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Micromechanical analysis of hydraulic fracturing in the toughness-dominated regime: implications to supercritical carbon dioxide fracturing

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Abstract

Supercritical carbon dioxide (SC-CO₂) is considered as an ideal non-aqueous fracturing fluid due to its superior properties of liquid-like density, gas-like viscosity, high compressibility, and diffusivity. This study aims to investigate the micromechanical behavior of SC-CO₂ fracturing in both intact and fractured rock samples by using a coupled fluid-solid discrete element method (DEM) model. A new numerical algorithm for hydraulic fracturing in the toughness-dominated regime is developed by assuming that the pressure in the whole fracture is uniform. This new numerical algorithm could achieve a much higher computational efficiency compared with the conventional hydromechanical scheme in DEM. Hydraulic fracturing cases using high-viscosity fracturing fluid are also performed for comparison. The results indicate that the fracture propagation induced by SC-CO₂ tends to be less smooth and continuous, more asymmetric, and tortuous compared to that induced by viscous fluid. Besides, the lowviscosity fluid like $SC-CO₂$ can lead to a lower breakdown pressure, and the fluid leak-off into the rock matrix can result in a lower breakdown pressure and higher fracture propagation pressure. The simulations also illustrate that SC-CO₂ fracturing tends to create a more complex and productive fracture network if the pre-existing natural fractures are involved. As a result, we can conclude that SC-CO₂ could be an alternative fracturing fluid to induce a more effective fracture network for hydrocarbon production.

Keywords Hydraulic fracturing . Supercritical carbon dioxide . DEM . Toughness-dominated . Tortuosity . Breakdown pressure

1 Introduction

Hydraulic fracturing has become a widespread reservoir stimulation technique for hydrocarbon production from low-

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permeability formations. However, conventional hydraulic fracturing with water-based fracturing fluid has some disadvantages, including environmental impacts caused by the flow-back water and the high cost of water usage. Furthermore, since gas productivity is greatly dependent on the water sensitivity of shale gas formations, water-based fracturing is not suitable for the water-sensitive clay-abundant formations [\[1](#page-15-0)–[4](#page-15-0)]. There is an urgent requirement to explore the feasibility of non-aqueous fracturing technologies to overcome these drawbacks. Supercritical carbon dioxide (SC- $CO₂$) has been considered as an ideal non-aqueous fracturing fluid due to its superior properties of liquid-like density, gas-like viscosity, high compressibility, and diffusivity [[1](#page-15-0)]. Furthermore, the use of $SC\text{-}CO₂$ could make shale a major host for the geological carbon storage [\[5](#page-15-0), [6\]](#page-15-0). In recent years, despite astounding progress made in understanding the mechanics of hydraulic fracturing following the boom of unconventional resource development, the mechanism of $SCCO₂$ fracturing has not been fully understood yet.

Hydraulic fracture is known to have complex multiscale behavior. The fracturing behavior is determined by the interplay of two energy dissipation mechanisms (viscosity and toughness) and two fluid storage mechanisms (storage and leak-off) [\[7](#page-15-0)]. These two sets of competing mechanism are associated with four limiting regimes: storage viscosity [[8\]](#page-15-0), storage toughness [[9\]](#page-15-0), leak-off viscosity [[10\]](#page-15-0), and leak-off toughness [[11\]](#page-15-0). Conventional hydraulic fracturing with water-based fracturing fluid usually approaches or even falls into the viscosity-dominated regime due to the relatively large viscosity of fracturing fluid, if the leak-off is not dominant. In such a regime, the energy is mainly dissipated by the viscous fluid flow inside the fracture, while the energy expended in fracturing solid medium is negligibly small. There is a gradual transition to the toughness-dominated regime with the decrease of fracturing fluid viscosity. Thus, the fracture propagation driven by $SC\text{-}CO₂$ likely corresponds to the toughnessdominated regime because of its gas-like viscosity [\[12\]](#page-15-0). Both the fracture profiles and pressure response can be vastly different among different hydraulic fracture propagation regimes [\[13](#page-15-0)–[15\]](#page-15-0).

Hydraulic fracture initiation and propagation are mainly affected by some key factors including rock properties, confining pressure, pressurization rate, and fluid type $[16–21]$ $[16–21]$ $[16–21]$ $[16–21]$ $[16–21]$. Experimental studies of hydraulic fracturing were carried out to investigate the fracturing behavior for different types of fluid. Besides the fluid viscosity, recent analyses showed that the composition and state of the fracturing fluid also play important roles [\[22](#page-15-0), [23](#page-15-0)]. Ishida et al. [[24](#page-15-0)–[26](#page-15-0)] conducted experiments using oil, water, liquid $CO₂$ (L-CO₂), and SC- $CO₂$ to fracture granite. It was observed that SC-CO₂ fracturing leads to the lowest breakdown pressure, the highest tortuosity [[27](#page-15-0)], and the most complex fracture network with many secondary branches due to the reduced viscosity and enhanced infiltration. A possible explanation can be that the lowviscosity fracturing fluid like $SC\text{-}CO₂$ could easily penetrate into defects in the rock and subsequently cause more fractures. The fracturing experiments in non-bedded sandstone using different viscous fracturing fluids showed similar results [\[28](#page-15-0)]. In addition, it is necessary for hydraulic fracture to activate and connect pre-existing natural fractures to generate a large fracture network [[29,](#page-16-0) [30\]](#page-16-0). The growth behavior of SC- $CO₂$ -induced fractures in layered shale [\[28,](#page-15-0) [31,](#page-16-0) [32\]](#page-16-0) and bedded sandstone [\[33](#page-16-0)] was explored in laboratory experiments. The results suggest that $SC\text{-}CO₂$ promotes the fracture complexity with secondary branches by activating microfractures. More shear-dominant fractures were also observed along the grain boundaries inclined in the direction of the maximum compressive stress when using low-viscosity fracturing fluid [\[26](#page-15-0), [27](#page-15-0)]. The key parameters for $SC\text{-}CO_2$ -related hydraulic fracturing experiments in the literature are summarized in Table 1. Those experimental studies showed some macroscale fracturing behavior induced by $SC\text{-}CO₂$, but the fracture growth behavior induced by $SC\text{-}CO₂$ in the naturally fractured rock needs further exploration and the micromechanical analysis to reveal the mechanism in depth is still lacking.

