RESEARCH PAPER

Efficient label-free CTC enrichment using novel elevated height chip chamber by vortex technology

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Abstract

Circulating tumor cells (CTCs) are essential evidence for monitoring tumor dynamics and efficient cancer therapy. Besides its clinical importance, it is a feasible method for liquid biopsy in oncological studies. Vortex is a label-free technique using inertial microfluidic technology to separate CTCs from blood with high throughput. A possible moderate separation efficiency arises due to the loss of particles from the vortex reservoir during particle-to-particle interactions. This study develops a novel three-dimensional reservoir geometry of the vortex with more efficient CTCs separation. In the proposed geometry, the height of the main channel is the same as the vortex chips and sets to 70 µm; however, the reservoir height increased to 100 µm. This design is called elevated height chip (EHC). The fnite element simulation method is employed for optimizing reservoir height and working Reynolds number (Re). To verify the functionality of the EHC, the constant height chip design was also fabricated and showed a good correlation with the previously reported result. Both chips are tested with polystyrene beads (20, 15, and 8 µm) and MCF-7 cells. The results show 40% and 20% increase in efficiency for polystyrene beads and MCF-7 cells, respectively. Finally, it is demonstrated that 45% efficiency and 80% purity per cycle are achieved using the new proposed structure, and it is increased to 80% efficiency and 80% purity for over three cycles repetition.

Keywords 3D vortex · Circulating tumor cells · CTC separation · Inertial focusing

1 Introduction

CTCs are cancerous cells that travel in the blood. They are disseminated from tumor tissue and can seed new tumors in the other tissues, called the metastases stage of cancer (Bankó et al. [2019](#page-11-0)). More than 90% of deaths from cancer are caused by metastases (Racila et al. [1998\)](#page-11-1). Thomas Ashworth, an Australian pathologist, frst suggested a theory that CTCs are a fundamental prerequisite for metastasis in the mid-nineteenth century (Rostami et al. [2019\)](#page-12-0). Moreover, scientists could also obtain the whole genome of a CTC with advances in next-generation sequencing (NGS) and single-cell sequencing (SCS) technologies and relate it to primary and metastatic tumors. CTC accumulation in the

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blood shows the aggressiveness and may provide insight into the patient's ongoing treatment response (Riethdorf et al. [2007](#page-11-2)). CTCs enumerating serves as liquid biopsy that could be an alternative to the cancer screening method, replacing more invasive and time-consuming techniques such as tumor biopsy (Bankó et al. [2019\)](#page-11-0). So, isolation and analysis of CTCs can provide a complete description of the metastatic progression to better indicate the efectiveness of targeted therapy and discover CTC properties (Shen et al. [2017](#page-12-1); Cho et al. [2018\)](#page-11-3). However, a few challenges include the purifcation of CTCs from other blood cells and isolating and extracting cells without destroying or losing DNA and RNA (Cheng et al. [2019](#page-11-4)).

Many existing technologies employ affinity-based capturing methods, using antibodies or aptamers binding to surface markers of cells (Karabacak et al. [2014](#page-11-5)). CellSearch (Menarini Silicon Biosystems), the frst and the unique FDA-approved CTC detection system, exploits epithelial molecular adhesion molecule (EpCAM) antibody in the enrichment step (Andreopoulou et al. [2012\)](#page-11-6). Monitoring the prognosis of tumor recurrence or metastasis in patients with

metastatic breast, prostate, or colorectal cancer is possible by this method (Castro-Giner and Aceto [2020](#page-11-7)). Although it is an efective clinical application device, it has a comparatively low detection sensitivity. It cannot diferentiate between heterogeneous tumor cells, which may involve the downregulation of target epithelial cell surface markers (Sollier et al. [2014;](#page-12-2) Che et al. [2016](#page-11-8)). This method requires cell fxation, preventing further analysis requiring viable cells for the isolated sample. CTCs are scarce cells that are less than 100 per mL of blood (Cheng et al. [2019](#page-11-4)). So, separation methods with low efficiency may cause faulty results in clinical data.

As an alternative to biophysical markers, physical properties such as size, density, deformability, optical and dielectric properties have been used (Cho et al. [2018](#page-11-3); Cheng et al. [2019\)](#page-11-4). For example, ISET (Vona et al. [2000](#page-12-3)) and ScreenCell (Desitter et al. [2011\)](#page-11-9) have used porous mechanical flters to separate CTCs. On the other hand, other isolation methods have been proposed, including hydrodynamic focusing with active separation chambers and positive or negative magnetic depletion (Cheng et al. [2019](#page-11-4)). These methods usually have drawbacks such as time-consuming preparation, clogging the filter pores, and low efficiency.

CTCs are generally larger than normal blood cells. Red Blood Cells (RBC) are the most numerous cells, about a few billion cells per mL with diameters between 6 and 8 μm (Liu et al. [2019\)](#page-11-10). White Blood Cells (WBC) are larger than RBCs, with an average diameter of $10-15 \mu m$, and there are few millions WBCs per 1 mL of blood (Liu et al. [2019](#page-11-10)). However, CTCs have an average diameter of about 15–20 µm (Liu et al. [2019\)](#page-11-10). The diference between cell diameters is essential for separating with inertial methods such as serpentine, spiral microchannel, vortex technology, etc. Despite the spiral channel's high efficiency for separation higher than 85% (Warkiani et al. [2014](#page-12-4)), the vortex separates the CTCs with high purity and high throughput in a low volume solution (Sollier et al. [2014\)](#page-12-2).

