Geodesic distance for right invariant Sobolev metrics of fractional order on the diffeomorphism group. II

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Abstract The geodesic distance vanishes on the group $\operatorname{Diff}_c(M)$ of compactly supported diffeomorphisms of a Riemannian manifold M of bounded geometry, for the right invariant weak Riemannian metric which is induced by the Sobolev metric H^s of order $0 \le s < \frac{1}{2}$ on the Lie algebra $\mathfrak{X}_c(M)$ of vector fields with compact support.

Keywords Diffeomorphism group · Geodesic distance · Sobolev metrics of non-integral order

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1 Introduction

In the article [1] we studied right invariant metrics on the group $\mathrm{Diff}_c(M)$ of compactly supported diffeomorphisms of a manifold M, which are induced by the Sobolev metric H^s of order s on the Lie algebra $\mathfrak{X}_c(M)$ of vector fields with compact support. We showed that for $M=S^1$ the geodesic distance on $\mathrm{Diff}(S^1)$ vanishes if and only if $s\leq \frac{1}{2}$. For other manifolds, we showed that the geodesic distance on $\mathrm{Diff}_c(M)$ vanishes for $M=\mathbb{R}\times N$, $s<\frac{1}{2}$ and for $M=S^1\times N$, $s\leq \frac{1}{2}$, with N being a compact Riemannian manifold.

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Now we are able to complement this result by: *The geodesic distance vanishes on* $\operatorname{Diff}_c(M)$ *for any Riemannian manifold M of bounded geometry, if* $0 \le s < \frac{1}{2}$.

We believe that this result holds also for $s = \frac{1}{2}$, but we were able to overcome the technical difficulties only for the manifold $M = S^1$, in [1]. We also believe that it is true for the regular groups $\mathrm{Diff}_{\mathcal{H}^\infty}(\mathbb{R}^n)$ and $\mathrm{Diff}_{\mathcal{S}}(\mathbb{R}^n)$ as treated in [8], and for all Virasoro groups, where we could prove it only for s = 0 in [2].

In Sect. 2, we review the definitions for Sobolev norms of fractional orders on diffeomorphism groups as presented in [1] and extend them to diffeomorphism groups of manifolds of bounded geometry. Section 3 is devoted to the main result.

2 Sobolev metrics H^s with $s \in \mathbb{R}$

2.1 Sobolev metrics H^s on \mathbb{R}^n

For $s \ge 0$ the Sobolev H^s -norm of an \mathbb{R}^n -valued function f on \mathbb{R}^n is defined as

$$||f||_{H^{s}(\mathbb{R}^{n})}^{2} = ||\mathcal{F}^{-1}(1+|\xi|^{2})^{\frac{s}{2}}\mathcal{F}f||_{L^{2}(\mathbb{R}^{n})}^{2}, \tag{1}$$

where \mathcal{F} is the Fourier transform

$$\mathcal{F}f(\xi) = (2\pi)^{-\frac{n}{2}} \int_{\mathbb{R}^n} e^{-i\langle x,\xi\rangle} f(x) \, dx,$$

and ξ is the independent variable in the frequency domain. An equivalent norm is given by

$$||f||_{\overline{H}^{s}(\mathbb{R}^{n})}^{2} = ||f||_{L^{2}(\mathbb{R}^{n})}^{2} + ||\xi|^{s} \mathcal{F} f||_{L^{2}(\mathbb{R}^{n})}^{2}.$$
(2)

The fact that both norms are equivalent is based on the inequality

$$\frac{1}{C} \left(1 + \sum_{j} |\xi_{j}|^{s} \right) \le \left(1 + \sum_{j} |\xi_{j}|^{2} \right)^{\frac{s}{2}} \le C \left(1 + \sum_{j} |\xi_{j}|^{s} \right),$$

holding for some constant C. For s > 1 this says that all ℓ^s -norms on \mathbb{R}^{n+1} are equivalent. But the inequality is true also for 0 < s < 1, even though the expression does not define a norm on \mathbb{R}^{n+1} . Using any of these norms we obtain the Sobolev spaces with non-integral s

$$H^{s}(\mathbb{R}^{n}) = \{ f \in L^{2}(\mathbb{R}^{n}) : ||f||_{H^{s}(\mathbb{R}^{n})} < \infty \}.$$

We will use the second version of the norm in the proof of the theorem, since it will make calculations easier.

