

# Room-temperature metal–insulator transition of MBE grown VO<sub>2</sub> film investigated by temperature dependent resistance and transmittance

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**Abstract** VO<sub>2</sub> films were grown on TiO<sub>2</sub> (001) substrate by a radio frequency (RF)-plasma assisted oxide molecular beam epitaxy. An excellent reversible metal-to-insulator (MIT) transition accompanied with an abrupt change in both resistivity and infrared transmittance was observed at room temperature (RT), which was much lower than the 341 K for bulk single crystal VO<sub>2</sub>. Remarkably, the MIT transition temperature ( $T_{MIT}$ ) deduced from resistivity-temperature curve was well consistent with that obtained from the temperature dependent IR transmittance. The lowered  $T_{MIT}$  was supposed to be originated from the internal stress induced by the interface lattice mismatch between VO<sub>2</sub> film and TiO<sub>2</sub> substrate, this assumption was supported by both Raman measurement and X-ray diffraction (XRD) 2theta peak shift. This achievement will potentially open up new opportunities for advanced applications of VO<sub>2</sub>-based devices where RT MIT is necessary.

## 1 Introduction

By virtue of its unique and fascinating properties, VO<sub>2</sub> was considered to be one of the key materials for a wide range

of energy-related applications [1, 2]. It undergoes an abrupt reversible phase transition, known as metal-to-insulator transition (MIT) or semiconductor-to-metal (SMT) first-order transition. At temperatures below the transition temperature ( $T_{MIT}$ ), VO<sub>2</sub> is in semiconducting state with monoclinic structure (space group P2<sub>1</sub>/c), in which the V atoms pair open an energy gap of 0.6 eV. While at temperatures above  $T_{MIT}$ , VO<sub>2</sub> is in metallic state with rutile-tetragonal structure (space group P4<sub>2</sub>/mnm), in which overlap between the Fermi level and the V<sub>3d</sub> band eliminates the band gap [3]. Most notably, this allotropic transition in crystal symmetry and electronic band structure, which can be triggered by some specific external stimuli such as temperature or voltage, was usually accompanied by an abrupt and dramatic change in physical properties. For example, the electrical resistance jumps up to four orders of magnitude and the optical transmission shows a distinct switching effect especially within the infrared wavelength region across the MIT boundary. With these unique and fascinating properties, VO<sub>2</sub> was considered to be one of the most promising energy-saving material for a wide range of energy-related applications, including smart windows for energy utilization, supercapacitors for energy storage, and thermoelectric generators for energy conversion. However, the intrinsic  $T_{MIT}$  for bulk single crystals VO<sub>2</sub> has been well accepted to be ~341 K, the specific applications will be seriously limited by the relatively rigid  $T_{MIT}$  [4]. Successfully modulating the  $T_{MIT}$  to room temperature (RT) has been a long-standing research topic. For quite a long time, intentionally impurity doping had been commonly attempted to achieve controllable  $T_{MIT}$  since both shifts towards lower and higher temperature can be realized by selecting appropriate dopant and ratios, i.e., doping with high-valent metal cations (such as W<sup>6+</sup>, Nb<sup>5+</sup>, Mo<sup>6+</sup>) were reported to lower the  $T_{MIT}$ , while doping with low-valent metal cations (such as

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Al<sup>3+</sup>, Ga<sup>3+</sup> and Cr<sup>3+</sup>) were expected to increase the  $T_{MIT}$  [4]. Nevertheless, the intentionally impurity doping generally suffered from the concomitant deterioration in change magnitude of electrical and optical characteristics as well as broadened hysteresis width, which was certainly undesirable for its practical applications as smart materials. Thus, the most critical challenge for VO<sub>2</sub> material lies in finding an effective and convenient way to reduce the  $T_{MIT}$  to room temperature while maintaining its abrupt MIT transition.

Recently, there has been a growing interest in modulating MIT behaviors of VO<sub>2</sub> films through internal stress since a relatively large  $T_{MIT}$  dependence on uniaxial stress along the *c* axis was proposed theoretically, which might be the most effective way to decrease the  $T_{MIT}$  due to its inherent advantages for many practical applications [5, 6]. Following this idea, the MIT process of the MBE grown VO<sub>2</sub> films was modulated continuously via the interfacial strain in our recent report [2]. However, there has been little investigation on the intimate correlation between electrical and optical properties across the MIT, which may shed light on the mechanism responsible for the specific MIT behavior in VO<sub>2</sub> film, for which no consensus has been reached.

