

PITT'S INEQUALITY AND LOGARITHMIC UNCERTAINTY PRINCIPLE FOR THE DUNKL TRANSFORM ON \mathbb{R}

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Abstract. We establish Pitt's inequality and deduce Beckner's logarithmic uncertainty principle for the Dunkl transform on \mathbb{R} . Also, we prove Stein–Weiss inequality for the Dunkl–Bessel potentials.

1. Introduction

In the Euclidean case Pitt's inequality for the Fourier transform [1,2], is given for $f \in \mathcal{S}(\mathbb{R}^n)$ (the Schwartz space) and $0 \leq \beta < n$, by

$$(1) \quad \int_{\mathbb{R}^n} |y|^{-\beta} |\widehat{f}(y)|^2 dy \\ \leq \pi^\beta \left[\Gamma\left(\frac{n-\beta}{4}\right) / \Gamma\left(\frac{n+\beta}{4}\right) \right]^2 \int_{\mathbb{R}^n} |x|^\beta |f(x)|^2 dx.$$

This inequality plays an important role for which some uncertainty principles hold. One of these uncertainty principles is the well-known Beckner's logarithmic uncertainty principle [1], that is, for every $f \in \mathcal{S}(\mathbb{R}^n)$,

$$(2) \quad \int_{\mathbb{R}^n} \ln(|x|) |f(x)|^2 dx + \int_{\mathbb{R}^n} \ln(|y|) |\widehat{f}(y)|^2 dy \\ \geq (\psi(n/4) - \ln \pi) \int_{\mathbb{R}^n} |f(x)|^2 dx,$$

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where

$$\psi(t) = \frac{d}{dt} [\ln \Gamma(t)].$$

Recently, Omri [9] has established an analogue of the Beckner logarithmic inequality for the Hankel transform.

In this paper, we consider $\alpha \geq -1/2$ and denote by $L_\alpha^p(\mathbb{R})$, $1 \leq p \leq \infty$, the space of measurable functions f on \mathbb{R} such that

$$\|f\|_{L_\alpha^p} := \begin{cases} \left(\int_{\mathbb{R}} |f(y)|^p |y|^{2\alpha+1} dy \right)^{1/p} < \infty, & 1 \leq p < \infty, \\ \operatorname{ess\,sup}_{y \in \mathbb{R}} |f(y)| < \infty, & p = \infty. \end{cases}$$

For $f \in L_\alpha^1(\mathbb{R})$ the Dunkl transform is defined (see [5]) by

$$\mathcal{F}_\alpha(f)(y) := \frac{1}{2^{\alpha+1}\Gamma(\alpha+1)} \int_{\mathbb{R}} E_\alpha(-ixy) f(x) |y|^{2\alpha+1} dy, \quad y \in \mathbb{R},$$

where $E_\alpha(-ixy)$ denotes the Dunkl kernel (for more details see the next section).

Many uncertainty principles have already been proved for the Dunkl transform, namely by Rösler [11] and Shimeno [12] who established (by two different methods) the Heisenberg–Pauli–Weyl inequality for the Dunkl transform. Kawazoe and Mejjaoli [8] gave some related versions of the uncertainty principle for the Dunkl transform (Cowling–Price’s theorem, Miyachi’s theorem, Beurling’s theorem and Donoho–Stark’s theorem). The author [15] proved a general form of the Heisenberg–Pauli–Weyl inequality for the Dunkl transform.

Building on the ideas of Beckner [1] and Omri [9] we show a Pitt’s inequality for the Dunkl transform (Theorem 1) and we deduce logarithmic uncertainty inequality for the Dunkl transform \mathcal{F}_α (Theorem 2). This inequality generalizes the Beckner’s logarithmic uncertainty inequality given by (2).

The Pitt’s inequality for the Dunkl transform also leads to a Stein–Weiss inequality in the Dunkl setting (Theorem 5).

This paper is organized as follows. In Section 2, we give Pitt’s inequality for the Dunkl transform \mathcal{F}_α and deduce Beckner’s logarithmic uncertainty inequality for \mathcal{F}_α . The last section is devoted to prove Stein–Weiss inequality in the Dunkl setting.

