

Homotopy analysis Sumudu transform method for time—fractional third order dispersive partial differential equation

Rishi Kumar Pandey¹ · Hradyesh Kumar Mishra¹

Received: 3 May 2016 / Accepted: 4 October 2016 /
Published online: 13 October 2016
© Springer Science+Business Media New York 2016

Abstract In this article, we apply the newly introduced numerical method which is a combination of Sumudu transforms and Homotopy analysis method for the solution of time fractional third order dispersive type PDE equations. It is also discussed generalized algorithm, absolute convergence and analytic result of the finite number of independent variables including time variable.

Keywords Dispersive partial differential equation · Homotopy analysis method · Homotopy analysis Sumudu transform method · Linear and nonlinear partial differential equation

Mathematics Subject Classification (2010) 26A33 · 34A08 · 60G22 · 65Gxx

1 Introduction

The Fractional calculus is as old as classical calculus because importance of this theory was marked as soon as the ideas of the classical calculus were born from the discussion of half derivative in epistle of Leibniz and L'Hôpital in the year 1695. Further, many mathematicians contributed on this theory and strengthened the notion

Communicated by: Helge Holden

✉ Hradyesh Kumar Mishra
hk.mishra@juet.ac.in

Rishi Kumar Pandey
rishipandey.9@rediffmail.com

¹ Department of Mathematics, Jaypee University of Engineering and Technology,
Guna, 473226 MP, India

of generalized order differential and integrals viz. Liouville, Euler, Fourier, Abel, Riemann, Weyl. Liouville took initial steps for the fractional order integration and published the series of papers (1832–1837). The Riemann–Liouville operator was the most popular among mathematicians who solved fractional order integration problems [1]. Evidently, up to 300 years mentioned theory was the asset of only pure mathematicians due to unavailability of geometrical and physical interpretation of fractional order differential and integral operators. Caputo [2] described useful formula for generalized order derivatives. Oldham and Spanier [1] discussed the initial framework of application in diffusion problem, classical calculus with proper explanation. Ross [3] presented the chronological development of this theory after completing his PhD in fractional calculus and also published a monograph [4]. In the consequence, Podlubny [5], Kilbas et al. [6], Anatoly et al. [7], Diethelm [8], Caponetto et al. [9], Samko et al. [10] introduced the generalized differential and integral operators in more precise form with existence and uniqueness of results in application. Now a days, enormous model and physical phenomena like anomalous diffusion equation theory [11], mechanics of non-Hamiltonian systems [12], theory of long range interaction [13], astrophysics [14], optics [15], mechanics of fractal media [16], plasma physics [17, 18], physical kinetics [19], quantum mechanics [20], chaotic dynamics [21], which cannot meaningfully describe without means of fractional operators. Because in dynamical systems, integer order derivatives only evaluate a fixed number of derivatives wherein fractional derivatives can evaluate the value for any arbitrary order of derivative correspond to real numbers. Payable to its incredible scope and relevance in many branches of science and engineering, an extensive attention has been shown to find the solution of differential and integral equations involving the fractional derivatives. Except the modelling approach of mentioned differential equations and its solution procedure, including efficiency of convergence, divergence or junctions solutions of the model are uniformly important in numerical evaluation analysis. In order to achieve more convenient and highly adorable results, numerous numerical methods have been proposed to solve the differential equations of fractional order. Some of semi-analytic/analytic methods or numerical methods are differential transform method [22–24], Variational iteration method (VIM) [25, 26], fractional variational iteration method (FVIM) [27], Wavelet Operational matrix method [28], generalized differential transform method [29], Fractional sub equation method [30], Homotopy perturbation method [31–34], Homotopy analysis method [35–38], Homotopy analysis transform method [39–43], Fractional differential transform and Modified Fractional differential transform method [44, 45], Homotopy analysis fractional Sumudu transform method (HAFSTM) [46].

In order to convert the complex linear and nonlinear form of fractional order partial differential equations into simpler algebraic form many type of fractional integral and differential transforms have been applied to gain the exact and approximate solutions of FPDE's [47, 48]. Kumar and his co-workers successfully applied homotopy analysis transform method which is cumulation of Laplace transform and homotopy analysis method for the solution of fractional Fornberg–Whitham equation arising in wave breaking [39], volterra integral equation [40], fractional wave equations [41],

coupled Boussinesq–Burger’s equations arise in propagation of shallow water waves [42], unidirectional propagation of long waves in dispersive media [43]. Watugala [49] introduced the Sumudu transform and some properties discussed by Weerakoon [50, 51]. Further, Belgacem [52–59] provide precise definition of Sumudu transform and also discussed better implementations for the solution of FDE’s, FPDE’s using many results, properties and relations, which enhances the literature of this transform. It can easily convert many fractional order linear and nonlinear partial differential equations in time domain without loss of generality for different type of included fractional operator viz. Caputo, Riemann–Liouville, Ritzs space, etc. Multi-stage HAM is introduced in [60] for Solving non-linear Riccati Differential Equations. Since the homotopy analysis method applied to solve in wide variety of linear and nonlinear partial differential equations such as some fractional order smoking model [61], Lorenz system [62], a class of partial differential equations [63], space– and time-fractional kdv equation [64], Foam Drainage Equation with Space– and Time–Fractional Derivatives [65] and so on. The disadvantage of perturbation method is to solve each iteration and convergence region is very less. ADM, VIM, provide weak convergent and not necessarily accurate always to exact solutions. DTM and FDTM, MFDTM require additional information and basic formula to evaluate the results. While HASTM is easily evaluate the nonlinear term with high accuracy due to independence of physical parameters and absolute convergence of series towards the exact solutions.

In this article we have applied Homotopy analysis Sumudu transform method to solve third- order fractional dispersive partial differential equations [66–71] included fractional derivative in caputo sense. The HASTM obtains semi analytic solutions in the form of series solutions. It is different from other transforms and semi analytic method, which does not require additional information except some initial and boundary conditions. It easily changes the original problem to lucid manner and then one can evaluate the result with high convergence and accuracy.

The article taxonomy is arranged as follows: In Section 2 rudimental definitions of fractional calculus and properties are discussed. The rudimental concept of HASTM is explained in Section 3. To demonstrate the method and advantages, three examples of fractional order dispersive partial differential equations are solved with discussion of convergence in Section 4. At the end concluding remark is presented in Section 5.

2 Basic definitions

Definition 2.1 A real function $f(t)$, $t > 0$, is said to be in the space C_μ , $\mu \in \mathbb{R}$ if there exists a real number $p (> \mu)$, such that $f(t) = t^p f_1(t)$, where $f_1(t) \in C[0, \infty)$, and it is said to be in the space C_μ^m iff $f^{(m)} \in C_\mu$, $m \in \mathbb{N}$.

Definition 2.2 The Riemann–Liouville Fractional integral operator of order $\alpha \geq 0$, of a function $f(t) \in C_\mu$, and $\mu \geq -1$ is defined as [72, 73]

$$J^\alpha f(t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t - \tau)^{\alpha-1} f(\tau) d\tau, \quad \alpha > 0, \quad x > 0 \text{ and } J^0 f(t) = f(t).$$

For the Riemann–Liouville fractional integral, we have

$$J^\alpha t^y = \frac{\Gamma(y + 1)}{\Gamma(y + \alpha + 1)} t^{\alpha+y}.$$

Definition 2.3 The fractional derivative of $f(t)$ in the Caputo sense is defined as [1]

$$D_t^\alpha f(t) = \begin{cases} J^{m-\alpha} D^n f(t), \\ \frac{1}{\Gamma(n-\alpha)} \int_0^t (t-\tau)^{m-\alpha-1} f^{(m)}(\tau) d\tau, \end{cases}$$

where $m - 1 < \alpha \leq m$, $m \in N$, $t > 0$.