In this study, VO<sub>2</sub> films with superior crystalline quality and ideal chemical stoichiometry were grown on TiO<sub>2</sub> (001) substrate by oxide molecular beam epitaxy (O-MBE). A systematical analysis were made on the crystal structure, morphology, electrical transport and transmittance within the infrared region (IR) of the VO<sub>2</sub> film. The results indicated that the RT MIT performance with narrow hysteresis width and large amplitude contrast in both IR optical transmittance and resistivity was realized, and the  $T_{MIT}$  deduced from resistivity–temperature curve was well consistent with the one obtained from the temperature dependent infrared (2 μm) transmittance. Our results will be favorable for exploit the advanced device applications based on VO<sub>2</sub> film working near RT [2, 7].

## 2 Experimental details

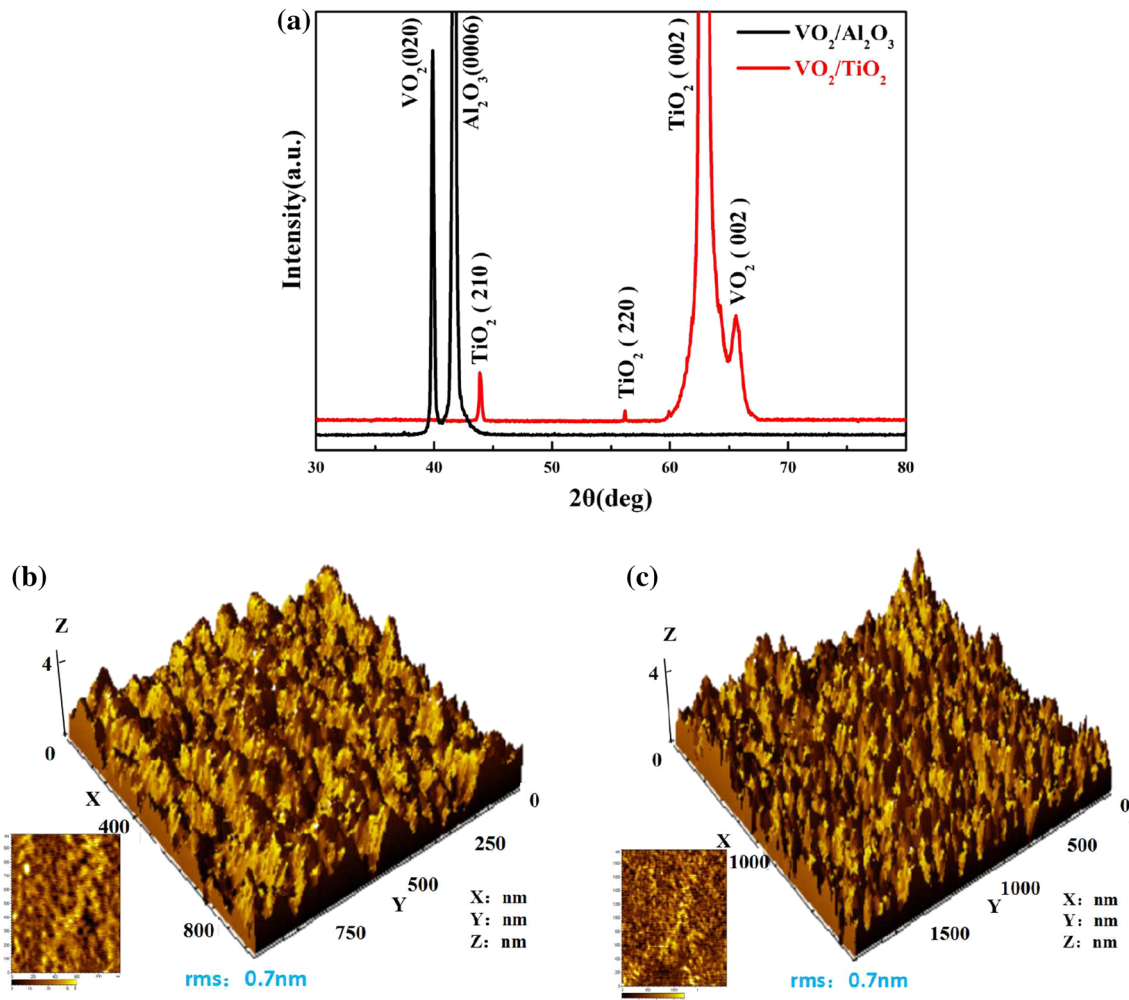
VO<sub>2</sub> films were grown on single crystal TiO<sub>2</sub> (001) substrate by a radio frequency (RF)-plasma assisted MBE with a base pressure better than  $5.5 \times 10^{-7}$  Pa. The details of the epitaxial film preparation have been reported elsewhere [2, 3]. The rutile TiO<sub>2</sub> single crystal is selected as substrate due to its stable thermal properties and the same crystal structure and space group as that of VO<sub>2</sub> film in high temperature tetragonal metallic phase. In addition, the measurements of electrical and optical properties across the MIT of VO<sub>2</sub> will not be affected considering the insulating and transparent characteristics of TiO<sub>2</sub> substrate. The TiO<sub>2</sub> (001) substrates with the dimensions of  $10 \times 10 \times 0.5$  mm<sup>3</sup>