Numerical simulation is a powerful tool to reproduce the fracturing process and help to understand the mechanism of

Table 1 Summary of the key parameters for SC-CO₂-related hydraulic fracturing experiments in the literature

Sample	Size (mm)	Fracturing fluid	Fracturing characteristics	Reference
Granite	Cubic, 190	Viscous oil \rightarrow Water	More wavelike, branched, and thinner fractures	$\lceil 24 \rceil$
Granite	Cubic, 170	$L-CO2$ \rightarrow SC-CO ₂	More three-dimensional fractures; lower breakdown pressure	$\lceil 25 \rceil$
Granite	Cubic, 170	Viscous oil \rightarrow Water \rightarrow L-CO ₂ \rightarrow SC-CO ₂	More three-dimensional, tortuous, and shear-dominant fractures; lower breakdown pressure	$\lceil 26 \rceil$
Granite	Cubic, 170	Viscous oil \rightarrow Water \rightarrow SC-CO ₂	More tortuous, branched, and shear-dominant fractures	[27]
Shale, non-bedded sandstone	Cubic, 200	Water \rightarrow L-CO ₂ \rightarrow SC-CO ₂	More branches, more likely to connect with natural fracture and bedding; lower breakdown pressure	[28]
Shale	Cylindrical, 85/170 (D/L)	Viscous oil \rightarrow Water \rightarrow L-CO ₂	More shear-dominant and branched fractures; larger fracture opening	$[31]$
Shale	Cylindrical, $100 \sim 150$, $100/\sim 200$ (D/L) Cubic, 300	Water \rightarrow CO ₂	More shear-dominant fractures; more complex networks; higher leak-off; lower breakdown pressure	$\lceil 32 \rceil$
Layered tight sandstone	Cubic, 300	X-linked guar \rightarrow Slick water \rightarrow SC-CO ₂	More branches; smaller fracture opening; higher leak-off; lower breakdown pressure	$\lceil 33 \rceil$

hydraulic fracturing. The injection fluid properties, such as compressibility, viscosity, and interfacial tension, are the main factors investigated through numerical modeling [\[34](#page-16-0)–[37](#page-16-0)]. Based on the cohesive zone model in conjunction with a poroelasticity model, it was found that fluid viscosity has more significant effect on the hydraulic conductivity and leak-off coefficient than fluid compressibility [[34\]](#page-16-0). Moreover, fluid viscosity could influence the fluid flow behavior and the evolution of the poroelastic stress around the borehole, resulting in the low-viscosity fluid filtrating into the fractures easily [[35,](#page-16-0) [37](#page-16-0)]. Meanwhile, the interfacial tension of fracturing fluid controls whether fluid invades the pore space at the borehole wall [\[37\]](#page-16-0). Thus, hydraulic fracturing induced by $SCCO₂$ could be more efficient due to its high compressibility, low viscosity, and low interfacial tension, particularly for the rock containing natural microcracks [[38\]](#page-16-0).

In this study, we applied a fully coupled fluid-solid discrete element method (DEM) model to study the micromechanical behavior of $SC-CO₂$ hydraulic fracturing. We developed a new numerical algorithm for hydraulic fracturing in the toughness-dominated regime. The algorithm can achieve a much higher computational efficiency than the traditional hydromechanical scheme for hydraulic fracturing modeling with low fluid viscosity. Based on the numerical algorithm, the comparison between hydraulic fracturing with $SC-CO₂$ in the toughness-dominated regime and that with high-viscosity fluid approaching the viscosity-dominated regime was made.

2 Numerical algorithm of hydraulic fracturing with DEM

2.1 New algorithm for the toughness-dominated regime

The numerical modeling is carried out by employing the discrete element code PFC2D [[39](#page-16-0)], in which the model is

assembled by bonded circular particles [\[40\]](#page-16-0). The linear parallel bond model is applied to mimic the rock deformation and failure behavior at microscale [[41\]](#page-16-0). The fluid-driven fracture propagation process is an emergent behavior of bonding breakage in either tensile or shear modes. Besides, there are additional joints represented by using the smooth joint model (SJM) to simulate the discontinuities in fractured rock mass [\[42](#page-16-0)]. The SJM can simulate the behavior of a smooth interface with user-defined macroscopic properties such as friction, cohesion, and dilation.

The fluid-solid coupling is implemented by applying a fluid network based on the porous structure formed by the particles [\[39](#page-16-0), [43](#page-16-0)]. As shown in Fig. 1 a, a series of enclosed domains are created by connecting the centers of neighboring particles while the fluid pressure in a domain is applied to the surface of surrounding particles. Therefore, there is a drag force (F_{fluid}) for each particle calculated from the resultant force of surrounding domain pressure (P) (see Fig. 1b). The aperture between two particles is assumed to be a flow channel, which is initially inactivated since the sample in this study is impermeable unless mentioned. If the normal stress or the shear stress exceeds the corresponding bond strength, the bond breaks, which results in the generation of a microcrack. Then, the corresponding flow channel is activated and the fluid flow is allowed (see Fig. 1c).