Vortex chip includes straight pre-focusing microchannels and comes with expanded regions called reservoirs. Vortex chip traps CTCs inside reservoirs with the rotating flow called micro-vortices, induced by inertial forces at high Reynolds number (i.e., 150–250) followed by sudden downstream expansion region. Inertial forces push cells toward vortex flow in the expansion region, where particles larger than the specifed threshold gain enough displacement to enter rotating streams. Vortex technology, frst was introduced by Di Carlo in 2010, is a robust approach that presents high purity, high throughput, label-free, viable, and high enrichment ratio separation method (i.e., $10^3 - 10^4$); however, the capturing efficiency is still a challenging issue (Hur et al. 2011). Despite the limited capture efficiency, clinical validations on 12 patients showed that vortex successfully isolated≥5 CTCs/mL from 75% of patients, while CellSearch

just detected 17% samples positive from the same population (Sollier et al. [2014](#page-12-2)). Comparison of the vortex literature works is shown in Table [1](#page-2-0), briefy.

Mach in 2012 showed 20% capturing efficiency and 40% purity of trapping 500 MCF-7 cells/mL spiked in diluted blood using 1000×300 µm reservoirs (Mach et al. [2011](#page-11-12)). Che et al. ([2013\)](#page-11-13) presented a high correlation between MCF-7 and 15 μ m beads with almost 20% capturing efficiency and more than 80% purity using 720×220 µm reservoirs. Also, they showed 47% capturing efficiency for 20 µm beads. Sollier et al. [\(2014](#page-12-2)), characterized vortex capturing behavior concerning fow rate. Maximum capturing efficiencies for both 19 and 15 μ m beads were 60% and 22%, respectively, occurring in 6 mL/min ($Re = 210$), while for suspended MCF-7 in PBS sample was 21% in 4 mL/ min (Re=160) (Sollier et al. [2014\)](#page-12-2). Che et al. attempted to enhance vortex capturing efficiency by optimizing the total number of serialized and parallelized reservoirs and interconnecting channel length. This design which is called Vortex HT earned 62% capture efficiency for 20 μ m beads and 22% for MCF-7 $10 \times$ diluted blood, which was increased to ~36% by reintroducing waste (Che et al. 2016). The high number of serialized reservoirs increases back pressure and deforms the chip causing its lifetime reduction. Also, since cells are elastic material and highly deformable, highpressure separation has a negative effect on efficiency, so increasing efficiency for particle samples is higher than cell samples in Vortex HT. Vortex HT device reached 80% capturing efficiency after seven sample reintroducing cycles; however, due to extended run time and other practical circumstances, 45% capturing efficiency after 3 cycles were obtained as practical sample cycle in Vortex BioScience System. Paiè et al. ([2017\)](#page-11-14) introduced an engineered reservoir with two parallel rotation orbits inside the cavity by inserting a separation wall along the particles' pathway to decrease particle–particle interaction between trapped particles. The device showed 65% efficiency for 20 μ m rigid polystyrene beads versus 50% for the standard vortex chip (Paiè et al. [2017](#page-11-14)). At the last commercialization step, the Vortex chip was fabricated on a diferent material (Lemaire et al. [2018\)](#page-11-15) due to the high deformability of PDMS. The new PMMA chip could reach 80% in the separation of MCF-7 cells with triple sample reprocessing.

However, all the literature works demonstrate the 2D design of the Vortex chip, and it has not been studied how the out-of-plane forces could improve the trapping mechanism. Hence in this work, we propose a new threedimensional cavity vertical fow device that enables unique capturing efficiency for MCF-7 cells. The proposed device has a lower number of serialized reservoirs to decrease the backpressure and deformation of the chip and with an enormous volume to increase the particle capacity of reservoirs. Also, its larger height along the expanded region causes

multi-stable orbits with reduced particle–particle interaction. In addition, to choose optimized design parameters, a semi-theoretical method combined with previously reported experimental phenomena was adopted. Separatrix boundary location, vortex shape within the cavity, the magnitude of shear rate, and trapping time were enlisted as optimization parameters. Moreover, cavity capacitance, capturing efficiency, and purity of the new design were measured for both polystyrene beads and MCF-7 cells, and the results were compared to the CHC vortex chip. The newly developed device shows good consistency for MCF-7 and polystyrene bead trapping behavior over a higher Reynolds number range, while the standard control vortex chip shows low capturing for MCF-7 cells.