were thoroughly cleaned to remove surface contamination and provide a atomically smooth surface for epitaxy deposition of VO<sub>2</sub> film. All the growth parameters, including the substrate temperature, the chamber pressure, the metallic vanadium evaporation rate and growth duration time, have been optimized for high crystalline quality and abrupt high magnitude MIT property. After deposition, the samples were cooled down to RT under the same oxygen pressure.

The crystallographic properties were determined by X-ray diffraction (XRD) using a LabXRD-6000 (CuK<sub>α1</sub>: $\lambda=0.154056$  nm). The surface morphology was examined by atomic force microscopy (AFM) with tapping mode using MI PicoScan 2500. The MIT properties were investigated by monitoring the change in both electrical resistivity and infrared transmittance across the MIT process, i.e., the samples were thermally cycled in the temperature ranging from 10 to 45 °C. The resistivity–temperature (*R*–*T*) behavior was investigated by measuring the electrical resistivity during heating and cooling process using the conventional electrical method by Keithley 2635 A source meter with 4-probe configuration. The temperature dependent infrared transmittance was studied by detecting the transmittance in the wavelength ranging from 1400 to 5000 nm using a FT-IR FTIR SPECTROMETER (Bruker EQUINDX 55), and the transmittance in 2 μm wavelength was recorded as a function of temperature. The films composition and valence state of V were investigated by X-ray photoelectron spectroscopy (XPS) analyses on Thermo Scientific ESCALAB250Xi system with an AlK<sub>α</sub> 1486.8 eV X-ray radiation source under a base pressure of  $3 \times 10^{-8}$  Pa. Before measurement, the sample was sputtered by Ar ion bombardment for 30 s for surface cleaning. The O1s binding energy line at 530 eV was taken as a reference for calibration and the composition was calculated from the peaks areas taking into account the sensitivity factors. Raman spectra measurements were carried out on DXR Raman Microscope with a 532 nm excitation laser source.

## 3 Results and discussion

Figure 1a shows the typical XRD patterns of the VO<sub>2</sub> film grown on single crystal TiO<sub>2</sub> (001) substrate by MBE under optimized conditions. For comparison, the XRD patterns of high quality VO<sub>2</sub> film (~60 nm) grown on *c*-plane sapphire substrate were also present in Fig. 1a, which exhibit MIT properties very close to bulk single crystal VO<sub>2</sub> material according to our previous experiments [8]. In addition to the dominant diffraction peaks corresponding to the TiO<sub>2</sub> substrate, the well-defined XRD diffraction peak corresponding to the VO<sub>2</sub> (002) was observed at 65.6° (JCPDS 76–0675) [9]. It should be noted that the growth of high quality pure phase VO<sub>2</sub>



**Fig. 1** **a** XRD patterns of  $\text{VO}_2$  film grown on  $\text{TiO}_2$  (001) substrate by MBE. For comparison, the XRD patterns of high quality  $\text{VO}_2$  film grown on c-plane sapphire substrate were also present in Fig. 1, which exhibit MIT properties very close to bulk single crys-

tal  $\text{VO}_2$  material. **b** and **c** show the surface morphology AFM images of as-grown  $\text{VO}_2$  film with the scanning area of ( $1 \times 1 \mu\text{m}^2$ ) and ( $2 \times 2 \mu\text{m}^2$ ), respectively

film was still rather challenging due to the complexity of the vanadium-oxygen system [10]. Here, however, the high quality pure monoclinic phase  $\text{VO}_2$  films with (002) preferred orientation have been achieved by MBE based on our previous optimization process [3, 8]. Comparing with the lower XRD profile, the only difference lies in the preferred orientation of  $\text{VO}_2$  films, which was no doubt resulted from the induction of  $\text{TiO}_2$  substrate. Because it was very likely that very thin  $\text{VO}_2$  films (around 30 nm for current case) below a critical thickness are constrained to have the same lattice structure as the tetragonal  $\text{TiO}_2$  (001) substrate. Here, the critical thickness refers to the thickness at which the misfit dislocations are generated to relax the film and form the equilibrium phase. It has been demonstrated that the phase transformation and strain relaxation phenomena follow totally