2. Logarithmic uncertainty principle for the Dunkl transform

The Dunkl operator Λ_α , $\alpha \geq -1/2$, associated with the reflection group \mathbb{Z}_2 on \mathbb{R} :

$$\Lambda_\alpha f(x) := \frac{d}{dx} f(x) + \frac{2\alpha + 1}{x} \left[\frac{f(x) - f(-x)}{2} \right],$$

is the operator devised by Dunkl in connection with a generalization of the classical theory of spherical harmonics.

For $\alpha \geq -1/2$ and $y \in \mathbb{R}$, the initial problem:

$$\Lambda_\alpha f(x) = iyf(x), \quad f(0) = 1,$$

has a unique analytic solution $E_\alpha(ixy)$ called Dunkl kernel [4,10] given by

$$E_\alpha(ixy) = j_\alpha(xy) + \frac{ixy}{2(\alpha + 1)} j_{\alpha+1}(xy), \quad x \in \mathbb{R},$$

where

$$j_\alpha(xy) := \Gamma(\alpha + 1) \sum_{n=0}^{\infty} \frac{(-1)^n (xy)^{2n}}{2^{2n} n! \Gamma(n + \alpha + 1)},$$

is the spherical Bessel function of order α (see [19]).

For $x, y \in \mathbb{R}$, the Dunkl kernel $E_\alpha(ixy)$ has the following Bochner-type representation (see [4,10])

$$E_\alpha(ixy) = \frac{\Gamma(\alpha + 1)}{\sqrt{\pi} \Gamma(\alpha + 1/2)} \int_{-1}^1 e^{ixyt} (1 - t^2)^{\alpha-1/2} (1 + t) dt.$$

In particular, we have

$$(3) \quad |E_\alpha(ixy)| \leq 1, \quad x, y \in \mathbb{R}.$$

The Dunkl kernel gives rise to an integral transform, which is called Dunkl transform on \mathbb{R} , and was introduced by Dunkl in [5], where already many basic properties were established. Dunkl's results were completed and extended later on by de Jeu [6]. The Dunkl transform of a function f in $L^1_\alpha(\mathbb{R})$, is

$$\mathcal{F}_\alpha(f)(y) := \frac{1}{2^{\alpha+1} \Gamma(\alpha + 1)} \int_{\mathbb{R}} E_\alpha(-ixy) f(x) |x|^{2\alpha+1} dx, \quad y \in \mathbb{R}.$$

We notice that $\mathcal{F}_{-1/2}$ agrees with the Fourier transform \mathcal{F} that is given by

$$\mathcal{F}(f)(x) := (2\pi)^{-1/2} \int_{\mathbb{R}} e^{-ixy} f(x) dx, \quad y \in \mathbb{R}.$$

The Dunkl transform of a function $f \in L^1_\alpha(\mathbb{R})$ could be computed via the associated Hankel transform \mathcal{H}_α that is

$$\mathcal{F}_\alpha(f)(y) = \mathcal{H}_\alpha(f_e)(|y|) - iy\mathcal{H}_{\alpha+1}\left(\frac{f_o}{\cdot}\right)(|y|), \quad y \in \mathbb{R},$$

where $f_e(x) = \frac{1}{2}(f(x) + f(-x))$, $f_o(x) = \frac{1}{2}(f(x) - f(-x))$ and

$$\mathcal{H}_\alpha(f_e)(|y|) := \frac{1}{2^\alpha\Gamma(\alpha + 1)} \int_0^\infty f_e(r)j_\alpha(|y|r) r^{2\alpha+1} dr.$$

More details for the Hankel transform are collected in [3,16].

Some of the properties of Dunkl transform \mathcal{F}_α are collected below (see [5,6]).

(a) The Dunkl transform \mathcal{F}_α is a topological isomorphism from $\mathcal{S}(\mathbb{R})$ onto itself, and from $\mathcal{S}'(\mathbb{R})$ onto itself.