The developed new algorithm for the toughnessdominated regime is shown in the flow chart in Fig. [2.](#page-3-0) In the toughness-dominated regime, the energy is mainly spent in fracturing the rock, and the energy dissipated by viscous flow is negligible. The pressure gradient along the fracture length is nearly zero. Thus, this algorithm assumes that the pressure in the whole fracture is uniform. The fluid volume inside the hydraulic fracture can be viewed as a whole, and the goal is to match the current opened fracture volume (V_c) with the injection fluid volume (V_f) . The change of fluid pressure is proportional to the volume difference as given in Eq. ([1](#page-3-0)). Note that the

Fig. 1 Fluid flow model based on DEM (modified from [\[35](#page-16-0), [43\]](#page-16-0)). a The enclosed domain formed by the neighboring particles. b The drag force on a particle calculated from the pressure of surrounding domains. c Fluid

flow in the activated channel after the corresponding bond breaks. P_1 , P_2 , P_3 , and P_4 are the surrounding domain pressures acting on the particle, and F_{fluid} is the resultant force of P_1 , P_2 , P_3 , and P_4

Fig. 2 Proposed numerical algorithm for the toughnessdominated regime

compressibility of $SC-CO₂$ is not considered here and the effect of compressibility is analyzed in Section [5](#page-10-0). The fracture volume is calculated from the aperture width (w) and the length of flow channel (l_p) as shown in Eqs. (2) and (3) . The apertures and current opened fracture volume are updated with the apparent fluid pressure acting on the surrounding particles. At each fluid time step, up to ten mechanical steps are executed except that the balance of the fluid volume and fracture volume is achieved (less than 1% difference).

$$
\Delta P = \alpha (V_f - V_c) / V_t \tag{1}
$$

$$
V_{\rm c} = \sum l_{\rm p}(w - w_0) \tag{2}
$$

$$
V_t = \sum l_p w \tag{3}
$$

where ΔP is the pressure increment at each mechanical time step, α is a coefficient proportional to the rock stiffness, V_t is the total fracture volume, w is the aperture width, w_0 is the assumed initial aperture width, and l_p is the length of the flow channel, which is the distance between two connecting domain centers.

2.2 Comparison with the conventional hydromechanical scheme

For the conventional hydromechanical scheme [\[39,](#page-16-0) [43](#page-16-0)], unequal fluid pressure accumulates in each domain as a result of pressure gradient along the fracture length. The fluid flow inside the fracture is calculated at each fluid time step (d_t) , and the change of fluid pressure for each fluid domain is given by the following equation:

$$
\Delta P = K_{\rm f} (\sum q d_{\rm t} - dV_{\rm d}) / V_{\rm d} \tag{4}
$$

where K_f is the fluid bulk modulus, V_d is the apparent volume of the domain, dV_d is the change of domain volume per fluid time step, and Σq is the total flow rate per fluid time step in the domain. The flow rate (q) in the channel is given as a function of the aperture (w) , the fluid viscosity (μ) , the pressure difference (ΔP) between the two domains, and the length of flow channel (l_p)

$$
q = \frac{w^3}{12\mu} \frac{\Delta P}{l_p} \tag{5}
$$

For both the conventional hydromechanical scheme and the developed new algorithm for the toughness-dominated regime, the aperture width (w) is updated with contact normal force as follows:

Compressive normal force :
$$
w = \frac{w_0 F_0}{F + F_0}
$$
 (6)

Zero normal force : $w = w_0 + w_{\text{mech}}$ (7)

where w_{mech} is the gap between two particles, F is the current contact normal force, and F_0 is the normal force at which the aperture decreases to 50% of the initial aperture (w_0) .

In the conventional hydromechanical scheme, since the fluid time step is proportional to the term of fluid viscosity over fracture aperture ($dt \propto \frac{\mu}{K_f w^2}$), the fluid time step could be very small for low-viscosity fluid such as $SCCO₂$. In contrast, fluid viscosity and fracture aperture are irrelevant to the fluid time step for the toughness-dominated regime and it can therefore be 2 orders of magnitude larger than that for the conventional hydromechanical scheme. Consequently, a much higher computational efficiency can be achieved based on the new proposed algorithm for hydraulic fracturing in the toughnessdominated regime.

3 Numerical model setup

3.1 Model description

A two-dimensional intact rock sample $(1.0 \text{ m} \times 1.0 \text{ m})$ with an injection point at the center is generated in this study. The particle radius is randomly distributed between 3 and 5 mm by using a given random seed number. The sample is subjected to two types of biaxial compressive stresses: one with 10 MPa in the x direction and 12 MPa in the y direction and the other with 10 MPa in the x direction and 15 MPa in the y direction. Young's modulus of the intact rock sample is 29.1 GPa, and Poisson's ratio is 0.30 based on the numerical simulation of unconfined compressive strength test. The uniaxial compressive strength and tensile strength of the intact rock sample are 22.3 MPa and 4.8 MPa, respectively. Through the biaxial compressive test, the intact rock friction angle and cohesion are 37.5° and 5.8 MPa, respectively. Besides the intact rock sample, a fractured rock sample which contains 18 natural fractures is also generated as shown in Fig. 3. The natural fractures obey the SJM, which allows slip and separation at particle contacts [[42](#page-16-0)]. The frictional stress on the weakly bonded natural fractures is assumed to obey Coulomb's frictional law. The initial aperture is 0.1 mm, which is 17 times larger than the assumed residual aperture. The fracturing fluid is not able to invade the natural fracture until the corresponding bond breaks and connects to the hydraulic fracture. The algorithm for the toughnessdominated regime is applied to the modeling cases with $SC-CO₂$ (fluid viscosity 0.06 cp), while the traditional hydromechanical coupling scheme is applied to the cases with high viscosity of 100 cp. The input parameters for the numerical model and the microscale parameters for the contact models are summarized in Tables [2](#page-5-0) and [3](#page-5-0), respectively. Hydraulic fracturing simulation cases in both intact and fractured rock samples are listed in Table [4](#page-6-0).