Definition 2.4 The Sumudu transform is defined over the set of functions [52–54]

$$A = \left\{ f(t) \mid \exists M, \tau_1, \tau_2 > 0, |f(t)| < M e^{\frac{|t|}{\tau_j}}, \text{ if } t \in (-1)^j \times [0, \infty) \right\},$$

by the following formula

$$\bar{f}(u) = \mathbb{S}[f(t)] = \int_0^\infty f(ut) e^{-t} dt, \quad u \in (-\tau_1, \tau_2).$$

Definition 2.5 The Sumudu transform of $f(t) = t^\alpha$ is defined as [53]

$$\mathbb{S}[t^\alpha] = \int_0^\infty e^{-t} t^\alpha dt = \Gamma(\alpha + 1) u^\alpha, \quad R(\alpha) > 0.$$

Definition 2.6 The Sumudu transform $\mathbb{S}[f(t)]$ of the Riemann -Liouville fractional integral is defined as [53]

$$\mathbb{S}[J^\alpha f(t)] = u^{-\alpha} F(u).$$

Definition 2.7 The Sumudu transform $\mathbb{S}[f(t)]$ of the Caputo fractional derivative is defined as [53]

$$\mathbb{S}[D_t^\alpha f(t)] = u^{-\alpha} \mathbb{S}[f(t)] - \sum_{k=0}^{m-1} u^{-\alpha+k} f^{(k)}(0^+), \quad \text{where } m - 1 < \alpha \leq m.$$

3 Solution by Homotopy analysis Sumudu transform method

To illustrate the rudimental conception of the HASTM for the fractional partial differential equation, we consider the linear third order dispersive partial differential equations in following manner:

$$D_t^{n\alpha} \xi(x_1, x_2, \dots, x_n, t) + \sum_{i=1}^n l_i \frac{\partial^3 \xi(x_1, x_2, \dots, x_n, t)}{\partial x_i^3} = G(x_1, x_2, \dots, x_n, t); \quad (1)$$

$\forall l_i, t > 0, \forall x_i \in R, n - 1 < \alpha \leq n$, and the $G(x_1, x_2, \dots, x_n, t)$ is the source function.

For simplicity, we ignore all initial and boundary conditions, which can be treated in a homogeneous way. Now the methodology consists of applying the Sumudu transform first on both sides of the Eq. 1, we get

$$\mathbb{S} [D_t^{n\alpha} \xi (x_1, x_2, \dots, x_n, t)] + \mathbb{S} \left[\sum_{i=1}^n l_i \frac{\partial^3 \xi (x_1, x_2, \dots, x_n, t)}{\partial x_i^3} \right] = \mathbb{S} [G (x_1, x_2, \dots, x_n, t)]; \tag{2}$$

Using the definition (2.7) differentiation property of the Sumudu transform

$$u^{-\alpha} \mathbb{S} [\xi (x_1, x_2, \dots, x_n, t)] - \sum_{k=0}^{n-1} \frac{\xi^{(k)} (0)}{u^{\alpha-k}} + \mathbb{S} \left[\sum_{i=1}^n l_i \frac{\partial^3 \xi (x_1, x_2, \dots, x_n, t)}{\partial x_i^3} \right] = \mathbb{S} [G (x_1, x_2, \dots, x_n, t)];$$

which gives

$$\mathbb{S} [\xi (x_1, x_2, \dots, x_n, t)] - \sum_{k=0}^{n-1} \frac{\xi^{(k)} (0)}{u^{-k}} + u^\alpha \mathbb{S} \left[\sum_{i=1}^n l_i \frac{\partial^3 \xi (x_1, x_2, \dots, x_n, t)}{\partial x_i^3} - G (x_1, x_2, \dots, x_n, t) \right] = 0; \tag{3}$$

we define nonlinear operator as

$$\mathbb{N} [\phi (x_1, x_2, \dots, x_n, t; p)] = \mathbb{S} [\phi (x_1, x_2, \dots, x_n, t; p)] - \sum_{k=0}^{n-1} \frac{\phi^{(k)} (0)}{u^{-k}} + u^\alpha \mathbb{S} \left[\sum_{i=1}^n l_i \frac{\partial^3 \phi (x_1, x_2, \dots, x_n, t; p)}{\partial x_i^3} - G (x_1, x_2, \dots, x_n, t; p) \right], \tag{4}$$

where $p \in [0, 1]$ be an embedding parameter and $\phi (x_1, x_2, \dots, x_n, t; p)$ is a real function of x_1, x_2, \dots, x_n, t and p .

We construct a homotopy as follow:

$$(1 - p) \mathbb{S} [\phi (x_1, x_2, \dots, x_n, t; p) - \xi_0 (x_1, x_2, \dots, x_n, t)] = p \hbar H (x_1, x_2, \dots, x_n, t) \mathbb{N} [\phi (x_1, x_2, \dots, x_n, t; p)]; \tag{5}$$

where \hbar is a nonzero auxiliary parameter and $H (x_1, x_2, \dots, x_n, t) \neq 0$. An auxiliary function $\xi_0 (x_1, x_2, \dots, x_n, t)$ is an initial guess of $\xi (x_1, x_2, \dots, x_n, t)$ and $\phi (x_1, x_2, \dots, x_n, t; p)$ is an unknown function. It is important that one has great freedom to choose auxiliary parameter in HASTM. Obviously, when $p = 0$ and $p = 1$ it holds

$$\phi (x_1, x_2, \dots, x_n, t; 0) = \xi_0 (x_1, x_2, \dots, x_n, t), \quad \phi (x_1, x_2, \dots, x_n, t; 1) = \xi (x_1, x_2, \dots, x_n, t). \tag{6}$$

Thus, as p increases from 0 to 1, the solution varies from initial guess $\xi_0 (x_1, x_2, \dots, x_n, t)$ to the solution $\xi (x_1, x_2, \dots, x_n, t)$. Now, expanding $\phi (x_1, x_2, \dots, x_n, t; p)$ on Taylor's series with respect to q , we get

$$\phi (x_1, x_2, \dots, x_n, t; p) = \xi_0 (x_1, x_2, \dots, x_n, t) + \sum_{m=1}^{\infty} p^m \xi_m (x_1, x_2, \dots, x_n, t), \tag{7}$$

where

$$\xi_m(x_1, x_2, \dots, x_n, t) = \frac{1}{\Gamma(m+1)} \left. \frac{\partial^m \phi(x_1, x_2, \dots, x_n, t; p)}{\partial p^m} \right|_{p=0}. \tag{8}$$

The convergence of the series solution (7) is controlled by \hbar . If the auxiliary linear operator, the initial guess, the auxiliary parameter \hbar and the auxiliary function are properly chosen, the series (7) converges at $p = 1$. Hence we obtain

$$\xi(x_1, x_2, \dots, x_n, t) = \xi_0(x_1, x_2, \dots, x_n, t) + \sum_{m=1}^{\infty} \xi_m(x_1, x_2, \dots, x_n, t), \tag{9}$$

which must be one of the solutions of original nonlinear equations. The above expression provides us with a relationship between the initial guess $\xi_0(x_1, x_2, \dots, x_n, t)$ and the exact solution $\xi(x_1, x_2, \dots, x_n, t)$ by means of the terms $\xi_m(x_1, x_2, \dots, x_n, t)$ ($m = 1, 2, 3, \dots$), which are still to be determined.