(b) For all $f \in L^1_\alpha(\mathbb{R})$, we have $\mathcal{F}_\alpha(f) \in L^\infty_\alpha(\mathbb{R})$, and

$$\|\mathcal{F}_\alpha(f)\|_{L^\infty_\alpha} \leq \frac{1}{2^{\alpha+1}\Gamma(\alpha + 1)} \|f\|_{L^1_\alpha}.$$

(c) *Inversion theorem:* Let $f \in L^1_\alpha(\mathbb{R})$ such that $\mathcal{F}_\alpha(f) \in L^1_\alpha(\mathbb{R})$. Then

$$f(x) = \mathcal{F}_\alpha(\mathcal{F}_\alpha(f))(-x) \quad \text{a.e. } x \in \mathbb{R}.$$

(d) *Plancherel theorem:* The Dunkl transform \mathcal{F}_α extends uniquely to an isometric isomorphism of $L^2_\alpha(\mathbb{R})$ onto itself. In particular, we have

$$(4) \quad \|f\|_{L^2_\alpha} = \|\mathcal{F}_\alpha(f)\|_{L^2_\alpha}, \quad f \in L^2_\alpha(\mathbb{R}).$$

In [9], Omri proved the following Pitt's inequality for the Hankel transform, that is, for $f \in \mathcal{S}_e(\mathbb{R})$ (the Schwartz space of even functions) and $0 \leq \beta < 2\alpha + 2$,

$$(5) \quad \int_0^\infty y^{-\beta} |\mathcal{H}_\alpha(f)(y)|^2 y^{2\alpha+1} dy \leq A_{\alpha,\beta} \int_0^\infty x^\beta |f(x)|^2 x^{2\alpha+1} dx,$$

where

$$(6) \quad A_{\alpha,\beta} = 2^{-\beta} \left[\Gamma\left(\frac{2\alpha + 2 - \beta}{4}\right) / \Gamma\left(\frac{2\alpha + 2 + \beta}{4}\right) \right]^2.$$

In the following, we extend the Pitt's inequalities (1) and (5) to a more general case.

THEOREM 1 (Pitt's inequality). Let $0 \leq \beta < 2\alpha + 2$ and let $f \in \mathcal{S}(\mathbb{R})$, then

$$\int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_{\alpha}(f)(y)|^2 |y|^{2\alpha+1} dy \leq M_{\alpha,\beta} \int_{\mathbb{R}} |x|^{\beta} |f(x)|^2 |x|^{2\alpha+1} dx,$$

with $M_{\alpha,\beta} = \sup(A_{\alpha,\beta}, A_{\alpha+1,\beta})$, where $A_{\alpha,\beta}$ is the constant given by (6).

PROOF. Let $0 \leq \beta < 2\alpha + 2$ and let $f \in \mathcal{S}(\mathbb{R})$, then

$$\begin{aligned} \int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_{\alpha}(f)(y)|^2 |y|^{2\alpha+1} dy &= \int_{\mathbb{R}} |y|^{-\beta} |\mathcal{H}_{\alpha}(f_e)(|y|)|^2 |y|^{2\alpha+1} dy \\ + \int_{\mathbb{R}} |y|^{-\beta} \left| \mathcal{H}_{\alpha+1} \left(\frac{f_o}{\cdot} \right) (|y|) \right|^2 |y|^{2\alpha+3} dy &\leq 2 \int_0^{\infty} y^{-\beta} |\mathcal{H}_{\alpha}(f_e)(y)|^2 y^{2\alpha+1} dy \\ &\quad + 2 \int_0^{\infty} y^{-\beta} \left| \mathcal{H}_{\alpha+1} \left(\frac{f_o}{\cdot} \right) (y) \right|^2 y^{2\alpha+3} dy. \end{aligned}$$

Then by (5) we obtain

$$\begin{aligned} \int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_{\alpha}(f)(y)|^2 |y|^{2\alpha+1} dy &\leq 2A_{\alpha,\beta} \int_0^{\infty} x^{\beta} |f_e(x)|^2 x^{2\alpha+1} dx \\ &\quad + 2A_{\alpha+1,\beta} \int_0^{\infty} x^{\beta} \left| \frac{f_o(x)}{x} \right|^2 x^{2\alpha+3} dx \\ &\leq 2 \sup(A_{\alpha,\beta}, A_{\alpha+1,\beta}) \int_0^{\infty} x^{\beta} [|f_e(x)|^2 + |f_o(x)|^2] x^{2\alpha+1} dx. \end{aligned}$$