To test the accuracy of the new numerical algorithm for the toughness-dominated regime, the simulation results are compared with the reference solution for a plane strain Khristianovic-Geertsma-de Klerk (KGD) hydraulic fracture [\[14](#page-15-0)]. It is first necessary to determine the regime of propagation. There are four limiting regimes as shown in Fig. [4:](#page-6-0) storage viscosity (M) [[8](#page-15-0)], storage toughness (K) [[9\]](#page-15-0), leak-off viscosity (M) [[10](#page-15-0)], and leak-off toughness (K) [\[11\]](#page-15-0). The fracture propagation regimes are determined by the dimensionless

Fig. 3 Pre-existing natural fractures in rock sample. a 18 fractures inserted in the model sample. b Close-up view of one natural fracture. The red lines represent pre-existing natural fractures, and the blue and gray thick lines in b indicate the contact normal

Table 2 The parameters for the numerical model

Model properties	Values
Model size (m)	1.0×1.0
Number of particles	17,511
Particle radius (mm)	$3 - 5$
Particle density $(kg/m3)$	3169
Compressive stress in x, S_{min} (MPa)	10
Compressive stress in v , S_{max} (MPa)	12/15
Young's modulus, $E(GPa)$	29.1
Poisson's ratio, ν	0.30
Uniaxial compressive strength, σ_c (MPa)	22.3
Uniaxial tensile strength, σ_t (MPa)	4.8
Friction angle, φ (°)	37.5
Cohesion, c (MPa)	5.8
Bulk modulus of rock, $K(GPa)$	8.3
Proportionality coefficient, α (MPa)	0.2
Assumed residual aperture, w_0 (m)	6e-6
Assumed joint initial aperture, w_{0j} (m)	$1e-4$
Injection rate, $Q \text{ (m}^2\text{/s)}$	$0.5e-4$
Injection fluid viscosity (cp)	0.06 and 100

time (τ) and the toughness parameter (K_m) as follows:

$$
\tau = \frac{E' C'^6}{\mu' Q^3} t \tag{8}
$$

$$
K_{\rm m} = \left(\frac{K^{'4}}{\mu' E^{'3} Q}\right)^{1/4} \tag{9}
$$

where Q is the fluid injection rate, t is the injection time, and μ , E, K, and C are the material parameters defined as follows:

Table 3 The microscale parameters for the contact models

Microscale input parameters	Values	
Young's modulus of the particle (GPa)	17	
Ratio of stiffness of the particle	2.5	
Friction coefficient of the particle, μ	0.5	
Young's modulus of the parallel bond (GPa)	17	
Ratio of stiffness of the parallel bond	2.5	
Tensile strength of the parallel bond (MPa)	8	
Cohesion of the parallel bond (MPa)	15	
Friction angle of the parallel bond $(°)$	30	
Normal stiffness of SJM (GPa/m)	120	
Shear stiffness of SJM (GPa/m)	60	
Tensile strength of SJM (MPa)	0.1	
Cohesion of SJM (MPa)	0.1	
Friction coefficient of SJM	0.5	

$$
\mu^{'} = 12\mu, E^{'} = \frac{E}{1 - v}, K^{'} = 4\left(\frac{2}{\pi}\right)^{1/2} K_{\rm IC}, C^{'} = 2C_{\rm L}
$$
 (10)

where μ is the fluid viscosity, E is Young's modulus, v is Poisson's ratio, K_{IC} is the mode I fracture toughness of the rock, and C_{L} is Carter's leak-off parameter.

In this study, K_m is the decisive parameter due to the absence of leak-off. As shown in $[14]$ $[14]$, the hydraulic fracture propagates in the toughness-dominated regime for $K_m > 4.80$ and in the viscosity-dominated regime for $K_m < 0.70$. When the mode I fracture toughness of the random model is 2.66 Pa· $m^{1/2}$, which is inversely calculated in the next section of this paper, the toughness parameter (K_m) calculated for the modeling cases with $SC-CO₂$ and 100 cp fluid are 6.99 and 1.09, respectively. Therefore, the fracture propagation driven by $SC-CO₂$ belongs to the toughness-dominated regime, while that induced by 100 cp fluid approaches the viscositydominated regime.

3.2 Model verification

In order to verify the hydraulic fracturing model in the toughness-dominated regime, simulation of hydraulic fracturing using $SC-CO₂$ was first performed using the cubic array particle assembly (Fig. [5](#page-7-0)). The particle radius is 5 mm, and Young's modulus is 35.1 GPa. The two applied stresses are 10 MPa in the x direction and 12 MPa in the y direction. Other input parameters are the same as those in Table 2.

As shown in [[14](#page-15-0)], the KGD solutions of fracture aperture (w) and fracture length (l) for the toughness-dominated regime are given as follows:

$$
w(\xi, t) = 0.6828 \left(\frac{K^2 Qt}{E^2}\right)^{1/3} \left(1 - \xi^2\right)^{1/2} \tag{11}
$$

$$
l(t) = 0.9324 \left(\frac{E'Qt'}{K'}\right)^{1/6}
$$
 (12)

where $\xi = x/l$ is the normalized coordinate along the fracture and x is the coordinate measured from the injection point.

The evolutions of fracture aperture and fracture length at $t =$ 0.9 s in the numerical model and KGD model are compared in Fig. [6.](#page-7-0) Note that the input of mode I fracture toughness of the regular model was 1.88 Pa·m^{1/2} in the KGD model in order to match numerical solution. The fracture length in the cubic array particle assembly is the distance between two fracture tips in the y direction. The fracture aperture at the injection point and the length in the KGD model are as follows: $w(0,0.9) = 7.52e-5$ m and $l(0.9) = 0.763$ m, while that calculated in the numerical model are 7.24e−5 m and 0.770 m, respectively. The fracture width is underestimated by 3.9%, and the fracture length is overestimated by 1.1%. The small discrepancy can be attributed

Table 4 Summary of hydraulic fracturing simulation cases

Case ID	Fracturing fluid	S_{\min}	$S_{\rm max}$	Pre-existing fractures	Leak- off
Case 1-S1	100 cp fluid	10	12	N ₀	N ₀
Case 1-S2	100 cp fluid	10	15	N ₀	N ₀
Case 2-S1	SC - $CO2$	10	12	N ₀	N ₀
$Case 2-S2$	SC - $CO2$	10	15	N ₀	N ₀
Case N1-S1	100 cp fluid	10	12	Yes	N ₀
$Case N1-S2$	100 cp fluid	10	15	Yes	N ₀
Case N2-S1	SC - $CO2$	10	12	Yes	N ₀
$Case N2-S2$	SC - $CO2$	10	15	Yes	N ₀
Case L1-S1	100 cp fluid	10	12	N ₀	Yes

to the gradual increase of rock modulus in the DEM modeling due to the opening of the fracture. Nevertheless, the results in Fig. [6](#page-7-0) indicate that the numerical algorithm for the toughnessdominated regime is relatively accurate.

