Complete self-shrinkers confined into some regions of the space

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Abstract We study geometric properties of complete non-compact bounded self-shrinkers and obtain natural restrictions that force these hypersurfaces to be compact. Furthermore, we observe that, to a certain extent, complete self-shrinkers intersect transversally a hyperplane through the origin. When such an intersection is compact, we deduce spectral information on the natural drifted Laplacian associated to the self-shrinker. These results go in the direction of verifying the validity of a conjecture by H.D. Cao concerning the polynomial volume growth of complete self-shrinkers. A finite strong maximum principle in case the self-shrinker is confined into a cylindrical product is also presented.

Keywords Bounded self-shrinkers · Hyperplane intersection · Weighted manifolds · Drifted Laplacian

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1 Introduction

By a self-shrinker "based at" $x_0 \in \mathbb{R}^{m+1}$ we mean a connected, isometrically immersed hypersurface $x: \Sigma^m \to \mathbb{R}^{m+1}$ whose mean curvature vector field **H** satisfies the equation

$$(x - x_0)^{\perp} = -\mathbf{H},$$

where $(\cdot)^{\perp}$ denotes the projection on the normal bundle of Σ . Note that we are using the convention

$$\mathbf{H} = \mathrm{tr}_{\Sigma}\mathbf{A},$$

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where the second fundamental form of the immersion is defined as the generalized Hessian

$$\mathbf{A} = Ddx.$$

With this convention, if Σ is oriented by the outer unit normal ν and we let

$$\mathbf{H} = H v$$
,

then Σ is mean-convex provided $H \leq 0$, and furthermore, the self-shrinker equation takes the scalar form

$$\langle x - x_0, v \rangle = -H.$$

In this paper we shall consider only self-shrinkers based at $0 \in \mathbb{R}^{m+1}$. Natural examples of complete, properly embedded self-shrinkers are the cylindrical products

$$\mathcal{C}_{\sqrt{k}}^{k,m-k} = \mathbb{S}_{\sqrt{k}}^k \times \mathbb{R}^{m-k}, \quad k = 0, \dots, m,$$
(1.1)

which include, as extreme cases, the sphere $\mathbb{S}_{\sqrt{m}}^m$ and all the hyperplanes through the origin of \mathbb{R}^{m+1} . Actually, according to a classification theorem by Colding and Minicozzi [9], these are the only complete, embedded and mean-convex self-shrinkers with extrinsic polynomial volume growth, i.e.,

$$\operatorname{vol}(\mathbb{B}_R^{m+1} \cap \Sigma) \leq CR^n$$

for some C > 0, $n \in \mathbb{N}$ and for every $R \gg 1$; here \mathbb{B}_R^{m+1} denotes the ball of radius R in the ambient Euclidean space.

We stress that it was conjectured by Cao [6], that every complete self-shrinker has extrinsic polynomial (Euclidean, in fact) volume growth. By a very interesting result due to Cheng and Zhou [8], that completes a previous work by Ding and Xin [10], this is equivalent to the fact that the immersion is proper. Thus, by way of example, if Cao Conjecture was true, then any complete self-shrinker in a ball of \mathbb{R}^{m+1} should be compact. In order to obtain indications on the validity of this conjecture, it is then relevant to understand which geometric constraints are imposed by the assumption that a complete self-shrinker is bounded and to obtain natural and general restrictions that force these hypersurfaces to be compact. For instance, we will prove the following results.

Theorem 1 Let $x : \Sigma^m \to \mathbb{B}^{m+1}_{R_0}(0) \subset \mathbb{R}^{m+1}$ be a complete self-shrinker.

- (a) Assume $|\mathbf{A}| \leq 1$. Then:
 - (a.1) $R_0 \ge \sup_{\Sigma} |\mathbf{H}| = \sqrt{m}$.
 - (a.2) If m = 2, then $\Sigma = \mathbb{S}^2_{\sqrt{2}}$.
 - (a.3) If $m \ge 3$ and Σ is non-compact, then Σ must be connected at infinity, i.e., it has only one end. Moreover, $|\mathbf{A}| < 1$, the universal cover $\tilde{\Sigma}$ enjoys the loops to infinity property along every ray [23], and every finitely generated subgroup of the fundamental group of Σ grows at most polynomially of order m.
- (b) Assume $\lim_{R\to\infty} \sup_{\Sigma\setminus B_R^{\Sigma}} |\mathbf{A}| < 1$. Then Σ is compact.
- (c) Assume $|\mathbf{A}| \in L^p(\Sigma)$, for some $p \ge m$. Then Σ is compact.

More generally, one can try to understand the geometry of self-shrinkers which are confined in a connected region bounded by some dilated cylinder $C_R^{k,m-k}$, $R \ge \sqrt{k}$. In this setting, as a preliminary and simple fact, we observe the validity of the following (finite) strong maximum principle.

Theorem 2 Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be a complete self-shrinker. Assume that $|\mathbf{H}| \leq \sqrt{k}$ and that $x(\Sigma)$ is confined inside the domain bounded by $\mathcal{C}_R^{k,m-k}$. If $x(\Sigma) \cap \mathcal{C}_R^{k,m-k} \neq \emptyset$ then:

(a) R = √k,
(b) x : Σ → C^{k,m-k}_{√k} is a Riemannian covering map. In particular, if k ≥ 2, then Σ = C^{k,m-k}_{√k} in the Riemannian sense.

Actually, when k = 0 and, hence, $C^{0,m}$ is a hyperplane through the origin, it is reasonable to expect that the self-shrinker cannot be located into one of the corresponding half-spaces. We are able to verify that, to a certain extent, this is in fact true. The next result can be considered as a weak half-space theorem for complete self-shrinkers.

Theorem 3 Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be a complete, self-shrinker. Assume that either one of the following assumptions is satisfied:

- (a) Σ has extrinsic polynomial volume growth (equivalently, Σ is properly immersed).
- (b) $|\mathbf{A}|^2 \in L^p(d\mathrm{vol}_f)$ with $|\mathbf{A}|^2 \le 1 + \frac{1}{p}$, for some p > 1.

Then, for every hyperplane Π through the origin of \mathbb{R}^{m+1} , Σ cannot be contained in one of the closed half-spaces determined by Π unless $\Sigma = \Pi$.

Accordingly, and in view of the strong maximum principle, it is also reasonable to assume that some transversal intersection between a self-shrinker and a hyperplane through the origin occurs. When such an intersection is compact, we can obtain information on the spectrum of the natural drifted Laplacian $\Delta_f = \Delta - \langle \nabla, \nabla f \rangle$, with $f = |x|^2/2$.