2 Vortex technology and device design

Vortex technology was developed as a microfuidic platform to quickly isolate intact CTCs from whole blood using inertial microfuidics and microscale vortices. Micro-vortices form as soon as a high inertia fuid inside a narrow microchannel suddenly jets into an expanded reservoir connected to another narrow microchannel (Fig. [1](#page-3-0)a). In the case of high fuid inertia, the relative low-pressure region forms nearby the cavity entrance while the high-pressure region forms near the cavity outlet. This pressure gradient detaches the streamlines near the outlet (trailing wall) and forms isolated circulating streams inside the cavity. Channel Reynolds number (Re_c) is a dimensionless parameter defined as a ratio of inertial to viscous forces.

$$
\text{Re}_{\text{c}} = \rho U_{\text{m}} D_{\text{h}} / \mu. \tag{1}
$$

Moreover, the Particle Reynolds number is defned as the following relation.

$$
\text{Re}_{\text{p}} = \text{Re}_{\text{c}} \bigg(\frac{a_{\text{p}}}{D_{\text{h}}} \bigg),\tag{2}
$$

where D_h is the characteristic channel diameter, ρ is the fluid density, U_m is the maximum velocity, μ is the dynamic viscosity, and a_p is particle diameter. At particle Reynolds number (Re_p) > 1 and a_p/D_h > 0.07 (i.e., when the particle diameter is noticeable compared to the channel diameter), inertial focusing takes place, and particles focus at the center of the broader face of a rectangular cross-section (For instance, a 10 μ m particle in a 40*70 μ m channel). The entering narrow microchannel focuses the particles in two stable equilibrium positions and prepares the particle to take the maximum pos-sible effect from the unexpected geometry changes (Fig. [1](#page-3-0)a). The shear gradient lift force directs particles toward the

Fig. 1 Chip design inspired from Vortex HT reservoirs; **a** trajectory of 8, 20 μ m microparticles in the EHC design. The reservoir height is 110 μ m, and the channel height is about 70 μ m. **b** Vortex HT reservoir (CHC) in the same fow rate with a constant height of about 70 μ m. Shear and wall-induced forces on the particles in **c–f** areas in red boxes are shown in **c**–**f**, respectively, to compare forces in the *Y* and *Z* directions. Shear-induced lift forces and wall-induced lift forces are shown with black and white arrows, respectively, in **c**–**f**

schematics. Microparticles have a vertical displacement in the proposal design. When they enter the upper part of the reservoir, their speed will be reduced soon, and the up straight path is close to continuing. Thus, they are forced to orbit in the vortex and cannot return to the main fow stream because of the high diference in fow speed between the two parts of the reservoir. Also, in the EHC design, the space for the particles is wider to orbit on the *z*-axis. Hence, there could be non-planar orbits that could carry the particles

channel wall while it repels by counteracting wall-induced lift force. The overall forces form an ordered train of particles at approximately 0.4 channel diameter away from the wall. Moreover, the wall force is abruptly eliminated when a particle enters the cavity (Fig. [1](#page-3-0)a). Therefore, the shear gradient lift force (F_s) accelerates the particles toward the vortex region with no resistant force except stokes drag. F_s can be defned as:

$$
F_{\rm s} = f_{\rm L} \rho U_{\rm m}^2 a^4 / H^2,\tag{3}
$$

where f_L is defined as a non-dimensional lift coefficient that is a function of the Reynolds number and the normalized cross-sectional position. Since F_s scales with a^4 (Di Carlo [2009\)](#page-11-18). Therefore, only larger than the threshold size particles may have enough lateral displacement to cross the separatrix boundary and enter the vortex region (Liu et al. [2016](#page-11-19)). When a particle enters the vortex region, it experiences inertial and momentum forces and forms a stable limit cycle orbit.

In the CHC vortex chip, the particles entering the chamber are afected with a balanced force containing shear gradient and wall-induced lift force from sidewalls. As the particle enters the reservoir, the wall-induced force is suddenly removed, and the particle experiences a net-force perpendicular to the fow, which leads it through the chamber. Meanwhile, a new balance would be formed, and the particle orbits in the chamber at a specifc height, as shown in Fig. [1b](#page-3-0). However, in the EHC design, there are more forces in touch with the particle. The conventional forces in CHC also play the dominant role in EHC design; however, new forces would appear and alter the mechanism efectively. Wall-induced forces are present from the sidewalls and the top and bottom walls due to the considerable size of target particles.

On the other hand, shear-induced forces contain horizontal and vertical components because of the 3D surface of the velocity profle in the channel cross-section. However, in the conventional CHC, the vertical force was not employed. In

the EHC design, the vertical force component causes a vertical movement of particles in the reservoir and an asymmetric shape in orbiting path. This out-of-plane net force makes the particle move more facile and easier to get out of the mainstream and keeps it more stable in the reservoir. This vertical movement presents a new space with a higher volume and lower high-pressure region near the walls. It enables the particle to change its altitude in the reservoir continuously. Hence, there would be more space in the EHC device for the particle to orbit than CHC, which is limited to a specifc orbit in the reservoir and reduces the particle–particle efect.

The separatrix region is one of the remarkable factors for evaluating each design at diferent fow rates. Separatrix is the largest vertical distance that a passing particle deviates from the main streamline without entering into the vortex. As shown in Fig. [3](#page-5-0)a, particles smaller than a threshold move upward in the flow, and the flow draws them into the outlet and finally exits from the reservoir. Particles that pass vertically from this zone would enter the orbit and return to the reservoir entrance. The flow rate at the most achievable efficiency is defined as critical Reynolds. After injection of all samples, a bufer solution with the same Re is injected into the channel as a substitution for sample solution to wash unmeant minor debris contamination while larger particles maintain their orbits (Khojah et al. [2017\)](#page-11-20). Later, controlled releasing of trapped particles is achieved by reducing the fow rate, eliminating the vortex, and providing uniform laminar fow inside the cavity to wash out trapped cells (Fig. [1\)](#page-3-0).