Define the vectors

$$\vec{\xi} = \{\xi_0(x_1, x_2, \dots, x_n, t), \xi_1(x_1, x_2, \dots, x_n, t), \xi_2(x_1, x_2, \dots, x_n, t), \dots, \xi_m(x_1, x_2, \dots, x_n, t)\}. \tag{10}$$

Differentiating the zero order deformation Eq. 5 m times with respect to embedding parameter p and then setting $p = 0$, and finally dividing them by $\Gamma(m+1)$ we obtain the m^{th} order deformation equation as follows:

$$\begin{aligned} & \mathbb{S} [\xi_m(x_1, x_2, \dots, x_n, t) - \chi_m \xi_{m-1}(x_1, x_2, \dots, x_n, t)] \\ & = \hbar H(x_1, x_2, \dots, x_n, t) R_m(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t). \end{aligned} \tag{11}$$

Operating the inverse Sumudu transform of both sides, we get

$$\begin{aligned} \xi_m(x_1, x_2, \dots, x_n, t) & = \chi_m \xi_{m-1}(x_1, x_2, \dots, x_n, t) \\ & + \mathbb{S}^{-1} \left[\hbar H(x_1, x_2, \dots, x_n, t) R_m(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t) \right], \end{aligned} \tag{12}$$

where

$$R_m(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t) = \frac{1}{\Gamma(m)} \left. \frac{\partial^{m-1} \varphi(x_1, x_2, \dots, x_n, t; p)}{\partial p^{m-1}} \right|_{p=0}, \tag{13}$$

and

$$\chi_m = \begin{cases} 0, & m \leq 1, \\ 1 & m > 1. \end{cases} \tag{14}$$

In our case

$$\begin{aligned} R_m(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t) & = D_t^{n\alpha} \xi_{m-1}(x_1, x_2, \dots, x_n, t) + \sum_{i=1}^n l_i \frac{\partial^3 \xi_{m-1}(x_1, x_2, \dots, x_n, t)}{\partial x_i^3} \\ & - (1 - \chi_m) G(x_1, x_2, \dots, x_n, t). \end{aligned} \tag{15}$$

In this way, it is easy to obtain $\xi_m(x_1, x_2, \dots, x_n, t)$ for $m \geq 1$, at M^{th} order, we have

$$\xi(x_1, x_2, \dots, x_n, t) = \sum_{m=0}^M \xi_m(x_1, x_2, \dots, x_n, t), \tag{16}$$

where $M \rightarrow \infty$, we obtain an accurate approximation of the original equation (1).

Theorem 3.1 (Convergence Theorem) *If the series (16) is converging for $M \rightarrow \infty$, where $\xi_m(x_1, x_2, \dots, x_n, t)$ is obtained by Eq. 12 and using the conditions (14) and (15). Then, it must be the exact solution of original discussed partial differential equation (1).*

Proof Let the series (16) be the convergent series then

$$\sum_{m=0}^{\infty} \xi_m(x_1, x_2, \dots, x_n, t) = \xi_0(x_1, x_2, \dots, x_n, t) + \sum_{m=1}^{\infty} \xi_m(x_1, x_2, \dots, x_n, t) = K(x_1, x_2, \dots, x_n, t). \tag{17}$$

Now we have $\lim_{M \rightarrow \infty} \xi_m(x_1, x_2, \dots, x_n, t) = 0$. Using definition of Eq. 11 we obtained

$$\begin{aligned} & \lim_{M \rightarrow \infty} \left[\hbar H(x_1, x_2, \dots, x_n, t) \sum_{m=1}^M R_m \left(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t \right) \right] \\ &= \lim_{M \rightarrow \infty} \left(\sum_{m=1}^M \mathbb{S} \left[\xi_m(x_1, x_2, \dots, x_n, t) - \chi_m \xi_{m-1}(x_1, x_2, \dots, x_n, t) \right] \right) \\ &= \lim_{M \rightarrow \infty} \left(\sum_{m=1}^M \mathbb{S} \left[\xi_m(x_1, x_2, \dots, x_n, t) - \chi_m \xi_{m-1}(x_1, x_2, \dots, x_n, t) \right] \right) \\ &= \mathbb{S} \left(\lim_{M \rightarrow \infty} \sum_{m=1}^M \left[\xi_m(x_1, x_2, \dots, x_n, t) - \chi_m \xi_{m-1}(x_1, x_2, \dots, x_n, t) \right] \right) \\ &= \mathbb{S} \left(\lim_{M \rightarrow \infty} \xi_M(x_1, x_2, \dots, x_n, t) \right) \\ &= 0. \end{aligned}$$

Since $\hbar \neq 0, H(x_1, x_2, \dots, x_n, t) \neq 0$, therefore $\sum_{m=1}^{\infty} R_m \left(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t \right) = 0$.

From (15)

$$\begin{aligned} & \sum_{m=1}^{\infty} R_m \left(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t \right) = \sum_{m=1}^{\infty} \left(D_t^{n\alpha} \xi_{m-1}(x_1, x_2, \dots, x_n, t) \right. \\ & \left. + \sum_{i=1}^n l_i \frac{\partial^3 \xi_{m-1}(x_1, x_2, \dots, x_n, t)}{\partial x_i^3} - (1 - \chi_m) G(x_1, x_2, \dots, x_n, t) \right) \\ & \sum_{m=1}^{\infty} R_m \left(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t \right) = \sum_{m=1}^{\infty} D_t^{n\alpha} \xi_{m-1}(x_1, x_2, \dots, x_n, t) \\ & + \sum_{m=1}^{\infty} \sum_{i=1}^n l_i \frac{\partial^3 \xi_{m-1}(x_1, x_2, \dots, x_n, t)}{\partial x_i^3} - \sum_{m=1}^{\infty} (1 - \chi_m) G(x_1, x_2, \dots, x_n, t) \end{aligned}$$

$$\begin{aligned} \sum_{m=1}^{\infty} R_m \left(\vec{\xi}_{m-1}, x_1, x_2, \dots, x_n, t \right) &= D_t^{n\alpha} \sum_{m=0}^{\infty} \xi_m(x_1, x_2, \dots, x_n, t) \\ &+ \sum_{i=1}^n l_i \frac{\partial^3 \sum_{m=0}^{\infty} \xi_m(x_1, x_2, \dots, x_n, t)}{\partial x_i^3} - G(x_1, x_2, \dots, x_n, t) \\ D_t^{n\alpha} K(x_1, x_2, \dots, x_n, t) &+ \sum_{i=1}^n l_i \frac{\partial^3 K(x_1, x_2, \dots, x_n, t)}{\partial x_i^3} - G(x_1, x_2, \dots, x_n, t) = 0. \end{aligned} \tag{18}$$

Above equation (18) shows that, $K(x_1, x_2, \dots, x_n, t)$ satisfies the original problem (1). □

4 Numerical illustrations

In this section we consider the time fractional dispersive partial differential equations to authenticate the method discussed in the previous section.

Example 4.1 We consider the linear time fractional KDV [71]

$$\xi_t^\alpha(x, t) + 2 \frac{\partial \xi(x, t)}{\partial x} + \frac{\partial^3 \xi(x, t)}{\partial x^3} = 0, \quad t > 0, \quad 0 < \alpha \leq 1, \tag{19}$$

subject to the initial condition

$$\xi(x, 0) = \text{Sin } x. \tag{20}$$

The exact solution at $\alpha = 1$ is given by

$$\xi(x, t) = \text{Sin}(x - t). \tag{21}$$

Applying the Sumudu transform of both sides in Eq. (19) and after using the definition (2.7) of Sumudu transform for fractional derivative, we get

$$\mathbb{S}[\xi(x, t)] + u^\alpha \mathbb{S} \left[2 \frac{\partial \xi(x, t)}{\partial x} + \frac{\partial^3 \xi(x, t)}{\partial x^3} \right] = 0, \quad t > 0. \tag{22}$$

The nonlinear operator is

$$N[\phi(x, t; p)] = \mathbb{S}[\phi(x, t; p)] + u^\alpha \mathbb{S} \left[2 \frac{\partial \phi(x, t; p)}{\partial x} + \frac{\partial^3 \phi(x, t; p)}{\partial x^3} \right] = 0, \quad t > 0, \quad 0 \leq p \leq 1, \tag{23}$$

and thus

$$R_m \left(\vec{\xi}_{m-1}, x, t \right) = \mathbb{S}[\xi_{m-1}(x, t)] + u^\alpha \mathbb{S} \left[2 \frac{\partial \xi_{m-1}(x, t)}{\partial x} + \frac{\partial^3 \xi_{m-1}(x, t)}{\partial x^3} \right] = 0, \quad t > 0. \tag{24}$$