Theorem 4 Let $i : \Sigma^m \hookrightarrow \mathbb{R}^{m+1}$ be a complete, embedded self-shrinker. Assume that, for some hyperplane $\Pi \approx \mathbb{R}^m$ through the origin, $\Sigma \cap \Pi = K$ is a compact (m-1)-dimensional submanifold. Then:

- (a) for every connected component Σ_1 of $\Sigma \setminus K$ (which is an open submanifold $\Sigma_1 \subset \Sigma$ with $\partial \Sigma_1 \subseteq K$) it holds $\lambda_1(-\Delta_f^{\Sigma_1}) \ge 1$.
- (b) If either Σ is compact or Σ has only one end, then there exists a compact connected component Σ₂ of Σ\K such that λ₁(−Δ^{Σ₂}_f) = 1.
- (c) If Σ_3 is an end of Σ with respect to K and

$$\operatorname{vol}\left(\Sigma_{3} \cap \mathbb{B}_{R}^{m+1}\right) = O\left(e^{\alpha R^{2}}\right), \text{ as } R \to +\infty,$$
$$\leq \alpha < 1/2, \text{ then } \lambda_{1}(-\Delta_{f}^{\Sigma_{3}}) = 1.$$

2 Some notations

for some 0

Throughout the paper we let

$$f = \frac{|x|^2}{2}$$

and we denote by $d \operatorname{vol}_f$ the corresponding weighted volume measure of Σ , i.e.,

$$d\operatorname{vol}_f = e^{-f} d\operatorname{vol}.$$

Thus, $\Sigma_f = (\Sigma, g, d \operatorname{vol}_f)$ is a smooth metric measure space. The weighted measure of the intrinsic geodesic ball $B_R^{\Sigma}(o) = \{p \in \Sigma : d_{\Sigma}(o, p) < R\}$ is given by

$$\operatorname{vol}_f(B_R^{\Sigma}) = \int\limits_{B_R^{\Sigma}} d\operatorname{vol}_f.$$

Note that, obviously,

$$\operatorname{vol}_f(B_R^{\Sigma}(o)) \leq \operatorname{vol}_f(\mathbb{B}_R(x(o)) \cap \Sigma),$$

where \mathbb{B}_R denotes the Euclidean ball.

There is a natural drifted Laplacian on Σ_f defined by

$$\Delta_f = e^f \operatorname{div}(e^{-f} \nabla) = \Delta - \langle \nabla, \nabla f \rangle.$$

It is symmetric on $L^2(dvol_f)$ and, since $\nabla f = x^T$, it can be expressed in the equivalent form

$$\Delta_f = \Delta - \langle \nabla, x^{\mathrm{T}} \rangle = \Delta - \langle \nabla, x \rangle,$$

where x^{T} denotes the tangential component of the immersion.

Recall also that the Bakry–Emery Ricci tensor of Σ_f is defined by

$$\operatorname{Ric}_{f} = \operatorname{Ric} + \operatorname{Hess}(f).$$

Using once again the self-shrinker equation we easily obtain the following very important estimate [21],

$$\operatorname{Ric}_{f} \ge 1 - |\mathbf{A}|^2 \tag{2.1}$$

where **A** denotes the second fundamental tensor of the immersion $x \colon \Sigma^m \to \mathbb{R}^{m+1}$. Indeed, by Gauss equations,

$$\operatorname{Ric} \geq \langle \mathbf{H}, \mathbf{A} \rangle - |\mathbf{A}|^2 g,$$

whereas, by the self-shrinker equation,

$$\operatorname{Hess}(f) = g + \langle x^{\perp}, \mathbf{A} \rangle = g - \langle \mathbf{H}, \mathbf{A} \rangle$$

3 A maximum principle

To begin with, we observe that if a complete self-shrinker with $|\mathbf{A}| \leq 1$ is contained in a ball and it is tangent to the boundary of this ball at a point, then it must be the standard sphere $\mathbb{S}_{\sqrt{m}}^{m}$. The analytic proof is a straightforward application of the maximum principle for subharmonic functions. Later on, in Sect. 4.2, we shall come back on this kind of arguments.

Proposition 5 Let $x : \Sigma^m \to \mathbb{R}^{m+1}$ be a complete bounded self-shrinker with $|\mathbf{H}| \le \sqrt{m}$. If there exist $x_0 \in \Sigma$ such that $|x|(x_0) = \sup_{\Sigma} |x|$, then $|x| \equiv \sqrt{m}$ and Σ is the standard sphere $\mathbb{S}^m_{\sqrt{m}}$.

Proof Recall that, [9, Lemma 3.20],

$$\Delta |x|^2 = 2(m - |\mathbf{H}|^2), \tag{3.1}$$

therefore, by assumption,

$$\Delta |x|^2 \ge 0.$$

Using the strong maximum principle we thus obtain $|x| \equiv c > 0$. This implies that $x : \Sigma^m \to C$ \mathbb{S}_{c}^{m} is a Riemannian covering projection, hence an isometry since \mathbb{S}_{c}^{m} is simply connected. In particular, by the self-shrinker equation, $c = \sqrt{m}$.

The above result can be deduced more geometrically via a suitable application of the usual touching principle. We adopt this viewpoint to obtain the following strong maximum principle for self-shrinkers. Recall that the oriented hypersurface $x: \Sigma^m \to \mathbb{R}^{m+1}$ is called mean-convex at $p \in \Sigma$ if $\mathbf{H}(p) = H(p)v$, where $H(p) \leq 0$ and v is the outward pointing unit normal at p.

Theorem 6 (Maximum principle). Let $\Omega \subset \mathbb{R}^{m+1}$ be a domain such that $i : \partial \Omega \hookrightarrow \mathbb{R}^{m+1}$ is a properly embedded self-shrinker. Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be a complete self-shrinker satisfying $x(\Sigma) \subseteq \overline{\Omega_{\lambda}}$ for some $\lambda > 0$, where $\Omega_{\lambda} = \lambda \Omega$ denotes the λ -dilation of Ω . Assume that $x(\Sigma) \cap \partial \Omega_{\lambda} \neq \emptyset$ and that, for each intersection point x(p), there exist a neighborhood $V \subset \mathbb{R}^{m+1}$ of x(p) and a neighborhood $W \subset \Sigma$ of p such that:

(i) $\partial \Omega \cap \lambda^{-1} V$ is mean convex

(ii) $\sup_{W} |\mathbf{H}_{\Sigma}| \leq \inf_{\lambda^{-1} V \cap \partial \Omega} |\mathbf{H}_{\partial \Omega}|.$

Then

- (a) $\lambda = 1$, (b) $\partial \Omega = S_{\sqrt{k}}^k \times \mathbb{R}^{m-k}$, for some $k \in \{0, ..., m\}$, (c) $x : \Sigma \to \partial \Omega$ is a Riemannian covering map.

In particular, if $\partial \Omega$ is simply connected (e.g., if k > 2 in (b)), then $\Sigma = \partial \Omega$ in the Riemannian sense.

A situation of special interest is obtained by choosing $\partial \Omega$ to be a cylindrical product shrinker $C_{k,m-k}^{k,m-k}$. Note that the case k = m is precisely the content of Proposition 5.

Corollary 7 Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be a complete self-shrinker. Assume that $|\mathbf{H}_{\Sigma}| \leq \sqrt{k}$ and that $x(\Sigma)$ is confined inside the solid cylinder bounded by $\mathcal{C}_R^{k,m-k} = \mathbb{S}_R^k \times \mathbb{R}^{m-k}$. If $x(\Sigma) \cap \mathcal{C}_{R}^{k,m-k} \neq \emptyset$ then

(a) $R = \sqrt{k}$, (b) $x : \Sigma \to C_{\sqrt{k}}^{k,m-k}$ is a Riemannian covering map.

In particular, if $k \ge 2$, then $\Sigma = C_{\sqrt{k}}^{k,m-k}$.

