RESEARCH ARTICLE

Ramping down a clinical 3 T scanner: a journey into MRI and MRS at 0.75 T

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Abstract

Object Lower-feld MR is reemerging as a viable, potentially cost-efective alternative to high-feld MR, thanks to advances in hardware, sequence design, and reconstruction over the past decades. Evaluation of lower feld strengths, however, is limited by the availability of lower-feld systems on the market and their considerable procurement costs. In this work, we demonstrate a low-cost, temporary alternative to purchasing a dedicated lower-feld MR system.

Materials and Methods By ramping down an existing clinical 3 T MRI system to 0.75 T, proton signals can be acquired using repurposed ¹³C transmit/receive hardware and the multi-nuclei spectrometer interface. We describe the ramp-down procedure and necessary software and hardware changes to the system.

Results Apart from presenting system characterization results, we show in vivo examples of cardiac cine imaging, abdominal two- and three-point Dixon-type water/fat separation, water/fat-separated MR Fingerprinting, and point-resolved spectroscopy. In addition, the ramp-down approach allows unique comparisons of, e.g., gradient fdelity of the same MR system operated at diferent feld strengths using the same receive chain, gradient coils, and amplifers.

Discussion Ramping down an existing MR system may be seen as a viable alternative for lower-feld MR research in groups that already own multi-nuclei hardware and can also serve as a testing platform for custom-made multi-nuclei transmit/ receive coils.

Keywords Magnetic resonance imaging · Cine MRI · Multiparametric MRI · Technology assessment · Field strength

Introduction

In recent years, magnetic resonance imaging at lower felds (0.1 T \leq B₀ \leq 1 T) has received renewed interest [[1–](#page-16-0)[7\]](#page-16-1). This development is driven in parts by the necessity to reduce procurement and operation costs of MR systems in healthcare markets. As the magnet accounts for more than 30% of the overall system cost $[8]$ $[8]$ $[8]$, lowering the static field has been pursued as a cost-cutting option in conjunction with

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the design of sealed cryo-systems requiring minimal liquid helium to reduce installation and operation expenses. Lowering the field strength is beneficial in relation to specifc absorption rate (SAR) limits, as higher transmit-feld strengths at lower static feld allow for reduced radiofrequency pulse durations and hence longer sampling windows for a given repetition time as well as permitting higher fip angles in balanced and spoiled feld-echo protocols. At the same time, longer readouts are also supported by the increase of T_2^* at lower field, which allow recouping some of the signal-to-noise ratio (SNR) losses by reducing the readout bandwidth [\[9](#page-17-1)]. Reduced T_1 times in conjunction with larger fip angles in balanced and spoiled feld-echo protocols further allow to reduce these losses [\[10](#page-17-2)]. Moreover, the impact of the magneto-hydrodynamic (MHD) efect [[11\]](#page-17-3) is reduced at lower field, which makes ECG-triggering or gating more reliable. Finally, from the perspective of patient comfort, acoustic noise is lowered due to reduced

Lorentz forces between the cryostat and the gradient coils [\[12\]](#page-17-4).

Since 0.5 T systems entered clinics in the 1980s, signifcant advances in magnet technology, gradient and radiofrequency hardware, pulse sequence design and image reconstruction have been achieved, now rendering lower-feld systems potentially competitive with current high-feld systems for several applications [[13\]](#page-17-5). For example using a 0.35 T MRI split-bore scanner integrated into a radiation therapy setup, the value of lower felds for cardiac imaging application was pointed out early [[14](#page-17-6), [15](#page-17-7)]. Meanwhile, a number of 0.55 T systems have been deployed at several sites. The frst of these installations, at the National Institutes of Health (NIH), was used to demonstrate excellent image quality for cardiac, body, lung, and interventional imaging [[2](#page-16-2)]. This and earlier work sparked our interest in evaluating lower-field imaging [[14](#page-17-6)[–18](#page-17-8)] for cardiac $[17–22]$ $[17–22]$ $[17–22]$ and quantitative body imaging applications $[23, 1]$ $[23, 1]$ [24](#page-17-12)]. To keep time and cost overheads to a minimum, it was decided to temporarily ramp down a clinical 3 T system with high-performance gradients and utilize the existing $13¹³C$ transmit/receive hardware to study proton MRI and MRS at 0.75 T. Contrary to similar ramp-downs by the NIH, University of Southern California, and the New York University, the 3 T body transmit/receive coil remained in the scanner, making our installation reversible and relatively quick to perform (1/2 day per ramp-down/-up) and by repurposing existing hardware, overall costs were reduced. In the present manuscript, we give a detailed account of the ramp-down procedure on both a hardware and software level, show system characterization results, and present several in vivo MRI and MRS use cases.

