



# Hydrothermal growth of one-dimensional Ce-doped TiO<sub>2</sub> nanostructures for solid-state DSSCs comprising Mg-doped CuCrO<sub>2</sub>

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## ABSTRACT

In this work, hydrothermal method was used to grow the one-dimensional Ce-doped TiO<sub>2</sub> nanostructures on the fluorine-doped tin oxide glass. The incorporation of cerium atoms into the TiO<sub>2</sub> lattice was investigated by structural and electrical analyses. The results showed that TiO<sub>2</sub> nanorod arrays had maximum electrical conductivity for 3 % of Ce dopant. The increase in the electrical conductivity can be ascribed to the formation of oxygen vacancy in TiO<sub>2</sub> nanostructures with Ce doping. The morphology of the as-grown nanorods revealed that the diameter of the TiO<sub>2</sub> nanorods increased with Ce doping concentration. Hydrothermal etching treatment was performed on the as-grown 3 % Ce-doped TiO<sub>2</sub> nanorod arrays in hydrochloric acid solution at 160 °C for different time durations. The obtained results clearly showed that the morphology change from the nanorods to the nanotubes occurred after 4 h of etching during hydrothermal treatment. In addition, sol–gel process was employed to synthesize 5 % Mg-doped CuCrO<sub>2</sub> nanoparticles with p-type conductivity as solid-state electrolyte. The solid-state DSSCs fabricated using hydrothermally treated 3 % Ce-doped TiO<sub>2</sub> nanorods displayed better photovoltaic performance than those made from the untreated ones. This is attributed to the larger specific surface of the hydrothermally treated nanorods compared with the untreated ones. The dye-loading measurement confirmed that the amounts of adsorbed dye on the surface of hydrothermally treated nanorods were higher than that of the untreated ones.

## Introduction

Among the different n-type semiconductors, titanium dioxide (TiO<sub>2</sub>) has been recognized as one of the most promising materials for use in photocatalysts, dye-

sensitized solar cells (DSSCs), lithium-ion batteries, and gas sensors due to its unique physical, optical, and chemical properties such as wide bandgap, good photocatalytic activity, chemical stability, and non-toxicity [1–3]. In addition, vertically aligned one-

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dimensional (1D) titanium dioxide nanostructures such as nanowires, nanorods, and nanotubes as photoanodes in DSSCs have received a great deal of attention because of their large surface area, high crystallinity, and excellent charge transport properties which effectively suppress the charge recombination at grain boundaries [4, 5]. Until now, 1D vertically aligned TiO<sub>2</sub> nanostructures have been grown by different methods. Wei et al. reported the growth of dense and vertically aligned TiO<sub>2</sub> nanowires on titanium substrate by direct oxidation of titanium foils [6]. Shankar et al. prepared highly ordered TiO<sub>2</sub> nanotube arrays by electrochemical anodization of Ti foils [7]. Foong et al. synthesized vertically aligned TiO<sub>2</sub> nanotube arrays by liquid-phase atomic layer deposition in nanoporous anodic alumina templates [8]. The growth of TiO<sub>2</sub> nanorods on a tantalum substrate by metal–organic chemical vapor deposition was also reported by Chen et al. [9]. Furthermore, single-crystalline rutile TiO<sub>2</sub> nanorod arrays on different substrates by hydrothermal method were recently reported [10]. In comparison with the other methods, the hydrothermal method has been identified as one of the simple, low-cost, and low-temperature techniques to grow the nanostructures [11]. Although, 1D TiO<sub>2</sub> nanostructures have been extensively investigated but the growth of TiO<sub>2</sub> nanorod arrays on fluorine-doped tin oxide conductive glass (FTO) substrates via hydrothermal method has just recently been reported [12].

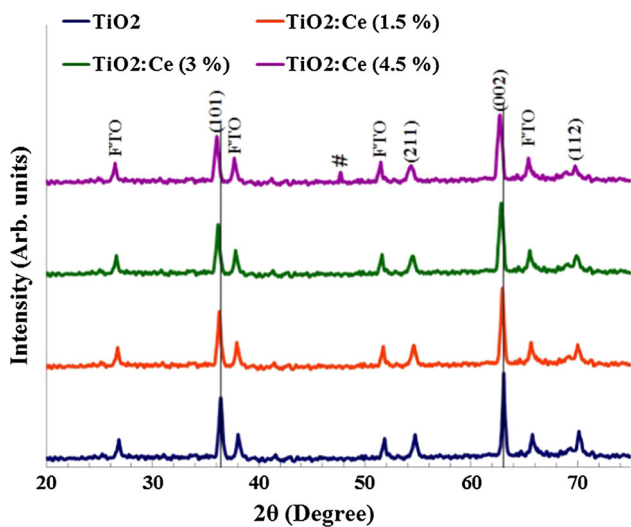
DSSCs have been considered as an important alternative to conventional Si-based solar cells, because of their high-energy conversion efficiency (>13 %), facile fabrication process, and low manufacturing cost [13–15]. However, major problems, such as leakage or evaporation of liquid electrolyte have limited the long-term operation and practical use of DSSCs [16]. To overcome these problems, considerable attempts have been made to substitute the liquid electrolyte with solid-state or gel-state electrolytes [17–19]. Unfortunately, the efficiency of fabricated solid-state DSSCs has been restricted due to pore-filling problems [20]. Pore-filling problems in solid-state DSSCs can be solved by one-dimensional TiO<sub>2</sub> nanorods/nanowires due to their direct pathways for penetration of solid-state electrolyte [21, 22]. Xu et al. fabricated high-efficiency solid-state DSSCs including TiO<sub>2</sub>-coated ZnO nanowire arrays and spiro-OMeTAD as photoanode and solid-state electrolyte, respectively [23]. Roh et al. synthesized

hierarchical anatase pine tree-like TiO<sub>2</sub> nanotubes by hydrothermal method for application in solid-state DSSCs [24]. Wang et al. fabricated solid-state DSSCs with ordered TiO<sub>2</sub> nanorods on FTO substrate as photoanodes [25]. In a previous work, the grown vertically aligned TiO<sub>2</sub> nanorods on FTO substrates were used as photoanode for fabrication of solid-state DSSCs comprising Zn-doped CuCrO<sub>2</sub> nanoparticles as solid-state electrolyte [26]. The use of TiO<sub>2</sub> nanorods instead of TiO<sub>2</sub> nanoparticles promoted the efficiency of fabricated solid-state DSSCs due to better penetration of solid-state electrolyte in space between nanorod arrays [27]. However, the efficiency improvement of fabricated devices was not significant due to insufficient dye adsorption on the surface of the nanorods [28, 29]. Vertically aligned TiO<sub>2</sub> nanotube arrays are superior to nanowires or nanorods because of their increased surface area for dye adsorption. Due to this reason, more efficient light harvesting for increasing the amount of photo-generated charge carriers can be occurred.

In the current work, we study the effect of Ce dopant concentration on structural properties of vertically aligned TiO<sub>2</sub> nanorod arrays. The growth of the TiO<sub>2</sub> nanorods on the FTO substrate is performed by hydrothermal method. The cerium is chosen as a dopant which can improve the n-type conductivity of the TiO<sub>2</sub> semiconductor through creation of oxygen vacancies [30–32]. It is well-known that the oxygen vacancies are the origin of the n-type conductivity in TiO<sub>2</sub> semiconductors [33, 34]. After preparation of Ce-doped TiO<sub>2</sub> nanorods, hydrothermal etching treatment is used for transforming them into Ce-doped TiO<sub>2</sub> nanotubes. We expect that the use of nanotubes instead of nanorods can improve the efficiency of constructed solid-state DSSCs as they have relatively higher surface area. Furthermore, in this way, a large amount of the solid-state electrolyte can fill into the space between TiO<sub>2</sub> nanostructures.

