

# Effect of NdAlO<sub>3</sub> on microstructure, dielectric properties and temperature-stable mechanism of (Sr, Ca, Nd)TiO<sub>3</sub> ceramics at microwave frequency

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**Abstract** Microwave dielectric ceramics  $(1-x)(\text{Sr}_{0.3}\text{Ca}_{0.427}\text{Nd}_{0.182})\text{TiO}_3-x\text{NdAlO}_3$  (abbreviated as SCNTA $x$  hereafter,  $0.1 \leq x \leq 0.4$ ) were prepared by conventional mixed oxide route, and their phase composition, microstructure and microwave dielectric properties were investigated as a function of the  $x$  value and sintering temperatures. A single tilted orthorhombic perovskite structure in space group Pnma was refined in the studied composition range. For microwave dielectric properties, the decreasing relative permittivity was strongly affected by the ionic polarizability of Nd<sup>3+</sup> and Al<sup>3+</sup> in SCNTA $x$  ceramic systems. Also, the quality factor of SCNTA $x$  solid solution had strongly depended on apparent densities and average grain sizes. As expected, the promising ceramic of SCNTA $x$  ( $x = 0.25$ ) sintered at 1520 °C for 4 h was found to possess good microwave dielectric properties: a relative permittivity ( $\epsilon_r$ ) of 55.6, a quality factor ( $Q \times f$ ) of 25,600 GHz (at 4.249 GHz) and a temperature coefficient of resonant frequency ( $\tau_f$ ) of 6.7 ppm/°C. Especially, the  $\tau_f$  values of SCNTA $x$  ceramics were not strongly depended on tolerance factor ( $t$ ) with increasing of the NdAlO<sub>3</sub> content, while these  $\tau_f$  values were essentially correlated with the B-site bond valence and octahedral tilting.

Wherein, either decreasing the B-site bond valence or increasing the octahedral tiltings ( $\theta$  and  $\varphi$ ) led to a decrease in  $\tau_f$  value for the present ceramics.

## 1 Introduction

With the rapid development of mobile and satellite communication, microwave electronic devices are required to be developed and fabricated for miniaturization and integration. Particularly, as the typical representative of perovskite paraelectric titanates (ABO<sub>3</sub> structure), CaTiO<sub>3</sub> has been attracted much attention as microwave dielectric because of its high relative permittivity ( $\sim 170$ ) [1, 2]. However, its dielectric loss ( $1/Q$ ) is rather high, compared with the complex perovskites. Thus, much work has been performed to control the microwave dielectric loss in CaTiO<sub>3</sub> through A-site substitution, such as  $\text{Ca}_{1-x}\text{Ln}_{2x/3}\text{TiO}_3$  (Ln = La, Nd and Sm,  $Q \times f \sim 8000\text{--}12,000$  GHz) [3–5]. But their temperature coefficient of resonant frequency ( $\tau_f \sim 190\text{--}440$  ppm/°C) are too large to apply those microwave frequency devices, so that much attention has to be paid to improve their temperature stability by mixing CaTiO<sub>3</sub>-based materials with those materials (ABO<sub>3</sub>) of negative  $\tau_f$  value. Among these candidates, CaTiO<sub>3</sub>-NdGaO<sub>3</sub> and CaTiO<sub>3</sub>-LaAlO<sub>3</sub> ceramic systems are attracting great interests due to their better microwave dielectric properties ( $\epsilon_r \sim 35\text{--}45$ ,  $Q \times f \geq 33,000$  GHz and  $\tau_f \approx 0$  ppm/°C) [6–8]. However, it is well known that a higher  $\epsilon_r$  can facilitate circuit miniaturization because the wavelength inside the dielectric resonators (DRs) is inversely proportional to the square root of its  $\epsilon_r$ . Although the good temperature stability and the lower dielectric loss can be obtained in those ceramics, the relative permittivity ( $\epsilon_r$ ) is not still high enough and can not meet the miniaturization for the high-

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demand DRs. Therefore, it should be an important issue to further enhance the  $\epsilon_r$  through the preparation of new dielectric ceramic materials.

