



New Method of Calculating a Multiplication by using the Generalized Bernstein-Vazirani Algorithm

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Abstract We present a new method of more speedily calculating a multiplication by using the generalized Bernstein-Vazirani algorithm and many parallel quantum systems. Given the set of real values $\{a_1, a_2, a_3, \dots, a_N\}$ and a function $g : \mathbf{R} \rightarrow \{0, 1\}$, we shall determine the following values $\{g(a_1), g(a_2), g(a_3), \dots, g(a_N)\}$ simultaneously. The speed of determining the values is shown to outperform the classical case by a factor of N . Next, we consider it as a number in binary representation; $M_1 = (g(a_1), g(a_2), g(a_3), \dots, g(a_N))$. By using M parallel quantum systems, we have M numbers in binary representation, simultaneously. The speed of obtaining the M numbers is shown to outperform the classical case by a factor of M . Finally, we calculate the product; $M_1 \times M_2 \times \dots \times M_M$. The speed of obtaining the product is shown to outperform the classical case by a factor of $N \times M$.

Keywords Quantum computation · Quantum algorithms

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1 Introduction

As for applications of the quantum theory, implementation of a quantum algorithm to solve Deutsch's problem [1–3] on a nuclear magnetic resonance quantum computer is reported firstly [4]. An implementation of the Deutsch-Jozsa algorithm on an ion-trap quantum computer is also reported [5]. There are several attempts to use single-photon two-qubit states for quantum computing. Oliveira *et al.* implements Deutsch's algorithm with polarization and transverse spatial modes of the electromagnetic field as qubits [6]. Single-photon Bell states are prepared and measured [7]. Also the decoherence-free implementation of Deutsch's algorithm is reported by using such a single-photon and by using two logical qubits [8]. More recently, a one-way based experimental implementation of Deutsch's algorithm is reported [9].

In 1993, the Bernstein-Vazirani algorithm was reported [10, 11]. It can be considered as an extended Deutsch-Jozsa algorithm. In 1994, Simon's algorithm was reported [12]. Implementation of a quantum algorithm to solve the Bernstein-Vazirani parity problem without entanglement on an ensemble quantum computer is reported [13]. Fiber-optics implementation of the Deutsch-Jozsa and Bernstein-Vazirani quantum algorithms with three qubits is discussed [14]. Quantum learning robust against noise is studied [15]. A quantum algorithm for approximating the influences of Boolean functions and its applications are recently reported [16]. Quantum computation with coherent spin states and the close Hadamard problem are also discussed [17]. Transport implementation of the Bernstein-Vazirani algorithm with ion qubits is more recently reported [18]. Quantum Gauss-Jordan elimination and simulation of accounting principles on quantum computers are discussed [19]. We mention that the dynamical analysis of Grover's search algorithm in arbitrarily high-dimensional search spaces is studied [20]. A method of computing many functions simultaneously by using many parallel quantum systems is reported [21].

On the other hand, the earliest quantum algorithm, the Deutsch-Jozsa algorithm, is representative to show that quantum computation is faster than the classical counterpart with a magnitude that grows exponentially with the number of qubits. In 2015, it was discussed that the Deutsch-Jozsa algorithm can be used for quantum key distribution [22]. In 2017, it was discussed that secure quantum key distribution based on Deutsch's algorithm using an entangled state [23]. Subsequently, a highly speedy secure quantum cryptography based on the Deutsch-Jozsa algorithm is proposed [24].

In this work we present a new method of more speedily calculating a multiplication by using the generalized Bernstein-Vazirani algorithm and many parallel quantum systems. Given the set of real values $\{a_1, a_2, a_3, \dots, a_N\}$ and a function $g : \mathbf{R} \rightarrow \{0, 1\}$, we shall determine the following values $\{g(a_1), g(a_2), g(a_3), \dots, g(a_N)\}$ simultaneously. The speed of determining the values is shown to outperform the classical case by a factor of N . Next, we consider it as a number in binary representation; $M_1 = (g(a_1), g(a_2), g(a_3), \dots, g(a_N))$. By using M parallel quantum systems, we have M numbers in binary representation, simultaneously. The speed of obtaining the M numbers is shown to outperform the classical case by a factor of M . Finally, we calculate the product; $M_1 \times M_2 \times \dots \times M_M$. The speed of obtaining the product is shown to outperform the classical case by a factor of $N \times M$.