Methods

The ramp-down was performed on a 3 T Philips Achieva scanner (Philips Healthcare, Best, the Netherlands) with a dual-amplifier gradient system delivering 200 T/m/s slew rate at a maximum gradient strength of 40 mT/m or 100 T/m/s at 80 mT/m.

Choice of feld strength

To use the 3 T multi-nuclei hardware for transmission and reception of the proton signal at the reduced feld strength $B₀$, the resonance frequency of the existing ¹³C transmit and receive coils had to match the Larmor frequency of protons. Using the gyromagnetic ratio of ¹³C ($\gamma_{13C} = 67.28$ rad MHz/T) and of ¹H ($\gamma_{1H} \approx 267.52$ rad MHz/T), the Larmor frequency is given by

$$
\gamma_{13C} \cdot 3T = \gamma_{1H} \cdot B_{0,}
$$

Fig. 1 Schematic of the ramp-down experiment. To receive proton signals at a reduced feld strength using existing 13C transmit/receive (*T*x/*R*x) hardware (**b**), a clinical 3 T system (**a**) was ramped down to 0.75 T (**c**). In both confgurations, the Larmor frequency is equal as the gyromagnetic ratio of ${}^{1}H$ is approximately four times larger than that of $13C$. As the scanner software remained configured for 3 T, pulse-programming software changes had to be made to account for

leading to

$$
B_0 = \frac{\gamma_{13C}}{\gamma_{1H}} \cdot 3T \approx 0.251 \cdot 3T \approx 0.75T.
$$

Hence, the target feld strength of 0.75 T was chosen. A schematic of the ramp-down procedure and required modifcations is shown in Fig. [1](#page-1-0).

Pulse‑programming changes

The system's software had to be changed to enable proton imaging at the reduced feld strength. From the software perspective, the scanner remained a 3 T transmitting and receiving at the 13C Larmor frequency. However, gradient felds were calculated using the proton gyromagnetic ratio instead to enable proton imaging at 0.75 T. The steps were as follows:

- (1) To calculate the correct gradient felds for imaging, the gyromagnetic ratio of ${}^{13}C$ was overwritten with the value for protons γ_{1H} .
- (2) After Larmor frequency determination, the frequency is checked to lie within a narrow range of the expected values based on the 3 T feld strength and the proton gyromagnetic ratio. These checks had to be disabled, otherwise the system would refuse to scan.
- (3) Preparation phases had to be adapted to use the local receive coils instead of the body coil. Pick-up coilbased preparation phases had to be disabled as these would have operated on the wrong resonance frequency.
- (4) Since parallel imaging (SENSE) and compressed sensing (CS-SENSE) scan features require a body coil reference scan for determining sensitivity maps, additional modifcations were required given that the former 3 T body coil was not operational:
	- a. The SENSE reference scan had to be adapted to run without a body coil. Here, the second acquisition that utilizes only the body coil for reception was adjusted to use a single receive coil element. These reference scans were only required to allow running SENSE/CS-SENSE accelerated scans and were automatically used by the vendor's reconstructor to produce images on the console. For ofine reconstruction, the acquired SENSE reference data were disregarded and coil sensitivities were estimated separately.
	- b. SENSE/CS-SENSE feature checks had to be overwritten to run without a body coil and with a non-SENSE-certifed coil.
-
- (5) For water/fat separation using multi-echo and multiacquisition protocols (DIXON), water/fat shift calculation had to be adapted to work with the resonance frequency determined during the preparation phases instead of using predefned feld strength-dependent settings.