Since

$$\int_0^{\infty} x^{\beta} [|f_e(x)|^2 + |f_o(x)|^2] x^{2\alpha+1} dx = \frac{1}{2} \int_{\mathbb{R}} |x|^{\beta} |f(x)|^2 |x|^{2\alpha+1} dx,$$

we obtain

$$\int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_{\alpha}(f)(y)|^2 |y|^{2\alpha+1} dy \leq \sup(A_{\alpha,\beta}, A_{\alpha+1,\beta}) \int_{\mathbb{R}} |x|^{\beta} |f(x)|^2 |x|^{2\alpha+1} dx,$$

which completes the proof. \square

In [9], Omri proved the following logarithmic uncertainty principle for the Hankel transform, that is, for $f \in \mathcal{S}_e(\mathbb{R})$,

$$(7) \quad \int_0^{\infty} \ln(x) |f(x)|^2 x^{2\alpha+1} dx + \int_0^{\infty} \ln(y) |\mathcal{H}_{\alpha}(f)(y)|^2 y^{2\alpha+1} dy$$

$$\geq \left[\psi \left(\frac{\alpha + 1}{2} \right) + \ln 2 \right] \int_0^\infty |f(x)|^2 x^{2\alpha+1} dx.$$

In the next, we extend the logarithmic uncertainty principles (2) and (7) to a more general case.

THEOREM 2 (logarithmic uncertainty inequality). *For every $f \in \mathcal{S}(\mathbb{R})$,*

$$\int_{\mathbb{R}} \ln |x| |f(x)|^2 |x|^{2\alpha+1} dx + \int_{\mathbb{R}} \ln |y| |\mathcal{F}_\alpha(f)(y)|^2 |y|^{2\alpha+1} dy \geq D_\alpha \|f\|_{L^2_\alpha}^2,$$

$$D_\alpha = \frac{1}{2} \left[\psi \left(\frac{\alpha + 1}{2} \right) + \psi \left(\frac{\alpha + 2}{2} \right) + 2 \ln 2 \right] - \frac{1}{2} \left| \psi \left(\frac{\alpha + 1}{2} \right) - \psi \left(\frac{\alpha + 2}{2} \right) \right|.$$

PROOF. Let $\varphi(x, \beta) = |x|^{\beta+2\alpha+1} |f(x)|^2$, then for every

$$\beta \in]-\alpha - 1, \alpha + 1[$$

and $0 < |x| < 1$, we have

$$\left| \frac{\partial}{\partial \beta} \varphi(x, \beta) \right| \leq |x|^\alpha |\ln |x|| |f(x)|^2.$$

However, for every real number σ such that $0 < \sigma < \alpha + 1$, the function $x \rightarrow |x|^\sigma |\ln |x||$ is bounded for $0 < |x| < 1$ so that

$$\int_{-1}^1 |x|^\alpha |\ln |x|| |f(x)|^2 dx < \infty.$$

In the same way, for every $\beta \in]-\alpha - 1, \alpha + 1[$ and $|x| > 1$, we have

$$\left| \frac{\partial}{\partial \beta} \varphi(x, \beta) \right| \leq |x|^{3\alpha+2} |\ln |x|| |f(x)|^2,$$

and since $f \in \mathcal{S}(\mathbb{R})$, then the function $x \rightarrow |x|^{3\alpha+2} |\ln |x|| |f(x)|^2$ is also integrable over $|x| \geq 1$. This justifies the differentiation under the integral sign and shows that for every $\beta \in]-\alpha - 1, \alpha + 1[$

$$\frac{\partial}{\partial \beta} \left(\int_{\mathbb{R}} |x|^\beta |f(x)|^2 |x|^{2\alpha+1} dx \right) = \int_{\mathbb{R}} |x|^\beta \ln |x| |f(x)|^2 |x|^{2\alpha+1} dx.$$

By the same way, one can easily see that for every $\beta \in]-\alpha - 1, \alpha + 1[$

$$\begin{aligned} & \frac{\partial}{\partial \beta} \left(\int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_\alpha(f)(y)|^2 |y|^{2\alpha+1} dy \right) \\ &= - \int_{\mathbb{R}} |y|^{-\beta} \ln |y| |\mathcal{F}_\alpha(f)(y)|^2 |y|^{2\alpha+1} dy. \end{aligned}$$