## Experimental methods

The vertically aligned Ce-doped TiO<sub>2</sub> nanorod arrays were grown directly on fluorine-doped tin oxide (FTO) glass substrate by hydrothermal method. Firstly, FTO glass substrates (1.5 × 2.0 cm<sup>2</sup>) were cleaned by ethanol and distilled water in an ultrasonic bath for 30 min to remove the possible impurities, were dried in air at 60 °C for 30 s and then



**Figure 1** X-ray diffraction patterns of undoped and Ce-doped TiO<sub>2</sub> nanorod arrays on FTO glass substrates. Peaks labeled as # are attributed to the CeO<sub>2</sub> structure. The solid lines are guide for the eyes.

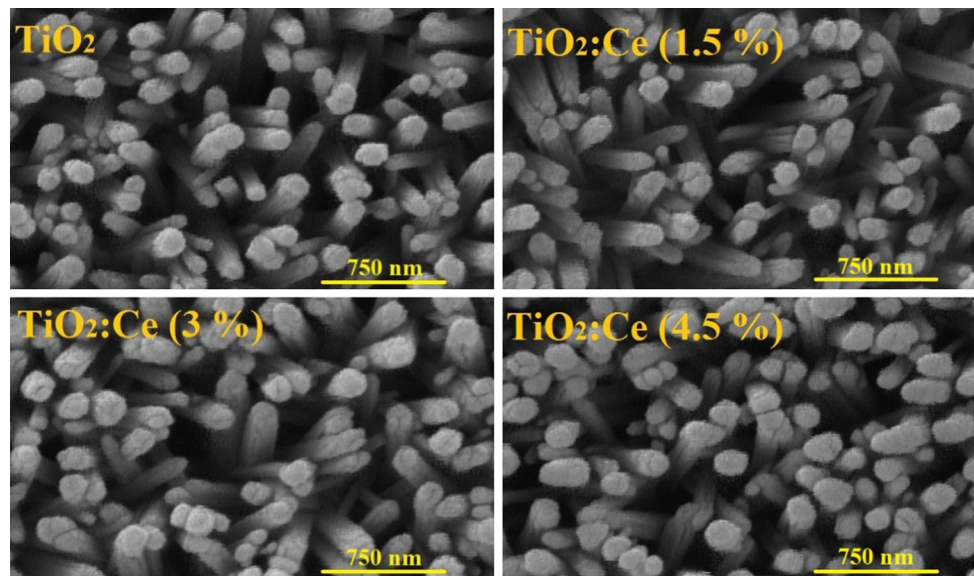
were put into a Teflon-lined stainless steel autoclave (100 ml volume). Then, 30 ml distilled water was mixed with 30 ml hydrochloric acid (HCl, Amertat Shimi Co., Iran) under magnetic stirring at room temperature. After stirring for 10 min, 1.1 ml titanium tetraisopropoxide (TTIP, Samchun Chemical Co., Korea) accompanied with different values of cerium (III) nitrate pentahydrate (Ce(NO<sub>3</sub>)<sub>3</sub>·5H<sub>2</sub>O, Chem-Lab NV Co., Belgium) as dopant source was added to the mentioned mixed solution and stirred vigorously for 20 min to obtain a clear and transparent solution. Afterwards, the resulting clear solution was transferred into the Teflon-lined stainless steel autoclave. The sealed autoclave was placed in a laboratory furnace (F35L-1200, Azar Furnace Co., Iran) at 160 °C for 12 h. At the end of the reaction, the autoclave was cooled down to room temperature naturally and the grown Ce-doped TiO<sub>2</sub> nanorod array on the FTO substrate was taken out, rinsed with distilled water and ethanol for several times to remove the residual reactants, and finally dried in air at 60 °C for 30 s. In order to convert the prepared nanorods to nanotubes, the as-grown TiO<sub>2</sub> nanorod arrays on FTO substrates were immersed in HCl (25 mL of 7 M aqueous solution) and hydrothermally treated at 160 °C for different time durations. The hydrothermally treated TiO<sub>2</sub> nanorods were rinsed with distilled water and ethanol for several times to remove the residual reactants, and finally dried in air at 60 °C for 30 s.

5 % Mg-doped CuCrO<sub>2</sub> nanoparticles were synthesized by sol-gel process. In this experiment, the solution was prepared by dissolving copper acetate monohydrate (Cu(CH<sub>3</sub>COO)<sub>2</sub>·H<sub>2</sub>O, Merck Co., Germany) and chromium nitrate nonahydrate (Cr(NO<sub>3</sub>)<sub>3</sub>·9H<sub>2</sub>O, Merck Co., Germany) as starting materials and magnesium nitrate hexahydrate (Mg(NO<sub>3</sub>)<sub>2</sub>·6H<sub>2</sub>O, Merck Co., Germany) as a dopant source in 2-methoxyethanol (CH<sub>3</sub>(CH<sub>2</sub>)<sub>2</sub>OOH, Merck Co., Germany) at 60 °C for 3 h using a magnetic stirrer. The molar ratio of copper ions to the sum of chromium and magnesium was fixed at 1.0 and the concentration of copper ions in the solution was equal to 0.7 M. The resultant transparent and deep green solution was then heated up to 110 °C and kept at this temperature for 3 h. The dried powder was annealed in air in the laboratory furnace at 900 °C for 1 h. After annealing, the prepared particles were milled in a planetary ball-milling machine (NARYA MPM-2\*250 H, Amin-Asia Co., Iran) at 400 rpm for 5 h in a zirconia vessel and with zirconia balls in ethanol to produce very fine particles. The weight ratio of ball to powder was 10:1.

The hydrothermally treated Ce-doped TiO<sub>2</sub> nanorod arrays on FTO glass were used as photoanode for DSSCs. Prior to dye sensitization, the hydrothermally treated Ce-doped TiO<sub>2</sub> nanorod arrays were heated up to 150 °C and kept at this temperature for 30 min. After cooling down to 80 °C, the Ce-doped TiO<sub>2</sub> electrodes were dipped into 0.5 mM solution of N719 dye (Dyesol Co., Australia) in ethanol for 24 h at room temperature and then washed with ethanol and dried in air. Subsequently, Mg-doped CuCrO<sub>2</sub> nanoparticles were ultrasonically dispersed in ethanol and deposited on the surface of the Ce-doped TiO<sub>2</sub> photoelectrode using a spin coater (SC-410 spincoater, Backer Viera Trading Co., Iran) at spin speed of 1000 rpm for 30 s. In addition, counter electrodes were prepared by spin-coating a 5 mM 2-propanol solution of H<sub>2</sub>PtCl<sub>6</sub> on FTO substrate at 500 rpm for 10 s and followed by annealing at 450 °C for 30 min. The solid-state DSSCs were assembled in a sandwich configuration with prepared Ce-doped TiO<sub>2</sub> photoanode and platinum-coated FTO counter electrode. The active area of the DSSCs was 0.2 cm<sup>2</sup>. Finally, to improve the interconnection between the nanoparticles and the penetration of a large amount of the nanoparticles into the photoanode, the constructed solid-state DSSCs were pressed by hydraulic press (769YP-15A, Tianjin Keqi High and New

**Table 1** Lattice parameters of the as-grown undoped and Ce-doped TiO<sub>2</sub> nanorods

Lattice parameter	TiO <sub>2</sub>	TiO <sub>2</sub> :Ce (1.5 %)	TiO <sub>2</sub> :Ce (3 %)	TiO <sub>2</sub> :Ce (4.5 %)
a (Å)	4.595	4.599	4.605	4.611
c (Å)	2.956	2.959	2.962	2.967

**Figure 2** Top-view FE-SEM images of the as-grown undoped and Ce-doped TiO<sub>2</sub> nanorod arrays with 1.5, 3, and 4.5 % Ce doping concentration.