Proof of Theorem 6 Let

$$\mathcal{O} = x^{-1}(\partial \Omega_{\lambda})$$

Since x is smooth and $\partial \Omega_{\lambda}$ is closed in \mathbb{R}^{m+1} , we have that \mathcal{O} is a closed subset of Σ . We claim that \mathcal{O} is also open so that, by a connectedness argument, $\mathcal{O} = \Sigma$, i.e., $x(\Sigma) \subset \partial \Omega_{\lambda}$. To this end, let $p \in \mathcal{O}$. Observe that, by the mean-convexity assumption (i), in a connected neighborhood $\lambda^{-1}U_{x(p)} \subset \partial \Omega$ it holds

$$\mathbf{H}_{\partial\Omega} = H_{\partial\Omega} \nu_{\partial\Omega},$$

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where $H_{\partial\Omega} \leq 0$ and $\nu_{\partial\Omega}$ denotes the exterior pointing unit normal to $\partial\Omega$. Moreover, the rescaling property of the mean curvature tells us that

$$H_{\partial\Omega_{\lambda}}(x(p)) = \lambda^{-1} H_{\partial\Omega}(\lambda^{-1}x(p)).$$

Whence, using the fact that $i : \partial \Omega \hookrightarrow \mathbb{R}^{m+1}$ is a self-shrinker, it is standard to deduce that either $H_{\partial \Omega_{\lambda}} \equiv 0$ in $U_{x(p)}$, or $H_{\partial \Omega_{\lambda}} < 0$ on $U_{x(p)}$; see e.g., the beginning of the proof of [21, Theorem 2]. In the first case, by assumption, we must have $\mathbf{H}_{\Sigma} = 0$ in a neighborhood of p in Σ and the result reduces to a well-known local maximum principle for minimal surfaces. Therefore, from now on, we assume

$$H_{\partial\Omega_{\lambda}} < 0$$
 in $U_{x(p)}$.

Since $x(\Sigma)$ lies inside $\overline{\Omega_{\lambda}}$, then $x(\Sigma)$ must intersect $\partial \Omega_{\lambda}$ tangentially at $p \in \mathcal{O}$ and

$$\nu_{\Sigma}(p) = \nu_{\partial \Omega_{\lambda}}(x(p))$$

the outward pointing unit normal to Ω_{λ} . It follows from the self-shrinker equations for $\partial \Omega$ and Σ , and the rescaling property of the mean curvature, that

$$\mathbf{H}_{\Sigma}(p) = \lambda^2 \mathbf{H}_{\partial \Omega_{\lambda}}(x(p)) = \lambda \mathbf{H}_{\partial \Omega}(\lambda^{-1}x(p)).$$

Combining this latter with assumption (ii) we get

$$\lambda^{2}|\mathbf{H}_{\partial\Omega_{\lambda}}(x(p))| = |\mathbf{H}_{\Sigma}(p)| \le |\mathbf{H}_{\partial\Omega}(\lambda^{-1}x(p))| = \lambda|\mathbf{H}_{\partial\Omega_{\lambda}}(x(p))|.$$

Thus

 $\lambda \leq 1.$

If we write, in a neighborhood of p:

$$\mathbf{H}_{\Sigma} = H_{\Sigma} \nu_{\Sigma} \quad \text{and} \quad \mathbf{H}_{\partial \Omega_{\lambda}} = H_{\partial \Omega_{\lambda}} \nu_{\partial \Omega_{\lambda}},$$

then, by mean convexity of $U_{x(p)}$, by the above equation at p, and by continuity, we have, in a neighborhood of p,

$$H_{\Sigma}, H_{\partial\Omega_{\lambda}}(x) < 0$$

and

$$H_{\Sigma} \ge H_{\partial\Omega}(\lambda^{-1}x) = \lambda H_{\partial\Omega_{\lambda}}(x) \ge H_{\partial\Omega_{\lambda}}(x).$$

We can now apply the usual touching principle and deduce that, actually, $x(\Sigma)$ and $\partial \Omega_{\lambda}$ coincide in a small neighborhood of p. This proves the claim and, as already remarked at the beginning of the proof, $x(\Sigma) \subseteq \partial \Omega_{\lambda}$.

Now, $x : \Sigma \to \partial \Omega_{\lambda}$ is a local isometry between complete manifolds, hence, it is a covering map. In particular, $x(\Sigma) = \partial \Omega_{\lambda}$, and from the equality

$$H_{\Sigma}(p) = \lambda^2 H_{\partial \Omega_{\lambda}}(x(p))$$

we deduce

$$H_{\partial\Omega_{\lambda}}(x) = H_{\Sigma} = \lambda^2 H_{\partial\Omega_{\lambda}}(x),$$

that is

$$\lambda = 1.$$

This shows that $x(\Sigma) = \partial \Omega$. Finally, by assumption (i), $\partial \Omega$ is a properly embedded self-shrinker satisfying $H_{\partial\Omega} \leq 0$ everywhere. Since properly immersed self-shrinkers have polynomial (actually Euclidean) volume growth [8,10], to complete the proof we apply a classification result by Colding and Minicozzi [9, Theorem 0.17].

4 Self-shrinkers in a ball

The aim of this section is to show that certain boundedness conditions on the norm of the second fundamental form prevent the existence of complete, non-compact, bounded self-shrinkers.

4.1 Estimate of the exterior radius

The sphere $\mathbb{S}_{\sqrt{m}}^{m}$ is a self-shrinker of constant mean curvature $-\sqrt{m}$ and contained in the compact ball $\mathbb{B}_{\sqrt{m}}^{m+1}(0)$. Our first remark is that if a complete self-shrinker with controlled intrinsic volume growth is contained in some ball $\mathbb{B}_{R_0}^{m+1}(0)$, then there is an obvious relation between the ray R_0 and the dimension m.

Proposition 8 Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be a complete non-compact self-shrinker whose intrinsic volume growth satisfies

$$R o rac{R}{\log \operatorname{vol}(B_R^{\Sigma})} \notin L^1(+\infty).$$

If $x(\Sigma) \subseteq \overline{\mathbb{B}}_{R_0}^{m+1}(0)$, then

$$R_0 \geq \sup_{\Sigma} |\mathbf{H}| \geq \sqrt{m}.$$

Proof Recall that, by the self-shrinker equation,

$$\Delta_f |x|^2 = 2(m - |x|^2).$$

On the other hand, since

$$c^{-1}d\operatorname{vol}_f \le d\operatorname{vol} \le cd\operatorname{vol}_f$$

for a large enough constant c > 1, then

$$\frac{R}{\log \operatorname{vol}_f(B_R^{\Sigma})} \notin L^1(+\infty)$$

and this implies that the weighted manifold Σ_f enjoys the weak maximum principle at infinity for the drifted Laplacian Δ_f [17,18]. Therefore

$$0 \ge 2\left(m - \sup_{\Sigma} |x|^2\right) \ge 2(m - R_0^2),$$

and the claimed lower estimate on R_0 follows. Now, from the self-shrinker equation we have

$$\sup_{\Sigma} |\mathbf{H}| \le |x| \le R_0.$$

Using this information into Eq. (3.1):

$$\Delta |x|^2 = 2(m - |\mathbf{H}|^2),$$

and noting also that the weak maximum principle at infinity for the Laplacian holds on Σ , we deduce

$$0 \ge 2\left(m - \sup_{\Sigma} |\mathbf{H}|^2\right).$$

This completes the proof.