The m^{th} - order deformation equation is given by

$$\mathbb{S}[\xi_m(x, t) - \chi_m \xi_{m-1}(x, t)] = \hbar H(x, t) R_m \left(\vec{\xi}_{m-1}, x, t \right).$$

Applying the inverse Sumudu transform, we have

$$\xi_m(x, t) = \chi_m \xi_{m-1}(x, t) + \mathbb{S}^{-1} \left[\hbar H(x, t) R_m \left(\vec{\xi}_{m-1}, x, t \right) \right]. \tag{25}$$

On solving above equation for $m = 1, 2, \dots$. For simplicity, we choose $H(x, t) = 1$,

$$\begin{aligned} \xi_1(x, t) &= \frac{t^\alpha \hbar \text{Cos } x}{\Gamma(1 + \alpha)}, \\ \xi_2(x, t) &= \frac{t^\alpha \hbar \text{Cos } x}{\Gamma(1 + \alpha)} + \frac{t^{2\alpha-1} \alpha \hbar^2 \text{Cos } x \Gamma(\alpha)}{\Gamma(1 + \alpha) \Gamma(2\alpha)} - \frac{t^{2\alpha} \hbar^2 \text{Sin } x}{\Gamma(1 + 2\alpha)}, \\ \xi_3(x, t) &= \frac{t^\alpha \hbar \text{Cos } x}{\Gamma(1 + \alpha)} + \frac{2t^{2\alpha-1} \alpha \hbar^2 \text{Cos } x \Gamma(\alpha)}{\Gamma(2\alpha) \Gamma(1 + \alpha)} - \frac{t^{3\alpha-2} \alpha \hbar^3 \text{Cos } x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} \\ &\quad + \frac{3t^{3\alpha-2} \alpha^2 \hbar^3 \text{Cos } x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} - \frac{3t^\alpha \hbar^3 \text{Cos } x}{\Gamma(1 + \alpha)} - \frac{t^{3\alpha-1} \alpha \hbar^3 \Gamma(\alpha) \text{Sin } x}{\Gamma(1 + \alpha) \Gamma(3\alpha)} \\ &\quad - \frac{2t^{2\alpha} \hbar^2 \text{Sin } x}{\Gamma(1 + 2\alpha)} - \frac{2t^{3\alpha-1} \alpha \hbar^3 \Gamma(2\alpha) \text{Sin } x}{\Gamma(1 + 2\alpha) \Gamma(3\alpha)}, \\ &\quad \vdots \\ &\quad \vdots \\ &\quad \vdots \end{aligned}$$

and so on.

Here, we consider the results upto $m = 10$ and rest of the components can be evaluated by iteration formula (24).

Therefore the solution of Eq. 19 is given by

$$\xi(x, t) = \xi_0(x, t) + \sum_{m=1}^{\infty} \xi_m(x, t). \tag{26}$$

At $\hbar = -1$ we obtained the following approximation:

$$\begin{aligned} \xi(x, t) &= \frac{-3t^\alpha \text{Cos } x}{\Gamma(1 + \alpha)} + \frac{3t^{2\alpha-1} \alpha \text{Cos } x \Gamma(\alpha)}{\Gamma(2\alpha) \Gamma(1 + \alpha)} + \frac{t^{3\alpha-2} \alpha \text{Cos } x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} \\ &\quad - \frac{2t^{3\alpha-2} \alpha^2 \text{Cos } x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} + \frac{t^{3\alpha} \text{Cos } x}{\Gamma(1 + 3\alpha)} + \text{Sin } x + \frac{t^{3\alpha-1} \alpha \Gamma(\alpha) \text{Sin } x}{\Gamma(3\alpha) \Gamma(1 + \alpha)} \\ &\quad - \frac{3t^{2\alpha} \text{Sin } x}{\Gamma(1 + 2\alpha)} + \frac{2t^{3\alpha-1} \alpha \Gamma(2\alpha) \text{Sin } x}{\Gamma(3\alpha) \Gamma(1 + 2\alpha)} + \dots \end{aligned} \tag{27}$$

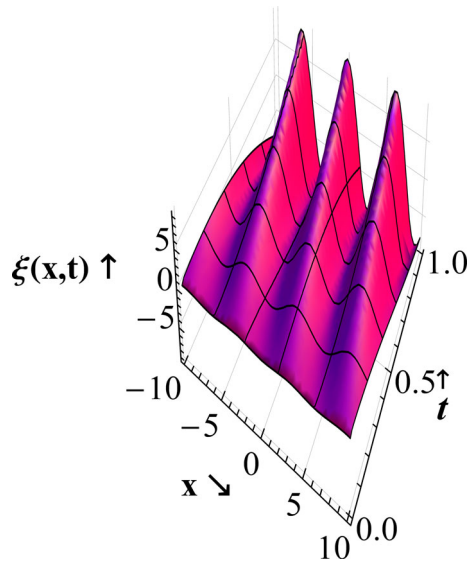
when $\alpha = 1$ Eq. 27 shows the similar results as [70] which is the exact solution of Eq. 19

$$\xi(x, t) = \text{Sin } x - \frac{1}{2} t^2 \text{Sin } x - t \text{Cos } x + \frac{1}{6} t^3 \text{Cos } x + \dots$$

After simplification we get Eq. 21.

Figures 1, 2, 3 and 4 show that the nature of fractional derivative and fluctuation changes from $\alpha = 0.9, 0.95, 1$ and exact solution at $\alpha = 1$.

Fig. 1 Plot of $\xi(x, t)$ w.r.t x and t at $\alpha = 0.9$



Example 4.2 Consider the linear time fractional KDV equation in one dimensional space

$$\xi_t^\alpha(x, t) + 3 \frac{\partial^3 \xi(x, t)}{\partial x^3} = 0, \quad t > 0, \quad 0 < x < 1, \quad 0 < \alpha \leq 1, \quad (28)$$

subject to the initial condition

$$\xi(x, 0) = \text{Cos } x, \quad 0 \leq x \leq 1. \quad (29)$$

Fig. 2 Plot of $\xi(x, t)$ w.r.t x and t at $\alpha = 0.95$

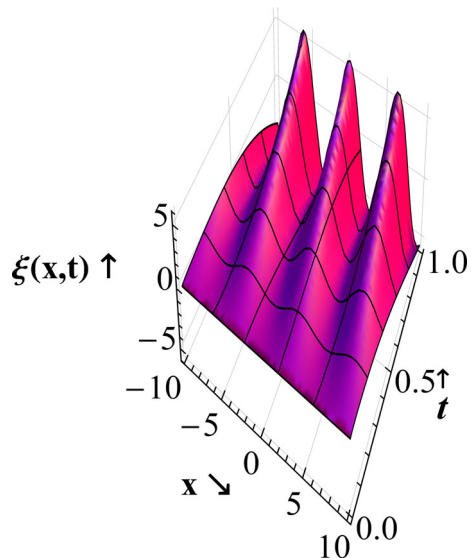
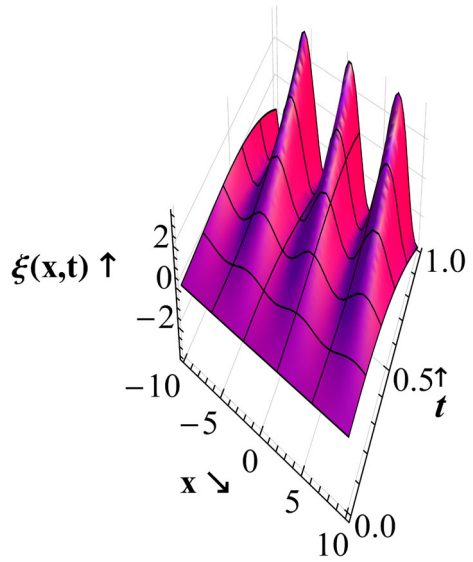