SrTiO<sub>3</sub> is a well-known incipient ferroelectric with cubic perovskite structure at room temperature, and it possesses a higher relative permittivity ( $\epsilon_r \sim 270$ ), a modest quality factor ( $Q \times f \sim 3,000$  GHz) and a larger temperature coefficient of resonant frequency ( $\tau_f \sim 1650$  ppm/°C) [2, 9], compared with that of CaTiO<sub>3</sub>. High relative permittivity makes it possible to noticeably miniaturize passive microwave components. Therefore, for practical application, some researchers have become increasingly interested in microwave dielectric properties of the SrTiO<sub>3</sub> series materials [10, 11]. In our previous work [12], the (Sr, Ca, Nd)TiO<sub>3</sub> ceramic systems were synthesized to investigate their structure and microwave dielectric properties, wherein the Sr<sup>2+</sup> substitutions for (Ca<sub>0.61</sub>Nd<sub>0.26</sub>)<sup>2+</sup> in (1 - x)SrTiO<sub>3-x</sub>Ca<sub>0.61</sub>Nd<sub>0.26</sub>TiO<sub>3</sub> (SCNTx) ceramics could form solid solution. Especially, the (Sr<sub>0.3</sub>Ca<sub>0.427</sub>Nd<sub>0.182</sub>)TiO<sub>3</sub> (SCNT) specimen sintered at 1400 °C for 4 h was an orthorhombic distorted perovskite-type microwave dielectric material with a moderate  $\epsilon_r \sim 145$  and a very high  $Q \times f \sim 7900$  GHz. In addition, a relatively lower positive  $\tau_f \sim 415.6$  ppm/°C, compared with that of SrTiO<sub>3</sub> and/or CaTiO<sub>3</sub>, can also be obtained for this sample. This means that the SCNTx ceramics can ensure the higher  $\epsilon_r$ , at the same time achieve a temperature-stable material by forming solid solutions with the microwave ceramics of high negative  $\tau_f$  values. On the other hand, NdAlO<sub>3</sub> belongs to a distorted perovskite structure at room temperature with low  $\epsilon_r \sim 22.3$ , but high  $Q \times f \sim 58,000$  and negative  $\tau_f \sim -33$  ppm/°C [13]. In that way, the higher positive  $\tau_f$  of SCNTx can be suppressed to low or near-zero  $\tau_f$  value with the addition of NdAlO<sub>3</sub>. Therefore, the aim of this work was to obtain a relative high  $\epsilon_r$  ( $\geq 55$ ), good  $Q \times f$  ( $\geq 25,000$  GHz) and near-zero  $\tau_f$  by synthesizing (1 - x)(Sr<sub>0.3</sub>Ca<sub>0.427</sub>Nd<sub>0.182</sub>)TiO<sub>3-x</sub>NdAlO<sub>3</sub> (referred to hereafter as SCNTAx hereafter, 0.1  $\leq x \leq$  0.4) ceramic systems. Moreover, the phase composition, microstructure, and sintering behavior of the SCNTAx ceramics were also investigated in composition range with  $x = 0.1$ –0.4. Specially, regulatory mechanism of the  $\tau_f$  values with the change of tolerance factor ( $t$ ), B-site bond valence ( $\bar{V}_B$ ) and octahedral tiltings ( $\theta$  and  $\varphi$ ) were discussed in details as well.

## 2 Experimental

Specimens with  $x = 0.1, 0.2, 0.25, 0.3$  and  $0.4$  of (1 - x)(Sr<sub>0.3</sub>Ca<sub>0.427</sub>Nd<sub>0.182</sub>)TiO<sub>3-x</sub>NdAlO<sub>3</sub> (SCNTAx) ceramics were prepared by the conventional solid state

reaction method using high-purity reagent-grade raw materials (more than 99.9 %) of CaCO<sub>3</sub>, SrCO<sub>3</sub>, Nd<sub>2</sub>O<sub>3</sub>, Al<sub>2</sub>O<sub>3</sub> and TiO<sub>2</sub>. These materials were mixed according to the desired compositions and ground in distilled water for 24 h with ZrO<sub>2</sub> balls. The mixtures were dried and calcined at 1150 °C ( $x = 0.1, 0.2, 0.25, 0.3$ ) and 1250 °C ( $x = 0.4$ ) for 2 h. The calcined powders were then sieved by using a 100 and 200 mesh screen. The sieved powders with 5 wt% PVA as a binder were pressed into pellets 11.5 mm in diameter and 5.5–6.0 mm in thickness under a pressure of 150 MPa. Additionally, these pellets were preheated at 650 °C for 2 h to expel the binder, then sintered at 1450–1620 °C for 4 h in air.

The crystalline structure of the prepared samples was analyzed by an X-ray diffractometer (XRD, D8-Advance, Bruker, Germany). After spraying the surfaces, the microstructures of these sintered specimens were examined by scanning electron microscopy (SEM, S-4800, Hitachi, Japan) coupled with energy dispersive spectroscopy (EDS). Additionally, the bulk density of the sintered ceramics was evaluated from the dimensions and weight. The relative permittivity ( $\epsilon_r$ ) was measured using the Hakki-Coleman dielectric resonator method with TE<sub>01 $\delta$</sub>  resonant modes (Agilent E5230C) [14], and the quality factor was characterized by  $Q \times f$  ( $Q = 1/\text{dielectric loss}$ ,  $f = \text{resonant frequency}$ ) value at  $3.8 = 6.1$  GHz, which were measured through the parallel plate method, as modified and improved by Courtney [15]. In addition, the temperature coefficient of the resonant frequency ( $\tau_f$ ) was measured by using a temperature chamber. Accordingly, these  $\tau_f$  values were calculated by the following equation (Eq. 1) [16]:

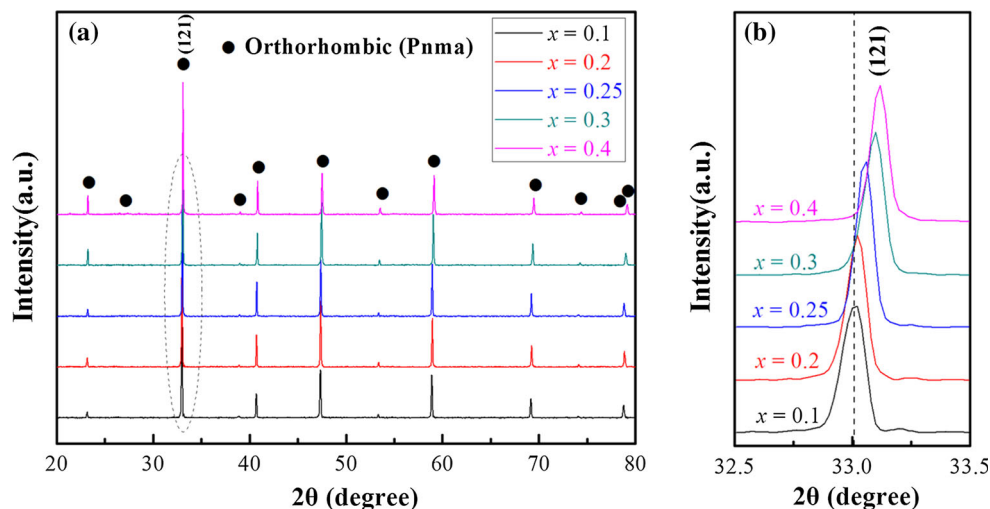
$$\tau_f = \frac{\Delta f_0}{f_0 \Delta T} = \frac{f_{75} - f_{25}}{f_{25} \times 50} \quad (1)$$

where  $f_{75}$  and  $f_{25}$  represent the resonant frequencies at 75 and 25 °C, respectively.

## 3 Results and discussions

The room-temperature X-ray diffraction (XRD) patterns recorded from the (1 - x)(Sr<sub>0.3</sub>Ca<sub>0.427</sub>Nd<sub>0.182</sub>)TiO<sub>3-x</sub>NdAlO<sub>3</sub> (SCNTAx, 0.1  $\leq x \leq$  0.4) ceramics sintered at 1550 °C for 4 h are illustrated in Fig. 1a. As shown, all peaks can be indexed based on CaTiO<sub>3</sub> (PDF#42-0423) with an orthorhombic Pbnm structure and no second phase is detected for all the specimens, indicating the formation of SCNTAx solid solution, even in the sample with  $x = 0.4$ . In addition, as shown in the Fig. 1b, the main diffraction peak (121) of the orthorhombic phase slightly shifts to a higher angle with increasing NdAlO<sub>3</sub> content. This is due to the fact that the ionic radii of Nd<sup>3+</sup> (1.27 Å, CN = 12) is slightly larger than that of (Sr<sub>0.3</sub>Ca<sub>0.427</sub>Nd<sub>0.182</sub>)<sup>2+</sup> (1.235 Å,

**Fig. 1** **a** X-ray diffraction patterns of SCNTA $x$  ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1550 °C for 4 h and **b** the magnified figure of (121) peak



CN = 12) at A-site, but the incorporation of the much smaller  $\text{Al}^{3+}$  (0.535 Å, CN = 6) in place of  $\text{Ti}^{4+}$  (0.605 Å, CN = 6) at B-site [17]. Thus, with an increase in the substitution of smaller  $\text{Al}^{3+}$ , all diffraction peaks gradually shift to the higher angles and the lattice parameters and unit cell volumes (Table 1) of the SCNTA $x$  ceramics gradually decrease as well, which obeys the general known Vegard's law.

Figure 2 shows the back-scattered electron images (BEIs) of as-sintered surfaces of the SCNTA $x$  ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1550 °C for 4 h. As shown in Fig. 2a–e, the color of all the grains is basically the same, which suggests that a single phase exists in the SCNTA $x$  ceramics within the studied composition range, and this is a well agreement with the results of the XRD patterns (see Fig. 1a). In addition, a small amount of porosity and very small grains of 1–4 μm can be observed in Fig. 2a–c, and the small average grain size will result in a large amount of grain boundary, causing the low relative densities of the ceramics. Moreover, the average grain sizes increase with increasing  $\text{NdAlO}_3$  content (see Fig. 2d, e), wherein a well-developed and uniform microstructure can be achieved by sintering the specimen with  $x = 0.3$  at 1550 °C, it also means that 1550 °C may be the

best densification sintering temperature for the SCNTA $x$  ( $x = 0.3$ ) sample.