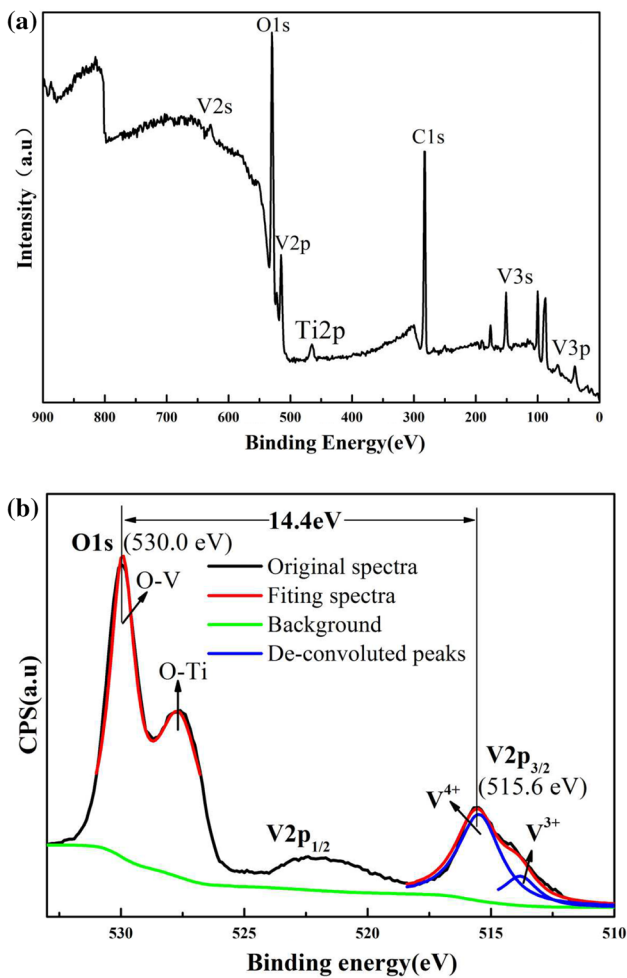
different kinetics in thin and thick layers [10]. In addition, according to Bragg equation:  $2d\sin\theta = n\lambda$ , where the  $d$  is crystal plane spacing,  $\theta$  is incident angle,  $n$  and  $\lambda$  are integers and constants, respectively. The slightly shift in  $\text{VO}_2$  (002)  $2\theta$  angle compared with the value for single crystal  $\text{VO}_2$  material (JCPDS 76-0675), indicating a smaller crystal plane spacing  $d$  and thus implying the existence a tensile stress in-plane mainly originated from the misfit between  $\text{VO}_2$  film and  $\text{TiO}_2$  substrate. Moreover, the surface morphology of the  $\text{VO}_2$  film is illustrated by AFM measurements. The AFM images with the scanned area of  $1 \times 1 \mu\text{m}^2$  and  $2 \times 2 \mu\text{m}^2$  were shown in the Fig. 1b, c, respectively. As can be seen from these images, the sample exhibit smooth and homogeneous surface morphology composed of uniformly and dense grains with spherical granular crystallites, and the same

root mean square (rms) was observed to be 0.7 nm for both images. While the uniform and dense nano columns are more clearly visible in the images of Fig. 1c.