Now let ϕ be the function defined on $]-\alpha - 1, \alpha + 1[$ by

$$\phi(\beta) = \int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_\alpha(f)(y)|^2 |y|^{2\alpha+1} dy - M_{\alpha,\beta} \int_{\mathbb{R}} |x|^\beta |f(x)|^2 |x|^{2\alpha+1} dx.$$

Then ϕ is well defined on $]-\alpha - 1, \alpha + 1[$, however Pitt's inequality implies that $\phi(\beta) \leq 0$ for every $0 \leq \beta < \alpha + 1$, and according to Plancherel's theorem, $\phi(0) = 0$. Since ϕ is differentiable at 0^+ , then it follows that $\phi'(0^+) \leq 0$, and by a basic calculus, we get

$$\begin{aligned} & \int_{\mathbb{R}} \ln |x| |f(x)|^2 |x|^{2\alpha+1} dx + \int_{\mathbb{R}} \ln |y| |\mathcal{F}_\alpha(f)(y)|^2 |y|^{2\alpha+1} dy \\ & \geq - \frac{\partial M_{\alpha,\beta}}{\partial \beta} \Big|_{\beta=0^+} \|f\|_{L_\alpha^2}^2. \end{aligned}$$

But

$$\begin{aligned} & \frac{\partial M_{\alpha,\beta}}{\partial \beta} \Big|_{\beta=0^+} = \frac{1}{2} \lim_{\beta \rightarrow 0^+} \frac{A_{\alpha,\beta} + A_{\alpha+1,\beta} - 2}{\beta} + \frac{1}{2} \lim_{\beta \rightarrow 0^+} \left| \frac{A_{\alpha,\beta} - A_{\alpha+1,\beta}}{\beta} \right| \\ &= \frac{1}{2} \left(\frac{\partial A_{\alpha,\beta}}{\partial \beta} + \frac{\partial A_{\alpha+1,\beta}}{\partial \beta} \right) \Big|_{\beta=0^+} + \frac{1}{2} \left| \frac{\partial A_{\alpha,\beta}}{\partial \beta} - \frac{\partial A_{\alpha+1,\beta}}{\partial \beta} \right| \Big|_{\beta=0^+} \\ &= -\frac{1}{2} \left[\psi \left(\frac{\alpha+1}{2} \right) + \psi \left(\frac{\alpha+2}{2} \right) + 2 \ln 2 \right] + \frac{1}{2} \left| \psi \left(\frac{\alpha+1}{2} \right) - \psi \left(\frac{\alpha+2}{2} \right) \right|, \end{aligned}$$

which completes the proof. \square

3. Stein–Weiss inequality for the Dunkl–Bessel potentials

The Dunkl transform allows us to define a generalized translation operators on $L_\alpha^2(\mathbb{R})$ by setting

$$\mathcal{F}_\alpha(\tau_x f)(y) = E_\alpha(ixy) \mathcal{F}_\alpha(f)(y), \quad y \in \mathbb{R}^d.$$

This is the definition of Thangavelu and Xu given in [17].

Note that from (3) and (4), the definition makes sense, and we have

$$\|\tau_x f\|_{L^2_\alpha} \leq \|f\|_{L^2_\alpha}, \quad f \in L^2_\alpha(\mathbb{R}).$$

Rösler [10] introduced the Dunkl translation operators for f in $L^1_\alpha(\mathbb{R})$, by

$$\tau_x f(y) = \begin{cases} \int_0^\pi \left[f_e((x, y)_\theta) + f_o((x, y)_\theta) \frac{x+y}{(x, y)_\theta} \right] d\nu_{x,y}(\theta), & (x, y) \neq (0, 0) \\ f(y), & x = 0, \end{cases}$$

where $(x, y)_\theta = \sqrt{x^2 + y^2 - 2|xy| \cos \theta}$ and

$$d\nu_{x,y}(\theta) = \frac{\Gamma(\alpha + 1)}{\sqrt{\pi} \Gamma(\alpha + \frac{1}{2})} [1 - \operatorname{sgn}(xy) \cos \theta] \sin^{2\alpha} \theta d\theta.$$

More details for the Dunkl translation operators are collected in [13,14].