Note that, by [25, Theorem 2.2] and inequality (2.1), a complete non-compact bounded self-shrinker $x : \Sigma^m \to \mathbb{R}^{m+1}$ with $|\mathbf{A}| \le 1$ satisfies the sharp estimate

$$\operatorname{vol}(B_R^{\Sigma}) \le CR^m.$$
 (4.1)

Moreover, since $|\mathbf{A}| \leq 1$, by the Cauchy–Schwarz inequality we have that $|\mathbf{H}|^2 \leq m$. We can hence specialize Proposition 8 to the following

Corollary 9 Let $x : \Sigma^m \to \mathbb{R}^{m+1}$ be a complete non-compact self-shrinker with $|\mathbf{A}| \le 1$. If $x(\Sigma) \subseteq \overline{\mathbb{B}}_{R_0}^{m+1}(0)$, then

$$R_0 \ge \sup_{\Sigma} |\mathbf{H}| = \sqrt{m}.$$

4.2 Bounded self-shrinkers with $|\mathbf{A}| \leq 1$

As a consequence of the strong maximum principle for the Laplace–Beltrami operator, we observed in Sect. 3 that, for a self-shrinker satisfying $|\mathbf{A}| \leq 1$, hence $|\mathbf{H}| \leq \sqrt{m}$, the norm of the immersion cannot attain a finite maximum unless the shrinker is a round sphere of radius \sqrt{m} . In particular, this applies to any compact self-shrinker with the same bound on the mean curvature. It is by now well understood that parabolicity is a good substitute of compactness. For two-dimensional shrinkers this property is implied by the above condition on the second fundamental form.

Theorem 10 Let $x : \Sigma^2 \to \mathbb{R}^3$ be a complete bounded self-shrinker with $|\mathbf{A}| \leq 1$. Then $\Sigma = \mathbb{S}^2_{\sqrt{2}}$.

Proof Since m = 2, we know from (4.1) that Σ has quadratic intrinsic volume growth, therefore it is parabolic (possibly compact); see e.g., [11]. As in Proposition 5, since $|\mathbf{H}| \le \sqrt{2}$, $|x|^2$ is a bounded subharmonic function and we obtain that $|x| \equiv \text{const.}$ This implies $\Sigma = \mathbb{S}^2_{\sqrt{2}}$.

In higher dimensions, the same control gives information on the topology at infinity of a bounded shrinker.

Theorem 11 Let $x : \Sigma^m \to \mathbb{R}^{m+1}$ be a complete non-compact bounded self-shrinker with $|\mathbf{A}| \leq 1$. Then Σ does not contain a line. In particular, Σ is connected at infinity, i.e., Σ has only one end.

Remark 12 Applying this result to the universal covering of Σ , and using [23–25], we also get the topological information collected in Theorem 1 stated in the Sect. 1.

Proof Assume by contradiction that Σ contains a line. By assumption and (2.1), we have that $Ric_f \ge 0$ with f bounded. Therefore, we can apply the Cheeger–Gromoll–Lichnerowicz splitting theorem [14], and obtain that Σ splits isometrically as the Riemannian product $(N^{m-1} \times \mathbb{R}, g_N + dt \otimes dt)$. Moreover f is constant along the line. Thus

$$\operatorname{Hess}(f)(\partial_t, \partial_t) = 0. \tag{4.2}$$

On the other hand, consider the Simons type equation, see [12, page 292], [9, Lemma 10.8],

$$\frac{1}{2}\Delta_f |\mathbf{A}|^2 + |\mathbf{A}|^2 (|\mathbf{A}|^2 - 1) = |D\mathbf{A}|^2.$$
(4.3)

Since, by assumption, $|\mathbf{A}| \leq 1$ then the strong maximum principle for the drifted Laplacian yields that either (a) $|\mathbf{A}| < 1$ or (b) $|\mathbf{A}| \equiv 1$, on Σ . In case (a), recalling (2.1), we deduce that

$$\operatorname{Ric}_{f}(\partial_{t}, \partial_{t}) = \operatorname{Hess}(f)(\partial_{t}, \partial_{t}) > 0 \text{ on } \Sigma,$$

contradicting (4.2). Suppose that (b) holds, namely, $|\mathbf{A}| \equiv 1$. Using again the Simons equation we get that \mathbf{A} is parallel. We can therefore apply a classification theorem by Lawson, [13, Theorem 4] and deduce that $x(\Sigma)$ is a cylindrical product $\mathbb{S}^k_{\sqrt{k}} \times \mathbb{R}^{m-k}$ with $k = 0, \ldots, m$. Since the self-shrinker is bounded, we conclude that $\Sigma = \mathbb{S}^m_{\sqrt{m}}$, contradicting the assumption that Σ is not compact.

4.3 Bounded self-shrinkers with $\limsup |\mathbf{A}| < 1$

In the two previous results we considered global bounds on the norm of the second fundamental form. The application of the Feller property for Δ_f in combination with the maximum principle at infinity enable us to prevent the existence of complete, non-compact, bounded self-shrinkers even in the case a pinching condition on $|\mathbf{A}|$ is required at infinity. Recall that the weighted manifold Σ_f is said to be Feller if, for some (hence any) smooth domain $\Omega \subset \subset \Sigma_f$ and $\lambda > 0$, the minimal solution h > 0 of the exterior boundary value problem

$$\begin{cases} \Delta_f h = \lambda h & \text{on } \Sigma \setminus \overline{\Omega} \\ h = 1 & \text{on } \partial \Omega \end{cases}$$

satisfies $h(x) \to 0$ as $x \to \infty$; see [20,4]. In particular, we obtain the following

Theorem 13 Let $x : \Sigma^m \to \mathbb{B}_{R_0}^{m+1}(0) \subset \mathbb{R}^{m+1}$ be a complete self-shrinker with $\lim_{R\to\infty} \sup_{\Sigma\setminus B_R^{\Sigma}} |A| < 1$. Then Σ is compact.

Remark 14 Suppose that Σ is compact. Then $\Sigma \setminus B_R^{\Sigma} = \emptyset$ for $R > \operatorname{diam}(\Sigma)$ and, therefore, $\lim_{R\to\infty} \sup_{\Sigma\setminus B_R^{\Sigma}} |A| = -\infty$, proving that the assumption of the theorem is automatically satisfied. Note also that, from a different perspective, the result states that a complete, non-compact, bounded self-shrinker must satisfy the asymptotic condition $\lim_{R\to\infty} \sup_{\Sigma\setminus B_R^{\Sigma}} |A| \ge 1$.

Proof First observe that, since $|\mathbf{A}| \in L^{\infty}(\Sigma)$ and $|\nabla f| = |x^{T}| \leq |x| < R_{0}$, we know by (2.1), Theorem 7 and Theorem 8 in [4] that M is both stochastically complete and Feller with respect to Δ_{f} . Furthermore, by (4.3) and our assumption, we have that $|\mathbf{A}|$ is a bounded nonnegative solution of

$$\Delta_f |\mathbf{A}|^2 \ge \lambda |\mathbf{A}|^2 \tag{4.4}$$

outside a smooth domain $\Omega \subset \Sigma$ and for a suitable $\lambda > 0$. An application of Theorem 2 in [4] permits to deduce that

$$|\mathbf{A}|(x) \to 0, \quad \text{as } x \to \infty.$$
 (4.5)

In the matter of this, note that the proof in [4] actually works for nonnegative solutions at infinity of inequalities of the form (4.4).