4 Numerical results

The cases of hydraulic fracturing with $SCCO₂$ in the toughness-dominated regime and with 100 cp fracturing fluid approaching the viscosity-dominated regime were described and compared in this section. All cases were run with the random particle assembly to better mimic the rock materials.

Fig. 4 The phase diagram of the fracture propagation regimes versus dimensionless time (τ) and toughness parameter (K_m) [[14](#page-15-0)]. M, K, M, and \widetilde{K} represent the storage viscosity regime, the storage toughness regime, the leak-off viscosity regime, and the leak-off toughness regime, respectively. The black dots indicate the fracture propagation regimes with the input parameters in Table [2](#page-5-0)

4.1 Hydraulic fracturing in intact rock sample

In the cases of intact rock samples, the hydraulic fractures driven by the two types of fracturing fluid extend along the direction of maximum principal stress as shown in Fig. [7](#page-8-0). The hydraulic fracture is tortuous, and the tortuosity can be used to evaluate fracture propagation pathways [\[27](#page-15-0)]. The tortuosity is defined as the total fracture length along a pathway (L_0) divided by the distance between the two ends of the fracture (L) in the reference area, and the total fracture length is calculated as the sum of the length of all microcracks generated in the model. When the stress difference is 2 MPa in case 1-S1 and case 2-S1, the results show that the tortuosity is 1.63 for SC-CO₂ and 1.19 for 100 cp fluid. The tendency in laboratory experiments using granite and shale is similar but with different values $[26, 27, 29]$ $[26, 27, 29]$ $[26, 27, 29]$ $[26, 27, 29]$ $[26, 27, 29]$ $[26, 27, 29]$. In the experiments using the granite $[26, 27, 29]$ $[26, 27, 29]$. 27], the fractures induced by SC-CO₂ have a tortuosity up to 1.109, while those induced by oil $(\sim 320 \text{ cp})$ have a tortuosity up to 1.062. In the experiments using shale [[29\]](#page-16-0), the measured total fracture length contains branches, and the tortuosity is 5.602 for L-CO₂ and 1.090 for oil $(\sim 270 \text{ cp})$. The cases with larger stress difference show similar results. The calculated tortuosity suggests that $SC\text{-}CO₂$ has potential to generate more complex fracture networks and is desirable to enhance the hydrocarbon production in the field operations.

The net pressure and aperture along the hydraulic fractures are shown in Figs. [8](#page-8-0) and [9.](#page-9-0) When comparing at the same injection fluid volume, hydraulic fracturing using more viscous fluid tends to induce a thicker and shorter fracture than that using $SC\text{-}CO₂$ in the toughness-dominated regime. The difference of fracture profiles can be explained by the distribution of fluid pressure in the fracture as shown in Figs. [8](#page-8-0) a and [9](#page-9-0) a. Uniform fluid pressure is applied throughout the SC-CO2-induced fracture surface, while large-pressure gradient exists in the viscous fluid–induced fracture, and the closer it is to the tip, the greater the pressure gradient. The influence of fluid pressure distribution is analyzed in Section [5.](#page-10-0)

Fig. 5 a Schematic of the hydraulic fracturing verification model at $t = 0.9$ s. b Close-up view of the cubic array particle assembly

Evolution of injection fluid pressure and the tip location on the y-axis for the two types of fracturing fluid are compared in Figs. [10](#page-9-0) and [11](#page-10-0). The fracture length growth with a series of jumps is observed in the cases using $SC\text{-}CO₂$, and the jumps cause the immediate drop of the injection fluid pressure in the initial state. Moreover, the fracture induced by $SC-CO₂$ is more asymmetric as shown in Figs. [10](#page-9-0) b and [11](#page-10-0) b. The plateaus of the curve indicate that the fracture growth of one wing could be temporarily arrested due to the local inherent heterogeneity, which may shift the position of maximum aperture from the injection point (Fig. [11](#page-10-0) b). By contrast, the hydraulic fracture driven by 100 cp fracturing fluid propagates more smoothly, continuously, and symmetrically over the flow

time. The results are in agreement with those observed in the hydraulic fracturing simulation in the toughness-dominated regime using the traditional hydromechanical coupling scheme [[44\]](#page-16-0). During the fracturing process, the ratio between the particle size and the accumulated fracture length becomes smaller as the hydraulic fracture propagates, resulting in a stable injection fluid pressure with a decreasing fluctuation. Moreover, the propagation pressure for $SCCO₂$ fracturing is much lower than that with high-viscosity fluid, and the former is slightly larger than the minimum principal stress (10 MPa). It is also observed that $SC\text{-}CO₂$ fracturing can reach a lower breakdown pressure (22.3 MPa for case 2-S1 and 22.1 MPa for case 2-S2) than the viscous fluid fracturing (24.1 MPa for

Fig. 6 Simulation results of hydraulic fracturing verification model with the cubic array particle assembly. a Comparison between the approximate solution and the numerical solution for fracture aperture. **b** Comparison between the approximate solution and the numerical solution for fracture length

Fig. 7 Fracture profiles induced by 100 cp fracturing fluid and SC-CO₂, a in case 1-S1 and case 2-S1 at $t = 1.5$ s and b in case 1-S2 and case 2-S2 at $t = 0.9$ s. The black line represents the hydraulic fracture, and the red line

shows the connection line of two ends. The numbers give the average tortuosity in the two directions for the fractures generated under different test conditions

case 1-S1 and 23.7 MPa for case 1-S2). However, the difference of breakdown pressure is less profound than that ob-served in previous experimental studies [[25](#page-15-0), [26\]](#page-15-0). The discrepancy may be attributed to the absence of circular borehole and fluid leak-off in this study, which is analyzed in Section [5](#page-10-0).