Let β be a real number such that $0 < \beta < 2\alpha + 2$. The Dunkl-type Riesz potentials $I_{\alpha,\beta} f$ are defined by (see [18]):

$$I_{\alpha,\beta} f(x) := (d_{\alpha,\beta})^{-1} \int_{\mathbb{R}} \tau_x(|\cdot|^{\beta-2\alpha-2})(-y) f(y) |y|^{2\alpha+1} dy, \quad f \in \mathcal{S}(\mathbb{R}), \quad x \in \mathbb{R},$$

where

$$d_{\alpha,\beta} := 2^{-\alpha-1+\beta} \frac{\Gamma(\beta/2)}{\Gamma(\alpha + 1 - \frac{\beta}{2})}.$$

Thangavelu and Xu established the following relation between the Dunkl-type Riesz potentials $I_{\alpha,\beta}$ and the Dunkl transform \mathcal{F}_α .

THEOREM 3 ([18], Proposition 4.1). *Let $0 < \beta < 2\alpha + 2$. The identity*

$$(8) \quad \mathcal{F}_\alpha(I_{\alpha,\beta} f)(y) = |y|^{-\beta} \mathcal{F}_\alpha(f)(y)$$

holds in the sense that

$$\int_{\mathbb{R}} I_{\alpha,\beta} f(x) g(x) |x|^{2\alpha+1} dx = \int_{\mathbb{R}} \mathcal{F}_\alpha(f)(y) |y|^{-\beta} \mathcal{F}_\alpha(g)(y) |y|^{2\alpha+1} dy,$$

whenever $f, g \in \mathcal{S}(\mathbb{R})$.

COROLLARY 1. *Let $\beta, \gamma > 0$ such that $\beta + \gamma < 2\alpha + 2$. Then for every $f \in \mathcal{S}(\mathbb{R})$, the Dunkl-type Riesz potential $I_{\alpha,\beta}$ satisfies the semigroup property*

$$(9) \quad I_{\alpha,\beta}(I_{\alpha,\gamma} f) = I_{\alpha,\beta+\gamma}(f).$$

THEOREM 4 ([7], Theorem 1.1). *Let $0 < \beta < 2\alpha + 2$ and let $1 < p < \frac{2\alpha+2}{\beta}$. Then the Dunkl-type Riesz potential $I_{\alpha,\beta}$ is bounded from $L_{\alpha}^p(\mathbb{R})$ into $L_{\alpha}^q(\mathbb{R})$ if and only if*

$$\frac{1}{p} - \frac{1}{q} = \frac{\beta}{2\alpha + 2}.$$

COROLLARY 2. *Let $f, g \in \mathcal{S}(\mathbb{R})$ and let $0 < \beta < 2\alpha + 2$. The Dunkl-type Riesz potential $I_{\alpha,\beta}$ satisfies the duality property*

$$(10) \quad \int_{\mathbb{R}} I_{\alpha,\beta} f(x) \overline{g(x)} |x|^{2\alpha+1} dx = \int_{\mathbb{R}} f(x) \overline{I_{\alpha,\beta} g(x)} |x|^{2\alpha+1} dx.$$

In particular, for every $f \in \mathcal{S}(\mathbb{R})$,

$$(11) \quad \int_{\mathbb{R}} f(x) \overline{I_{\alpha,\beta} f(x)} |x|^{2\alpha+1} dx = \int_{\mathbb{R}} |I_{\alpha,\beta/2} f(x)|^2 |x|^{2\alpha+1} dx.$$