On the other hand, using the self-shrinker equation, we compute

$$\operatorname{Hess}(f) = g - \langle \mathbf{H}, \mathbf{A} \rangle.$$

By (4.5) having fixed any ray $\gamma : [0, +\infty) \to \Sigma$, we have

$$\frac{\mathrm{d}^2}{\mathrm{d}t^2}(f \circ \gamma)(t) = \mathrm{Hess}(f)(\dot{\gamma}, \dot{\gamma}) \ge \frac{1}{2}, \quad \text{for} \quad t \gg 1.$$

It follows by integration that $|x|^2 \to +\infty$ along γ and, therefore, $x(\Sigma)$ is unbounded. Contradiction.

4.4 Bounded self-shrinkers with $|\mathbf{A}| \in L^{p \ge m}$

In the next result we switch from L^{∞} to L^{p} conditions on the norm of the second fundamental form. In particular, we show that complete bounded self-shrinkers with finite total curvature must be compact.

Theorem 15 Let $x : \Sigma^m \to \mathbb{R}^{m+1}$ be a complete, bounded self-shrinker satisfying $|\mathbf{A}| \in L^p(d\text{vol})$, for some $p \ge m$. Then Σ is compact.

Proof By contradiction, suppose that Σ is complete and non-compact. To illustrate the argument, let us first consider the case p = m. Since f is bounded and $|\mathbf{H}| \in L^m(\Sigma)$, it is standard to obtain that Σ enjoys the weighted L^2 -Sobolev inequality

$$\left(\int \varphi^{\frac{2m}{m-2}} d\operatorname{vol}_f\right)^{\frac{m-2}{m}} \leq S \int |\nabla \varphi|^2 d\operatorname{vol}_f,$$

for some constant S > 0 and for every $\varphi \in C_c^{\infty}(\Sigma)$. Indeed, first we can absorb the mean curvature term in the Sobolev inequality by Michael and Simon [15], outside a large compact set, then, according to [5], we can extend the resulting Sobolev inequality to all of Σ and, finally, we note that, since *f* is bounded,

$$c^{-1}d\operatorname{vol}_f \leq d\operatorname{vol} \leq c d\operatorname{vol}_f$$

for a large enough constant c > 1.

Now we recall that, using (4.3) and the Kato inequality, we have that the second fundamental form of the self-shrinker satisfies the Simons-type inequality

$$\Delta_f |\mathbf{A}| + |\mathbf{A}|^3 \ge 0.$$

Since $|\mathbf{A}| \in L^m(d \operatorname{vol}_f)$, combining the PDE with the weighted Sobolev inequality gives the Anderson-type decay estimate

$$\sup_{\Sigma \setminus B_{E}^{\Sigma}(o)} |\mathbf{A}| = o(R^{-1}), \quad \text{as} \quad R \to +\infty.$$
(4.6)

This follows, e.g., by adapting to the weighted setting the arguments in [16].

From this uniform estimate it is now standard to get that the immersion x is proper, thus contradicting the assumption that $x(\Sigma)$ is a bounded subset of \mathbb{R}^{m+1} . In fact, we have the following general result that, in the setting of minimal submanifolds of the Euclidean space, traces back to a paper by Anderson [1]; see also Remark 17 below.

Lemma 16 Let $x: (\Sigma^m, g) \to \mathbb{R}^{m+1}$ be a complete, non-compact hypersurface satisfying (4.6). Then x is proper and Σ has finite topological type, i.e., there exists a smooth compact subset $\Omega \subset \subset \Sigma$ such that $\Sigma \setminus \Omega$ is diffeomorphic to the half-cylinder $\partial \Omega \times [0, +\infty)$.

As a matter of fact, the uniform decay condition (4.6) on the second fundamental form, as well as the corresponding structure Lemma, are even too much strong for the desired conclusion to hold. This is illustrated in the next reasonings where we assume the general condition $p \ge m$.

Again, by contradiction, suppose that Σ is complete and non-compact. Since f is bounded, by the self-shrinker equation we get $|\mathbf{H}| \in L^{\infty}$. Whence, we obtain that Σ enjoys the weighted L^2 -Sobolev inequality (with potential term)

$$\left(\int_{\Sigma} \varphi^{\frac{2m}{m-2}} d\operatorname{vol}_f\right)^{\frac{m}{m}} \leq A \int_{\Sigma} |\nabla \varphi|^2 d\operatorname{vol}_f + B \int_{\Sigma} \varphi^2 d\operatorname{vol}_f,$$

for some constants A, B > 0 and for every $\varphi \in C_c^{\infty}(\Sigma)$. Since $|\mathbf{A}|$ is a solution of the semilinear equation

$$\Delta_f |\mathbf{A}| + |\mathbf{A}|^3 \ge 0,$$

and $|\mathbf{A}| \in L^p(d\text{vol}_f) = L^p(d\text{vol})$ for some $p \ge m$, we deduce that (see e.g., [16])

$$\sup_{\Sigma \setminus B_R^{\Sigma}} |\mathbf{A}| = o(1), \quad \text{as} \quad R \to +\infty.$$
(4.7)

Reasoning exactly as in the last part of the proof of Theorem 13 this leads to the fact that $x(\Sigma)$ is unbounded, yielding a contradiction.

Remark 17 The decay assumption (4.6) in Lemma 16 can be considerably relaxed. This was established in [3] where the authors used the notion of tamed submanifolds. We are grateful to Pacelli Bessa for having pointed out this fact to us.

5 Self-shrinkers and hyperplanes through the origin

5.1 Self-shrinkers in a half-space

It is reasonable that a complete self-shrinker has a certain homogeneous distribution around $0 \in \mathbb{R}^{m+1}$ and, therefore, it should intersect every hyperplane through the origin. For compact self-shrinkers this property is easily verified. In fact, more is true. It was proved in Theorem 7.3 of [24] that if the distance between two properly immersed self-shrinkers (either compact or not) is realized, then the self-shrinkers must intersect. In particular, a compact self-shrinker must intersect every hyperplane through the origin, as claimed. Moreover, the intersection must be non-tangential by maximum principle considerations. Summarizing, a compact self-shrinker cannot be contained in one of the half-spaces determined by a hyperplane through the origin. Needless to say, exactly the same proof works for a complete self-shrinker with

polynomial volume growth because, according to [8], it is properly immersed. We are going to recover the same conclusion by using more direct and analytic arguments that are suitable for a generalization to the complete (non-necessarily proper) setting.

Theorem 18 Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be a compact self-shrinker. Then, for every hyperplane Π through the origin of $\mathbb{R}^{m+1}, x(\Sigma)$ cannot be contained in one of the closed halfspaces determined by Π .

Proof Recall that, for a self-shrinker,

$$\Delta_f x = -x.$$

See e.g., [9, Lemma 3.20]. Therefore, if Π has normal equation

$$\Pi: \quad L(y) := \sum_{j=1}^{m+1} a_j y^j = 0, \tag{5.1}$$

we have that the self-shrinker satisfies also

$$\Delta_f L(x) = -L(x). \tag{5.2}$$

Whence, it follows easily that $x(\Sigma)$ cannot be contained in one of the closed half-spaces determined by Π . Indeed, otherwise, we would have that either $L(x) \ge 0$ or $L(x) \le 0$. Without loss of generality, suppose that $L(x) \ge 0$. Then, by the above equation, L(x) would be an *f*-superharmonic function on the compact manifold Σ . By the maximum principle $L \equiv \text{const}$ and by Eq. (5.2) $L \equiv 0$. This means that $x(\Sigma) \subseteq \Pi$ and, by geodesic completeness, $x(\Sigma) = \Pi$. This is clearly impossible because Σ is compact.