2.2 Sobolev metrics for Riemannian manifolds of bounded geometry

Following [13, Section 7.2.1] we will now introduce the spaces $H^s(M)$ on a manifold M. If M is not compact we equip M with a Riemannian metric g of bounded geometry which exists by [5]. This means that

- (I) The injectivity radius of (M, g) is positive.
- (B_{∞}) Each iterated covariant derivative of the curvature is uniformly *g*-bounded: $\|\nabla^i R\|_{\mathcal{E}} < C_i$ for $i = 0, 1, 2, \ldots$



The following is a compilation of special cases of results collected in [3, Chapter 1], who treats Sobolev spaces only for integral order.

Proposition [4,6,10] *If* (M, g) *satisfies* (I) *and* (B_{∞}) *then the following holds:*

- (1) (M, g) is complete.
- (2) There exists $\varepsilon_0 > 0$ such that for each $\varepsilon \in (0, \varepsilon_0)$ there is a countable cover of M by geodesic balls $B_{\varepsilon}(x_{\alpha})$ such that the cover of M by the balls $B_{2\varepsilon}(x_{\alpha})$ is still uniformly locally finite.
- (3) Moreover, there exists a partition of unity $1 = \sum_{\alpha} \rho_{\alpha}$ on M such that $\rho_{\alpha} \geq 0$, $\rho_{\alpha} \in C_{c}^{\infty}(M)$, $supp(\rho_{\alpha}) \subset B_{2\varepsilon}(x_{\alpha})$, and $|D_{u}^{\beta}\rho_{\alpha}| < C_{\beta}$ where u are normal (Riemann exponential) coordinates in $B_{2\varepsilon}(x_{\alpha})$.
- (4) In each $B_{2\varepsilon}(x_{\alpha})$, in normal coordinates, we have $|D_u^{\beta}g_{ij}| < C'_{\beta}$, $|D_u^{\beta}g^{ij}| < C''_{\beta}$, and $|D_u^{\beta}\Gamma^m_{ij}| < C'''_{\beta}$, where all constants are independent of α .

We can now define the H^s -norm of a function f on M:

$$||f||_{H^{s}(M,g)}^{2} = \sum_{\alpha=0}^{\infty} ||(\rho_{\alpha}f) \circ \exp_{x_{\alpha}}||_{H^{s}(\mathbb{R}^{n})}^{2}$$

$$= \sum_{\alpha=0}^{\infty} ||\mathcal{F}^{-1}(1+|\xi|^{2})^{\frac{s}{2}} \mathcal{F}((\rho_{\alpha}f) \circ \exp_{x_{\alpha}})||_{L^{2}(\mathbb{R}^{n})}^{2}.$$

If M is compact the sum is finite. Changing the charts or the partition of unity leads to equivalent norms by the proposition above, see [13, Theorem 7.2.3]. For integer s we get norms which are equivalent to the Sobolev norms treated in [3, Chapter 2]. The norms depend on the choice of the Riemann metric g. This dependence is worked out in detail in [3].

For vector fields we use the trivialization of the tangent bundle that is induced by the coordinate charts and define the norm in each coordinate as above. This leads to a (up to equivalence) well-defined H^s -norm on the Lie algebra $\mathfrak{X}_c(M)$.