The numerical results of case 2-S1 and the approximate solutions using the KGD model at $t = 1.5$ s are shown in Fig. [12.](#page-10-0) The fracture aperture at the injection point and the distance between two fracture tips in the y direction using the KGD model were $w(0,1.5) = 10.40e-5$ m and $l(1.5) = 0.919$ m, while those calculated by the numerical model are 10.64e−5 m and 0.683 m, respectively. The fracture aperture at the injection point is overestimated by 2.3%, while the distance between the two tips in the y direction is underestimated by 25.7%. It is obvious that there is a relatively large discrepancy, though the input of mode I

fracture toughness (2.66 $\text{Pa} \cdot \text{m}^{1/2}$) is attempted to match the numerical solutions as well as possible. The apparent toughness (K_{IC}) scales with the tensile strength of the contact bond $(\sigma_t^{'})$ and the particle radius (R) as shown in [\[41,](#page-16-0) [45\]](#page-16-0) is as follows: $K_{\text{IC}} \propto \sigma'_t \sqrt{R}$. However, due to the effect of internal heterogeneity, the mode I fracture toughness in case 2-S1 is larger than that in the regular model (1.88 Pa·m^{1/2}) regardless of the negative effect of decreasing the particle size. The possible explanation is that the fracture propagation in the toughness-dominated regime is more sensitive to the local inherent heterogeneity of the sample, resulting in less smooth fracture growth and more asymmetric and tortuous fractures as aforementioned. In real rock mass, the contact between particles might not be aligned with the direction of the maximum or minimum principal stresses, and the maximum horizontal stress may make some contribution to the con-

Fig. 8 a The net pressure of hydraulic fractures in case 1-S1 and case 2-S1 at $t = 1.5$ s. b The aperture of hydraulic fractures in case 1-S1 and case 2-S1 at $t = 1.5$ s. Both the color scale and the thickness of the lines indicate the net pressure in **a** and fracture width in **b**

Fig. 9 a The net pressure of hydraulic fractures in case 1-S2 and case 2-S2 at $t = 0.9$ s. b The aperture of hydraulic fractures in case 1-S2 and case 2-S2 at $t = 0.9$ s. Both the color scale and the thickness of the lines indicate the net pressure in a and fracture width in **b**

tact force along the path of the hydraulic fracture. It is harder to fracture a random solid medium, and the hydraulic fracture propagation path is more tortuous. There is an increase of the effective toughness and a decrease of direct fracture length for the random model [\[44](#page-16-0)]. If the tortuosity is taken into consideration, the total accumulated fracture length is 1.099 m in case 2-S1. Thus, the fracture in the numerical solution is longer, thinner, and more tortuous than the one predicted by the KGD theoretical solution.

4.2 Hydraulic fracturing in fractured rock sample

It is widely accepted that rock mass contains pre-existing natural fractures and the propagation of hydraulic fracture could be greatly affected by those natural fractures or geological discontinuities [\[46](#page-16-0)–[49](#page-16-0)]. As a fluid-driven fracture approaches the preexisting fractures, the stress state around the fracture tip is disturbed, and this may cause complex fracture interaction behavior [[50](#page-16-0)–[52\]](#page-16-0). In this study, four cases of hydraulic fracturing simulations were conducted to explore the fracture propagation behavior in rocks with pre-existing natural fractures.

Fig. [13](#page-11-0) shows the evolution of injection fluid pressure using different fracturing fluids in the fractured rock sample for two horizontal stress differences. It is observed that the breakdown pressure and propagation pressure in $SC\text{-}CO₂$ fracturing are much lower compared to those in viscous fluid fracturing. This finding is similar to that observed in the intact rock sample. However, the pre-existing natural fractures cause the increase of breakdown pressure in the fractured rock sample compared to that in the intact rock sample. Note that the effect of natural fractures near wellbore on the breakdown pressure may be smaller in the real case because of the alternative path in the third dimension. Figs. [14](#page-11-0) and [15](#page-12-0) display the hydraulic fracture profiles when the injection volumes of the two fracturing fluid are the same. Fracture propagation deviates from the

Fig. 10 a The history of injection fluid pressure in case 1-S1 and case 2-S1 at $t = 1.5$ s. b The evolution of tip location on the y-axis in case 1-S1 and case 2-S1 at $t = 1.5$ s

Fig. 11 a The history of injection fluid pressure in case 1-S2 and case 2-S2 at $t = 0.9$ s. b The evolution of tip location on the y-axis in case 1-S2 and case 2-S2 at $t = 0.9$ s

direction of the maximum confining stress through the interaction with natural fractures when the stress difference is low (see Fig. [14](#page-11-0)). When the stress difference is high, the hydraulic fracture induced by viscous fluid tends to propagate toward the preexisting fractures. Once reaching the pre-existing fractures, the hydraulic fracture prefers to propagate along the pre-existing fractures as shown in Fig. [15](#page-12-0) a. The asymmetric fractures driven by $SC-CO₂$ first propagate along the pre-existing natural fracture and then branch off and turn to the direction of the maxi-mum confining stress (see Fig. [15](#page-12-0) b).