A similar conclusion can be obtained for complete self-shrinkers $x: \Sigma^m \to \mathbb{R}^{m+1}$ with a controlled extrinsic geometry. By way of example, suppose that

$$|x| + |\mathbf{A}(p)| \le \sqrt{1 + r(p)^2},$$
 (5.3)

where $r(p) = d_{\Sigma}(p, o)$. Then, for every hyperplane Π through the origin, if $x(\Sigma)$ lies on one side of Π , then

$$\operatorname{dist}_{\mathbb{R}^{m+1}}(\Pi, x(\Sigma)) = 0$$

and the distance is not attained, unless $x(\Sigma) = \Pi$.

Indeed, note that, in light of (2.1), condition (5.3) implies

$$Ric_f \ge -C(1+r^2), \quad |\nabla f| = |x^{\mathrm{T}}| \le \sqrt{1+r^2}.$$

Then, according to Corollary 5.3 in [19], for every $u \in C^2(\Sigma)$ with $\inf_{\Sigma} u = u_* > -\infty$ there exists a sequence $\{p_n\} \subset \Sigma$ along which

$$u(p_n) < u_* + \frac{1}{n}, \quad |\nabla u|(p_n) < \frac{1}{n}, \quad \Delta_f u(p_n) > -\frac{1}{n}.$$

Now, as in the compact case, if $x(\Sigma)$ lies on one side of Π , we can assume that $L(x) \ge 0$ where L(y) is defined in (5.1). Evaluating (5.2) along $\{p_n\}$ we deduce that $\inf_{\Sigma} L(x) = 0$, as desired. The second conclusion is a consequence of the strong minimum principle for positive super-solutions of $\Delta_f + 1$.

In the next theorem we point out natural geometric conditions that permit to recover the full conclusion of the compact case.

Theorem 19 Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be a complete, non-compact self-shrinker. Assume that either one of the following assumptions is satisfied:

- (a) Σ has (extrinsic) polynomial volume growth.
- (b) $\operatorname{vol}_f(B_R^{\Sigma}) = O(R^2)$ as $R \to \infty$. (c) $|\mathbf{A}|^2 \in L^p(d\operatorname{vol}_f)$ and $|\mathbf{A}|^2 \le 1 + \frac{1}{p}$, for some p > 1.

Then, for every hyperplane Π through the origin, if $x(\Sigma)$ lies on one side of Π , then $x(\Sigma) = \Pi.$

Proof We shall use extensively the notation introduced so far. In particular, the hyperplane Π is described by the normal equation (5.1) and the function L(x) satisfies Eq. (5.2).

Assume we are in the assumptions of (a). Since Σ has polynomial volume growth, then $\operatorname{vol}_f(\Sigma) < +\infty$ and Σ_f is parabolic with respect to the drifted Laplacian Δ_f . Using the above notation, assume without loss of generality that L(x) > 0. By Eq. (5.2) we see that $L(x) \ge 0$ is f-superharmonic, hence it is constant by f-parabolicity. The desired conclusion now follows as in the proof of Theorem 18. Case (b) is completely similar. Assume now that the assumptions in (c) are satisfied. Let $x(\Sigma) \neq \Pi$ and, by contradiction, suppose that $x(\Sigma)$ is contained in a half-space determined by Π . Then, by the strong minimum principle, we can assume that L(x) > 0 is a solution of

$$\Delta_f L + L = 0.$$

Since

$$p(|\mathbf{A}|^2 - 1) \le 1,$$

for some p > 1, we obtain

$$\Delta_f L + p(|\mathbf{A}|^2 - 1)L \le 0.$$

Combining this latter with the Simons-type inequality [obtained from (4.3)]

$$|\mathbf{A}| \{\Delta_f |\mathbf{A}| + |\mathbf{A}|(|\mathbf{A}|^2 - 1)\} \ge |D\mathbf{A}|^2 - |\nabla|\mathbf{A}||^2 \ge 0,$$

and applying Theorem 8 in [21], we conclude that either $|\mathbf{A}| \equiv 1$ or $|\mathbf{A}| \equiv 0$. Using this information into the Simons-type equality (4.3)

$$\frac{1}{2}\Delta_f |\mathbf{A}|^2 + |\mathbf{A}|^2 (|\mathbf{A}|^2 - 1) = |D\mathbf{A}|^2$$

gives that $|D\mathbf{A}| \equiv 0$ and by Lawson classification theorem $x(\Sigma) = \mathbb{S}^{k}_{\sqrt{k}} \times \mathbb{R}^{m-k}$, with $0 \le k \le m$. Since $x(\Sigma)$ must lie on one side of Π we necessarily have k = 0, i.e., $x(\Sigma) = \Pi$, contradiction.

5.2 Bottom of the spectrum of the drifted Laplacian

Once we have understood that, to a certain extent, complete self-shrinkers intersect transversally a hyperplane through the origin, we are going to deduce spectral information on the drifted Laplacian whenever the intersection is compact, and some (extrinsic) volume growth condition is satisfied.

The intuition for the general result contained in Theorem 22 relies on the following two examples. Recall that, by definition, the bottom of the spectrum of $-\Delta_f$ on a domain $\Omega \subseteq \Sigma$,

with Dirichlet boundary conditions, is defined by

$$\lambda_1(-\Delta_f^{\Omega}) = \inf_{v \in C_c^{\infty}(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla v|^2 d\operatorname{vol}_f}{\int_{\Omega} v^2 d\operatorname{vol}_f}.$$

The bottom of the spectrum λ_1 is an eigenvalue of $-\Delta_f$ if there exists a function $u \in \text{Dom}(-\Delta_f^{\Omega})$ such that

$$-\Delta_f u = \lambda_1 u$$
 on Ω ,

where

$$\operatorname{Dom}(-\Delta_f^{\Omega}) = \{ u \in W_0^{1,2}(\Omega, d\operatorname{vol}_f) : \Delta_f u \in L^2(\Omega, d\operatorname{vol}_f) \}$$

is the domain of (the Friedrichs extension of) $-\Delta_f$ originally defined on $C_c^{\infty}(\Omega)$. For future purposes, we also recall that if Σ is complete and $\partial\Omega$ is compact then,

$$u \in W^{1,2}(\Omega, d\operatorname{vol}_f)$$
 and $u = 0$ on $\partial \Omega \Rightarrow u \in W^{1,2}_0(\Omega, d\operatorname{vol}_f)$.

Indeed, the interesting case occurs when Ω is non-compact, i.e., an exterior domain, in the complete manifold Σ . Let $0 \le \phi_R \le 1$ be the standard family of cut-off functions supported in the ball B_{2R}^{Σ} , satisfying $\phi_R = 1$ on B_R^{Σ} and such that $|\nabla \phi_R| \le 2/R$. Then, $u_R = u\phi_R \in W_0^{1,2}(\Omega)$ and it is easy to verify that $u_R \to u$ in $W^{1,2}(\Omega, d\text{vol}_f)$, as $R \to \infty$.