Technology Co., China) under constant pressure (5 MPa) at room temperature for 30 min.

The shape and surface morphology of the as-grown nanorods and Mg-doped CuCrO<sub>2</sub> nanoparticles were characterized by a Hitachi S-4100 field-emission scanning electron microscope and a Tescan mira 3 xmu, respectively. The length of the vertically aligned Ce-doped TiO<sub>2</sub> nanorods was also determined by cross-sectional field-emission scanning electron microscopy (FE-SEM). In addition, the transmission electron microscopy (TEM) was carried out using a Zeiss EM10C microscope operating with an accelerating voltage of 100 kV to estimate the size of the nanoparticles. The crystalline structure and growth orientation of the as-grown undoped and doped TiO<sub>2</sub> nanorods were also characterized by X-ray powder diffraction (XRD) analysis with Cu K $\alpha$  radiation ( $\lambda = 1.5406 \text{ \AA}$ ) in the Bragg angle ranging between 10° and 80°. The hot-probe and Hall effect measurements were carried out to study the electrical properties of the as-grown Ce-doped TiO<sub>2</sub> nanorods. The magnetic field in the Hall effect measurements was adjusted at 400 mT. The current density–voltage (J–V) characteristics were

determined with a computer-controlled I–V tracer (IV-25, Sharif Solar, Iran) under the illumination of an AM1.5G solar simulator (SIM-1000, Sharif Solar, Iran) to determine the photovoltaic performance of the assembled DSSCs. The light intensity was adjusted to 100 mW/cm<sup>2</sup> using a reference Si solar cell. To estimate the amount of adsorbed N719 dye on the surface of the TiO<sub>2</sub> nanorod arrays, the sensitized Ce-doped TiO<sub>2</sub> photoanodes were immersed into 5 mL of 0.1 M NaOH aqueous solution to desorb dye molecules from the TiO<sub>2</sub> photoanodes. Then, a UV–Vis spectrophotometer (AvaSpec-3648, Avantes, Netherlands) was used to measure the UV–Vis absorption spectra of the resulting solutions. Photoluminescence (PL) spectra of the as-grown Ce-doped TiO<sub>2</sub> nanorods were determined by a fluorescence spectrophotometer (Cary Eclipse, Varian, Australia) using the 300-nm line of a Xe lamp as the excitation source at room temperature. In order to investigate the charge recombination rate in the prepared DSSCs, the open-circuit voltage decay experiment was carried out by recording the decay of the open-circuit voltage ( $V_{OC}$ ) after turning off the illumination, at steady state.

**Table 2** Electrical properties of the undoped and Ce-doped TiO<sub>2</sub> nanorods

Sample	TiO <sub>2</sub>	TiO <sub>2</sub> :Ce (1.5 %)	TiO <sub>2</sub> :Ce (3 %)	TiO <sub>2</sub> :Ce (4.5 %)
Charge carriers concentration (cm <sup>-3</sup> )	$3.1 \times 10^{17}$	$8.3 \times 10^{17}$	$6.6 \times 10^{18}$	$1.4 \times 10^{18}$
Hall mobility (cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )	5.3	5.1	4.6	4.2
Conductivity (S cm <sup>-1</sup> )	0.26	0.68	4.8	0.94

## Results and discussion

In order to study the effect of Ce dopant concentration on the crystalline structure of nanorods, X-ray diffraction measurements were first performed on the prepared nanorods. The XRD patterns of the as-grown undoped and Ce-doped TiO<sub>2</sub> nanorods in Fig. 1 exclusively show that the as-synthesized TiO<sub>2</sub> nanorods without annealing are in rutile phase (JCPDS card No. 00-001-1292) [35]. Previously, Liao et al. showed that the formation of anatase, rutile, and brookite phases of TiO<sub>2</sub> during the hydrothermal method strongly depends on the starting materials [36]. Using the distilled water, hydrochloric acid, and titanium tetraisopropoxide as the starting materials, the expected phase is pure rutile, in agreement with the experimental result of Liu et al. [12]. The obtained results also verify that up to 3 % of Ce-doped TiO<sub>2</sub>, there is no trace of any secondary phase or clusters such as cerium oxide and binary titanium cerium phases, confirming the achievement of the single-phase as-grown TiO<sub>2</sub> nanorods. This signifies that the Ti ions are successfully replaced by Ce ions in the crystalline structure. Furthermore, the sharp and intense peaks in the XRD patterns indicate that the as-grown nanorods are highly crystalline but in polycrystalline phase. However with 4.5 % of Ce doping, the appearance of an extra peak at 47.44 can be attributed to the CeO<sub>2</sub> (JCPDS Card No. 00-002-1306) [37]. As can be seen from Fig. 1, as the cerium doping percentage increases, the position of the diffraction peaks shifts to lower angles, their intensity decreases, and the full-width at half-maximum (FWHM) of the TiO<sub>2</sub> diffraction peaks increases gradually. The observed decrease in the crystalline quality of the TiO<sub>2</sub> nanorods can be attributed to the formation of lattice disorder and strain induced by substitution of larger Ce<sup>4+</sup> (0.96 Å) in the place of Ti<sup>4+</sup> (0.60 Å) due to the increase of the lattice parameters [38]. This trend (the slight shift in the peak positions and the variation of the

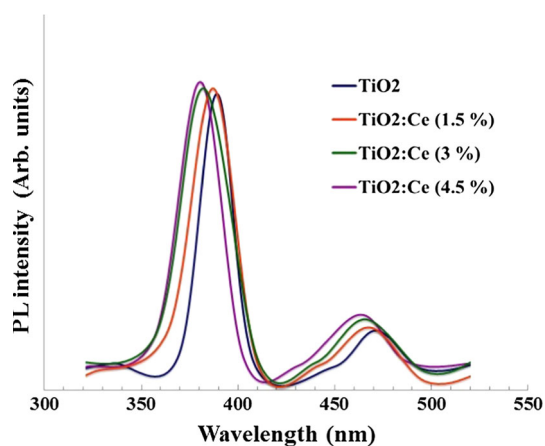
corresponding lattice constant) is consistent with the Vegard's law, which states that the lattice parameters of the Ti<sub>1-x</sub>Ce<sub>x</sub>O<sub>2</sub> ternary alloy can be linearly varied with Ce content from titanium oxide ( $x = 0$ ) to cerium oxide ( $x = 1$ ) [39]. The Lattice parameters of the as-grown nanorods obtained from the XRD patterns are presented in Table 1. Using Debye–Scherrer's equation [40],

$$d_{hkl} = \frac{0.9\lambda}{(\beta \cos(\theta))}$$

in which,  $d_{hkl}$  represents the size of the crystallites,  $\lambda$  is the wavelength of the radiation (1.5406 Å),  $\beta$  is FWHM in radians and  $\theta$  is half of the scattering angle, the size of the crystallites along (002) direction can be estimated to be about 30.1, 25.6, 20.3, and 16.7 nm for undoped, 1.5, 3, and 4.5 % Ce-doped TiO<sub>2</sub>, respectively.