We used two microfluidic chips and tested for CTC separation from the bufer. The geometry of CHC originates from maintaining the layout of the Vortex chip. This design consists of 16 parallel channels (i.e., separated by a 1 mm gap) with an arrangement of six reservoirs along each channel. The reason behind using six reservoirs in each channel is the high efect of pressure resistance in the separation behavior of each reservoir. Sollier's results (Sollier et al. [2014;](#page-12-2) Lemaire et al. [2018\)](#page-11-15) show that the fexible PDMS chip has lower capturing efficiency than rigid PMMA chips. Furthermore, it shows that the last four reservoirs of PDMS chips (since they experience low pressure) have around 70–80% capturing capacity. So, we chose six reservoirs to optimize the design. The dimension of each reservoir is $900 \times 300 \mu$ m, and the main channel has 70 µm height (*h*) and 40 µm width (*w*). In EHC, the height of the reservoir is increased to 110 µm, while other parameters were considered the same as the CHC design (Fig. [1\)](#page-3-0).

3 Materials and methods

3.1 Chemicals

A negative photoresist, SU-8 2050 (MicroChem, MA, USA), was used to make the mold by UV lithography on a 4-in.

Si wafer. Sylgard 184 Polydimethylsiloxane (PDMS) was procured from Dow Corning (Midland, MI, USA). Glycol Methyl Ether Acetate (PGMEA) is used as SU-developer. Standard microbeads from Phosphorex (Hopkinton, MA, USA) in three diferent average diameters of 8, 15, and 20 μm. MCF-7 cells were cultured in DMEM medium; fetal bovine serum (FBS) and penicillin–streptomycin (Invitrogen) were acquired. These chemicals were purchased from Sigma-Aldrich (Taufkirchen, Germany).

3.2 Device fabrication

Two microfuidic chips (i.e., CHC and EHC) were fabricated using SU-8 mold and PDMS for replica molding. The SU-8 master mold was made on a silicon wafer using the conventional UV lithography method. CHC device was fabricated based on the procedure mentioned in Dhar et al. ([2015](#page-11-21)) by applying and spinning SU-8 photoresist on a Si wafer to obtain 70 µm resist thickness and following with pre-baking of photoresist. UV-light was then exposed to the photoresist followed by post-exposure-bake, and fnally, SU-8 was developed using Propylene Glycol Methyl Ether Acetate (PGMEA) as shown in Fig. [2a](#page-5-1) step 1 to 4.

The EHC chip was fabricated using two-step lithography, as shown in Fig. [2](#page-5-1). In the frst step, a 70 μm layer of SU-8 was spin-coated on a Silicon wafer. The frst design step, which contained the chip's geometry, inlet, and outlet, and 16 microchannels were fabricated on the substrate exposed to the UV light and then developed. In the second step, a 110 μm layer of SU-8 was spin-coated on the previous geometry, followed by a second mask alignment using the former geometry, then baked and exposed to UV light. After developing, the microfuidic mold was ready to use.

The PDMS chips were fabricated using replica molding. PDMS and the curing agent were mixed with a ratio of 10:1, respectively. The fnal mixture was applied to the master mold and degassed, then placed in the oven at 80 °C for 2 h to cure. The cured PDMS was peeled off from the master mold; formerly, inlet and outlet holes were punched by a 2 mm biopsy punch. After exposure to oxygen plasma for 30 s (Henniker HPT-100, NH, USA), the PDMS and a glass sheet were bound.

3.3 Sample preparation

Fluorescent polystyrene particles were utilized with mean diameters of 8, 15, and 20 μ m. 0.1% Tween 20 also was added to particle suspensions to prevent particle aggregation. Two experiments were designed to show the diference between the two chips with various concentrations of solutions. In the frst experiment, 20 µm particles with diferent concentrations were injected into both chips to compare the performance. In the other experiment, solutions containing

Fig. 2 Device fabrication **a** steps of SU-8 mold fabrication with multiple heights, **b** the frst mask geometry for 70 µm height region, **c** the second mask geometry for exposing on 110 µm layer in the EHC vortex chip and **d** exhibits the geometry of CHC (vortex HT) chip with 70 µm height

Fig. 3 Vortex morphology and separatrix boundary. **a** Streamline of fuid in reservoir area in Re=130 and separatrix boundary. The separatrix width is shown by S. **b** 2D plot of Re and separatrix boundary

in diferent heights of the reservoir; **c** maximum of shear rate in the diferent heights of the reservoir with 50, 150, 350, and 550 µm distances from the leading wall

8 µm and 20 µm were injected into both chips. As mentioned before, deionized water was used for the buffer solution in all the samples.