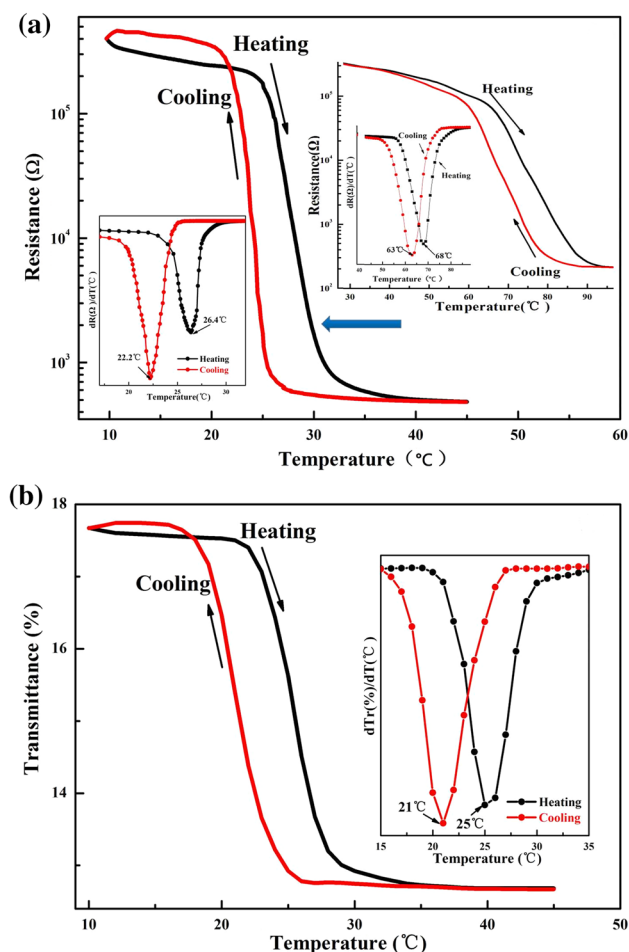
The XPS measurements were carried out to identify the valence state of V in the VO<sub>2</sub>/TiO<sub>2</sub> structure (Fig. 2). The XPS survey spectra was shown in Fig. 2a, all the elements can be well identified in the survey spectra according to previous reports in literature [11, 12]. The detected signal from C1<sub>s</sub> could be attributed to unintentionally contamination from ambient during the growth and measurement process since C element was ubiquitous element in the environment. To identify the precise valence state of V in the oxide layer, the enlarged high-resolution XPS spectra focusing on V<sub>2p<sub>3/2</sub></sub> and O1<sub>s</sub> with binding energy between 510 and 533 eV was shown in Fig. 2b. The O1<sub>s</sub> peak could be de-convoluted into two peaks, only the one at 530.0 eV

related to vanadium oxides, whereas the other peak at 527.8 eV comes from titanium oxide. Meanwhile, the typical two-peak structure (V<sub>2p<sub>3/2</sub></sub> at 515.6 eV and V<sub>2p<sub>1/2</sub></sub> at 523.3 eV) due to the spin-orbit splitting was observed for the V<sub>2p</sub> signal. According to Silversmit's method, the binding energy span between O1<sub>s</sub> and V<sub>2p<sub>3/2</sub></sub> core level could be utilized as a criterion to identify the different vanadium oxidation states, since it differs remarkably for VO<sub>x</sub> with disparate valence state [11]. The calculated binding energy span from Fig. 2b was 14.4 eV, which was in rational agreement with the value of 14.35 eV for stoichiometric VO<sub>2</sub> epitaxial layer [13], suggesting that the valence state of V in as-grown VO<sub>2</sub> film was principally composed of V<sup>4+</sup> [11]. Moreover, the slightly asymmetric feature of the V<sub>2p<sub>3/2</sub></sub> peak suggests the existence of trace amount of V<sup>3+</sup> in the VO<sub>2</sub> film. After subtracting the Shirley background (cyan) using Gaussian Lorentzian sum function by XPSPEAK 4.1 software, the V<sub>2p<sub>3/2</sub></sub> peak could be de-convoluted into two peaks, i.e., V<sup>4+</sup> peak at 515.6 eV and V<sup>3+</sup> peak at 513.8 eV, no other valence state of V was detected. The oxygen deficiency during high vacuum growth might be responsible for the formation of V<sup>3+</sup>. It should be noted that the commonly observed V<sup>5+</sup> peak at ~517.0 eV, which are usually caused by the surface oxidation as a slightly oxygen rich composition generally exist only on the surfaces (the top several nanometers), was not detected in current case. It has been confirmed that the MIT characteristics might be seriously degraded by the existence of a small amount of vanadium oxides with different stoichiometric composition rather than VO<sub>2</sub>, such as V<sub>2</sub>O<sub>3</sub> or V<sub>2</sub>O<sub>5</sub>, which may exist at the surface or grain boundaries of VO<sub>2</sub> grains, although may not be detected by the XRD measurement [14, 15]. So the ideal stoichiometry composition was imperative for reliable material properties of VO<sub>2</sub> film. Here, the ratio of oxygen to vanadium can be roughly estimated to be 1.91 from the integrated area of I<sub>V<sup>4+</sup></sub> and I<sub>V<sup>3+</sup></sub> by the empirical formula:  $(2I_{V^{4+}} + 1.5I_{V^{3+}})/(I_{V^{4+}} + I_{V^{3+}})$  [14], indicating the nearly stoichiometrical O/V ratio inside the prepared VO<sub>2</sub> film. Therefore, the super crystalline quality and ideal chemical stoichiometry of our MBE grown VO<sub>2</sub> film were clearly evidenced by XRD, AFM and XPS investigation.