2 The Generalized Bernstein-Vazirani Algorithm

Let us suppose that we are given the following sequence of real values

$$a_1, a_2, a_3, \dots, a_N. \quad (1)$$

Let us now introduce the function

$$g : \mathbf{R} \rightarrow \{0, 1\}. \tag{2}$$

One step is of determining the following values

$$g(a_1), g(a_2), g(a_3), \dots, g(a_N). \tag{3}$$

Recall that in the classical case, we need N queries, that is, N separate evaluations of the function (2). In our quantum algorithm, we shall require a single query. Suppose now that we introduce another function

$$f : \{0, 1\}^N \rightarrow \{0, 1\} \tag{4}$$

which is a function with a N -bit domain and a 1-bit range. We construct the following function

$$\begin{aligned} f(x) &= g(a) \cdot x = \sum_{i=1}^N g(a_i)x_i \pmod{2} \\ &= g(a_1)x_1 \oplus g(a_2)x_2 \oplus g(a_3)x_3 \oplus \dots \oplus g(a_N)x_N \\ x_i &\in \{0, 1\}, g(a_i) \in \{0, 1\}, a_i \in \mathbf{R} \end{aligned} \tag{5}$$

where a_i is a real value. Here $g(a)$ symbolizes

$$g(a_1)g(a_2) \dots g(a_N). \tag{6}$$

Let us follow the quantum states through the algorithm. The input state is

$$|\psi_0\rangle = |0\rangle^{\otimes N} |1\rangle \tag{7}$$

where $|0\rangle^{\otimes N} = \overbrace{|0\rangle \otimes |0\rangle \otimes \dots \otimes |0\rangle}^N$. After the componentwise Hadamard transforms on the state (7)

$$\overbrace{H|0\rangle \otimes H|0\rangle \otimes \dots \otimes H|0\rangle}^N \otimes H|1\rangle \tag{8}$$

we have

$$|\psi_1\rangle = \sum_{x \in \{0,1\}^N} \frac{|x\rangle}{\sqrt{2^N}} \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]. \tag{9}$$

Next, the function f is evaluated using

$$U_f : |x, y\rangle \rightarrow |x, y \oplus f(x)\rangle \tag{10}$$

giving

$$|\psi_2\rangle = \pm \sum_x \frac{(-1)^{f(x)}|x\rangle}{\sqrt{2^N}} \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]. \tag{11}$$

Here $y \oplus f(x)$ is the bitwise XOR (exclusive OR) of y and $f(x)$. By checking the cases $x = 0$ and $x = 1$ separately, we see that for a single qubit

$$H|x\rangle = \sum_z (-1)^{xz} |z\rangle / \sqrt{2}. \tag{12}$$

Thus

$$\begin{aligned}
 &H^{\otimes N}|x_1, \dots, x_N\rangle \\
 &= \frac{\sum_{z_1, \dots, z_N} (-1)^{x_1 z_1 + \dots + x_N z_N} |z_1, \dots, z_N\rangle}{\sqrt{2^N}}.
 \end{aligned} \tag{13}$$

This can be summarized more succinctly in the very useful equation

$$H^{\otimes N}|x\rangle = \frac{\sum_z (-1)^{x \cdot z} |z\rangle}{\sqrt{2^N}} \tag{14}$$

where $x \cdot z$ is the bitwise inner product of x and z , modulo 2. Using the (14) and (11), we can now evaluate $H^{\otimes N}|\psi_2\rangle = |\psi_3\rangle$

$$|\psi_3\rangle = \pm \sum_z \sum_x \frac{(-1)^{x \cdot z + f(x)} |z\rangle}{2^N} \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]. \tag{15}$$

Thus

$$|\psi_3\rangle = \pm \sum_z \sum_x \frac{(-1)^{x \cdot z + g(a) \cdot x} |z\rangle}{2^N} \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]. \tag{16}$$

Because we have

$$\sum_x (-1)^x = 0 \tag{17}$$

we can see that

$$\sum_x (-1)^{x \cdot z + g(a) \cdot x} = 2^N \delta_{g(a), z}. \tag{18}$$

Therefore, the sum is zero if $z \neq g(a)$ and is 2^N if $z = g(a)$. Thus

$$\begin{aligned}
 |\psi_3\rangle &= \pm \sum_z \sum_x \frac{(-1)^{x \cdot z + g(a) \cdot x} |z\rangle}{2^N} \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right] \\
 &= \pm \sum_z \frac{2^N \delta_{g(a), z} |z\rangle}{2^N} \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right] \\
 &= \pm |g(a)\rangle \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right] \\
 &= \pm |g(a_1)g(a_2) \dots g(a_N)\rangle \left[\frac{|0\rangle - |1\rangle}{\sqrt{2}} \right]
 \end{aligned} \tag{19}$$

from which

$$|g(a_1)g(a_2) \dots g(a_N)\rangle. \tag{20}$$

can be obtained. That is to say, if we measure $|g(a_1)g(a_2) \dots g(a_N)\rangle$ then we can retrieve the following values

$$g(a_1), g(a_2), g(a_3), \dots, g(a_N) \tag{21}$$

using a single query. All we have to do is of performing one quantum measurement.