Due to the low solubility of the cerium ions in TiO<sub>2</sub>, at higher levels of doping, only a few Ce<sup>4+</sup> ions penetrate into the TiO<sub>2</sub> lattice and the segregation of the dopants at the grain boundaries prevents the crystallite growth during the synthesizing process. The decrease in the size of the TiO<sub>2</sub> crystallites with Ce doping is mainly attributed to this phenomenon. Our results are in agreement with the other literature studies [41, 42].

Figure 2 shows typical FE-SEM images of the as-grown crystalline rutile TiO<sub>2</sub> nanorod arrays on FTO substrates with different concentrations of the dopant. The surface morphology of the TiO<sub>2</sub> nanorods reveals that the entire surface of the FTO substrate is uniformly covered by highly ordered nanorods. Furthermore, as expected, the as-grown undoped and Ce-doped TiO<sub>2</sub> nanorods have tetragonal shape with square top facets due to the tetragonal crystal structure. Moreover, the cross-sectional FE-SEM images confirm that the as-grown nanorods are perpendicular to the substrate, and that there are no side branches. The top-view FE-SEM images show that the average area of the rectangular cross section of the synthesized nanorods gradually increases with



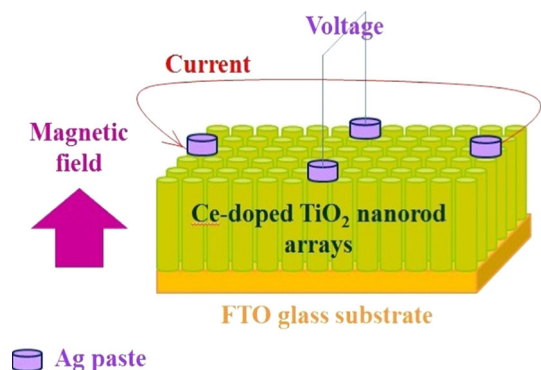
**Figure 3** Photoluminescence spectra of Ce-doped TiO<sub>2</sub> nanorods.

the increasing concentration of the Ce dopant from 0 to 4.5 %. A similar trend in the diameter of the Fe-doped ZnO nanorods was observed by Xiao et al. [43]. The morphology changes of the one-dimensional ZnO nanostructures were elucidated by the dislocation-driven growth mechanism due to the variation of the initial solution concentration [44]. The increase in the diameter of the as-grown TiO<sub>2</sub> nanorod arrays with Ce doping can be understood by this growth mechanism. Cerium doping has a similar effect on the supersaturation level of the initial growth solutions.

In order to investigate the electrical properties (charge carrier concentration, mobility, and conductivity) of the as-grown undoped and Ce-doped TiO<sub>2</sub> nanorods, Hall effect measurement was performed. The electrical properties of the vertically aligned TiO<sub>2</sub> nanorods with different Ce doping concentrations have been tabulated in Table 2. Hot-probe measurement confirmed the n-type conductivity of the as-grown nanorods. The n-type conductivity of TiO<sub>2</sub> structure is mainly attributed to the donor-type defects such as titanium interstitials and oxygen vacancies [33, 34]. When TiO<sub>2</sub> is doped with an element which has a higher oxidation state than Ti<sup>+</sup> [4], the electrical conductivity improvement originates from the presence of excess free electrons [45]. On the other hand, the increase in the electrical conductivity of the doped TiO<sub>2</sub> with tri- or tetravalent metals is ascribed to the oxygen vacancies in the structures. Iwaszuk and Nolan showed that Ce doping decreases the formation energy of oxygen vacancies in TiO<sub>2</sub> [31]. Similar trend is obtained for Ce-doped TiO<sub>2</sub> structure by Albuquerque [32]. Their results indicate

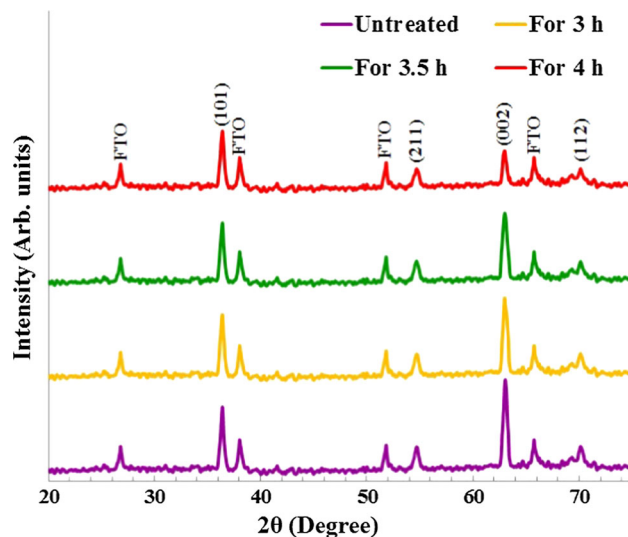
that the energy levels of oxygen vacancies in metal oxide semiconductors are affected by doping. In order to investigate the oxygen vacancy in the as-grown undoped and Ce-doped TiO<sub>2</sub> nanorods, the PL measurement was performed. Figure 3 shows the room temperature PL spectra of the as-grown nanorod arrays in the wavelength range of 320–520 nm. As can be seen from this figure, the nanorods exhibit two strong PL signals centered at approximately 370 and 470 nm. The first emission peak that appeared at ~370 nm (3.36 eV) is due to the emission of bandgap transition. The second emission peak around 470 nm (2.65 eV) results from the oxygen vacancies in TiO<sub>2</sub> structure [46], which are about 2.66 eV above the valence band and 0.7 eV below the conduction band of TiO<sub>2</sub>. There are three main steps during the PL process. First, the electrons are excited to the conduction band from the top of the valence band. Next, some of the excited electrons are captured by different levels of sub-bands related to oxygen vacancies. These oxygen vacancies act as radiative centers. Finally, the trapped electrons in oxygen vacancies are recombined with the holes in the valence band to give rise to PL signals. Thus, the larger the concentration of oxygen vacancies, the stronger the PL signal. As shown in Fig. 3, the PL intensity of Ce-doped TiO<sub>2</sub> nanorods increases with Ce doping concentration which confirms that a large amount of oxygen vacancy has been introduced into the TiO<sub>2</sub> structures by Ce substitution. Moreover, the emission peaks' position of Ce-doped TiO<sub>2</sub> nanorods shows a blue-shift compared to that of the undoped TiO<sub>2</sub> nanorods, which is attributed to the shift of the optical bandgap in Ce-doped TiO<sub>2</sub> nanorods due to the Burstein–Moss effect. Therefore, Ce doping facilitates the formation of oxygen vacancies in doped TiO<sub>2</sub>. In this way, two electrons released by formation of the oxygen vacancy play a significant role in increasing the electrical conductivity of the Ce-doped TiO<sub>2</sub>.