The MIT performance of the as-grown VO<sub>2</sub> sample was investigated by the temperature dependent electrical resistivity and IR optical transmittance measurements across the MIT process. As shown in Fig. 3, an excellent reversible MIT phase transition accompanied with an abrupt change in both electrical resistivity and IR optical transmittance was observed, and the smooth transition profiles suggesting a first-order reversible MIT behavior for the VO<sub>2</sub> film grown on TiO<sub>2</sub> (001) substrate. The differential d(R)/d(T) versus temperature curves for heating and cooling branches were shown in the left bottom of Fig. 3a, and the T<sub>MIT</sub> can be determined from the Gaussian fit of the d(R)/dT curves



**Fig. 2** **a** XPS survey spectra with binding energy in the range of 0-900 eV for VO<sub>2</sub>/TiO<sub>2</sub> structure. **b** The enlarged high-resolution spectra with binding energy in the range of 510–533 eV and the fitting results for sample VO<sub>2</sub>/TiO<sub>2</sub> heterostructure. After subtracting the Shirley background (cyan curve), the V<sub>2p<sub>3/2</sub></sub> peak could be de-convoluted into V<sup>4+</sup> and V<sup>3+</sup> peaks

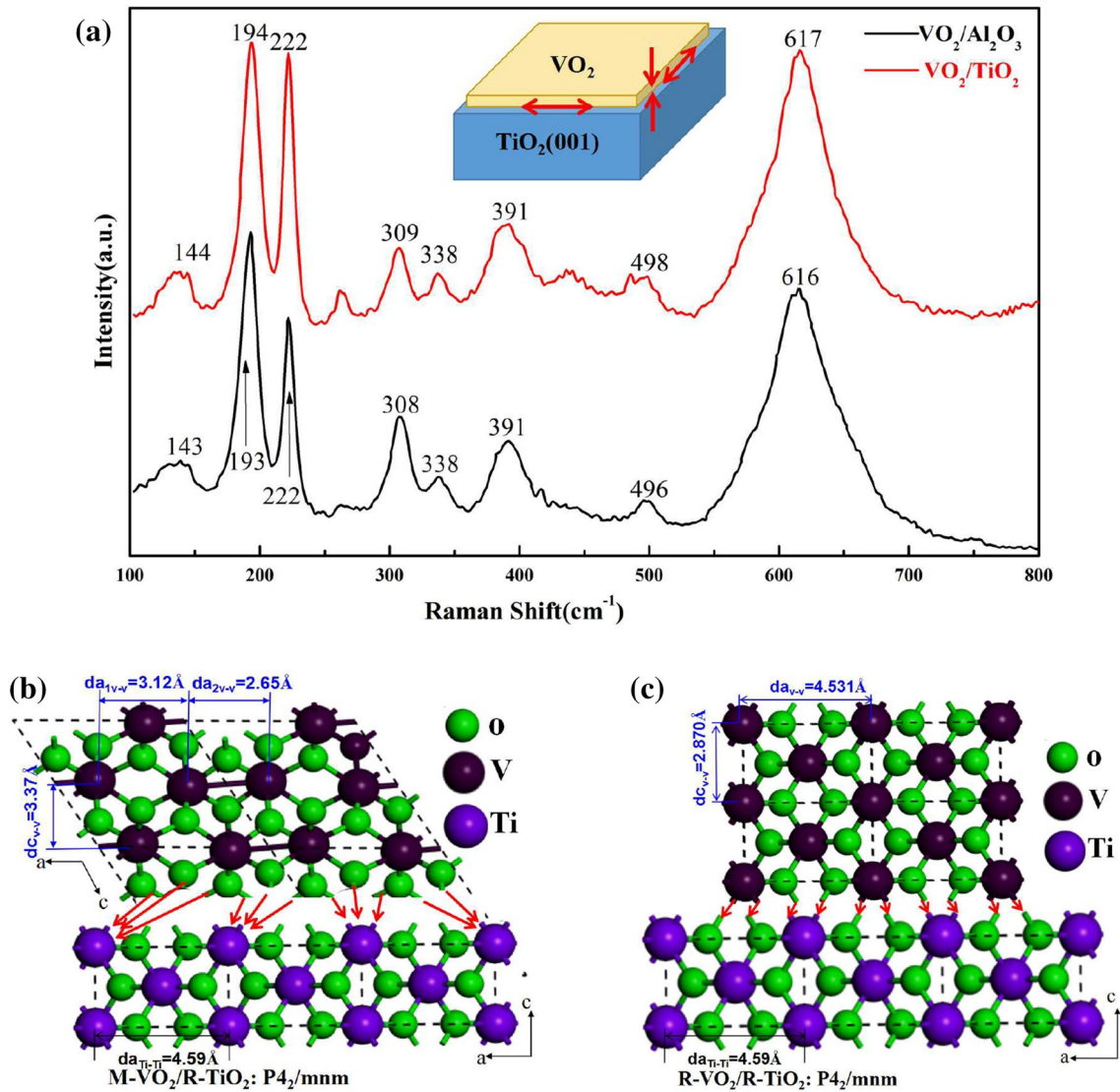


**Fig. 3** **a** Thermal hysteresis loops of resistance for the VO<sub>2</sub> film grown on TiO<sub>2</sub>(001) substrate by MBE. For comparison, the thermal hysteresis loops for the VO<sub>2</sub> film grown on sapphire substrate were shown on the *top right corner*. The differential  $d(R)/dT$  versus temperature curves were shown in the corresponding insets. **b** Thermal hysteresis loops of IR transmission at 2 μm wavelength for VO<sub>2</sub> film grown on TiO<sub>2</sub> (001) substrate. The differential  $d(Tr)/dT$  versus temperature curves were shown in the *inset* to determine the  $T_{MIT}$

(at which the value reach its extremum). For comparison, the result for the VO<sub>2</sub> film grown on c-plane sapphire substrate, which exhibit MIT properties very close to bulk single crystal VO<sub>2</sub> material according to our previous report [8], were also shown in the top right corner of Fig. 3a. The MIT properties can be characterized by the following parameters:  $T_{MIT}$  (defined as peak temperature for both heating and cooling runs),  $\Delta H$  (MIT thermal hysteresis width, defined as the difference of  $T_{MIT}$  for heating and cooling branch),  $\Delta T$  (MIT sharpness, characterized by the FWHM of the derivative curve of  $dR/dT$ -T plot), as well as  $\Delta A$  (MIT amplitude, defined as ratio of resistivity in insulator phase to that in metallic phase in the heating process). So it can be seen that the  $T_{MIT}$  show remarkable shift towards lower temperature, i.e., 68 °C for previous VO<sub>2</sub>