Generally, the tolerance factor ( $t$ ) is believed to be closely related with the temperature coefficient of resonant frequency ( $\tau_f$ ) in  $\text{ABO}_3$  perovskite solid solutions, and the  $\tau_f$  value increases linearly with an increasing of the tolerance factor if  $t < 1$  [18, 19]. Wherein, the tolerance factor has been proposed and also defined as [20, 21]:

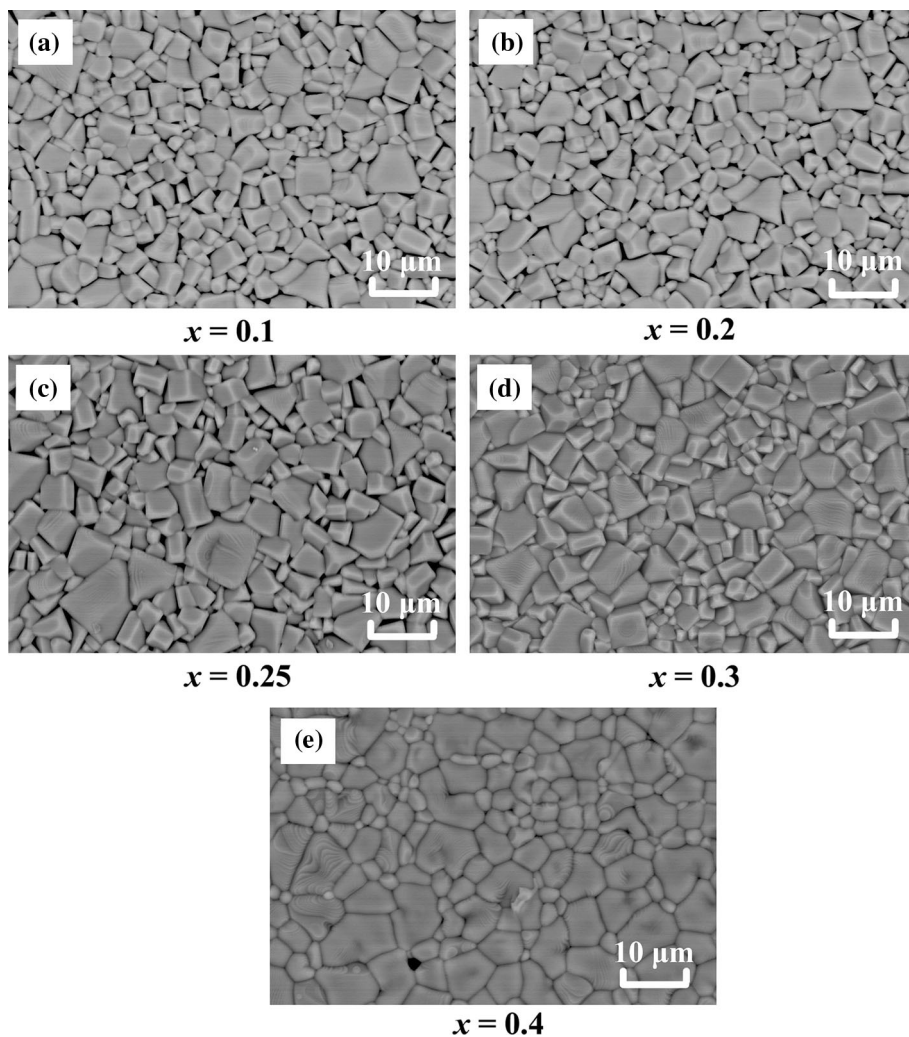
$$t = \frac{R_O + R_A}{\sqrt{2}(R_O + R_B)}, \quad (2)$$

where  $R_A$  is the radius of the A-site ion,  $R_B$  is the radius of B-site ion, and  $R_O$  is the radius of oxygen ion. However, our previous work (not be published) showed that the  $\tau_f$  value in  $(1-x)(\text{Sr}_{0.3}\text{Ca}_{0.427}\text{Nd}_{0.182})\text{TiO}_{3-x}\text{SmAlO}_3$  (SCNSTA $x$ ,  $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1480 °C for 4 h gradually decreased, while tolerance factor ( $t$ ) of the SCNSTA $x$  ceramics increased from 0.9307 to 0.9409 with increasing  $x$  value. This is due to the fact that the similarity of ionic radii between  $\text{Sm}^{3+}$  (1.24 Å, CN = 12) and  $(\text{Sr}_{0.3}\text{Ca}_{0.427}\text{Nd}_{0.182})^{2+}$  (1.235 Å, CN = 12) at A-site is not enough to change the bending of B–O–B bond in adjacent  $\text{BO}_6$  octahedra, as shown in Fig. 3a. When the size of A-site substitution is small enough, the tolerance