Example 20 Consider the self-shrinker sphere $\mathbb{S}_{\sqrt{m}}^m$. Then, each hyperplane Π through the origin divides $\mathbb{S}_{\sqrt{m}}^m$ into half-spheres isometric to $\mathbb{S}_{\sqrt{m}}^m = \mathbb{S}_{\sqrt{m}}^m \cap \{y_{m+1} > 0\}$. Since $f(x) \equiv m/2$, it holds

$$\lambda_1 \left(-\Delta_f^{+ \mathbb{S}_{\sqrt{m}}^m} \right) = \lambda_1 \left(-\Delta_f^{+ \mathbb{S}_{\sqrt{m}}^m} \right) = \frac{1}{m} \lambda_1 \left(-\Delta_1^{+ \mathbb{S}_1^m} \right) = 1;$$

see e.g., [7].

Example 21 Consider the self-shrinker cylinder $C = \mathbb{S}_{\sqrt{m-1}}^{m-1} \times \mathbb{R}$. Then the hyperplane $\Pi = \{y_{m+1} = 0\}$ intersects C along the sphere $\mathbb{S}_{\sqrt{m-1}}^{m-1}$ and divides C into two half-cylinders isometric to $C_+ = \mathbb{S}_{\sqrt{m-1}}^{m-1} \times \mathbb{R}_+$. These are the ends of Σ . We claim that

$$\lambda_1\left(-\Delta_f^{\mathcal{C}_+}\right) = 1.$$

Indeed, since

$$f = \frac{|x|^2}{2} = \frac{m-1}{2} + \frac{x_{m+1}^2}{2},$$

we have the decomposition

$$\Delta_f^{\mathcal{C}_+} = \Delta_{f}^{\mathbb{S}_{\sqrt{m-1}}^{m-1}} + \Delta_{t^2/2}^{\mathbb{R}_+}$$

and, therefore,

$$\begin{split} \lambda_1(-\Delta_f^{\mathcal{C}_+}) &= \lambda_1 \left(-\Delta_{\sqrt{m-1}}^{\mathbb{S}_{\sqrt{m-1}}^{m-1}} \right) + \lambda_1 \left(-\Delta_{t^2/2}^{\mathbb{R}_+} \right) \\ &= 0 + \lambda_1 \left(-\Delta_{t^2/2}^{\mathbb{R}_+} \right). \end{split}$$

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Now, the Ornstein–Uhlenbeck operator $\Delta_{t^2/2}^{\mathbb{R}}$ on $(\mathbb{R}_+, e^{-t^2/2}dt)$ satisfies

$$\lambda_1\left(-\Delta_{t^2/2}^{\mathbb{R}_+}\right) = 1.$$

See e.g., the lecture notes [22] for the basic theory and more advanced topics on the Ornstein–Uhlenbeck operator and its semigroup. Indeed, u(t) = t is a smooth, positive function on \mathbb{R}_+ satisfying

$$\Delta_{t^2/2}^{\mathbb{R}_+} u = u'' - tu' = -u \tag{5.4}$$

so that, by (the weighted version of) Barta's theorem [2],

$$\lambda_1\left(\Delta_{t^2/2}^{\mathbb{R}_+}\right) \ge \inf_{\mathbb{R}_+} \frac{-\Delta_{t^2/2}^{\mathbb{R}_+}u}{u} = 1.$$

On the other hand, $u \in W^{1,2}(\mathbb{R}_+, e^{-t^2/2}dt)$, therefore, by (5.4), $\Delta_{t^2/2}^{\mathbb{R}_+} u \in L^2(\mathbb{R}_+, e^{-t^2/2}dt)$. Furthermore, u(0) = 0. It follows that $u \in \text{Dom}(-\Delta_f^{\mathbb{R}_+})$ is also a Dirichlet eigenfunction of the Ornstein–Uhlenbeck operator on \mathbb{R}_+ .

Abstracting from the previous examples we are now ready to state the following general result.

Theorem 22 Let $i : \Sigma^m \hookrightarrow \mathbb{R}^{m+1}$ be a complete, embedded self-shrinker. Assume that, for some hyperplane $\Pi \approx \mathbb{R}^m$ through the origin, $\Sigma \cap \Pi = K$ is a compact (m-1)-dimensional submanifold. Then:

(a) for every connected component Σ_1 of $\Sigma \setminus K$ (which is an open submanifold $\Sigma_1 \subset \Sigma$ with $\partial \Sigma_1 \subseteq K$) it holds

$$\lambda_1\left(-\Delta_f^{\Sigma_1}\right) \ge 1.$$

(b) If either Σ is compact or Σ has only one end, then there exists a bounded connected component Σ₂ of Σ\K such that

$$\lambda_1\left(-\Delta_f^{\Sigma_2}\right) = 1.$$

(c) If Σ_3 is an end of Σ with respect to K with extrinsic volume growth

$$\operatorname{vol}\left(\Sigma_{3} \cap \mathbb{B}_{R}^{m+1}\right) = O\left(e^{\alpha R^{2}}\right), \quad \text{as} \quad R \to +\infty,$$
(5.5)

for some $0 \le \alpha < 1/2$, then

$$\lambda_1\left(-\Delta_f^{\Sigma_3}\right) = 1.$$

Remark 23 The conclusion in (a) holds regardless of the fact that the intersection K is compact.

Remark 24 Note that condition (5.5) in (c) is actually equivalent to the (only apparently less general) polynomial volume growth condition. Indeed, it is easy to see that (5.5) implies that $vol_f(\Sigma_3) < +\infty$ (see Lemma 25 below) and minor changes to the proofs of Theorem 2.2 in [10], and of Theorem 2.1 and Theorem 4.1 in [8] show that the equivalences in [8] can be localized to a given end. In particular, under assumption (5.5), Σ_3 is proper and of extrinsic

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polynomial (Euclidean) volume growth. For the sake of completeness we sketch out here the proof of the fact that a properly immersed end has Euclidean volume growth. The proofs of the remaining implications can be easily adapted from the original ones. Suppose that $\tilde{\Sigma}$ is a properly immersed end of a complete noncompact self-shrinker $x : \Sigma^m \to \mathbb{R}^{m+1}$. To prove that $\tilde{\Sigma}$ must have Euclidean extrinsic volume growth observe that, since $\partial \tilde{\Sigma}$ is compact and properly immersed we can find a regular value r_0 such that $\{p \in \tilde{\Sigma} : |x(p)| = r_0\}$ does not intersect $\partial \tilde{\Sigma}$. Then we can define for $r > r_0$ the set $D_r := \{p \in \tilde{\Sigma} : r_0 < |x(p)| < r\}$. Since the immersion is proper, letting $h = \frac{|x|^2}{4}$, we can define for $t > 0, r > r_0$,

$$I(t) = \frac{1}{t^{\frac{m}{2}}} \int_{\overline{D}_r} e^{-\frac{h}{t}} d\text{vol.}$$

Since on a self-shrinker

$$|\nabla h|^2 - h \le 0$$

$$\Delta_h h + h \le \frac{m}{2}$$

we obtain that, if $t \ge 1$,

$$I'(t) \leq -t^{-\frac{m}{2}-1} \int_{\overline{D}_r} \operatorname{div}\left(e^{-\frac{h}{t}} \nabla h\right) d\operatorname{vol}.$$