2.3 Sobolev metrics on $Diff_c(M)$

A positive definite weak inner product on $\mathfrak{X}_c(M)$ can be extended to a right-invariant weak Riemannian metric on $\mathrm{Diff}_c(M)$. In detail, given $\varphi \in \mathrm{Diff}_c(M)$ and $X, Y \in T_{\varphi}\mathrm{Diff}_c(M)$ we define

$$G^s_{\varphi}(X,Y) = \langle X \circ \varphi^{-1}, Y \circ \varphi^{-1} \rangle_{H^s(M)}.$$

We are interested solely in questions of vanishing and non-vanishing of geodesic distance. These properties are invariant under changes to equivalent inner products, since equivalent inner products on the Lie algebra

$$\frac{1}{C}\langle X, Y \rangle_1 \le \langle X, Y \rangle_2 \le C\langle X, Y \rangle_1$$

imply that the geodesic distances will be equivalent metrics

$$\frac{1}{C} \mathrm{dist}_1(\varphi, \psi) \leq \mathrm{dist}_2(\varphi, \psi) \leq C \mathrm{dist}_1(\varphi, \psi).$$

Therefore the ambiguity—dependence on the charts and the partition of unity—in the definition of the H^s -norm is of no concern to us.



3 Vanishing geodesic distance

Theorem 3.1 (Vanishing geodesic distance) *The Sobolev metric of order s induces vanishing geodesic distance on Diff*_c(M) *if*:

• $0 \le s < \frac{1}{2}$ and M is any Riemannian manifold of bounded geometry.

This means that any two diffeomorphisms in the same connected component of $Diff_c(M)$ can be connected by a path of arbitrarily short G^s -length.

In the proof of the theorem we shall make use of the following lemma from [1].

Lemma 3.2 [1, Lemma 3.2] Let $\varphi \in Diff_c(\mathbb{R})$ be a diffeomorphism satisfying $\varphi(x) \geq x$ and let T > 0 be fixed. Then for each $0 \leq s < \frac{1}{2}$ and $\varepsilon > 0$ there exists a time-dependent vector field $u_{\mathbb{R}}^{\varepsilon}$ of the form

$$u_{\mathbb{R}}^{\varepsilon}(t,x) = 1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]} * G_{\varepsilon}(x),$$

with $f, g \in C^{\infty}([0, T])$, such that its flow $\varphi^{\varepsilon}(t, x)$ satisfies—independently of ε —the properties $\varphi^{\varepsilon}(0, x) = x$, $\varphi^{\varepsilon}(T, x) = \varphi(x)$ and whose H^{s} -length is smaller than ε , i.e.,

$$Len(\varphi^{\varepsilon}) = \int_{0}^{T} \|u_{\mathbb{R}}^{\varepsilon}(t,\cdot)\|_{H^{s}} dt \leq C \|f^{\varepsilon} - g^{\varepsilon}\|_{\infty} \leq \varepsilon.$$

Furthermore $\{t: f^{\varepsilon}(t) < g^{\varepsilon}(t)\} \subseteq supp(\varphi)$ and there exists a limit function $h \in C^{\infty}([0,T])$, such that $f^{\varepsilon} \to h$ and $g^{\varepsilon} \to h$ for $\varepsilon \to 0$ and the convergence is uniform in t.

Here, $G_{\varepsilon}(x) = \frac{1}{\varepsilon}G_1(\frac{x}{\varepsilon})$ is a smoothing kernel, defined via a smooth bump function G_1 with compact support.

Proof of Theorem 3.1 Consider the connected component $\mathrm{Diff}_0(M)$ of Id, i.e. those diffeomorphisms of $\mathrm{Diff}_c(M)$, for which there exists at least one path, joining them to the identity. Denote by $\mathrm{Diff}_c(M)^{L=0}$ the set of all diffeomorphisms φ that can be reached from the identity by curves of arbitrarily short length, i.e., for each $\varepsilon>0$ there exists a curve from Id to φ with length smaller than ε .

Claim A. $\operatorname{Diff}_c(M)^{L=0}$ is a normal subgroup of $\operatorname{Diff}_0(M)$. Claim B. $\operatorname{Diff}_c(M)^{L=0}$ is a non-trivial subgroup of $\operatorname{Diff}_0(M)$.