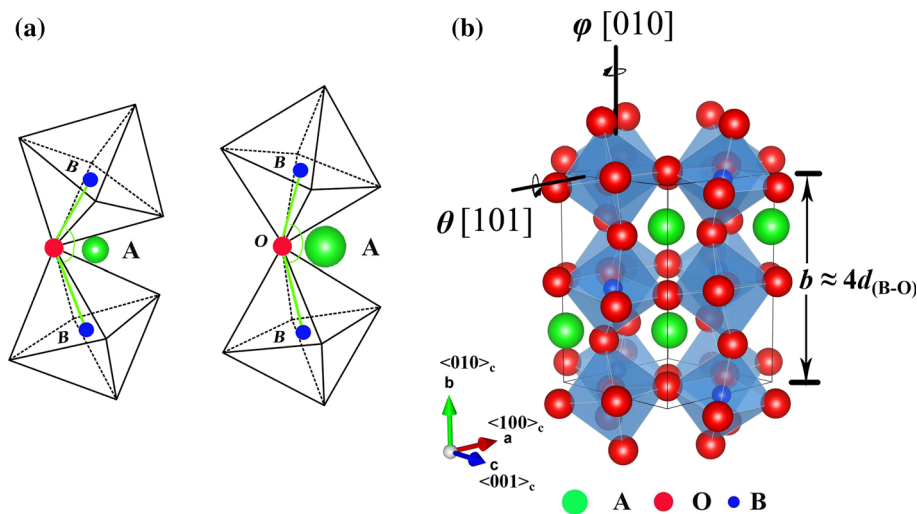
**Table 1** B-site bond valence and lattice parameters of the SCNTA $x$  ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1550 °C for 4 h

Composition	$a$ (Å)	$b$ (Å)	$c$ (Å)	Cell volume (Å <sup>3</sup> )	$\bar{R}_{\text{B-O}}$ (Å)	$d_{\text{B-O}}$ (Å)	$b$ (Å)	$\nu_{\text{B-O}}$	$\bar{V}_{\text{B}}$
$x = 0.1$	5.4274	7.6669	5.4234	225.67	1.7986	1.9167	0.37	0.7267	4.3602
$x = 0.2$	5.4245	7.6658	5.4202	225.39	1.7822	1.9164	0.37	0.6958	4.1748
$x = 0.25$	5.4236	7.6618	5.4123	224.91	1.7740	1.9155	0.37	0.6822	4.0932
$x = 0.3$	5.4175	7.6538	5.4082	224.25	1.7658	1.9135	0.37	0.6711	4.0266
$x = 0.4$	5.4092	7.6521	5.4120	224.01	1.7494	1.9130	0.37	0.6426	3.8556

**Fig. 2** Back-scattered electron images of as-sintered surfaces of the SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1550 °C for 4 h



**Fig. 3 a** The relationship between large and/or small A-site ions and the bending of B–O–B bond in adjacent oxygen octahedra, **b** the distorted orthorhombic perovskite with a space group of Pnma derived through the octahedral tilting  $\theta$  and  $\varphi$



factor ( $t$ ) can effectively reflect the structural stability in  $ABO_3$  structure. Moreover, for SCNSTAx ceramic systems, the increasing  $t$  is attributed to the incorporation of

smaller  $Al^{3+}$  ( $0.535 \text{ \AA}$ , CN = 6) in place of  $Ti^{4+}$  ( $0.605 \text{ \AA}$ , CN = 6) at B-site, but the substitution of  $Al^{3+}$  for B-site just results in an increase in the intensity of oxygen and

B-site ions at single  $\text{BO}_6$  octahedra, which hardly impacts on the thermal stability of ceramic systems. An increase in tolerance factor ( $t$ ) merely indicates that the volume of the  $\text{BO}_6$  octahedron is better matched to the size of the  $\text{AO}_{12}$  polyhedron. In the present work, for the SCNTAx ceramics, the Eq. (2) becomes:

$$t = \frac{R_{\text{O}^{2-}} + 0.427(1-x)R_{\text{Ca}^{2+}} + 0.182(1-x)R_{\text{Nd}^{3+}} + 0.3(1-x)R_{\text{Sr}^{2+}} + xR_{\text{Nd}^{3+}}}{\sqrt{2}[R_{\text{O}^{2-}} + (1-x)R_{\text{Ti}^{4+}} + xR_{\text{Al}^{3+}}]} \quad (3)$$

By calculating, the tolerance factor ( $t$ ) gradually increases from 0.9317 at  $x = 0.1$  to 0.9452 at  $x = 0.4$ , while the  $\tau_f$  value also gradually declines according to the analysis below. These results are similar with that of SCNTAx ceramics, and it also illustrates that the  $\tau_f$  values are not strongly depended on tolerance factor ( $t$ ) with increasing  $x$  value in SCNTAx ceramic systems. Thus, we consider that, without the phase transitions [22, 23], the  $\tau_f$  value is closely related with the the size of A-site ions substitution and the bending of B–O–B bonds. It can be concluded that the increasement of the bending in B–O–B bonds means the higher stability and lower symmetry in  $\text{ABO}_3$  perovskite structure. Also, it also indicates that the rattling effect in single  $\text{BO}_6$  octahedra decreases, giving rise to a decrease in  $\tau_f$  value.