Fig. 3 Plot of $\xi(x, t)$ w.r.t x and t at $\alpha = 1$



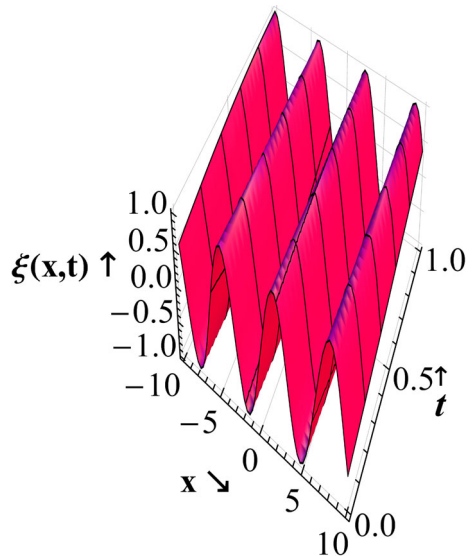
The exact solution at $\alpha = 1$ is given by

$$\xi(x, t) = \text{Cos}(x + 3t). \tag{30}$$

Applying the Sumudu transform of both sides in Eq. 28 and after using the definition (2.7) of Sumudu transform for fractional derivative, we get

$$\mathbb{S}[\xi(x, t)] + u^\alpha \mathbb{S}\left[3 \frac{\partial^3 \xi(x, t)}{\partial x^3}\right] = 0, \quad t > 0. \tag{31}$$

Fig. 4 Plot of Exact Solution of $\xi(x, t)$ w.r.t x and t



The nonlinear operator is

$$N[\phi(x, t; p)] = \mathbb{S}[\phi(x, t; p)] + 3u^\alpha \mathbb{S}\left[\frac{\partial^3 \phi(x, t; p)}{\partial x^3}\right] = 0, \quad t > 0, \quad 0 \leq p \leq 1, \tag{32}$$

and thus

$$R_m(\vec{\xi}_{m-1}, x, t) = \mathbb{S}[\xi_{m-1}(x, t)] + 3u^\alpha \mathbb{S}\left[\frac{\partial^3 \xi_{m-1}(x, t)}{\partial x^3}\right] = 0, \quad t > 0. \tag{33}$$

The m^{th} – order deformation equation is given by

$$\mathbb{S}[\xi_m(x, t) - \chi_m \xi_{m-1}(x, t)] = \hbar H(x, t) R_m(\vec{\xi}_{m-1}(x, t)).$$

Applying the inverse Sumudu transform, we have

$$\xi_m(x, t) = \chi_m \xi_{m-1}(x, t) + \mathbb{S}^{-1}\left[\hbar H(x, t) R_m(\vec{\xi}_{m-1}(x, t))\right]. \tag{34}$$

On solving above equation for $m = 1, 2, \dots$. For simplicity, we choose $H(x, t) = 1$

$$\begin{aligned} \xi_1(x, t) &= \frac{3t^\alpha \hbar \sin x}{\Gamma(1 + \alpha)}, \\ \xi_2(x, t) &= -\frac{9t^{2\alpha} \hbar^2 \cos x}{\Gamma(1 + 2\alpha)} + \frac{3t^\alpha \hbar \sin x}{\Gamma(1 + \alpha)} + \frac{3t^{2\alpha-1} \alpha \hbar^2 \sin x \Gamma(\alpha)}{\Gamma(1 + \alpha) \Gamma(2\alpha)}, \\ \xi_3(x, t) &= -\frac{9t^{3\alpha-1} \alpha \hbar^3 \cos x \Gamma(\alpha)}{\Gamma(1 + \alpha) \Gamma(3\alpha)} - \frac{18t^{2\alpha} \hbar^2 \cos x}{\Gamma(1 + 2\alpha)} \\ &\quad - \frac{18t^{3\alpha-1} \alpha \hbar^3 \cos x \Gamma(2\alpha)}{\Gamma(1 + 2\alpha) \Gamma(3\alpha)} + \frac{3t^\alpha \hbar \sin x}{\Gamma(1 + \alpha)} \\ &\quad + \frac{6t^{2\alpha-1} \alpha \hbar^2 \sin x \Gamma(\alpha)}{\Gamma(1 + \alpha) \Gamma(2\alpha)} - \frac{3t^{3\alpha-2} \alpha \hbar^3 \sin x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} \\ &\quad + \frac{6t^{3\alpha-2} \alpha^2 \hbar^3 \sin x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} - \frac{27t^{3\alpha} \hbar^3 \sin x}{\Gamma(3\alpha + 1)}, \\ &\quad \vdots \\ &\quad \vdots \\ &\quad \vdots \end{aligned}$$

and so on.

Here, we consider the results upto $m = 10$ and rest of the components can evaluate by iteration formula (34).

Therefore the solution of Eq. 28 is given by

$$\xi(x, t) = \xi_0(x, t) + \sum_{m=1}^{\infty} \xi_m(x, t). \tag{35}$$

At $\hbar = -1$ we obtained the following approximation:

$$\begin{aligned} \xi(x, t) = & \frac{-3t^\alpha \text{Cos } x}{\Gamma(1 + \alpha)} + \frac{3t^{2\alpha-1}\alpha \text{Cos } x \Gamma(\alpha)}{\Gamma(2\alpha) \Gamma(1 + \alpha)} + \frac{t^{3\alpha-2}\alpha \text{Cos } x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} \\ & - \frac{2t^{3\alpha-2}\alpha^2 \text{Cos } x \Gamma(\alpha) \Gamma(2\alpha - 1)}{\Gamma(2\alpha) \Gamma(1 + \alpha) \Gamma(3\alpha - 1)} + \frac{t^{3\alpha} \text{Cos } x}{\Gamma(1 + 3\alpha)} + \text{Sin } x + \frac{t^{3\alpha-1}\alpha \Gamma(\alpha) \text{Sin } x}{\Gamma(3\alpha) \Gamma(1 + \alpha)} \\ & - \frac{3t^{2\alpha} \text{Sin } x}{\Gamma(1 + 2\alpha)} + \frac{2t^{3\alpha-1}\alpha \Gamma(2\alpha) \text{Sin } x}{\Gamma(3\alpha) \Gamma(1 + 2\alpha)} + \dots \end{aligned} \tag{36}$$

when $\alpha = 1$ Eq. 36 shows the similar results as [70] which is the exact solution of Eq. 28.

$$\xi(x, t) = \text{Sin } x - \frac{1}{2}t^2 \text{Sin } x - t \text{Cos } x + \frac{1}{6}t^3 \text{Cos } x + \dots$$

After simplification we can get Eq. 30.

As discussed in previous example Figs. 5, 6, 7 and 8 show that the nature of fractional derivative and fluctuation changes for $\alpha = 0.9, 0.95, 1$ and exact solution at $\alpha = 1$.