The speed of determining N values improves by a factor of N as compared to the classical counterpart. Notice that we recover the Bernstein-Vazirani algorithm when $g : a_i \rightarrow a_i$.

3 Calculating a Multiplication by using the Generalized Bernstein-Vazirani Algorithm

We present a new method of more speedy calculating a multiplication by using many parallel quantum systems. By using M parallel quantum systems, we can compute M functions g^1, g^2, \dots, g^M simultaneously.

Let us suppose that we are given the following another sequence of real values

$$b_1, b_2, b_3, \dots, b_N. \tag{22}$$

Let us now introduce the function

$$g^2 : \mathbf{R} \rightarrow \{0, 1\}. \tag{23}$$

We can determine the following values by using the generalized Bernstein-Vazirani algorithm

$$g^2(b_1), g^2(b_2), g^2(b_3), \dots, g^2(b_N). \tag{24}$$

By using M parallel quantum systems, we can retrieve the following values

$$\begin{aligned} &g^1(a_1), g^1(a_2), g^1(a_3), \dots, g^1(a_N) \\ &g^2(b_1), g^2(b_2), g^2(b_3), \dots, g^2(b_N) \\ &\dots \\ &g^M(c_1), g^M(c_2), g^M(c_3), \dots, g^M(c_N). \end{aligned} \tag{25}$$

In the case, we measure the following quantum state

$$\begin{aligned} &|g^1(a_1)g^1(a_2) \dots g^1(a_N)\rangle \otimes \\ &|g^2(b_1)g^2(b_2) \dots g^2(b_N)\rangle \otimes \\ &\dots \otimes |g^M(c_1)g^M(c_2) \dots g^M(c_N)\rangle. \end{aligned} \tag{26}$$

All we have to do is of performing one quantum measurement.

We consider them as numbers in binary representation

$$\begin{aligned} M_1 &= (g^1(a_1), g^1(a_2), g^1(a_3), \dots, g^1(a_N)) \\ M_2 &= (g^2(b_1), g^2(b_2), g^2(b_3), \dots, g^2(b_N)) \\ &\dots \\ M_M &= (g^M(c_1), g^M(c_2), g^M(c_3), \dots, g^M(c_N)). \end{aligned} \tag{27}$$

Therefore, by using M parallel quantum systems, we have M numbers in binary representation, simultaneously. The speed of obtaining the M numbers is shown to outperform the classical case by a factor of M . Finally, we calculate the product; $M_1 \times M_2 \times \dots \times M_M$. The speed of obtaining the product is shown to outperform the classical case by a factor of $N \times M$.

As an example, if $N = 2$ and $M = 3$, we may have

$$(g^1(a_1), g^1(a_2)) = (0, 1) = 1 \tag{28}$$

$$(g^2(b_1), g^2(b_2)) = (1, 0) = 2 \tag{29}$$

$$(g^3(c_1), g^3(c_2)) = (1, 1) = 3 \tag{30}$$

and we have

$$(g^1(a_1), g^1(a_2)) \times (g^2(b_1), g^2(b_2)) \\ \times (g^3(c_1), g^3(c_2)) = 1 \times 2 \times 3 = 6. \quad (31)$$

An experimental evidence is very interesting and it is a further investigation.

4 Conclusions

In conclusion, we have presented a new method of more speedily calculating a multiplication by using the generalized Bernstein-Vazirani algorithm and many parallel quantum systems. Given the set of real values $\{a_1, a_2, a_3, \dots, a_N\}$ and a function $g : \mathbf{R} \rightarrow \{0, 1\}$, we shall have determined the following values $\{g(a_1), g(a_2), g(a_3), \dots, g(a_N)\}$ simultaneously. The speed of determining the values has been shown to outperform the classical case by a factor of N . Next, we have considered it as a number in binary representation; $M_1 = (g(a_1), g(a_2), g(a_3), \dots, g(a_N))$. By using M parallel quantum systems, we have had M numbers in binary representation, simultaneously. The speed of obtaining the M numbers has been shown to outperform the classical case by a factor of M . Finally, we have calculated the product; $M_1 \times M_2 \times \dots \times M_M$. The speed of obtaining the product has been shown to outperform the classical case by a factor of $N \times M$.

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