At a regular value r of |x|, for $t \ge 1$, by Stokes' Theorem we have thus

$$I'(t) \leq -t^{-\frac{m}{2}-1} \left| \int_{\{|x|=r\}} \left\langle e^{-\frac{h}{t}} \nabla h, \frac{\nabla h}{|\nabla h|} \right\rangle d\text{vol} \right|$$
$$- \int_{\{|x|=r_0\}} \left\langle e^{-\frac{h}{t}} \nabla h, \frac{\nabla h}{|\nabla h|} \right\rangle d\text{vol} \right|$$
$$\leq t^{-\frac{m}{2}-1} \int_{\{|x|=r_0\}} e^{-\frac{h}{t}} |\nabla h| d\text{vol}.$$

Integrating on $[1, r^2]$, with $r^2 > r_0^2 \ge 1$, we get

$$e^{-\frac{1}{4}}r^{-m}\int_{\overline{D}_r}d\text{vol} \leq \int_{\overline{D}_r}e^{-h}d\text{vol} + \int_{1}^{r^2}t^{-\frac{m}{2}-1}e^{-\frac{r_0^2}{4r}}dt\int_{\{|x|=r_0\}}|\nabla h|d\text{vol}.$$
 (5.6)

Proceeding now as in [8] we can conclude that, for any positive integer N, we have

$$\int_{\overline{D}_{r+N}} e^{-h} d\operatorname{vol} \le \left[\prod_{i=0}^{N} \frac{1}{1 - e^{-(r+i)}} \right] \left[\int_{\overline{D}_{r-1}} e^{-h} d\operatorname{vol} + e^{-r} \int_{1}^{r^{2}} t^{-\frac{m}{2} - 1} e^{-\frac{r^{2}}{4t}} dt \frac{r_{0}}{2} \operatorname{vol}_{m-1}(\{|x| = r_{0}\}) \right].$$

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This implies that $\int_{\tilde{\Sigma}} e^{-h} d\text{vol} < +\infty$ and the desired Euclidean extrinsic volume growth of $\tilde{\Sigma}$ follows from (5.6).

Proof of Theorem 22 Let Π be represented by the normal equation

$$\Pi : L(y) := \sum_{j=1}^{m+1} a_j y^j = 0.$$

Recall that, for every self-shrinker,

 $\Delta_f x = -x.$

It follows that

$$\Delta_f L(x) + L(x) = 0$$
, on Σ .

In particular, this equation holds on Σ_1 . Moreover, since Σ_1 is contained in one of the open halfspaces determined by Π , then either L < 0 or L > 0 on Σ_1 . Thus, up to changing the sign of L, we can assume L > 0 and using (the weighted version of) Barta's theorem we deduce

$$\lambda_1\left(-\Delta_f^{\Sigma_1}\right) \ge \inf_{\Sigma_1} \frac{-\Delta_f L}{L} = 1$$

This proves (a).

Suppose now that Σ is non-compact and has only one end. We claim that there exists a compact connected component Σ_2 of $\Sigma \setminus K$. In this case, since L = 0 on $\partial \Sigma_2 \subseteq K$, we deduce that L is an eigenfunction of $\Delta_f^{\Sigma_2}$ corresponding to the eigenvalue +1. When combined with (a) this clearly implies that $\lambda_1(-\Delta_f^{\Sigma_2}) = 1$, completing the proof of (b). To prove the claim, we first observe that $\Sigma \setminus K$ cannot be connected. Indeed, by contradiction, suppose the contrary. Then Σ must be contained in one of the closed half-spaces determined by Π and intersects Π tangentially along K. Without loss of generality, we can assume that $L(x) \ge 0$ on Σ and L(x) = 0 on K. Since $\Delta_f L(x) = -L(x) \le 0$ on Σ , by the strong minimum principle we get $L(x) \equiv 0$ on Σ , i.e., $\Sigma \subseteq \Pi$. Actually, $\Sigma = \Pi$ by geodesic completeness and this clearly prevents $K = \Sigma \cap \Pi$ to be compact, contradiction. Thus, $\Sigma \setminus K$ has at least two connected components. Since we are assuming that Σ has one end, at most one of them can be unbounded. We therefore find a bounded component $\Sigma_2 \subseteq \Sigma$ of $\Sigma \setminus K$, as claimed.

It remains to prove (c). The argument is completely similar to the above. According to (a), $\lambda_1(-\Delta_f^{\Sigma_3}) \ge 1$ and $L(x) \ge 0$ is a solution of

$$\begin{cases} \Delta_f L(x) + L(x) = 0, & \text{on } \Sigma_3 \\ L = 0, & \text{on } \partial \Sigma_3 \subseteq K. \end{cases}$$

To conclude that, in fact, $\lambda_1(-\Delta_f^{\Sigma_3}) = 1$ it suffices to show that $L \in \text{Dom}(\Sigma_3)$. Since L = 0 on the compact boundary $\partial \Sigma_3$, we have to show that $L \in W^{1,2}(\Sigma_3, d\text{vol}_f)$. To this aim, we simply note that

$$\frac{|L(x)|}{\sqrt{\sum a_j^2}} \le |x|,$$

and

$$\frac{|\nabla L(x)|}{\sqrt{\sum a_j^2}} \le 1.$$

Therefore, we can apply the next simple lemma. This proves (c) and completes the proof of the theorem. $\hfill \Box$

Lemma 25 Let $x: \Sigma^m \to \mathbb{R}^{m+1}$ be any hypersurface satisfying

$$\operatorname{vol}\left(\Sigma \cap \mathbb{B}_{R}^{m+1}\right) = O\left(e^{\alpha R^{2}}\right), \quad \text{as} \quad R \to +\infty.$$

for some $0 \le \alpha < 1/2$. Then, for every polynomial $\mathcal{P}(t)$ and for every $0 \le \beta < 1/2 - \alpha$,

$$\mathcal{P}(|x|)e^{\beta|x|^2} \in L^1(d\mathrm{vol}_f).$$

Proof Note that, by assumption, there exists t > 1 such that

$$\frac{1}{2}-t^2\alpha-\beta>0.$$

Now, we simply compute

$$\begin{split} \int_{\Sigma} |x|^{p} e^{\beta |x|^{2}} d\operatorname{vol}_{f} &= \int_{\Sigma} |x|^{p} e^{-\left(\frac{1}{2} - \beta\right)|x|^{2}} d\operatorname{vol} \\ &= C_{1} + C_{2} \sum_{n=0}^{+\infty} \int_{\Sigma \cap \left(\mathbb{B}_{t^{n+1}}^{m+1} \setminus \mathbb{B}_{t^{n}}^{m+1}\right)} |x|^{p} e^{-\left(\frac{1}{2} - \beta\right)|x|^{2}} d\operatorname{vol} \\ &\leq C_{1} + C_{2} \sum_{n=0}^{+\infty} t^{pn+p} e^{-\left(\frac{1}{2} - \beta\right)t^{2n}} \operatorname{vol}\left(\Sigma \cap \mathbb{B}_{t^{n+1}}^{m+1}\right) \\ &\leq C_{1} + C_{2} \sum_{n=0}^{+\infty} t^{pn+p} e^{-\left(\frac{1}{2} - t^{2}\alpha - \beta\right)t^{2n}} \\ &< +\infty. \end{split}$$

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