Example 4.3 Now, we consider the linear time fractional KDV equation in two dimensional space [71]

$$\xi_t^\alpha(x, y, t) + 2 \frac{\partial^3 \xi(x, y, t)}{\partial x^3} + \frac{\partial^3 \xi(x, y, t)}{\partial y^3} = 0, \quad t > 0, \quad 0 < \alpha \leq 1, \tag{37}$$

subject to the initial condition

$$\xi(x, y, 0) = \text{Cos}(x + y). \tag{38}$$

Fig. 5 Plot of $\xi(x, t)$ w.r.t x and t at $\alpha = 0.9$

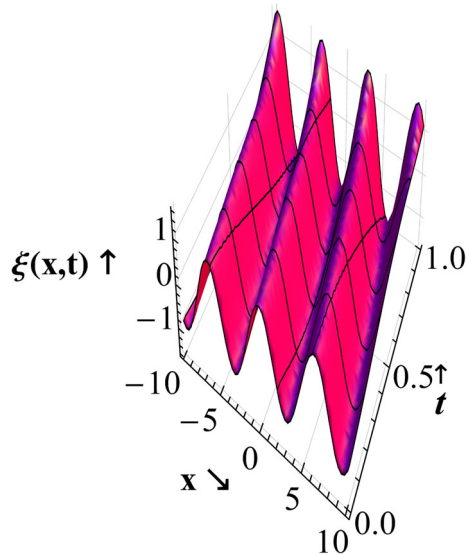
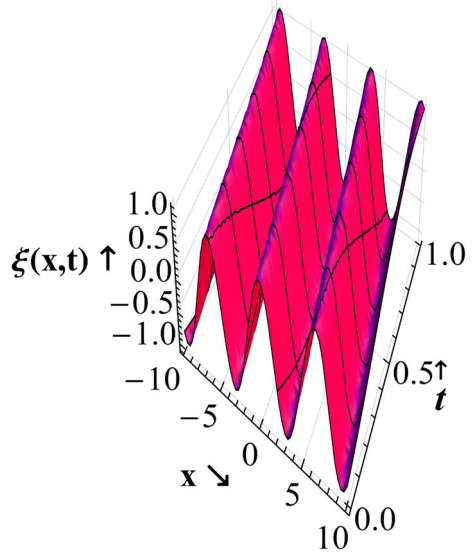


Fig. 6 Plot of $\xi(x, t)$ w.r.t x and t at $\alpha = 0.95$



The exact solution at $\alpha = 1$ is given by

$$\xi(x, y, t) = \text{Sin}(x + y + 2t). \tag{39}$$

Applying the Sumudu transform of both sides in Eq. 37 and after using the definition (2.7) of Sumudu transform for fractional derivative, we get

$$\mathbb{S}[\xi(x, y, t)] + u^\alpha \mathbb{S}\left[2 \frac{\partial^3 \xi(x, y, t)}{\partial x^3} + \frac{\partial^3 \xi(x, y, t)}{\partial y^3}\right] = 0, t > 0. \tag{40}$$

Fig. 7 Plot of $\xi(x, t)$ w.r.t x and t at $\alpha = 1$

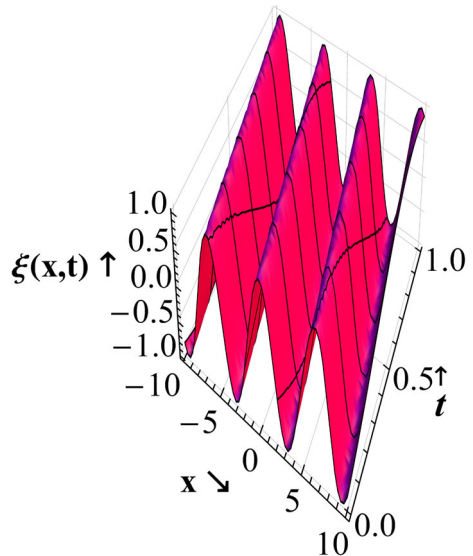
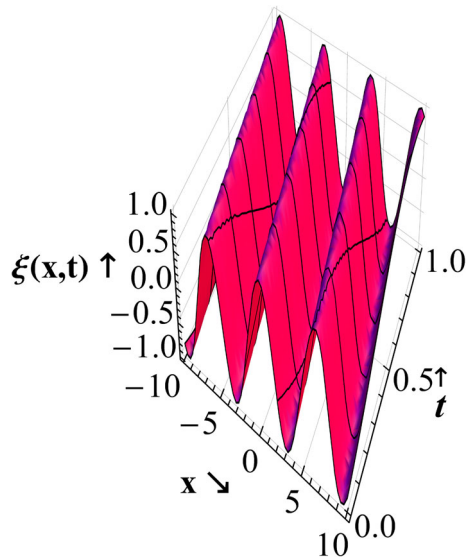


Fig. 8 Plot of Exact Solution of $\xi(x, t)$ w.r.t x and t



The nonlinear operator is

$$N[\phi(x, y, t; p)] = \mathbb{S}[\phi(x, y, t; p)] + u^\alpha \mathbb{S} \left[2 \frac{\partial^3 \phi(x, y, t; p)}{\partial x^3} + \frac{\partial^3 \phi(x, y, t; p)}{\partial y^3} \right] = 0, \quad t > 0, 0 \leq p \leq 1, \tag{41}$$

and thus

$$R_m \left(\vec{\xi}_{m-1}, x, y, t \right) = \mathbb{S}[\xi_{m-1}(x, y, t)] + u^\alpha \mathbb{S} \left[2 \frac{\partial^3 \xi_{m-1}(x, y, t)}{\partial x^3} + \frac{\partial^3 \xi_{m-1}(x, y, t)}{\partial y^3} \right] = 0, \quad t > 0. \tag{42}$$

The m^{th} – order deformation equation is given by

$$\mathbb{S}[\xi_m(x, y, t) - \chi_m \xi_{m-1}(x, y, t)] = \hbar H(x, y, t) R_m \left(\vec{\xi}_{m-1}(x, y, t) \right).$$

Applying the inverse Sumudu transform, we have

$$\xi_m(x, y, t) = \chi_m \xi_{m-1}(x, y, t) + \mathbb{S}^{-1} \left[\hbar H(x, y, t) R_m \left(\vec{\xi}_{m-1}(x, y, t) \right) \right]. \tag{43}$$

On solving above equation for $m = 1, 2, \dots$. For simplicity, we choose $H(x, t) = 1$,

$$\xi_1(x, y, t) = \frac{2t^\alpha \hbar \sin(x+y)}{\Gamma(1+\alpha)},$$

$$\xi_2(x, y, t) = -\frac{4t^{2\alpha} \hbar^2 \cos(x+y)}{\Gamma(1+2\alpha)} + \frac{2t^\alpha \hbar \sin(x+y)}{\Gamma(1+\alpha)} + \frac{2t^{2\alpha-1} \alpha \hbar^2 \Gamma(\alpha) \sin(x+y)}{\Gamma(1+2\alpha) \Gamma(1+\alpha)},$$

$$\begin{aligned} \xi_3(x, t) = & -\frac{4t^{3\alpha-1}\alpha\hbar^3\text{Cos}(x+y)\Gamma(\alpha)}{\Gamma(1+\alpha)} - \frac{8t^{2\alpha}\hbar^2\text{Cos}(x+y)}{\Gamma(1+2\alpha)} \\ & - \frac{8t^{3\alpha-1}\alpha\hbar^3\text{Cos}(x+y)\Gamma(2\alpha)}{\Gamma(3\alpha)\Gamma(1+2\alpha)} + \frac{2t^\alpha\hbar\text{Sin}(x+y)}{\Gamma(1+\alpha)} \\ & + \frac{4t^{2\alpha-1}\alpha\hbar^2\Gamma(\alpha)\text{Sin}(x+y)}{\Gamma(1+\alpha)\Gamma(2\alpha)} - \frac{2t^{3\alpha-2}\alpha\hbar^3\Gamma(\alpha)\Gamma(2\alpha-1)\text{Sin}(x+y)}{\Gamma(1+\alpha)\Gamma(3\alpha)\Gamma(3\alpha-1)} \\ & + \frac{4t^{3\alpha-2}\alpha\hbar^3\Gamma(\alpha)\Gamma(2\alpha-1)\text{Sin}(x+y)}{\Gamma(1+\alpha)\Gamma(2\alpha)\Gamma(3\alpha-1)} - \frac{8t^{3\alpha}\hbar^3\text{Sin}(x+y)}{\Gamma(1+3\alpha)} \\ & \vdots \\ & \vdots \\ & \vdots \end{aligned}$$

and so on.

Here, we consider the results upto $m = 10$ and rest of the components can be evaluated by iteration formula (43).