By [12] or [7], the group $\operatorname{Diff}_0(M)$ is simple. Thus claims A and B imply $\operatorname{Diff}_c(M)^{L=0} = \operatorname{Diff}_0(M)$, which proves the theorem.

The proof of claim A can be found in [1, Theorem 3.1] and works without change in the case of M being an arbitrary manifold and hence we will not repeat it here. It remains to show that $\mathrm{Diff}_c(M)^{L=0}$ contains a diffeomorphism $\varphi \neq \mathrm{Id}$.

We shall first prove claim B for $M=\mathbb{R}^n$ and then show how to extend the arguments to arbitrary manifolds. Choose a diffeomorphism $\varphi_{\mathbb{R}}\in \mathrm{Diff}_c(\mathbb{R})$ with $\varphi_{\mathbb{R}}(x)\geq x$ and $\mathrm{supp}(\varphi_{\mathbb{R}})\subseteq [1,\infty)$. Then let

$$u_{\mathbb{R}}^{\varepsilon}(t,x) := 1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]} * G_{\varepsilon}(x)$$

be the family of vector fields constructed in Lemma 3.2, whose flows at time T equal $\varphi_{\mathbb{R}}$. We extend the vector field $u_{\mathbb{R}}^{\varepsilon}$ to a vector field $u_{\mathbb{R}}^{\varepsilon}$ on \mathbb{R}^n via

$$u_{\mathbb{R}^n}^{\varepsilon}(t, x_1, \dots, x_n) := (u_{\mathbb{R}}^{\varepsilon}(t, |x|), 0, \dots, 0).$$



The flow of this vector field is given by

$$\varphi_{\mathbb{R}^n}^{\varepsilon}(t, x_1, \dots, x_n) = (\varphi_{\mathbb{R}}^{\varepsilon}(t, |x|), x_2, \dots, x_n),$$

where $\varphi_{\mathbb{R}}^{\varepsilon}$ is the flow of $u_{\mathbb{R}}^{\varepsilon}$. In particular we see that at time t=T

$$\varphi_{\mathbb{P}^n}^{\varepsilon}(t,x_1,\ldots,x_n)=(\varphi_{\mathbb{R}}(|x|),x_2,\ldots,x_n),$$

the flow is independent of ε . So it remains to show that for the length of the path $\varphi_{\mathbb{R}^n}^{\varepsilon}(t,\cdot)$ we have

$$\operatorname{Len}(\varphi_{\mathbb{R}^n}^{\varepsilon}) \to 0 \text{ as } \varepsilon \to 0.$$

We can estimate the length of this path via

$$\operatorname{Len}(\varphi_{\mathbb{R}^{n}}^{\varepsilon})^{2} = \left(\int_{0}^{T} \|u_{\mathbb{R}^{n}}^{\varepsilon}(t,.)\|_{H^{s}(\mathbb{R}^{n})} dt\right)^{2} \leq T \int_{0}^{T} \|u_{\mathbb{R}^{n}}^{\varepsilon}(t,.)\|_{H^{s}(\mathbb{R}^{n})}^{2} dt$$

$$= T \int_{0}^{T} \|u_{\mathbb{R}}^{\varepsilon}(t,|\cdot|)\|_{H^{s}(\mathbb{R}^{n})}^{2} dt = T \int_{0}^{T} \|1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]} * G_{\varepsilon}(|x|)\|_{H^{s}(\mathbb{R}^{n})}^{2} dt$$

$$\leq C(G_{1},T) \int_{0}^{T} \|1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)\|_{H^{s}(\mathbb{R}^{n})}^{2} dt,$$

