From the relationship between dielectric constant ε_r and temperature, we can see that all the samples have a hightemperature dielectric anisotropy peak at high temperature, and ferroelectric phase transit to paraelectric phase at the dielectric anomalous peak. The frequency dependence at low temperature is not very obvious. With the increase of temperature, the frequency dependence becomes more and more obvious. With the increase of temperature, the frequency dependence decreases. At the high temperature dielectric anomalous peak temperature T_m , the dielectric

Fig. 5 The relative dielectric constant ε_r and dielectric loss tand of the BF–BT–*x*BGa ceramics at room temperature

constant decreases with increasing frequency and moves to the high temperature as the frequency increases. Ferroelectrics can be divided into normal ferroelectrics, diffused ferroelectrics and relaxation ferroelectrics. Normal ferroelectrics show a sharp phase transition peak; diffuse ferroelectrics show a broad peak. The relaxor ferroelectrics show a wide peak and the Curie point moves towards the high temperature as the frequency increases. It can be seen from the figure that as *x* increases, the sample changes from normal ferroelectric to relaxed ferroelectrics gradually. On the other hand, it can be seen from the relationship between the dielectric loss and the temperature that the dielectric loss of the sample is low at low temperature and the frequency dependence is not obvious. With the increase of temperature, when it reaches a certain temperature, the dielectric loss increases sharply. The dielectric loss at low frequency is larger than that at high frequency, and the frequency dependence is obvious. Figure [7f](#page-6-0) shows the BF–BT–*x*BGa system of lead-free piezoelectric ceramic Curie temperature with composition *x* changes. We can see from the figure, the Curie temperature increases with the increase of *x* and then decreases from 434 to 404 °C. The Curie temperature T_c decline may be due to the doping of Ga in the grain boundary precipitation, resulting in the

dielectric properties of ceramics decreases and reducing the Curie temperature (T_c) .

Figure [8](#page-7-21)a shows the depolarization curves of BF–BT–*x*BGa (*x*=0, 0.01, 0.02, 0.03, 0.04) system leadfree piezoelectric ceramics. We used the method of *ex situ* measurement d_{33} to measure the depolarization temperature. As is shown in the figure, the piezoelectric constant d_{33} of samples with different compositions changes little with the increase of temperature at low temperature. When the temperature increases further, d_{33} decreases gradually with the increase of temperature. When it reaches a certain temperature, the d_{33} decreases rapidly. This temperature is the critical temperature (depolarization temperature). Figure [8](#page-7-21)b is a diagram of the variation of the depolarization temperature of the lead-free piezoelectric ceramics in the BF–BT–*x*BGa system with the change of *x*. From the diagram we can see that with the increase of *x*, the depolarization increases first and reaches the maximum of $(T_d=422$ °C at $x=0.02$). The $d₃₃$ of the same component point decreases with the increase of temperature. Piezoelectric ceramics only have polarized electric field to polarize, so that the electrical domain is aligned in the direction of polarization electric field, or most of the domains are aligned in the direction of polarized electric field. With the increase of temperature, the electric domain rearranged by the polarized electric field has a tendency to revert to the original disorderly arrangement, which makes the ceramic depolarize. The lower the stability of the domain after polarization, the lower the depolarization temperature. On the other hand, the higher the stability of the domain after polarization, the higher the depolarization temperature. It is found that the doping of Ga is beneficial to the improvement of the domain stability of the samples after polarization.

Fig. 7 The temperature dependence of ε_r and tan δ for BF–BT– *x*BGa ceramics with BiGaO₃ content at 1, 10, 100 kHz, **a** *x*=0, **b** *x*=0.01, **c** *x*=0.02, **d** *x*=0.03, **e** *x*=0.04, **f** Curie temperature (T_c) as a function of x for the BF–BT–*x*BGa ceramics

4 Conclusions

In this paper, the effects of Ga doping on the phase structure, morphology, piezoelectric ferroelectric properties and temperature stability of BF–BT–*x*BGa system lead-free piezoelectric ceramics have been studied systematically. Proper amount of Ga doping can increase the piezoelectric constant d_{33} , electromechanical coupling coefficient k_p and mechanical quality factor Q_m , reduce dielectric loss, and get *d*33=170 pC/N, *k*p=0.306, *Q*^m =42.877, tanδ=0.053 at $x = 0.02$. When $x = 0.02$ at $E = 60$ kV/cm, the ceramics

obtained the coercive field $(E_c=22.059 \text{ kV/cm})$, and the remanent polarization get $P_r = 25 \mu C/cm^2$. The temperature makes the lead-free piezoelectric ceramics of BF–BT–*x*BGa system change from normal ferroelectric to diffused ferroelectrics, and the temperature of depolarization increases firstly and then decreases, which indicates that proper Ga can improve the temperature stability of samples. These excellent properties show that BF–BT–*x*BGa lead free piezoelectric ceramic is a potential lead free piezoelectric material.

Fig. 8 a The piezoelectric constant d_{33} of BF–BT–*x*BGa lead-free piezoelectric ceramics as a function of temperature for doping different BGa contents. **b** The depolarization temperature T_d as a function of *x* for the BF–BT–*x*BGa ceramics

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