film on c-plane sapphire substrate while RT (26 °C) for current VO<sub>2</sub> film on TiO<sub>2</sub> (001) substrate. While other MIT parameters (including  $\Delta H$ ,  $\Delta T$ , and  $\Delta A$ ) remains almost on the same orders of magnitude as that in the top right corner of Fig. 3a, suggesting the distinct abrupt MIT performance was maintained perfectly during the decrease of  $T_{MIT}$  to RT. The relatively large MIT amplitude of exceeding three orders of magnitude and a narrow hysteresis width will be especially imperative for switching-type applications. Generally extremely demanding processing conditions are required for VO<sub>2</sub> films with lowered  $T_{MIT}$  and abrupt MIT behavior [8]. So the MIT investigation further confirming the high crystalline quality and phase purity of current VO<sub>2</sub> film, which was in well agreement with above XRD, AFM and XPS analysis. Moreover, the IR transmittance at 2 μm wavelength was recorded during heating and cooling process as shown in Fig. 3b, and the differential  $d(Tr)/dT$  ( $Tr$  stands for transmittance) versus temperature curves were shown in the insets of Fig. 3b. The  $T_{MIT}$  determined from the Gaussian fit of the  $d(Tr)/dT$  curves were shown in the insets of Fig. 3b. The  $T_{MIT}$  determined from the Gaussian fit of the  $d(Tr)/dT$  curves was 25 °C. Remarkably, the rather small hysteresis width was comparable with the high quality stoichiometric VO<sub>2</sub> film grown on c-plane sapphire substrate, which will be especially favorable for rapid response to environmental temperature as thermochromic smart window application. Therefore, the  $T_{MIT}$  and hysteresis width deduced from R–T curve (Fig. 3a) was well consistent with the one obtained from the temperature dependent infrared (2 μm) transmittance (Fig. 3b).

According to theoretical prediction, the observed remarkable decrease of  $T_{MIT}$  should be closely related with the internal stress induced by the interface lattice mismatch between TiO<sub>2</sub> substrate and VO<sub>2</sub> film [2]. Herein, the Raman spectra measurements were carried out since the line shape of Raman spectrum was extremely sensitive to the residual stress inside the films. For comparison, the Raman spectra for high quality VO<sub>2</sub> film with fully relaxed stress were also present with black line, which exhibit peaks very close to bulk single crystal VO<sub>2</sub> material according to our previous experiments [8]. As can be seen from Fig. 4, all the main peaks can be well indexed to the monoclinic VO<sub>2</sub>, especially the Raman peak at 617 cm<sup>-1</sup> (which can be regarded as the signal of M–VO<sub>2</sub>) was quite pronounced, indicating that the high quality single phase VO<sub>2</sub> films were obtained on TiO<sub>2</sub> substrate [16]. Moreover, a slight red shift in the main peaks indicating an in-plane tensile stress due to the lattice mismatch, which was also supported by the XRD 2θ peak shift as mentioned above. As expected from Poisson effect (shown in the inset of Fig. 4a, as the VO<sub>2</sub> film was stretched in plane, it would tend to be compressed out-of-plane, i.e., in the direction perpendicular to the direction of expansion. Fig. 4b, c show the schematic diagram of microscopic structure of the epitaxy