Also, MCF-7 cells were cultured in a DMEM medium. The medium was supplemented with 10% fetal bovine serum and 1% penicillin–streptomycin incubated at 37 °C under 5% $CO₂$. Cells were passaged once they reached 80% confluence, then fxed by 1% glutaraldehyde and spiked in PBS solution. A high concentration of the sample (with 2000 particles or cells) was used for saturation tests, and a low concentration (50–150 particles or cells per mL) was used for capture efficiency and purity tests.

3.4 Experimental setup and methods

An adjustable syringe pump was used to control fuid fow. In all the experiments, a digital camera with a high shutter speed (Canon D700) was mounted on the microscope (Nikon Ti-U) to analyze the trapping and releasing processes. The sample on the microscope was illuminated by a mercury arc lamp with a $10 \times$ and $4 \times$ objective lens. The minimum shutter speed of 1/4000 was used to investigate vortex formation and the particle motion in reservoirs. The fuorescence images were captured by a CCD Coolsnap HQ2 camera (Roper Scientifc, Evry, France), then processed by Zen2 software.

Before injecting particle suspensions, the chip was primed with DI water with an optimized flow rate for each chip to form the vortex streamlines and eliminate microbubbles in reservoirs. Promptly, the Particle suspension was injected into the channel to separate from the entire solution. Then, DI water was pumped into the channel again for 30 s to ensure all particles had entered the device. In the last step, by decreasing the flow rate $(Re=5)$, trapped particles were released from the reservoir and collected in a 96 well-plate for fnal analysis. Samples were mixed by vortex shaker to avoid particles to sediment before each step to preserve a uniform distribution of particles.

For the capacity test, 6 mL solution of 2000 particles per mL of 20 µm size particles was used. On the other hand, two solutions with concentrations of 50 and 150 particles/ mL of 20 µm size particles were used for efficiency analysis. Homogeneous mixtures of 8, 15, and 20 μ m particles were utilized to study the purity. Also, the same procedure was used to quantify the capacity of the MCF-7 cells in the device, and 6 mL of a sample of 100 cells/mL was used to analyze trapping efficiency in both chips.

3.5 Simulation

Both designs were simulated by COMSOL Multiphysics 5.2 (COMSOL, Burlington, MA) to investigate the streamlines' stability and velocity distribution in the micro vortex fow.

Laminar fow physics was applied in a single-phase fow regime. Water was used as the carrier fuid with a density (ρ) of 1000 kg/m³ and dynamic viscosity (μ) of 10⁻³ kg/m s. The effect of reservoir height was investigated by simulating diferent reservoirs heights (i.e., 70, 90, 110,130, 150, and 190 μ m). The microchannels connected to the reservoir contained a 500 µm length for each inlet and outlet. The flow was introduced to the channel by applying a fully-developed flow to satisfy the practical situation. The inlet flow rate was altered from 10 to 500 µl/min for a single channel in diferent steps to observe the best situation of the velocity, separatrix region, shear gradient force, and the vortex fow geometry results for the optimization. The fow Reynolds was measured by evaluating the mean channel velocity. Also, the results were validated by comparing the vortex shape in the fuorescent images of orbiting microparticles and the streamlines in the simulated reservoir at the same fow rate.

4 Results and discussion

4.1 Simulation and optimization

Vortex simulations in previous studies were performed using two diferent methods, including direct numerical simulation (DNS) and conventional CFD. DNS uses Lattice–Boltzmann Method (LBM) (Haddadi and Carlo [2017](#page-11-22); Haddadi et al. [2018\)](#page-11-16) and uses complex computations, but is not sufficiently consistent with the experimental results for separating CTC from blood due to the complex pathway of particles considering their interaction. Also, DNS is a time-consuming simulation process. Most importantly, particle–particle interaction is not considered in these methods (Jiang et al. [2018\)](#page-11-23). So numerical solutions are not reliable to obtain precise results.

The analytical CFD method was adopted in this work, and parameters were investigated in the conventional simulation method. We studied four parameters to obtain optimum cavity geometry in two groups. The frst group relates to the physical variables of the vortex, including the vortex morphology and the separatrix width. The second parameters group employed force-dependent variables, including shear gradient lift and residence time of particles. Vortex formation starts at a low Reynolds number (i.e., Re=30) from the beginning of the reservoir, called the leading wall, in the form of an incomplete triangular vortex cell. Increasing fow rate increases the Reynolds number of flow, and the vortex shape progressively expands and flls the cavity (Fig. S1). A complete fow cell is formed when a vortex covers the whole length of the cavity.

The effect of cavity aspect ratio and Re on the number of entrapped particles has been investigated in previous works (Haddadi and Carlo [2017](#page-11-22); Paiè et al. [2017;](#page-11-14) Haddadi

et al. [2018](#page-11-16)). It was shown that for each reservoir aspect ratio (defined as λ), the number of entrapped particles is maximized at a specifc Re, called Remax. Haddadi studied the *λ* efect on vortex formation and showed that the number of particles inside the cavity increases by increasing the cavity length, and Re_{max} shifts to higher values while the vortex center shifts longitudinally (Haddadi and Carlo [2017](#page-11-22); Haddadi et al. [2018](#page-11-16)). Similarly, the entrapped particles increase by increasing the reservoir width; however, no signifcant change in Remax (vertical shift) would be observed (Haddadi and Carlo 2017). While Re_{max} increases the flow rate, more advanced test equipment and stronger channel bonding would be needed for testing. This work shows that by increasing the height of the cavity, the chip can increase the particle trapping at a lower Reynolds number, as shown in Fig. [1.](#page-3-0) The obtained results in Sect. 4.4 show the efficiency improvement of this modifed geometry. To obtain optimum reservoir height and proper Reynolds number, their efect on residence time, shear gradient lift, and separation boundary were examined.