Hardware modifcations

Ramp‑down

The 3 T magnet was ramped down three times between 2019 and 2021. As the critical temperature of the super-conductor (SC) increases with decreasing feld strength and self-heating is reduced due to the lower current density at lower magnetic feld, the SC heater had to be tested to ensure that the magnet could be safely taken off field. Hence, a quench test was performed at 0.75 T during the initial ramp-down trial. Using the 3 T quench system, the feld strength could be reduced to a value below 10 mT within 2 min as compared with 20 s for the product feld strength. This was deemed safe given that only specially trained personnel were allowed to enter the MR room.

Re‑optimization of shim irons

The 3 T system utilizes two sets of shim irons, (1), a factory shim set, which is bolted to the cryostat, and (2), a removable shim set, which can be optimized on-site using a staticfeld camera and a simulation tool provided by the vendor. For the 0.75 T configuration, the removable shim irons were re-optimized to compensate both the intrinsic inhomogeneity of the magnet at 0.75 T as well as the feld generated by the fxed 3 T factory shims. A main feld homogeneity of 1.3 ppm root-mean-squared (RMS), equal to 42 Hz RMS, was achieved within a 40 cm sphere. The homogeneity in the standard $50 \times 50x45$ cm³ volume was 4.7 ppm RMS (150 Hz RMS) and, thus, considerably lower than the relative specification for 3 T (\leq 3 ppm RMS), due to limited space on the shim rails to compensate for the bolted shim set. In absolute frequency terms, the lower-feld confguration provided a roughly 2.5-fold better feld homogeneity within the standard $50 \times 50 \times 45$ cm³ volume compared to the minimum requirements of the system at 3 T (absolute: \leq 383 Hz RMS). Only within a narrow 20 cm sphere could the system be shimmed within industry requirements $(\leq 0.04$ ppm) providing a true fourfold increase in feld homogeneity.

RF coils

The 3 T proton body coil could not be used and remained connected inside the scanner in a de-tuned state. For transmission and reception, an existing 13 C jacket double Helmholtz pair (Clinical MR Solutions, Brookfeld, WI, USA) was employed as a transmitter coil, while an existing $13C$ four-channel cardiac receive array (Clinical MR Solutions, Brookfeld, WI, USA) with two posterior and two anterior rectangular elements (7.6 cm x 18.3 cm) was used for signal reception. Both coils pre-existed and were used in $13¹³C$ experiments at 3 T before. The transmit coil was tested and calibrated on a standard load phantom. The maximum transmit-field strength $(B1+)$ was limited to 45 mT. To ensure safe operation in the in vivo setting, electromagnetic (EM) felds produced by coils and cables were measured and checked using a dedicated EM exposure acquisition system (EASY4MRI, Schmid & Partner Engineering AG, Zurich, Switzerland).

Costs of operation

The following provides an approximate list of costs to pursue the ramp-down experiment. Approximate costs are listed in Swiss Francs (CHF).

- **Static feld probes (12.5 k CHF):** Magnetic feld measurement device to map out the static feld to optimize shim iron placement.
- **Shim rails (5 k CHF):** A set of shim rails were purchased to simplify switching between feld strengths. With these, the 3 T shim rails could be removed after ramp-down of the magnet and replaced by the dedicated lower-feld shim sets, before bringing the system to its fnal 0.75 T main feld strength and vice versa.
- Ramping down and ramping up (12 work hours of Philips engineers per ramp-down cycle): After the initial rampdown test, each successive ramp-down took approximately six hours for one service engineer and additional six hours for ramping up.
- **Helium (~ 50 l per ramp-down):** Including the quench test in the initial ramp-down, a total of 150 l of Helium was purchased over the course of the three ramp-downs. As Helium is captured during ramp-down, boil-off is minimal and 50 l/ramp-down should be seen as an upper bound.
- **Transmit and receive coils (0 CHF):** No direct costs arose from the purchase of T/R hardware as the existing $3 T$ ¹³C hardware could be repurposed.
- **Indirect Costs:** Work hours by our scientists to make scanning possible cannot be directly quantifed and can hardly be disentangled from the research work that has been performed on the system.