On the other hand, the temperature coefficient of the resonant frequency ( $\tau_f$ ) strongly depends on B-site bond valence or octahedral tilting for  $\text{ABO}_3$  perovskite structure [24–27]. The bond valences of B-site ( $\bar{V}_B$ ) for the SCNTAx ceramics are calculated using the following equations [24, 28]:

$$\bar{V}_i = \sum v_{ij} \quad (4)$$

$$v_{ij} = \exp\left[\frac{\bar{R}_{ij} - d_{ij}}{b}\right], \quad (5)$$

where  $\bar{R}_{ij}$  is the average bond valence parameter,  $d_{ij}$  is the length of a bond between atom  $i$  and  $j$ , and  $b$  is a universal constant (0.37 Å). In this work, the SCNTAx is an distorted orthorhombic perovskite with a space group of Pnma on a  $\sqrt{2} \times \sqrt{2} \times 2$  cell in the investigational range with  $x = 0.1$ – $0.4$ , and its  $d_{\text{B-O}}$  is approximately equal to  $\frac{1}{4}b$  (lattice parameter) in an unit cell volume [29], as shown in Fig. 3b. Additionally, replacing  $\text{Ti}^{4+}$  with  $\text{Al}^{3+}$  in the SCNTAx ceramics will lead to a variation in the atomic interactions, which results in a change in bond valence of this material. Table 1 summarizes the B-site bond valences calculated using Eqs. (4), (5) and the relevant lattice

parameters for the present SCNTAx ceramics. Meanwhile, Fig. 4 displays the  $\tau_f$  value of the SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1550 °C for 4 h as a function of B-site bond valences. As can be seen from Fig. 4, the  $\tau_f$  value gradually shifts to near zero or negative direction with a decrease in  $\bar{V}_B$  as  $x$  value increases from 0.1 to 0.4.

According to the bond valence theory [28], the shorter the bond length is, the stronger the bond energy become, which also explains that the rattling effect of B-site cation is decreased, thus causing  $\tau_f$  value to decrease. Moreover, the tilting of octahedra is another important parameters related to the  $\tau_f$  value of  $\text{ABO}_3$  perovskites that significantly affect the stability of crystal structure. The orthorhombic Pnma in these SCNTAx systems gives a tilt system of  $a^-b^+a^-$  by using Glazer's notation [21], meaning that there are two approximately equal anti-phase tilts about the  $a$  axis and an in-phase tilt about the  $b$  axis. It can also be conceived as the result of tiltings ( $\theta$  and  $\varphi$ ) about the [101] and [010] orientations of the regular octahedra (see Fig. 3b). Additionally, the unit cell parameters are related to the octahedral tiltings and bond valence parameters through the following relationships [30]:

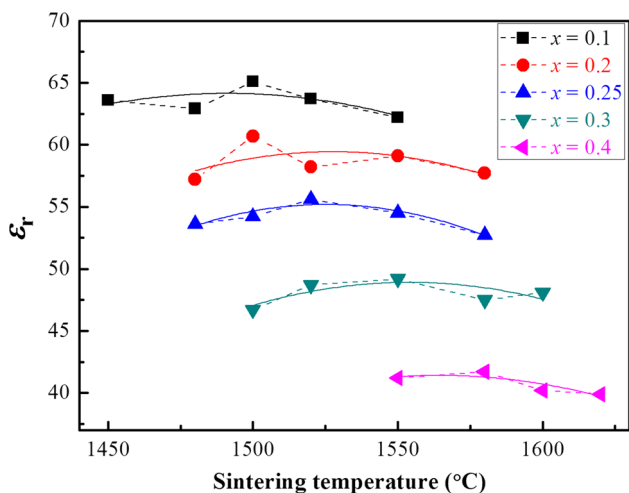
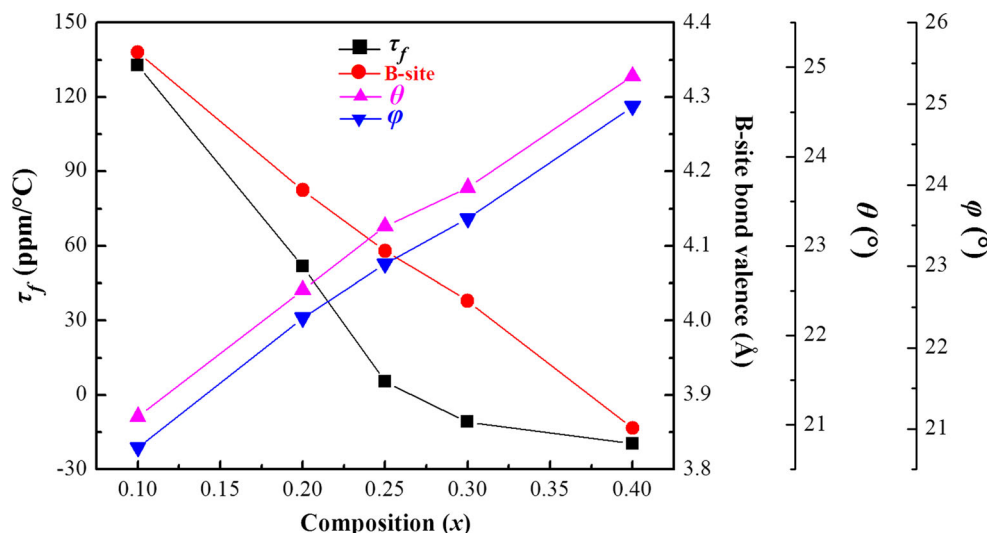
$$\cos \theta = 1 - \frac{|2\sqrt{2} \times \bar{R}_{\text{B-O}} - a|}{2\sqrt{2} \times \bar{R}_{\text{B-O}}} \quad (6)$$