In our simulations, the reservoir shows distinguished vortex geometry due to the reservoir height at the same fow Re. As shown in Fig. S1, the vortex cells reach the critical Reynolds depending on the chamber height. The Re_{max} for 70, 110, and 150 µm reservoirs were obtained at 200, 150, and 100 Re respectively. The results represent the vortex cell forming and progress by streamlines compared to experimental images. Simulation results were validated by the vortex shape of the fuorescent beads orbiting the chamber in experiments and the obtained streamlines. Figure [1](#page-3-0) and Fig. S1 demonstrate some evidence of the verifed geometries of simulated models. However, the CHC device could be more explicit proof of the orbiting particles in the 2D ring and closer to the simulated model because of its inplane orbits. On the other hand, the EHC formed 3D orbits for particles and poses a diferent shape than the 2D simulated streamlines and remains close to the 3D streamline in simulation. Overall, the simulation results were verifed by the geometrical properties of the vortex comparing to the experiments results.

The separatrix width is the minimum lateral displacement of particles needed to enter vortex fow and be trapped (Fig. [3a](#page-5-0)) (Haddadi et al. [2018](#page-11-16)). Increasing the reservoir height would decrease the separatrix width (Fig. [4b](#page-7-0)), and the probability of trapping particles increases. Figure S1 demonstrates vortex fow formation at diferent channel heights, and it shows that vortex flow forms at a lower flow rate with increasing channel height. Residence time is considered when a particle moves across the reservoir from the leading wall to the trailing wall. Having a more significant residence time (the lower particle speed) results in experiencing higher total shear gradient lift force, higher lateral motion of the particle toward the core of the vortex, and a higher probability of trapping. As a result, increasing channel height increases particle traveling time and trapping probability.

Another critical parameter is the shear gradient lift force (F_s) . This force is directly related to the shear rate (i.e., Eq. (3) (3) ; the magnitude of Fs could be estimated by analyzing the shear rate. When a particle approaches the trailing wall, the shear rate induced force on the particle will reduce. Figure [3a](#page-5-0) shows four diferent positions of shear rate measurement, and Fig. [3c](#page-5-0) shows these measurements for diferent reservoir heights. Simulation results show that 110 µm height has a higher shear rate than other heights. Decreasing

Fig. 4 Tracking pathway of 20 µm fuorescent particles in CHC and EHC (new chip); **a**, **b** Re=80, **c**, **d** Re=150, **e**, **f** Re=170

the shear rate at heights higher than 110 µm may be due to the reduction of the overpressure and the stagnation of the fow. This reduces vertical fow magnitude in the *Z* direction compared to the lateral flow. Conclusively, the optimized height is chosen as $He = 110 \mu m$ because it shows higher Fs and has smaller separatrix width and higher traveling time regarding the formation of complete vortex in lower Re.

4.2 Stable orbits

Flow condition forms a limit cycle or stable orbits for specifc particle size (Haddadi and Carlo [2017](#page-11-22); Shen et al. [2018a](#page-12-5), [b](#page-12-6)). Simulation of particle trajectories shows that particles spiral towards a unique limit cycle orbit (Shen et al. [2018a](#page-12-5), [b\)](#page-12-6). The motion of an isolated particle on the limit cycle trajectory continues without deviation unless hydrodynamic interaction with other particles perturbs it. An equal amount of centrifuge and shear-induced lift force make the stable particle orbit. Collisions between multiple entrapped particles generate deviation from stable limit cycle orbit, and due to random mass exchange, particles may exit the cavity. This depletion depends on the number of particles inside the cavities. Particles may leave the limit cycle by collision with one another or because of lift force at inertial fow, which generates a lateral migration towards the recirculating fow. As the reservoir volume increases by increasing the height, more stable cycles appear and entrap more particles in different levels of *Z*-direction, which we call 3D vortex orbits.

The experiment results are in agreement with our simulations and indicate that elevating reservoirs height increases the reservoir capacity for particle trapping until an optimum height. The designed EHC and the CHC chips were tested for Re of 80 (Fig. [4a](#page-7-0), b), 150 (Fig. [4](#page-7-0)c, d), and 170 (Fig. [4](#page-7-0)e, f). The vortex flow formation is a validated simulation result of Fig. S.1. EHC streamlines in Re=170 seem denser, and particles experience smaller radius trajectories than $Re = 150$; the balanced forces in the vortex containing centrifugal and the shear gradient force might be unbalanced. Hence, the particle could exit the vortex trajectory, and the efficiency decrease (Jiang et al. [2018\)](#page-11-23). The following section's result shows EHC chip has a higher trapping efficiency in $Re = 15$. A comparison of CHC and EHC forescent imaging shows the number of stable rotation circuits has increased and are more efficient in EHC design. Furthermore, separatrix width is reduced, which is consistent with simulation results.