where the last estimate follows from

$$\begin{split} &\|\mathbf{1}_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]} * G_{\varepsilon}(|x|)\|_{H^{s}(\mathbb{R}^{n})}^{2} \\ &= \int_{\mathbb{R}^{n}} (1 + |\xi|^{2s}) \left[\mathcal{F}\left(\mathbf{1}_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)\right)(\xi) \right]^{2} \left[\mathcal{F}\left(G_{\varepsilon}(|\cdot|)\right)(\xi) \right]^{2} \, \, \mathrm{d}\xi \\ &= \int_{\mathbb{R}^{n}} (1 + |\xi|^{2s}) \left[\mathcal{F}\left(\mathbf{1}_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)\right)(\xi) \right]^{2} \left[\mathcal{F}\left(G_{1}(|\cdot|)\right)(\varepsilon\xi) \right]^{2} \, \, \mathrm{d}\xi \\ &\leq \|\mathcal{F}G_{1}(|\cdot|)\|_{L^{\infty}}^{2} \cdot \|\mathbf{1}_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)\|_{H^{s}(\mathbb{R}^{n})}^{2}. \end{split}$$

Hence it is sufficient to show that

$$\|1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)\|_{H^{s}(\mathbb{R}^{n})} \to 0$$
 as $\varepsilon \to 0$ uniformly in t .

To compute the H^s -norm of $1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)$ we first Fourier transform it. The Fourier transform of a radially symmetric function $v(|\cdot|) \in L^1(\mathbb{R}^n)$ is again radially symmetric and given by the following formula, see [11, Theorem 3.3],

$$(\mathcal{F}v(|\cdot|))(\xi) = 2\pi |\xi|^{1-n/2} \int_{0}^{\infty} J_{n/2-1}(2\pi |\xi|s) v(s) s^{n/2} ds,$$

with $J_{n/2-1}$ denoting the Bessel function of order $\frac{n}{2}-1$. To simplify notation we will omit the dependence of the vector field $1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)$ on t and ε . Changing coordinates, this becomes

$$(\mathcal{F}1_{[f,g]}(|\cdot|))(\xi) = (2\pi)^{-n/2} |\xi|^{-n} \int_{2\pi f|\xi|}^{2\pi g|\xi|} J_{n/2-1}(s) s^{n/2} \, \mathrm{d}s.$$

This integral can be evaluated explicitly using the following integral identity for Bessel functions from [9, (10.22.1)]

$$\int z^{\nu+1} J_{\nu}(z) \, \mathrm{d}z = z^{\nu+1} J_{\nu+1}(z), \quad \nu \neq -\frac{1}{2}.$$

This gives us

$$(\mathcal{F}1_{[f,g]}(|\cdot|))(\xi) = |\xi|^{-n/2} \left(J_{n/2}(2\pi g|\xi|)g^{n/2} - J_{n/2}(2\pi f|\xi|)f^{n/2} \right).$$

The H^s -norm of $1_{[f,g]}(|\cdot|)$ is given by

$$\|1_{[f,g]}(|\cdot|)\|_{H^{s}(\mathbb{R}^{n})}^{2} = \int_{\mathbb{R}^{n}} (1+|\xi|^{2s}) \mathcal{F}1_{[f,g]}(|\cdot|)(\xi)^{2} d\xi.$$

We will only consider the term involving $|\xi|^{2s}$, since the L^2 -term can be estimated in the same way by setting s = 0. Transforming to polar coordinates we obtain

$$\begin{split} \int\limits_{\mathbb{R}^n} |\xi|^{2s} \left(\mathcal{F} \mathbf{1}_{[f,g]}(|\cdot|)(\xi) \right)^2 \ \mathrm{d}\xi &= \int\limits_{\mathbb{R}^n} |\xi|^{2s-n} \left(J_{n/2}(2\pi g |\xi|) g^{n/2} - J_{n/2}(2\pi f |\xi|) f^{n/2} \right)^2 \ \mathrm{d}\xi \\ &= \mathrm{Vol}(S^{n-1}) \int\limits_0^\infty r^{2s-1} \left(J_{n/2}(2\pi g r) g^{n/2} - J_{n/2}(2\pi f r) f^{n/2} \right)^2 \ \mathrm{d}r. \end{split}$$