$$\cos \varphi = 1 - \frac{|4 \times \bar{R}_{\text{B-O}} - b|}{4 \times \bar{R}_{\text{B-O}}}, \quad (7)$$

where  $\bar{R}_{ij}$  is the same as that of Eq. (2). Figure 4 also shows that the  $\text{NdAlO}_3$ -addition increases octahedral tiltings ( $\theta$  and  $\varphi$ ) from 21.09° to 24.90°, and from 20.77° to 24.97°, affecting a reduction in  $\tau_f$  value of the SCNTAx ceramics from 132.8 to  $-19.5$  ppm/°C. These results can be attributed to the fact that the increasing octahedral tilting results in a corresponding increase in the restoring force for recovering the tilting, and this stronger restoring force can inhibit the increase of heat effect and suppress the ionic rattling of A- and B-sites in a certain degree [27, 31], then the  $\tau_f$  value accordingly decreases. Therefore, the temperature coefficient of the resonant frequency ( $\tau_f$ ) can be effectively evaluated by the B-site bond valences and octahedral tilting for these SCNTAx ( $0.1 \leq x \leq 0.4$ ) perovskite ceramics.

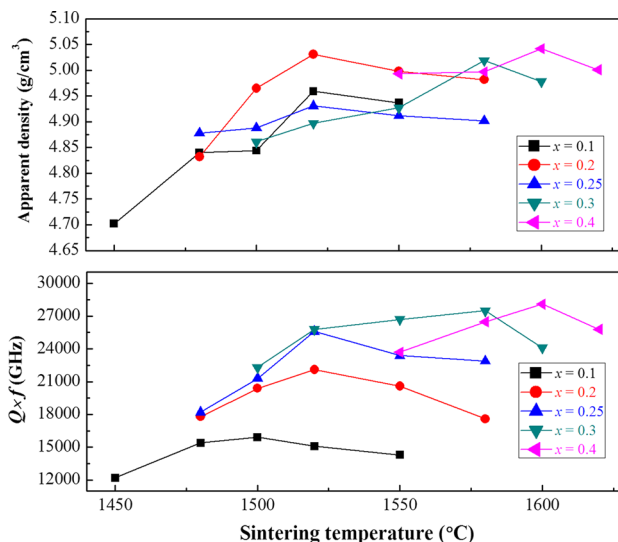
Figure 5 illustrates the relative permittivity ( $\epsilon_r$ ) of SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at various

**Fig. 4**  $\tau_f$  value of the SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1550 °C for 4 h as a function of B-site bond valence and octahedral tiltings ( $\theta$  and  $\varphi$ )



**Fig. 5**  $\epsilon_r$  of the SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics as functions of various sintering temperatures for 4 h

sintering temperatures (1450–1620 °C) for 4 h with different  $x$  value. It can be found that the  $\epsilon_r$  decreases steadily from 62.2 to 41.2 with the increase of  $\text{NdAlO}_3$  content when sintering temperature is 1550 °C. We approve that the relative permittivity is strongly dependent on the ionic polarizability for  $\text{ABO}_3$  perovskite solid solutions materials [25, 32]. For the present SCNTAx ( $0.1 \leq x \leq 0.4$ ) solid solutions, although  $(\text{Sr}_{0.3}\text{Ca}_{0.427}\text{Nd}_{0.182})^{2+}$  has a similar ionic polarizability ( $3.53 \text{ \AA}^3$ ) than that of  $\text{Nd}^{3+}$  ( $5.01 \text{ \AA}^3$ ), the ionic polarizability of  $\text{Al}^{3+}$  ( $0.79 \text{ \AA}^3$ ), however, is much lower than that of  $\text{Ti}^{4+}$  ( $2.93 \text{ \AA}^3$ ). Therefore, the  $\epsilon_r$  decreases gradually with an increasing of  $x$  value on account of the fact that the integral dielectric polarizability of  $\text{NdAlO}_3$  ( $11.83 \text{ \AA}^3$ ) is smaller than that of SCNT ( $12.49 \text{ \AA}^3$ ) [33]. In addition, as shown, the  $\epsilon_r$  are not affected greatly by sintering temperatures in the same