**Fig. 4** **a** The Raman spectra for VO<sub>2</sub> film grown on TiO<sub>2</sub>(001) substrate by MBE (red line). For comparison, the Raman spectra for high quality VO<sub>2</sub> film with fully relaxed stress were also present with black

line. **b** and **c** show the schematic diagram of microscopic structure of the VO<sub>2</sub>(002)/TiO<sub>2</sub>(001) interface with VO<sub>2</sub> in monoclinic ( $T < T_{MIT}$ ) and tetragonal ( $T > T_{MIT}$ ) structure, respectively. (Color figure online)

VO<sub>2</sub> (002)/TiO<sub>2</sub> (001) interface with VO<sub>2</sub> in monoclinic ( $T < T_{MIT}$ ) and tetragonal ( $T > T_{MIT}$ ) structure. And the lattice parameters of VO<sub>2</sub> and TiO<sub>2</sub>, together with the corresponding lattice mismatch between them were listed in Table 1. Since the lattice parameters of VO<sub>2</sub> are smaller than those of TiO<sub>2</sub> in both forms, the in-plane tensile stress will be generated for the epitaxially grown VO<sub>2</sub> films on TiO<sub>2</sub>(001) substrate due to the lattice mismatch, further leading to compression along c-axis according to the well-known Poisson's relation. This was in well consistent with above Raman analysis. Furthermore, as the temperature is lower than the  $T_{MIT}$  Fig. 4b, a very large mismatch up to 25% across the VO<sub>2</sub><sup>monoclinic</sup>/TiO<sub>2</sub><sup>tetragonal</sup>

interface could be obtained. While as the temperature exceeds the  $T_{MIT}$  (Fig. 4c), only a tiny lattice mismatch exist across the VO<sub>2</sub><sup>tetragonal</sup>/TiO<sub>2</sub><sup>tetragonal</sup> interface due to the similar tetragonal structure of rutile phase VO<sub>2</sub> and TiO<sub>2</sub>. The reduced V<sup>4+</sup>-V<sup>4+</sup> distance in the crystal structure results in direct overlapping of d orbitals, which stabilized the metallic phase of the rutile structure. In other words, the residual compressive stresses can counteract the transition energy barriers and stimulate the occurrence of the MIT sharply. In addition, the trace amount of V<sup>3+</sup> caused by the oxygen deficiency during high vacuum growth might also partially contributes to the decrease of  $T_{MIT}$  [2].

**Table 1** The lattice parameters of VO<sub>2</sub> and TiO<sub>2</sub> in monoclinic and rutile form, together with the corresponding lattice mismatch between them

Material	Space group	Lattice parameter			Lattice mismatch				
		a (Å)	b (Å)	c (Å)	β (°C)	a (%)	b (%)	c (%)	β (%)
VO <sub>2</sub>	Monoclinic P <sub>21</sub> /c	5.74	4.52	5.38	122.6	25.05	1.525	81.75	36.22
VO <sub>2</sub>	Rutile P4 <sub>2</sub> / mmm	4.53	4.53	2.87	90	1.307	1.307	3.041	0
TiO <sub>2</sub>		4.59	4.59	2.96	90				

## 4 Conclusion

In summary, the VO<sub>2</sub> film with superior crystalline quality and ideal chemical stoichiometry was grown on TiO<sub>2</sub> (001) substrates. An excellent RT reversible MIT performance was observed with an abrupt change in both resistivity and infrared optical transmittance. Remarkably, the T<sub>MIT</sub> deduced from R–T curve was well consistent with the one obtained from the temperature dependent infrared (2 μm) transmittance. Based on the Raman measurement and microstructure model, the RT MIT properties was mainly attributed to the residual compressive stresses imposed by the TiO<sub>2</sub> (001) substrates. Our achievement will potentially open new opportunities for advanced applications of VO<sub>2</sub>-based devices.

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## Compliance with Ethical Standards

**Conflict of interest** The authors declare that they have no potential conflict of interest to this work.

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