It is obvious that $SC-CO₂$ induces a much longer fracture path. The number of accumulated microcracks is obtained to evaluate the ability to stimulate natural fractures as listed in Table [5](#page-12-0). Compared with viscous fracturing fluid, more microcracks corresponding [to](https://cn.bing.com/dict/search?q=to&FORM=BDVSP6&mkt=zh-cn) the broken parallel bond are generated using $SC-CO₂$ at a low and high stress difference. When the stress difference is relatively high, $SCCO₂$ fracturing

Fig. 12 Simulation results of hydraulic fracturing for random models. a Comparison between the approximate solution and the numerical solution for hydraulic fracture aperture. b Comparison between the approximate

generates 168 microcracks in the pre-existing natural fractures, and 117 of them connect to the hydraulic fracture. However, in the case with viscous fluid, only 109 microcracks are generated and 57 of them connect to the hydraulic fracture. The results indicate that $SC\text{-}CO$ ₂ fracturing tends to achieve a more complex and productive fracture network than fracturing with highviscosity fluid. The results are consistent with the findings in the laboratory experiments [\[28,](#page-15-0) [33](#page-16-0)].

5 Discussion

5.1 Breakdown pressure

A fluid-driven fracture in rock usually propagates in mode I [\[7](#page-15-0)]. In the case of a two-dimensional fracture in an infinite plane as shown in Fig. [16](#page-13-0) a, the stress intensity factor at the

solution and the numerical solution for the distance in the y direction between two fracture tips

Fig. 13 The history of injection fluid pressure of a case N1-S1 and case N2-S1 and b case N1-S2 and case N2-S2

fracture tip is as follows [[53](#page-16-0)]:

$$
K_{1} = \frac{\sqrt{2l}}{\sqrt{\pi}} \int_{0}^{\frac{1}{2l}} \frac{p(\xi)}{\sqrt{\frac{1}{4}l^{2}-\xi^{2}}} d\xi
$$
\n(13)

$$
K_{\rm I} = \frac{F}{\sqrt{\frac{1}{2}\pi l}}\tag{14}
$$

$$
K_{\rm I} = p\sqrt{\frac{1}{2}\pi l} \tag{15}
$$

where $p(\xi)$ is the pressure applied on the fracture surface and l is the fracture length. The fracture propagates once the mode I stress intensity factor at the fracture tip is equal to the fracture toughness.

In two particular cases, case 1 with a pair of point loads (F) was applied to the fracture (see Fig. [16](#page-13-0)b) and case 2 with the uniform pressure load (p) acting throughout the whole fracture (see Fig. $16c$ $16c$) [[54](#page-16-0)]. Then, Eq. (13) can be simplified into the following two expressions:

Fig. 14 Fracture profiles at $t = 2.0$ s. a Case N1-S1. b Case N2-S1. The black and green lines represent the hydraulic fracture and the pre-existing fracture, respectively. The blue and red lines indicate the active

microcracks (in pre-existing fractures) connecting and disconnecting to the hydraulic fracture, respectively

Fig. 15 Fracture profiles at $t = 0.9$ s. a Case N1-S2. b Case N2-S2

sample. Therefore, when the fracture toughness and the finite length are respectively $K_{\text{IC}} = 2.66 \text{ Pa} \cdot \text{m}^{1/2}$ and $l = 0.032 \text{ m}$, the back-calculated pressure is 11.9 MPa based on Eq. ([15](#page-11-0)). Taking the minimum principal pressure (10 MPa) into consideration, the estimated injection pressure is about 21.9 MPa when the fracture initiates. Then, the breakdown pressure is 21.9 MPa or larger, since the fracture initiation could occur before the breakdown [\[19](#page-15-0)]. The above results for the breakdown pressure are consistent with the numerical results. Besides, because of the absence of a circular borehole, the effect of stress difference on the breakdown pressure is minimized. The existence of preset fracture with a finite length accounts for the small difference in the breakdown pressure in these four cases (case 1-S1, case 1-S2, case 2-S1, and case 2-S2).

To investigate the influence of fluid leak-off, an additional viscous fluid fracturing case (case L1-S1) considering fluid leak-off with a low stress difference was carried out, and a comparison with zero fluid leak-off case (case 1-S1) was made. The simulation parameters of case L1-S1 are the same as those of case 1-S1 except that the fluid flow into the rock matrix is turned on in case L1-S1. The injection fluid pressure and fracture length with respect to flow time are shown in

Fig. [17](#page-13-0) a. The breakdown pressure in case L1-S1 and case 1-S1 is 21.0 MPa and 24.1 MPa, respectively. There is a decrease of breakdown pressure if the fluid leak-off is involved. The fluid infiltration before fracture initiation causes an additional pore pressure around the borehole [\[35](#page-16-0), [37\]](#page-16-0) (see Fig. [17](#page-13-0) b). The pore pressure reduces the effective stress and effective toughness around the borehole, resulting in a lower initiation pressure. However, the additional energy dissipation due to the viscous fluid infiltration leads to a higher fracture propagation pressure after the breakdown. It is noted that the additional energy dissipation by fluid infiltration could become less significant if a smaller fluid viscosity is applied.

5.2 Fracture propagation process

The fluid pressure acting on the fracture surface varies depending on the regimes of hydraulic fracture propagation [[13](#page-15-0), [14\]](#page-15-0). The fracture propagations in the viscositydominated regime and the toughness-dominated regime are in a state close to Fig. [16](#page-13-0) b and c, respectively [\[26](#page-15-0)]. Fig. [18](#page-14-0) displays the minimum principal stress contour after the hydraulic fracture propagation in case 1-S1 and case 2-S1. The black line represents the traces of the hydraulic fracture.

Fig. 16 Fracture growth from a viewpoint of mode I stress intensity factor at the fracture tip under a distributed load, b point load, and c uniform load

Case 1-S1 using viscous fluid shows a much larger variation of stress from the injection point to the tip than case 2-S1 with $SCCO₂$ due to the difference of pressure distribution inside the fracture. With the uniform fluid pressure in the fracture for case 2-S1, the information of injection point and the ability to guide the geometrical evolution of the fracture tip could be lost. Besides, the fracture growth in the toughness-dominated regime is more sensitive to the inherent heterogeneity of rock [\[44\]](#page-16-0). Thus, the fracture propagation direction in case 2-S1 is more variable, resulting in larger fracture tortuosity and asymmetry. It should be noted that the simulation cases are performed in two-dimensional in this study. In reality, the more alternative paths in the third dimension have a more significant impact on the fracture profiles when using different fracturing fluids, and $SC\text{-}CO₂$ has potential to generate more complex fracture networks.