4.3 Saturation reservoir capacity

Although particle entrapment inside cavities can be controlled by the fow rate (Re), it also depends on the size of particles and saturation reservoir capacity. Expanding the volume of the reservoirs is a way to increase the number of trapped particles. We injected 20 mL of 20 µm beads and MCF-7 cells containing 2000 cells or particles at various flow rates (Re) to investigate the particle entrapment process. Lemaire showed that the reservoir capacity changes gradually in a single channel longitudinally because of PDMS deformation, and so the last reservoir has the most probability to trap the cells (Lemaire et al. [2018](#page-11-15)). We analyzed reservoir capacity by taking images from three diferent regions beginning, middle, and the end of the channel by A, B, and C, as shown in Fig. S.2, respectively. For instance, Fig. [5a](#page-8-0) and b show trapped beads in the frst and middle reservoirs.

Figure [5c](#page-8-0) shows the number of particles trapped, and as we can see, the reservoir capacity in our proposed chip (i.e., EHC chip) is higher than the CHC. It is noteworthy that the CHC design reaches the critical Re in higher fow rates because of its reservoir geometry than the EHC design. The

c Number of entrapped 20 µm and MCF-7 cells in each reservoir for EHC and CHC at the middle section in diferent Re

comparison between the two approaches in our tests clears that the reservoir's 3D geometry has a remarkable efect on the trapping mechanism. So, altering the geometry tends the reservoir to separate more particles/cells at lower Re. The EHC design represents the peak performance at about $Re = 150$ compared to the CHC device at $Re = 200$ for 20 μ m microspheres.

Maximum reservoir capacity takes place at $(Re_{\text{max}}=150)$ for the proposed chip (EHC) and at (Remax=200) for the CHC, which is consistent with simulation results. The new chip traps more particles at section C than the CHC and has a larger quantity for particle saturation. Sollier showed that the trapping efficiency for particles with different sizes is diferent at various Re (Sollier et al. [2014\)](#page-12-2). So, both chips were tested with 15 μ m and 20 μ m particles and MCF-7 cell line, and the result is shown in Fig. [5c](#page-8-0), and Fig. S.3, which is in agreement with the results reported by Sollier. The maximum number of 20 μ m and 15 μ m trapped beads for the CHC are 55 and 32, respectively, at Re=200, while the maximum number of MCF-7 trapped cells is 20 at $Re = 150$, and it happens because of the deformation of cell and its efect on particle–particle collision (Haddadi et al. [2018\)](#page-11-16).

However, trapped beads and cell samples show the same behavior in the proposed chip (EHC, Fig. [5](#page-8-0)c). The maximum number of trapped 20 µm and 15 µm particles and MCF-7 cells at $Re = 150$ are 60, 23, and 50, respectively (Fig. S4). So, the trapping peak of cells and microparticles are achieved at the same Re number in this work, which is crucial for cell separation from the blood. More stable orbits and lower particle–particle collision could be the reason for this improvement.

4.4 Capture efficiency and purity

Purity and enrichment ratio are the main advantages of the vortical fow separation (Sollier et al. [2014](#page-12-2); Che et al. [2016](#page-11-8)). The separation mechanism is extremely size-dependent (Wang et al. [2016](#page-12-7)). The induced lift force is theoretically unable to entrap the smaller particles from the main fow. In the previous papers (Mach et al. [2011;](#page-11-12) Sollier et al. [2014](#page-12-2); Che et al. [2016](#page-11-8)), many smaller particles were observed orbiting in the reservoirs because of the high concentration of particles/cells in the fow. Meanwhile, the washing step would exit them from the reservoirs. They orbit in the vortex region unstably, and afterward, most of them are thrown out of the vortex (Che et al. [2016](#page-11-8), Haddadi and Carlo [2017,](#page-11-22) Haddadi et al. [2018;](#page-11-16) Shen et al. [2018a,](#page-12-5) [b\)](#page-12-6). Haddadi showed that the removal efect for smaller particles might be the efect of particle–particle interaction (Haddadi and Carlo [2017](#page-11-22)). Altogether, the presence of smaller particles is low in this method.

Capture efficiency is the most critical factor for CTC separation. From the beginning of vortex technology, most of the studies were aimed to improve capture efficiency. These studies contained altering channel width (Mach et al. [2011](#page-11-12); Sollier et al. [2014;](#page-12-2) Dhar et al. [2015](#page-11-21); Haddadi and Carlo [2017](#page-11-22)), channel height (Zhou et al. [2013](#page-12-8); Sollier et al. [2014](#page-12-2)), channel serialization (Hur et al. [2011](#page-11-11); Che et al. [2016](#page-11-8)), changing the reservoir aspect ratio (Haddadi and Carlo [2017](#page-11-22); Shen et al. [2018a](#page-12-5), [b\)](#page-12-6), changing the outlet direction (Wang et al. [2013](#page-12-9); Rastogi et al. [2021](#page-11-17)), and changing the chip material (Lemaire et al. [2018;](#page-11-15) Volpe et al. [2019](#page-12-10)). The Vortex HT obtained the best clinical result, lower than 60%, after twice recycling the outlet (Che et al. [2016\)](#page-11-8). On the other hand, the last and the best vortex cell separation device, VTX-1, was proposed using a rigid plastic chip and achieved 80% capture efficiency (Lemaire et al. [2018](#page-11-15)).