Therefore the solution of Eq. 19 is given by

$$\xi(x, y, t) = \xi_0(x, y, t) + \sum_{m=1}^{\infty} \xi_m(x, y, t) . \tag{44}$$

At $\hbar = -1$ we obtain the following approximation:

$$\begin{aligned} \xi(x, y, t) = & \text{Cos}(x+y) + \frac{4t^{3\alpha-1}\alpha\text{Cos}(x+y)\Gamma(\alpha)}{\Gamma(3\alpha)\Gamma(1+\alpha)} - \frac{12t^{2\alpha}\text{Cos}(x+y)}{\Gamma(1+2\alpha)} \\ & + \frac{8t^{3\alpha-1}\alpha\text{Cos}(x+y)\Gamma(2\alpha)}{\Gamma(3\alpha)\Gamma(1+2\alpha)} - \frac{6t^\alpha\text{Sin}(x+y)}{\Gamma(1+\alpha)} + \frac{6t^{2\alpha-1}\alpha\Gamma(\alpha)\text{Sin}(x+y)}{\Gamma(2\alpha)\Gamma(1+\alpha)} \\ & + \frac{2t^{3\alpha-2}\alpha\Gamma(\alpha)\Gamma(2\alpha-1)\text{Sin}(x+y)}{\Gamma(2\alpha)\Gamma(1+\alpha)\Gamma(3\alpha-1)} + \frac{4t^{3\alpha-2}\alpha^2\Gamma(\alpha)\Gamma(2\alpha-1)\text{Sin}(x+y)}{\Gamma(2\alpha)\Gamma(1+\alpha)\Gamma(3\alpha-1)} \\ & + \frac{8t^{3\alpha}\text{Sin}(x+y)}{\Gamma(3\alpha+1)} \end{aligned} \tag{45}$$

when $\alpha = 1$ Eq. 45 shows the similar results as [70] which is the exact solution of Eq. 37

$$\xi(x, y, t) = \text{Cos}(x+y) - 2t^2\text{Cos}(x+y) - 2t\text{Sin}(x+y) + \frac{4}{3}t^3\text{Sin}(x+y) + \dots$$

After simplification we can get Eq. 39.

5 Conclusion

Here, we have applied HASTM for solving fractional third order dispersive partial differential equations. It is shown that HASTM is an effective alternate tool for the evaluation of linear and nonlinear partial differential equations, which is not require any physical perturbation quantity, lucid to understand. Thus, we can conclude that

the transform method with HAM is very effective and highly accurate for any lower and higher order fractional partial differential equations.

References

1. Oldham, K.B., Spanier, J.: *The Fractional Calculus: Theory and Application of Differentiation and Integration to Arbitrary Order*. Academic Press, California (1974)
2. Caputo, M.: Linear models of dissipation whose Q is almost frequency independent. Part II. *Geophys. J. R. Astron. Soc.* **13**, 529–39 (1967)
3. Ross, B.: The development of fractional calculus 1695–1900. *Hist. Math.* **4**, 75–89 (1977)
4. Miller, K., Ross, B.: *An Introduction to the Fractional Calculus and Fractional Differential Equations*. Wiley, New York (1993)
5. Podlubny, I.: *Fractional Differential Equations*. Academic Press, New York (1999)
6. Kilbas, A.A., Srivastava, H.M., Trujillo, J.J.: *Theory and Applications of Fractional Differential Equations*. Elsevier, Amsterdam (2006)
7. Kilbas, A.A., Srivastava, H.M., Trujillo, J.J.: *Theory and Applications of Fractional Differential Equations*. Jan Van Mill, North–Holland (2006)
8. Diethelm, K.: *The Analysis of Fractional Differential Equations*. Springer, New York (2010)
9. Caponetto, R., Dongola, G., Fortuna, L., Petras, I.: *Fractional Order Systems: Modeling and Control Applications*. WSPC, Singapore (2010)
10. Samko, S.G., Kilbas, A.A., Marichev, O.I.: *Fractional Integration and Derivatives: Theory and Application*. G & B, Amsterdam (1993)
11. Uchaikin, V.V.: Self-similar anomalous diffusion and levy stable laws. *Phys. Usp.* **26**(8), 821–849 (2003)
12. Tarasov, V.E.: Fractional generalization of gradient and Hamiltonian systems. *J. Phys. A Math. Gen.* **38**(26), 5929–5943 (2005)
13. Laskin, N., Zaslavsky, G.M.: Nonlinear fractional dynamics of lattice with long-range interaction. *Phys. A Stat. Mech. Appl.* **368**(1), 38–54 (2006)
14. Tarasov, V.E.: Gravitational field of fractal distribution of particles. *Celest. Mech. Dyn. Astron.* **94**(1), 1–15 (2006)
15. Fujioka, J.: Lagrangian structure and Hamiltonian conservation in fractional optical solitons. *Commun. Fract. Calc.* **1**(1), 1–14 (2010)
16. Tarasov, V.E.: *Fractional Dynamics: Applications of Fractional Calculus to Dynamics of Particles, Fields and Media*. Springer, Heidelberg (2010)
17. El-Wakil, S.A., Abulwafa, E.M., El-Shewy, E.K., Mahmoud, A.A.: Time-fractional KdV equation for electron-acoustic waves in plasma of cold electron and two different temperature isothermal ions. *Astrophys. Space Sci.* **333**(1), 269–267 (2011)
18. El-Wakil, S.A., Abulwafa, E.M., El-Shewy, E.K., Mahmoud, A.A.: Time-fractional study of electron acoustic solitary waves in plasma of cold electron and two isothermal ions. *J. Plasma Phys.* **78**(6), 641–649 (2012)
19. Petr, I.: *Fractional-Order Nonlinear Systems Modelin. Analysis and Simulation*. Higher Education Press and Springer, Beijing (2011)
20. Klafter, J., Lim, S.C., Metzler, R. (eds.): *Fractional Dynamics: Recent Advances*. World Scientific, Singapore (2012)
21. Baleanu, D., Tenreiro Machado, J.A., Luo, A.C.J. (eds.): *Fractional Dynamics and Control*. Springer, New York (2012)
22. Zhou, J.K.: *Differential Transformation Applications for Electrical Circuits*. Wuhan Huazhong University Press (1986)
23. Chen, C.K., Ho, S.H.: Solving partial differential equations by two-dimensional differential transform method. *Appl. Math. Comput.* **106**(2–3), 171–179 (1999)
24. Jang, M.J., Cheng, C.L., Liu, Y.C.: Two-dimensional differential transform for partial differential equations. *Appl. Math. Comput.* **121**(2–3), 261–270 (2001)
25. Yang, S., Xiao, A., Su, H.: Convergence of the variational iteration method for solving multi-order fractional differential equations. *Comput. Math. Appl.* **60**(10), 2871–2879 (2010)