In addition to oxygen vacancies and doping concentration, the electrical conductivity is significantly affected by the crystallinity of the prepared nanorods [26]. As can be seen from Table 2, when a small amount of cerium is incorporated into the TiO<sub>2</sub> lattice, the electrical conductivity starts to increase. With the substitution of cerium for titanium, the oxygen vacancy in the TiO<sub>2</sub> structure begins to increase and then two extra free electrons are created from each



**Figure 4** Schematic image of Hall effect measurement from the as-grown nanorod arrays.

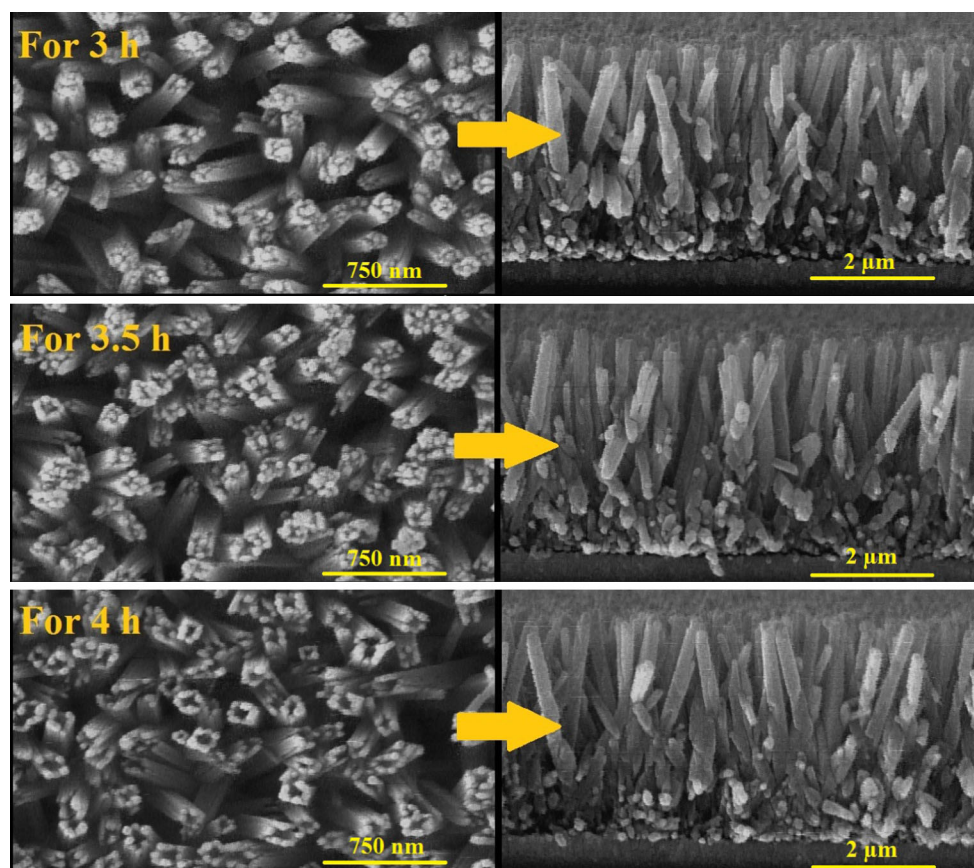
oxygen vacancy. Due to this process, concentration of charge carriers increases with Ce content up to 3 % but further increase in the value of cerium content decreases the conductivity of the TiO<sub>2</sub> nanorods. The reduction of conductivity is attributed to the crystallinity deterioration of TiO<sub>2</sub> nanorods and appearance of an additional phase (see Fig. 1). Because of the limited solubility of Ce inside TiO<sub>2</sub> structure, at higher doping concentrations (>3 at %), some cerium atoms cannot occupy the correct places inside TiO<sub>2</sub> lattice and produce an extra phase (CeO<sub>2</sub>) in the crystallite boundaries, which degrades the crystal structure of TiO<sub>2</sub> nanorods. Subsequently, the increase of cerium dopant concentration does not produce more charge carriers. Moreover, the mobility gradually decreases with the dopant content. In typical semiconductor materials, the decrease in Hall mobility is mainly due to different scattering mechanisms such as optical and acoustic phonon scattering mechanisms (also called lattice scattering), ionized and neutral impurity scattering, and grain boundary and defect scattering. In our study, the characterization of all samples has been carried out at room temperature in the same conditions. Thus, the phonon scattering does not contribute to the reduction of the Hall mobility. As mentioned above, with increasing Ce dopant content, the crystallinity of TiO<sub>2</sub> structure is deteriorated and the excess phase is formed (see Fig. 1). Therefore, the defect scattering originated from crystallinity deterioration of TiO<sub>2</sub> structure and the formation of doping oxide in the crystallite boundaries reduces the Hall mobility. These factors are also dominant in the reduction of electrical conductivity. The electrical conductivity ( $\sigma$ ) is found to be influenced by the charge carrier concentration ( $n$ ) and mobility of the charge carriers ( $\mu$ ),



**Figure 5** XRD patterns of 3 % Ce-doped TiO<sub>2</sub> nanorod arrays after hydrothermal treatment for different time durations.

according to the formula  $\sigma = n\mu e$  [47]. The results indicate that only appropriate amount of Ce can improve the electrical properties. The Hall effect measurement shows that the charge carrier concentration, mobility of the charge carriers, and the conductivity of the as-grown 3 % Ce-doped TiO<sub>2</sub> nanorods are about  $6.6 \times 10^{18} \text{ cm}^{-3}$ ,  $4.6 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ , and  $4.8 \text{ S cm}^{-1}$ , respectively. Schematic image of Hall effect measurement from the as-grown nanorod arrays is illustrated in Fig. 4.

A hydrothermal treatment was performed in the same Teflon-lined autoclave to convert the 3 % Ce-doped TiO<sub>2</sub> nanorods to nanotubes and the influence of hydrothermal treatment time on the morphology transformation was studied. The morphology change from the nanorods to the nanotubes takes place by the acid etching during hydrothermal treatment because the etching rate along (001) direction is faster than (101) direction [48, 49]. The preferential etching direction of the TiO<sub>2</sub> nanorods is elucidated by surface energy theory [50]. The surface energy is a main factor which determines the stability and reactivity of the surface [50]. In the rutile TiO<sub>2</sub> nanorods, the (101) and (001) surfaces have the lowest and highest surface energy values, respectively. Therefore, according to the surface energy theory, (101) facet is the most stable surface with lowest etching rate, whereas (001) facet is the most reactive surface with highest etching rate. Pan et al. and Zeng et al. showed that core and sidewall of the TiO<sub>2</sub> nanorods were composed of the (001) and (101) crystallites, respectively, and during



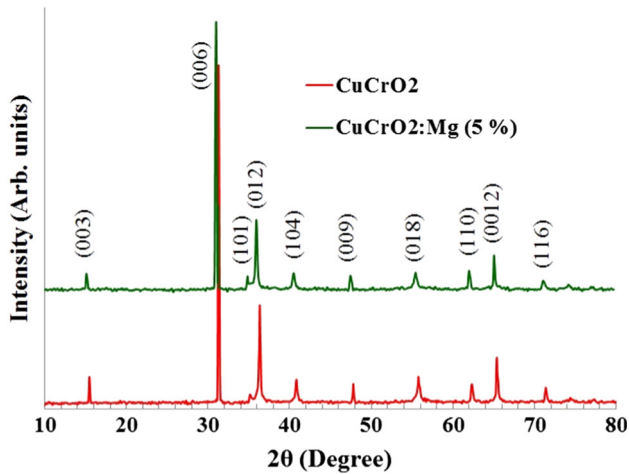
**Figure 6** Top-view and cross-sectional FE-SEM images of 3 % Ce-doped TiO<sub>2</sub> nanorod arrays after hydrothermal treatment for different time durations.