PROOF. Since $f \in \mathcal{S}(\mathbb{R})$ then f belongs to $L_{\alpha}^2(\mathbb{R}) \cap L_{\alpha}^{\frac{2\alpha+2}{\alpha+1+\beta}}(\mathbb{R})$, so that according to Theorem 4, $I_{\alpha,\beta} f$ belongs to $L_{\alpha}^2(\mathbb{R})$. Therefore by Plancherel formula (4), we deduce that

$$\begin{aligned} \int_{\mathbb{R}} I_{\alpha,\beta} f(x) \overline{g(x)} |x|^{2\alpha+1} dx &= \int_{\mathbb{R}} |y|^{-\beta} \mathcal{F}_{\alpha}(f)(y) \overline{\mathcal{F}_{\alpha}(g)(y)} |y|^{2\alpha+1} dy \\ &= \int_{\mathbb{R}} \mathcal{F}_{\alpha}(f)(y) \overline{\mathcal{F}_{\alpha}(I_{\alpha,\beta} g)(y)} |y|^{2\alpha+1} dy = \int_{\mathbb{R}} f(x) \overline{I_{\alpha,\beta} g(x)} |x|^{2\alpha+1} dx. \end{aligned}$$

The relation (11) is an immediate consequence of the semigroup and duality properties given by (9) and (10). \square

In the Euclidean case [1] and in the Hankel setting [9], Pitt's inequality is derived from Stein–Weiss inequality. In the following we show the opposite.

THEOREM 5 (Stein–Weiss inequality). *Let $0 < \beta < 2\alpha + 2$. Then for every $f \in \mathcal{S}(\mathbb{R})$ we have*

$$\int_{\mathbb{R}} \int_{\mathbb{R}} f(x) \tau_x(|\cdot|^{\beta-2\alpha-2}) (-y) \overline{f(y)} (|x| |y|)^{-\beta/2+2\alpha+1} dx dy \leq B_{\alpha,\beta} \|f\|_{L_{\alpha}^2},$$

where

$$B_{\alpha,\beta} = \frac{\Gamma(\beta/2) M_{\alpha,\beta}}{2^{\alpha-\beta+1} \Gamma(\alpha + 1 - \frac{\beta}{2})}.$$

PROOF. According to (8) and Plancherel formula (4), we have

$$\int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_{\alpha}(f)(y)|^2 |y|^{2\alpha+1} dy = \int_{\mathbb{R}} |I_{\alpha,\beta/2}f(x)|^2 |x|^{2\alpha+1} dx.$$

By using relation (11) we deduce that

$$\begin{aligned} \int_{\mathbb{R}} |y|^{-\beta} |\mathcal{F}_{\alpha}(f)(y)|^2 |y|^{2\alpha+1} dy &= \int_{\mathbb{R}} f(x) \overline{I_{\alpha,\beta}f(x)} |x|^{2\alpha+1} dx \\ &= (d_{\alpha,\beta})^{-1} \int_{\mathbb{R}} f(x) \left[\int_{\mathbb{R}} \tau_x(|\cdot|^{\beta-2\alpha-2}) (-y) \overline{f(y)} |y|^{2\alpha+1} dy \right] |x|^{2\alpha+1} dx. \end{aligned}$$

However, since $f \in \mathcal{S}(\mathbb{R})$ then f and $I_{\alpha,\beta}f$ belong to $L_{\alpha}^2(\mathbb{R})$. Therefore by Hölder's inequality, we deduce that

$$\begin{aligned} \int_{\mathbb{R}} \int_{\mathbb{R}} |f(x) \tau_x(|\cdot|^{\beta-2\alpha-2}) (-y) \overline{f(y)}| (|x||y|)^{2\alpha+1} dy dx \\ \leq d_{\alpha,\beta} \|f\|_{L_{\alpha}^2} \|I_{\alpha,\beta}f\|_{L_{\alpha}^2} < \infty. \end{aligned}$$

Hence, by Fubini's theorem and Theorem 1, we deduce that

$$\begin{aligned} \int_{\mathbb{R}} \int_{\mathbb{R}} f(x) \tau_x(|\cdot|^{\beta-2\alpha-2}) (-y) \overline{f(y)} (|x||y|)^{2\alpha+1} dx dy \\ \leq \frac{\Gamma(\beta/2)M_{\alpha,\beta}}{2^{\alpha-\beta+1}\Gamma(\alpha+1-\frac{\beta}{2})} \int_{\mathbb{R}} |x|^{\beta} |f(x)|^2 |x|^{2\alpha+1} dx. \end{aligned}$$

Replacing f by $|x|^{-\beta/2}f$ in the preceding inequality we obtain the desired result. \square

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