We believe that in previous works there was not enough attention paid to the reservoirs' third axis expansion. As described in Sect. [2,](#page-2-1) the effect of the vertical net force is remarkable due to the result of our simulations and experiments. The particle passes through the reservoir with a velocity approximately equal to the fow (about 4.5 m/s). Therefore, the particles have just few millisecond time to cross the mainstream borders and enter the orbiting streamlines. Afecting the vertical force in the EHC chamber could move the particle more efectively than the 2D mechanism. That is due to lower velocity of flow in the expanded height of the chamber compared to the channel. Nevertheless, in the CHC design, the Re number of the fow is higher, and the passing time is less than the proposed design. Moreover, the orbiting particles in the chambers are not just in a 2D plane and have more space to orbit without much interaction with others, as shown in Fig. [4](#page-7-0). In the CHC design the particles are focused around a ring at a middle height of the reservoir and it saturates fast (Hur et al. [2011](#page-11-11)).

To measure capturing efficiency, three different concentrations of 20 µm particle were prepared. 300, 600, and 900 particles were diluted in 6 mL of water. Figure [6](#page-10-0)a, b show the last reservoirs of both chips with their trapped particles for diferent samples (Lemaire et al. [2018](#page-11-15)), and it shows that the new EHC traps more beads than CHC. A similar test by MCF-7 samples (600 cells/mL) was run. In Fig. [6](#page-10-0)c, the yield capacity of both chips is shown for diferent concentrations of beads and one concentration of MCF-7 cells. The new EHC and CHC efficiency for MCF-7 cells is measured at 45% and 19%, respectively. Also, for 20 µm beads, we obtained the separation efficiency of 75% and 55% for the EHC and the CHC, respectively. We believe that this considerable increase in efficiency is due to the larger orbiting space, lower possibility of collision, and better trapping elements.

For increasing the separation efficiency, recycling the chip's outlet and reintroducing it into the device is done twice (Che et al. [2016](#page-11-8)). In Fig. [6d](#page-10-0), the cyclic behavior of EHC and CHCs is shown. CHC reaches 45% efficiency for cell samples in 3

Fig. 6 Comparison of the number of particles and cells trapped in both chips. **a** Sample: 6 mL of 50 particles/mL. **b** Sample: 6 mL of 150 particles/mL. **c** Shows the number of captured beads, with the two devices varying the number of processed beads. Note that the new device always presents a higher efficiency than CHC. **d** Effi-

ciency of separation and serial processing on samples to further increase total cell capture. **e** Comparison of purity between two designs. Both chips represent approximately the same purity, similar to the reports in the literature

cycles of process, while the new chip efficiency has about 45% efficiency for only one cycle, and it reaches 80% after 3 cycles. Therefore, the EHC would save processing time besides the high efficiency.

To evaluate the purity of the EHC, a test was conducted using the Re_{max} of each chip (i.e., Re_{max} = 200 for CHC and Re_{max} = 150 for EHC). 6 mL homogeneous mixture of 8, 15, and 20 μ m particles with concentration of 100 particles per mL were used. We observed a reduction in the large particles outflow due to collisions while the high number of large particles removes smaller particles from the reservoir. Purity was calculated by dividing the number of separated 20 µm by the whole number of separated particles. As expected, the ratio of 20 µm to the total count of trapped particles in the EHC is approximately equal to the CHC. The purity of the EHC and CHC are around 80% and 78%, respectively (Fig. [6](#page-10-0)e).

5 Conclusion

Vortex technology has been widely used in size-based sorting of rare cells, being able to process the sample in a short time and with good trapping efficiency, high concentration of the sample, and high purity at very low volume. Nonetheless, the interparticle interactions and high fuid velocity are undoubtedly obstacles that do not allow all cells to enter and stay in the vortex currents. Also, increasing serialized reservoirs cause higher injection pressure that is not suitable for separating deformable cells. In this work, we examined the impact of diferent reservoirs height layouts on the efficiency of vortex-based systems, preferring geometries that encourage the splitting of the trapped particles and reduce the rate of vortex fow and increase the particle trapping with decreasing serialized reservoirs and injection pressure. In addition, the asymmetric structure of the channel causes the formation of a diferent stable of rotation orbits and lowers the possibility of impact between particles and their release, which was proved in our simulation results and experimental tests. Furthermore, parameters such as shear gradient lift force, separatrix boundary width, and residence time have been investigated in this work. The Re_{max} for the conventional device was about 200, while the new device presented in this work is 150. Lowering the optimal rate will reduce the speed of particles in the reservoir area and reduce their collisions. Furthermore, the simulation shows an increase in shear gradient force as the primary function of capturing particles. Therefore, the saturation number of particles for each reservoir in the new chip increases. Also, the constant separation efficiency of 45% for cells and particles was achieved. Finally, the efficiency increased to 80% after 3 cycles of separation.

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Declarations

Conflict of interest The authors declare that they have no confict of interest.

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