the hydrothermal etching process, the (101) sidewalls undergo slower etching rate than the (001) cores [48, 49]. Comparison of XRD patterns of 3 % Ce-doped TiO<sub>2</sub> nanorods before and after hydrothermal treatment shows that the intensity of (002) peak suddenly decreases with hydrothermal treatment process (see Fig. 5). The etching occurs along [001] direction which covers both (001) and (002) facets. These results confirm that the etching rate along (001) direction of TiO<sub>2</sub> nanorods is faster than other directions. Our results are consistent with the recent reports for undoped 1D TiO<sub>2</sub> nanostructures [48, 49]. Figure 6 exhibits the FE-SEM images of 3 % Ce-doped TiO<sub>2</sub> nanorods after hydrothermal treatment for different time durations. During hydrothermal treatment at short time duration (<4 h), the TiO<sub>2</sub> nanorods are insufficiently etched and only a small portion of the nanorod tips is etched. When the hydrothermal treatment process is performed for 4 h, the central portion of the TiO<sub>2</sub> nanorods is successfully etched by the etching solution, and the as-

grown vertically aligned 3 % Ce-doped TiO<sub>2</sub> nanorods are converted into ordered nanotubes. During hydrothermal treatment at the long time duration (>4 h), over-etching of the nanorods and the peeling-off of the nanorod arrays from the FTO substrate occurred.

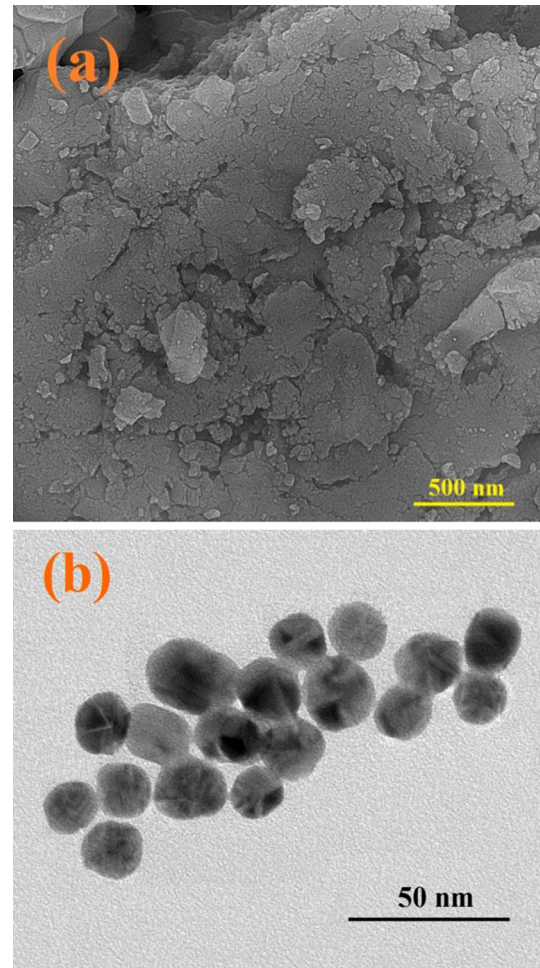
The XRD pattern of the prepared nanoparticles (after annealing at 900 °C for 1 h) is shown in Fig. 7. For comparison, the XRD spectrum of undoped CuCrO<sub>2</sub> nanoparticles (prepared in a similar condition) is also presented. All peaks are indexed to a delafossite structure of CuCrO<sub>2</sub> and consistent with JCPDS Card No. 01-089-6744 [26]. The XRD results show that there is no extra peak corresponding to any impurity phase such as Mg and MgO, which indicates that magnesium ions have been homogeneously incorporated into the CuCrO<sub>2</sub> lattice during the preparation process. As can be seen from Fig. 7, the intensity of the diffraction peaks decreases after Mg-doping, indicating that Mg-doping can lead to reduction in the crystalline quality of CuCrO<sub>2</sub>





**Figure 7** X-ray diffraction patterns of undoped and Mg-doped  $\text{CuCrO}_2$  nanoparticles.

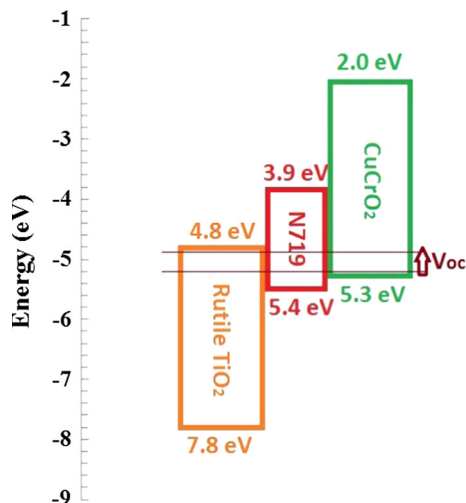
nanopowder. Furthermore, the addition of Mg dopant into  $\text{CuCrO}_2$  structure increases the full-width at half-maximum of the XRD peaks and shifts their position toward lower angles [26]. The decrease in crystal quality may be due to the appearance of stress because of the difference in ionic size of Cr ( $r_{\text{Cr}^{3+}} = 0.063$  nm) and Mg ( $r_{\text{Mg}^{2+}} = 0.066$  nm). The theoretical calculations and experimental studies have shown that the conductivity of the  $\text{CuCrO}_2$  structure is markedly increased when the small amount of divalent cations such as  $\text{Ni}^{2+}$ ,  $\text{Ca}^{2+}$ , and  $\text{Mg}^{2+}$  are substituted in trivalent  $\text{Cr}^{3+}$  site due to the creation of free holes [51, 52]. This increase in p-type conductivity is linked to  $\text{Cu}^{1+}/\text{Cu}^{2+}$  hole mechanism [51]. In  $\text{CuCrO}_2$  delafossite structure, the nominal valence of Cu and Cr is 1+ and 3+, respectively, and therefore three holes are required for charge compensation of the Cr vacancy. When one neutral Mg (with a valence of 2+) occupies a Cr site, the Mg donates two electrons to compensate holes generated by the Cr vacancy with one uncompensated hole. As a result,  $\text{Cu}^{1+}$  is oxidized to  $\text{Cu}^{2+}$ , creating a hole on Cu ion. This mechanism improves the p-type conductivity of the  $\text{CuCrO}_2$  structure. The average crystallite size of the synthesized nanoparticles is determined from the most intense diffraction peak (006) using the Debye–Scherrer’s equation—about 37.7 nm. According to this value, it is realized that the nanoparticles are relatively large and for a better penetration of the nanoparticles in the space between the nanorods, their size should be reduced. The typical FE-SEM and TEM images of the prepared 5 %



**Figure 8** a FE-SEM and b TEM images of the synthesized Mg-doped  $\text{CuCrO}_2$  nanoparticles after ball-milling for 5 h.