**Fig. 6** **a** Apparent density and **b** room-temperature  $Q \times f$  value of the SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics as a function of sintering temperature for 4 h

$x$  value, and a maximum  $\epsilon_r$  value of  $\sim 55.6$  is obtained for the specimen using SCNTAx ( $x = 0.25$ ) sintered at 1520 °C for 4 h. Moreover, these ranges in sintering temperature of the SCNTAx ceramics increase from 1450–1550 °C at  $x = 0.1$  to 1550–1620 °C at  $x = 0.4$  because  $\text{NdAlO}_3$  requires a higher sintering temperature of  $\sim 1650$  °C [13].

Figure 6a shows the apparent density of SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at different sintering temperature for 4 h. It can be observed that, for the same  $x$  value, the specimens' apparent densities initially increase with increasing sintering temperature and reach a maximum, then slightly decrease at a further increasing sintering temperature. On the other hand, the room-temperature  $Q \times f$  values of the

SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics sintered at 1450–1620 °C for 4 h are demonstrated in Fig. 6b. As reported, the  $Q \times f$  value strongly depends on the extrinsic factors, such as second phase, density, and grain size [34, 35]. In these SCNTAx ceramic systems, the effects of the secondary phase can be neglected, because the secondary phase in SCNTAx solid solution is not detected through the investigated composition range with  $x = 0.1$ –0.4, as confirmed in Fig. 1. As expected, an ordinary behavior is found that the  $Q \times f$  value is inside the range exhibited by the end members SCNT with a  $Q \times f$  value of 7900 GHz and NdAlO<sub>3</sub> with a  $Q \times f$  value of 58,000 GHz. In addition, the variation in  $Q \times f$  value of SCNTAx (except  $x = 0.1$ ) ceramics is basically consistent to that of the apparent density with the increasing sintering temperature, and the higher density means the lower porosity, producing a decrease in dielectric loss of the ceramic systems [36]. As a result, the  $Q \times f$  value is also mainly controlled by the microstructure in the present ceramics. Although the  $Q \times f$  value increases gradually with an increase in  $x$  value due to the higher  $Q \times f$  value of NdAlO<sub>3</sub>, a slow increase in  $Q \times f$  values for the samples with  $x = 0.3$  and 0.4 is shown in Fig. 6b. It can be explained by that large grains in dielectric ceramics usually contain more defects, such as dislocations, planar defects, and even the dislocation networks, which can lead to a decrease in the  $Q \times f$  value to some extent [31, 37]. Therefore, the increasing extent of  $Q \times f$  value in SCNTAx ( $0.1 \leq x \leq 0.4$ ) ceramics is decreased due to the increasing average grain sizes, as shown in Fig. 2a–e.

## 4 Conclusions

The microstructure, phase composition and microwave dielectric properties of  $(1-x)(\text{Sr}_{0.3}\text{Ca}_{0.427}\text{Nd}_{0.182})\text{TiO}_3-x\text{NdAlO}_3$  (SCNTAx,  $0.1 \leq x \leq 0.4$ ) ceramics were investigated in this work. The XRD results showed that the single orthorhombic perovskite solid solution could be formed in the composition range of  $x = 0.1$ –0.4. Also, a temperature-stable mechanism of these SCNTAx ceramics at microwave frequency were also explained through the tolerance factor ( $t$ ), B-site bond valence and octahedral tiltings. Wherein, the  $\tau_f$  values of SCNTAx ceramics were not strongly depended on tolerance factor ( $t < 1$ ) with increasing of the NdAlO<sub>3</sub> content, but on the bending of B–O–B bond at adjacent BO<sub>6</sub> octahedron for ABO<sub>3</sub> perovskite structure. On the other hand, these  $\tau_f$  values were essentially correlated with the B-site bond valence and octahedral tilting, and either decreasing the B-site bond valence or increasing the octahedral tiltings ( $\theta$  and  $\varphi$ ) led to a decrease in  $\tau_f$  value for these ceramic systems. Moreover,  $\varepsilon_r$  decreased with increasing  $x$  value due to the smaller dielectric polarizability of NdAlO<sub>3</sub> than that of SCNT, and the  $Q \times f$  value in SCNTAx solid solution had

strongly depended on the apparent density and average grain sizes. As a result, excellent microwave dielectric properties of  $\varepsilon_r \sim 55.6$ ,  $Q \times f \sim 25,600$  GHz (at 4.249 GHz) and  $\tau_f \sim 6.7$  ppm/°C were obtained for SCNTAx ( $x = 0.25$ ) ceramic sintered at 1520 °C for 4 h, which is a promising candidate for miniaturization components in developing wireless communication systems.

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