The above integral is non-zero only for those t, where $f^{\varepsilon}(t) \neq g^{\varepsilon}(t)$. From Lemma 3.2 and our assumptions on $\varphi_{\mathbb{R}}$ we know that

$$\{t : f^{\varepsilon}(t) < g^{\varepsilon}(t)\} \subseteq \operatorname{supp}(\varphi_{\mathbb{R}}) \subseteq [1, \infty).$$

Thus both $f^{\varepsilon}(t)$ and $g^{\varepsilon}(t)$ are different and away from 0 and we can evaluate the above integral using the identity [9, (10.22.57)],

$$\int\limits_{0}^{\infty} \frac{J_{\mu}(at)J_{\nu}(at)}{t^{\lambda}} \ \mathrm{d}t = \frac{\left(\frac{1}{2}a\right)^{\lambda-1} \Gamma\left(\frac{\mu}{2} + \frac{\nu}{2} - \frac{\lambda}{2} + \frac{1}{2}\right) \Gamma\left(\lambda\right)}{2\Gamma\left(\frac{\lambda}{2} + \frac{\nu}{2} - \frac{\mu}{2} + \frac{1}{2}\right) \Gamma\left(\frac{\lambda}{2} + \frac{\mu}{2} - \frac{\nu}{2} + \frac{1}{2}\right) \Gamma\left(\frac{\lambda}{2} + \frac{\mu}{2} + \frac{\nu}{2} + \frac{1}{2}\right)},$$

which holds for $Re(\mu + \nu + 1) > Re\lambda > 0$ and the identity [9, (10.22.56)],

$$\int\limits_{0}^{\infty} \frac{J_{\mu}(at)J_{\nu}(bt)}{t^{\lambda}} \ \mathrm{d}t = \frac{a^{\mu}\Gamma\left(\frac{\nu}{2} + \frac{\mu}{2} - \frac{\lambda}{2} + \frac{1}{2}\right)}{2^{\lambda}b^{\mu - \lambda + 1}\Gamma\left(\frac{\nu}{2} - \frac{\mu}{2} + \frac{\lambda}{2} + \frac{1}{2}\right)} \mathbf{F}\left(\frac{\nu}{2} + \frac{\mu}{2} - \frac{\lambda}{2} + \frac{1}{2}, \frac{\mu}{2} - \frac{\nu}{2} - \frac{\lambda}{2} + \frac{1}{2}; \mu + 1; \frac{a^{2}}{b^{2}}\right),$$

which holds for 0 < a < b and $\text{Re}(\mu + \nu + 1) > \text{Re}\lambda > -1$. Here $\mathbf{F}(a, b; c; d)$ is the regularized hypergeometric function. Using these identities with $\lambda = 1 - 2s$, $\mu = \nu = \frac{n}{2}$, $a = 2\pi f$ and $b = 2\pi g$ we obtain

$$\int_{0}^{\infty} r^{2s-1} J_{n/2} (2\pi f r)^{2} dr = \frac{1}{2} (\pi f)^{-2s} \frac{\Gamma(\frac{n}{2} + s) \Gamma(1 - 2s)}{\Gamma(1 - s)^{2} \Gamma(\frac{n}{2} + 1 - s)}$$

and

$$\int_{0}^{\infty} r^{2s-1} J_{n/2}(2\pi f r) J_{n/2}(2\pi g r) dr = \frac{1}{2} (\pi g)^{-2s} \left(\frac{f}{g} \right)^{n/2} \frac{\Gamma\left(\frac{n}{2} + s\right)}{\Gamma(1-s)} \mathbf{F}\left(\frac{n}{2} + s, s; \frac{n}{2} + 1; \frac{f^{2}}{g^{2}}\right).$$