26. Wazwaz, A.M.: The variational iteration method for analytic treatment for linear and nonlinear ODEs. *Appl. Math. Comput.* **212**(1), 120–134 (2009)
27. Saravanan, A., Magesh, N.: An efficient computational technique for solving the Fokker–Planck equation with space and time fractional derivatives. *J. King Saud Univ. Sci.* **28**, 160–166 (2016)
28. Balaji, S.: Legendre wavelet operational matrix method for solution of fractional order Riccati differential equation. *J. Egyptian Math. Soc.* **23**(2), 263–270 (2015)
29. Erturk, V.S., Momani, S., Odibat, Z.: Application of generalized differential transform method to multi-order fractional differential equations. *Commun. Nonlinear Sci. Numer. Simul.* **13**(8), 1642–1654 (2008)
30. Zhao, J., Tang, B., Kumar, S., Hou, Y.: The extended fractional sub-equation method for nonlinear fractional differential equations. *Math. Probl. Eng.* **2012**(11), 924956 (2012)
31. He, J.H.: Asymptotology by homotopy perturbation method. *Appl. Math. Comput.* **156**(3), 591–596 (2004)
32. He, J.H.: The homotopy perturbation method for non-linear oscillators with discontinuities. *Appl. Math. Comput.* **151**(1), 287–292 (2004)
33. He, J.H.: A coupling method of a homotopy technique and a perturbation technique for nonlinear problems. *Int. J. Non-Linear Mech.* **35**(1), 37–43 (2000)
34. He, J.H.: Application of homotopy perturbation method to nonlinear wave equations. *Chaos Solit Fractals.* **26**(3), 695–700 (2000)
35. Liao, S.J.: The proposed homotopy analysis technique for the solution of nonlinear problems. Ph.D. Thesis, Shanghai, Jiao Tong University (1992)
36. Liao, S.J.: *Beyond Perturbation Introduction to the Homotopy Analysis Method*. Chapman & Hall CRC Press, Washington DC (2004)
37. Liao, S.J.: *Homotopy Analysis Method in Nonlinear Differential Equations*. Springer, Heidelberg (2012)
38. Liao, S.J.: *Advances in the Homotopy Analysis Method*. World Scientific Press, Singapore (2014)
39. Kumar, S.: An analytical algorithm for nonlinear fractional Fornberg–Whitham equation arising in wave breaking based on a new iterative method. *Alexand. Eng. J.* **53**(1), 225–231 (2014)
40. Kumar, S., Singh, J., Kumar, D., Kapoor, S.: New homotopy analysis transform algorithm to solve volterra integral equation. *Ain Shams Eng. J.* **5**(1), 243–246 (2014)
41. Yin, X.B., Kumar, S., Kumar, D.: A modified homotopy analysis method for solution of fractional wave equations. *Adv. Mech. Eng.* **7**(12), 1–8 (2015)
42. Kumar, S., Kumar, A., Baleanu, D.: Two analytical methods for time-fractional nonlinear coupled Boussinesq–Burger’s equations arise in propagation of shallow water waves. *Nonlinear Dyn.* **85**(2), 699–715 (2016)
43. Kumar, S., Kumar, D., Singh, J.: Fractional modelling arising in unidirectional propagation of long waves in dispersive media. *Adv. Nonlinear Anal.* (2016). doi:[10.1515/anona-2013-0033](https://doi.org/10.1515/anona-2013-0033). ISSN (Online) 2191–950X, ISSN (Print) 2191–9496
44. Arikoglu, A., Ozkol, I.: Solution of a fractional differential equations by using differential transform method. *Chaos Solit. Fractals.* **34**(4), 1473–1481 (2007)
45. Erturk, V.S., Momani, S.: Solving systems of fractional differential equations using differential transform method. *J. Comput. Appl. Math.* **215**(1), 142–151 (2008)
46. Pandey, R.K., Mishra, H.K.: Homotopy analysis fractional Sumudu transform method for time-fractional fourth order differential equations with variable coefficients. *Am. J. Numer. Anal.* **3**(3), 52–64 (2015)
47. Waleed, H.: Solving nth-order integro-differential equations using the combined laplace transform–Adomian decomposition method. *Am. J. Appl. Math.* **4**(6), 882–886 (2013)
48. Haghghi, A.R., Dadvand, A., Ghejlo, H.H.: Solution of the fractional diffusion equation with the Riesz fractional derivative using McCormack method. *Commun. Adv. Comput. Sci. App.* **2014**(11), 00024 (2012)
49. Watugala, G.K.: Sumudu transform—a new integral transform to solve differential equations and control engineering problems. *Math. Eng. Ind.* **6**(4), 319–329 (1998)
50. Weerakoon, S.: Application of Sumudu transform to partial differential equations. *Int. J. Math. Educ. Sci. Tech.* **25**(2), 277–283 (1994)
51. Weerakoon, S.: Complex inversion formula for Sumudu transform. *Int. J. Math. Educ. Sci. Tech.* **29**(4), 618–621 (1998)

52. Belgacem, F.B.M., Karaballi, A.A., Kalla, S.L.: Analytical investigations of the Sumudu transform and applications to integral production equations. *Math. Probl. Eng.* **2003**(3), 103–118 (2003)
53. Belgacem, F.B.M., Karaballi, A.A.: Sumudu transform fundamental properties investigations and applications. *J. Appl. Math. Stoch. Anal.* **2006**(23), 91083 (2006)
54. Belgacem, F.B.M.: Introducing and analysing deeper Sumudu properties. *Nonlinear Studies* **13**(1), 23–41 (2006)
55. Hussain, M.G.M., Belgacem, F.B.M.: Transient solutions of Maxwell's equations based on Sumudu transform. *Prog. Electromagn. Res.* **74**, 273–289 (2007)
56. Belgacem, F.B.M.: Sumudu applications to Maxwell's equations. *Prog. Electromagn. Res.* **5**(4), 355–360 (2009)
57. Belgacem, F.B.M.: Sumudu transform applications to bessel functions and equations. *Appl. Math. Sci.* **4**(74), 3665–3686 (2010)
58. Gupta, V.G., Sharma, B., Belgacem, F.B.M.: On the solutions of generalized fractional kinetic equations. *Appl. Math. Sci.* **5**(19), 899–910 (2011)
59. Katatbeh, Q.D., Belgacem, F.B.M.: Applications of the Sumudu transform to fractional differential equations. *Nonlinear Studies* **18**(1), 99–112 (2011)
60. Guerrero, F., Santonja, F.J., Villanueva, R.J.: Solving a model for the evolution of smoking habit in Spain with homotopy analysis method. *Nonlinear Anal. Real World Appl.* **14**(1), 549–558 (2013)
61. Jafari, H., Firoozjaee, M.: A multistage homotopy analysis method for solving non-linear Riccati differential equations. *IJRRAS* **4**(2), 128–132 (2010)
62. Alomari, A.K., Noorani, M.S.M., Nazar, R., Li, C.P.: Homotopy analysis method for solving fractional Lorenz system. *Commun. Nonlinear Sci. Numer. Simulat.* **15**(7), 1864–1872 (2010)
63. Elsaid, A.: Homotopy analysis method for solving a class of fractional partial differential equations. *Commun. Nonlinear Sci. Numer. Simulat.* **16**(9), 3655–3664 (2011)
64. Mohyud-Din, S.T., Yildrm, A., Yikl, E.: Homotopy analysis method for space- and time-fractional kdv equation. *Int. J. Numer. Method H.* **22**(7), 928–941 (2012)
65. Fadravi, H.H., Nik, H.S., Buzhabadi, R.: Homotopy analysis method for solving foam drainage equation with space- and time-fractional derivatives. *Int. J. Differ. Equ.* **2011**(12), 237045 (2011)
66. Djidjeli, K., Twizell, E.H.: Global extrapolations of numerical methods for solving a third-order dispersive partial differential equations. *Int. J. Comput. Math.* **41**(1–2), 81–89 (1999)
67. Twizell, E.H.: *Computational Methods for Partial Differential Equations*. Ellis Horwood New York: Chichester and John Wiley and Sons (1984)
68. Lewson, J.D., Morris, J.L.I.: The extrapolation of first-order methods for parabolic partial differential equations. I. *SIAM J. Numer. Anal.* **15**(6), 1212–12124 (1978)
69. Mengzhaoh, Q.: Difference scheme for the dispersive equation. *Computing* **31**, 261–261 (1983)
70. Wazwaz, A.M.: An analytic study on the third-order dispersive partial differential equation. *Appl. Math. Comput.* **142**(2–3), 511–520 (2003)
71. Kanth, A.S.V.R., Aruna, K.: Solution of fractional third-order dispersive partial differential equations. *J. Egyptian Math. Soc.* **2**(3), 190–199 (2015)
72. Luchko, Y., Gorenflo, R.: An operational method for solving fractional differential equations with the Caputo derivatives. *Acta. Math. Vietnamica.* **24**(2), 207–233 (1999)
73. Moustafa, O.L.: On the Cauchy problem for some fractional order partial differential equations. *Chaos Solit Fractals.* **18**(1), 135–140 (2003)