Based on Eq. (15) (15) (15) , K_I increases with the fracture length (*l*) if the injection pressure (p) is fixed. In reality, the injection fluid pressure decreases once the K_I at the fracture tip is equal to K_{IC} and the fracture tip extends. After each fracture extension, the pressure in the fracture needs to accumulate again before the next extension. Thus, the

fracture growth induced by low-viscosity fluid like SC- $CO₂$ is less smooth and continuous with a series of jumps as observed in Figs. [10](#page-9-0) and [11.](#page-10-0)

5.3 Effect of fluid compressibility

 $SCCO₂$ has higher compressibility than traditional hydraulic fracturing fluids such as slick water. The relationship between $SC-CO₂$ density and pressure is almost linear for the value greater than 35 MPa when the temperature is constant [\[55](#page-16-0)–[57\]](#page-16-0). Based on the assumption of linear-compressibility law, Wang et al. [\[56](#page-16-0), [57\]](#page-16-0) found that the compressibility affects fracture growth only at the very beginning period in the storage-dominated regime $(M \text{ and } K \text{ in Fig. 4})$ $(M \text{ and } K \text{ in Fig. 4})$ $(M \text{ and } K \text{ in Fig. 4})$, and a slightly greater fracture length is observed due to the fluid expansion. Besides, the influence of fluid compressibility can be neglected in the permeable rock when the fluid leak-off dominates (\tilde{M} and \tilde{K} in Fig. [4\)](#page-6-0). The numerical study by Zhou and Burbey [\[34\]](#page-16-0) indicated that the effect of fluid viscosity can dominate and mask the effect of fluid compressibility. Therefore, the compressibility of $SC\text{-}CO₂$ is the minor factor

Fig. 17 a The history of injection pressure and the distance between two tips in the y direction in case 1-S1 and case L1-S1 at $t = 1.5$ s. **b** The distribution of pore pressure before fracture initiation. The gray hollow

circles represent particles with dry pore space, and the filled circles show those with saturated pore spaces. The filled color indicates the magnitude of the locally averaged pore pressure at the positions of particles

Fig. 18 The minimum principal stress contour after the hydraulic fracture propagation in a case 1-S1 and b case 2-S1. The black line represents the traces of the hydraulic fracture

and it is valid that the compressibility is not considered in this study.

6 Conclusions

In this study, a new DEM-based numerical algorithm is developed to investigate the micromechanical behavior of $SCCO₂$ fracturing in both intact and fractured rock samples, and hydraulic fracturing cases using 100 cp fracturing fluid are also performed for comparison. In the toughness-dominated regime, the pressure gradient inside the fracture can be ignored and the pressure is treated as uniform. The new algorithm shows high accuracy based on a comparison with the reference solution for a plane strain KGD hydraulic fracture. Moreover, it has much higher computational efficiency, which is more than 100 times than the traditional hydromechanical scheme for hydraulic fracturing with low fluid viscosity. The following conclusions can be drawn from this numerical study.

- 1. The hydraulic fracture propagation in intact rock driven by $SC-CO₂$ in the toughness-dominated regime tends to be less smooth and continuous, more asymmetric, and tortuous than that driven by high-viscosity fluid.
- 2. The fluid pressure acting on the fracture surface varies for the high- and low-viscosity fracturing fluids. The fluid pressure in the whole fracture driven by $SC-CO₂$ is uniform, while large pressure gradient exists in the fracture induced by viscous fluid.
- 3. In the cases using $SC\text{-}CO₂$, the uniform pressure inside the fracture causes a much smaller variation of stress and makes the information of injection point and the ability to

guide the geometrical evolution of the fracture tip lost. Therefore, the fracture propagation direction is more variable in $SC-CO₂$ fracturing, resulting in a more tortuous and asymmetric fracture path.

- 4. For $SC-CO₂$ fracturing, the injection fluid pressure decreases once the K_I at the fracture tip is equal to K_{IC} and the fracture tip extends. After each fracture extension, the pressure in the fracture needs to accumulate again before the next extension. Thus, the fracture growth induced by low-viscosity fluid like $SC\text{-}CO₂$ is less smooth and continuous.
- 5. The breakdown pressure increases with the viscosity of fracturing fluid. The difference of breakdown pressure is less profound in this study. The existence of the preset fracture with a finite length eases the influence of principal stress and fluid viscosity. In reality, the low-viscosity fluid like $SC-CO₂$ can lead to a much lower breakdown pressure with the circular borehole and the fluid leak-off involved.
- 6. In the fractured rock mass, the fracture propagation could deviate from the direction of the maximum confining stress under the low stress difference. Compared with viscous fluid, $SC\text{-}CO₂$ tends to cause a more complex fracture network with a longer main fracture in the rock matrix and more activated microcracks in the pre-existing natural fractures.
- 7. The fluid infiltration before fracture initiation could reduce the effective stress and the effective toughness around the borehole, resulting in a lower initiation pressure. However, the additional energy dissipation due to viscous fluid infiltration could lead to a higher fracture propagation pressure after the breakdown.

It is also worth mentioning that there are some limitations in this study and future work is required. The influence of leak-off on the breakdown pressure and propagation pressure is discussed using viscous fluid in this study. For further investigation on the fluid infiltration of $SCCO₂$, the improvement of the new algorithm for the toughness-dominated regime is required. The high injection rate and low viscosity of $SC-CO₂$ can lead to high Reynolds numbers and potential turbulence in the fracture [\[58](#page-16-0)], which could be considered in the future work.

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