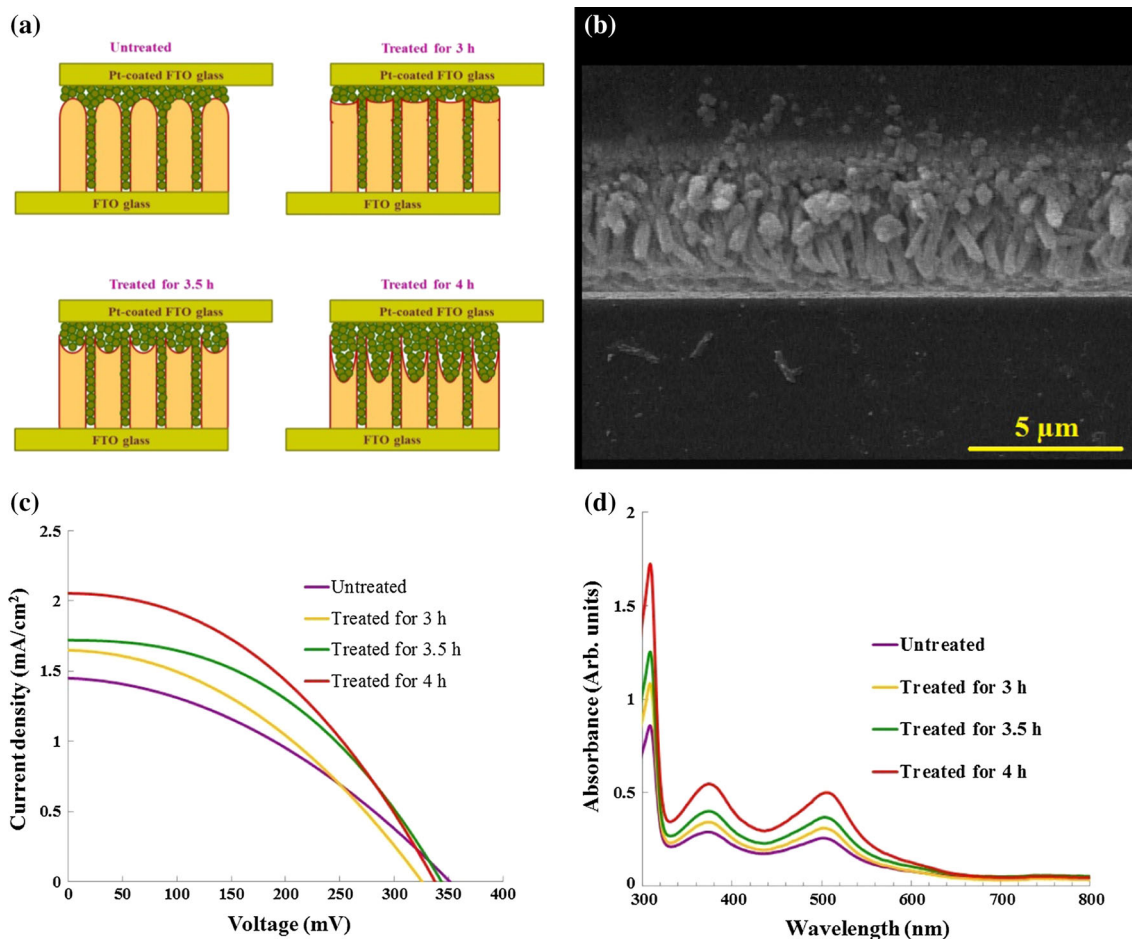
Mg-doped  $\text{CuCrO}_2$  nanoparticles after ball-milling for 5 h are illustrated in Fig. 8. This figure shows that the nanoparticles have spherical shape with uniform size which has been reduced to about 17 nm by milling.

In order to investigate the effect of hydrothermal treatment on the performance of the solid-state DSSCs, different types of the hydrothermally treated nanorods were used as photoanode. This type of solid-state DSSC consists of dye-sensitized nanostructure metal oxide semiconductor (such as 1D  $\text{TiO}_2$  nanostructure) on transparent conductive oxide glass, whose pore volume is filled with a transparent p-type semiconductor (such as  $\text{CuCrO}_2$  nanoparticles). Schematic energy diagram of the solid-state DSSC components used in this study is shown in Fig. 9. Under sunlight irradiation, a thin layer of dye molecules at the interface between the n- and p-type



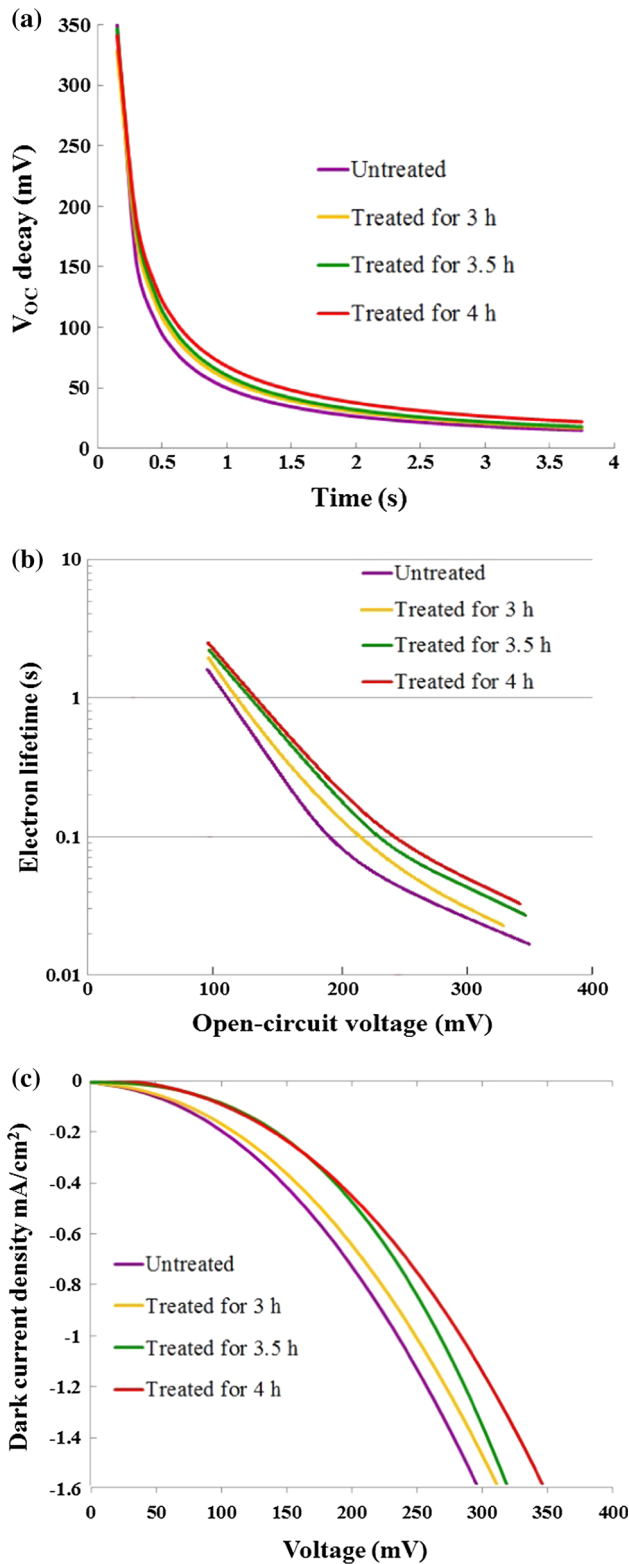
**Figure 9** Schematic energy diagram of the solid-state DSSC device components used in this study. The mentioned values were extracted from ref. [53–55].

semiconductors injects electrons into the conduction band of the n-type semiconductor and then is regenerated by injecting holes into the valence band of the p-type semiconductor. The open-circuit voltage generated by this type of solid-state DSSC under illumination corresponds to the difference between the Fermi level of electrons in the n-type semiconductor and Fermi level of holes in the p-type semiconductor. Schematic images of the constructed solid-state DSSCs based on 1D nanostructures (nanorods and nanotubes) are displayed in Fig. 10a. A cross-sectional SEM image is also shown in Fig. 10b. It shows conductive glass substrates, nanorods, and nanoparticles which to some extent, penetrated into the space between nanorods. The photocurrent density–voltage ( $J$ – $V$ ) curves of DSSCs based on vertically aligned 3 % Ce-doped  $\text{TiO}_2$  nanorods and



**Figure 10** **a** Schematic images and **b** cross-sectional FE-SEM image of the constructed solid-state DSSCs based on 1D nanostructures, **c** current density–voltage characteristics of the solid-state DSSCs made from hydrothermally treated 3 % Ce-

doped  $\text{TiO}_2$  nanorods under simulated AM1.5G solar light ( $100 \text{ mW/cm}^2$ ), and **d** UV–Vis absorption spectra of the desorbed N719 dye molecules from the photoanodes.