Putting it together results in

$$\begin{split} \int\limits_{\mathbb{R}^{n}} |\xi|^{2s} (\mathcal{F}1_{[f,g]}(|\cdot|))(\xi)^{2} \, \mathrm{d}\xi &= \mathrm{Vol}(S^{n-1}) \Bigg(\frac{f^{-2s} + g^{-2s}}{2\pi^{2s}} \frac{\Gamma\left(\frac{n}{2} + s\right) \Gamma(1 - 2s)}{\Gamma(1 - s)^{2} \Gamma\left(\frac{n}{2} + 1 - s\right)} \\ &- \frac{g^{-2s}}{\pi^{2s}} \frac{f^{n/2}}{g^{n/2}} \frac{\Gamma\left(\frac{n}{2} + s\right)}{\Gamma(1 - s)} \mathbf{F}\left(\frac{n}{2} + s, s; \frac{n}{2} + 1; \frac{f^{2}}{g^{2}}\right) \Bigg). \end{split}$$

In the limit $\varepsilon \to 0$ we know from Lemma 3.2 that $f^{\varepsilon}(t) \to h(t)$ and $g^{\varepsilon}(t) \to h(t)$ uniformly in t on [0, T] and hence $\frac{f^{\varepsilon}(t)}{g^{\varepsilon}(t)} \to 1$. For the regularized hypergeometric function $\mathbf{F}(a, b; c; d)$ at d = 1 we have the identity [9, (15.4.20)]

$$\mathbf{F}(a,b;c;1) = \frac{\Gamma(c-a-b)}{\Gamma(c-a)\Gamma(c-b)},$$

for Re(c-a-b) > 0. Applying the identity with $a = \frac{n}{2} + s$, b = s and $c = \frac{n}{2} + 1$ we get

$$\mathbf{F}\left(\frac{n}{2}+s,s;\frac{n}{2}+1;1\right) = \frac{\Gamma(1-2s)}{\Gamma(1-s)\Gamma\left(\frac{n}{2}+1-s\right)}.$$

Using the continuity of the hypergeometric function it follows that

$$\int_{\mathbb{D}^n} |\xi|^{2s} \left(\mathcal{F} \mathbf{1}_{[f,g]}(|\cdot|))(\xi) \right)^2 d\xi \to 0,$$

as $\varepsilon \to 0$ and the convergence is uniform in t. This concludes the proof that

$$\|1_{[f^{\varepsilon}(t),g^{\varepsilon}(t)]}(|\cdot|)\|_{H^{s}(\mathbb{D}^{n})} \to 0$$
 as $\varepsilon \to 0$ uniformly in t ,

and hence we have established claim B for $Diff_c(\mathbb{R}^n)$.

To prove this result for an arbitrary manifold M of bounded geometry we choose a partition of unity (τ_j) such that $\tau_0 \equiv 1$ on some open subset $U \subset M$, where normal coordinates centred at $x_0 \in M$ are defined. If $\varphi_{\mathbb{R}}$ is chosen with sufficiently small support, then the vector field $u_{\mathbb{R}^n}^\varepsilon$ has support in $\exp_{x_0}(U)$ and we can define the vector field $u_M^\varepsilon := (\exp_{x_0}^{-1})^* u_{\mathbb{R}^n}^\varepsilon$ on M. This vector field generates a path $\varphi_M^\varepsilon(t,\cdot) \in \mathrm{Diff}_0(M)$ with an endpoint $\varphi_M^\varepsilon(T,\cdot) = \varphi_M(\cdot)$ that does not depend on ε with arbitrarily small H^s -length since

$$\operatorname{Len}(\varphi_M^{\varepsilon}) \leq C_1(\tau) \int_0^T \|u_M^{\varepsilon}\|_{H^s(M,\tau)} dt = C_1(\tau) \int_0^T \|\exp_{x_0}^*(\tau_0.u_M^{\varepsilon})\|_{H^s(\mathbb{R}^n)} dt$$
$$= C_1(\tau) \int_0^T \|u_{\mathbb{R}^n}^{\varepsilon}\|_{H^s(\mathbb{R}^n)} dt.$$

Thus we can reduce the case of arbitrary manifolds to \mathbb{R}^n and this concludes the proof.

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