◀ **Figure 11** a Open-circuit voltage decay for different solid-state DSSCs and b calculated electron lifetime as a function of open-circuit voltage. c Dark current density–voltage characteristics of the solid-state DSSCs made from hydrothermally treated 3 % Ce-doped TiO<sub>2</sub> nanorods.

power conversion efficiency (PCE) of 0.18 % is obtained for the 3 % Ce-doped TiO<sub>2</sub> nanorod-based DSSC. After hydrothermal treatment for 3 and 3.5 h, the fabricated DSSCs with 3 % Ce-doped TiO<sub>2</sub> nanorod arrays show  $J_{sc}$  of 1.65 and 1.72 mA/cm<sup>2</sup>,  $V_{oc}$  of 328 and 344 mV, and PCE of 0.21 and 0.23 %, respectively. Finally, obvious improvement in  $J_{sc}$  2.05 mA/cm<sup>2</sup>, and the efficiency of 0.29 % is observed after hydrothermal treatment at 160 °C for 4 h. In comparison with untreated Ce-doped TiO<sub>2</sub> nanorod arrays, there is an observable enhancement in the short-circuit current density (from 1.45 to 2.05 mA/cm<sup>2</sup>) as a result of 4 h hydrothermal treatment. Many parameters can influence the  $J_{sc}$  of the DSSCs, but in our case only one of them has significant effect. As can be seen from the FE-SEM images in Fig. 6, the tip of the nanorods is etched with hydrothermal treatment and then is converted to the nanotubes with the increasing treatment duration. The formation of large specific surface area for dye loading on the inner and outer sides of the nanotubes can be mentioned as the main reason for  $J_{sc}$  improvement. When more dye molecules are adsorbed on the surface of the photoanode, a large amount of electron–hole pair can be generated and participate in the photocurrent improvement. The UV–Vis absorption spectra verify that the amounts of adsorbed dye on the surface of the hydrothermally treated photoanodes are higher than that of untreated one (Fig. 10d).

Open-circuit photovoltage decay measurement is used to determine the electron lifetime in DSSCs. Zaban et al. clearly showed that the rate of open-circuit photovoltage decay is inversely proportional to the electron lifetime in DSSCs [56]. The relationship between the photovoltage decay and electron lifetime is given by

$$\tau_n = \frac{k_B T}{q} \left( \frac{dV_{oc}}{dt} \right)^{-1},$$

where  $k_B$ ,  $T$ , and  $q$  are the Boltzmann constant, absolute temperature, and positive elementary charge, respectively. In order to measure the  $V_{oc}$

nanotubes as photoanodes are shown in Fig. 10c. A short-circuit current density ( $J_{sc}$ ) of 1.45 mA/cm<sup>2</sup>, an open-circuit voltage ( $V_{oc}$ ) of 351 mV, and an overall

decay, the simulated solar light is turned off and the decay of the photovoltage is recorded. Figure 11a shows the  $V_{OC}$  decay of the fabricated DSSCs. As can be seen, the decay of the photovoltage for hydrothermally treated photoanodes is slower than that of untreated one which indicates that the electron lifetime increases with hydrothermal treatment. The electron lifetimes determined from the  $V_{OC}$  decay measurements are about 15.4 ms for untreated photoanode and 20.7 ms (treated for 3 h), 25.9 ms (treated for 3.5 h), and 31.0 ms (treated for 4 h) for hydrothermally treated photoanodes at open-circuit voltage (see Fig. 11b). The highest lifetime of 31.0 ms has been recorded from the DSSC with 4 h hydrothermally treated Ce-doped  $TiO_2$  nanorods (3 %), which suggests that the treatment of nanorods has an impressive influence on dye adsorption, electron generation, and electron transport through  $TiO_2$  nanorods. The increase in the electron lifetime by hydrothermal treatment indicates that the number of defect sites and recombination centers has been reduced. This deduction could be confirmed by the dark J–V characteristics in Fig. 11c. The dark current in DSSCs is attributed to the recombination process at the photoanode/electrolyte interface. The 4 h hydrothermally treated Ce-doped  $TiO_2$ -based solid-state DSSC shows the lowest value of  $J_{SC}$ . The smallest dark current for this device specifies that the recombination at the  $TiO_2$ /Mg-doped  $CuCrO_2$  interface has been suppressed and so, the electron transportation along conducting channels (hydrothermally treated nanorods) is more efficient.

To the best of our knowledge, these solar cells are the best reported to date. In 2007, the photovoltaic parameters of the solid-state DSSCs utilizing  $CuAlO_2$  nanoparticles in  $TiO_2$ /N719 dye/ $CuAlO_2$  structure have been reported by Bandara et al. with  $J_{SC}$  of 0.08 mA/cm<sup>2</sup> [57]. Lee et al. prepared n- $TiO_2$ /p-NiO heterojunction electrodes for all-solid-state dye-sensitized solar cells with overall energy conversion efficiency of about 0.025 % [58]. Furthermore, Yuhas et al. have observed photovoltaic properties in FTO/ $ZnO$  nanowire/ $Cu_2O$  nanoparticles structure with energy conversion efficiency of about 0.053 % using a  $TiO_2$  blocking layer [59]. We hope that our results will help to extend the potential applications of the delafossite-type semiconductors as new materials for all-oxide solid-state DSSCs in the future.

## Conclusions

In summary, one-dimensional Ce-doped  $TiO_2$  nanostructures were successfully grown on FTO glass by hydrothermal method. The influence of the Ce doping concentration on the structural properties and electrical characteristics of the  $TiO_2$  nanorods were investigated. The highest electrical conductivity was obtained for 3 % Ce-doped  $TiO_2$  nanorod arrays. The increase in the electrical conductivity is the result of the oxygen vacancy formation in  $TiO_2$  nanostructures with Ce doping. Hydrothermal etching treatment for 4 h in acid solution created nanotubes from the nanorods. Sol–gel process was also employed to synthesize 5 % Mg-doped  $CuCrO_2$  nanoparticles with p-type conductivity as solid-state electrolyte. The photovoltaic performance of the fabricated solid-state DSSCs with hydrothermally treated Ce-doped  $TiO_2$  nanorods were improved due to the large specific surface area. The dye-loading measurement confirmed that the amount of adsorbed dye on the surface of hydrothermally treated photoanodes is higher than that of untreated one. By the techniques employed in this study, the highest efficiency of about 0.3 % was obtained for the solid-state DSSCs based on delafossite nanoparticles.

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