Total: We estimate a total investment of approximately 100 k CHF over the whole project including own work hours.

Acquisition, data processing and reconstruction

The ramped-down system was used as a test platform for a variety of scan techniques and hardware tests. Apart from mapping static, transmit, and receive felds, gradient modulation transfer functions and audio response functions were obtained. In vivo data were acquired on healthy subjects in accordance with institutional and ethical guidelines and upon informed consent. A detailed list of scan parameters can be found in Table [1.](#page-4-0)

Reconstruction using the vendor's reconstructor was possible for most of the acquired data producing images for planning and initial data evaluation on the console. As no coil sensitivity prescans were available, coil-combined reconstructions showed artifacts from phase diferences between the receive elements. Hence, all raw data including scan metadata were exported using ReconFrame (GyroTools LLC, Winterthur, Switzerland) and reconstructed ofine with MRecon (GyroTools LLC, Winterthur, Switzerland) in MATLAB 2020b (Mathworks, Natick, Massachusetts, USA) and the Berkley Advanced Reconstruction Toolbox (BART; Berkley, California, USA) [\[25](#page-17-13)].

Unless otherwise stated, coil sensitivities *S* were determined using BART's ecalib [[26](#page-17-14)] implementation with the option 'm1' to produce only the frst sensitivity map and 'c0' to disable cropping of the maps. To homogenize image contrast, the magnitude of the Roemer coil-combined images [[27\]](#page-17-15) was blurred by a 20 px Gaussian flter to capture primarily the low-frequency contrast variations produced by the receive coils. This provided a fake quadrature body coil (QBC) image *Q*, which was used to regularize the coil combination on a per-pixel basis. Per-pixel, the coil-combined image *I* was calculated using [[28\]](#page-17-16)

$$
I = (S^{\dagger}S + \lambda Q^p)^{-1}S^{\dagger}i,
$$

where *i* is the vector of single-coil reconstructions, \dagger is the Hermitian conjugate. The power $p = 1$ and the regularization parameter $\lambda = 0.3$, where chosen manually to compromise between background signal suppression and contrast homogenization.

All spiral scans were deblurred using multi-frequency interpolation with 19 equidistant demodulation frequencies between -300 and $+150$ Hz (25 Hz steps) [\[29](#page-17-17)]. B0 maps for deblurring were either obtained from the data itself (B_0) Mapping and Dixon Multi-Acquisition scans) or prescribed (MRF scan).

Reconstruction code and raw data for all imaging experiments are available for download on gitlab [https://gitlab.](https://gitlab.ethz.ch/ibt-cmr-public/recon-0.75t-mri) [ethz.ch/ibt-cmr-public/recon-0.75t-mri.](https://gitlab.ethz.ch/ibt-cmr-public/recon-0.75t-mri)

System characterization

B₀ mapping

Main field (B_0) mapping was performed in a transversal and coronal, abdominal slice using a dual echo time, spoiled gradient-recalled echo (GRE) sequence with a spiral readout. Spiral deblurring and feld-map estimation was performed in a self-consistent, iterative fashion. In each iteration, the estimated B0 map was frst unwrapped in 2D and then blurred using an adaptive 2D Gaussian flter with decreasing flter strength from 60 to 10 px in 10 steps followed by 10, 8, 6, 4, 2, 1 px in subsequent iterations. Due to the strong signal attenuation in the transversal orientation, the fake QBC image was calculated using a reduced Gaussian flter of 10 px for better contrast homogenization.

DREAM B1+mapping

For B1+mapping, the DREAM sequence $[30]$ $[30]$ was used in both transversal and coronal angulations.

Receive feld mapping

Approximate receive feld maps were obtained from the dual-TE B0 mapping scan by reconstructing each coil image separately and dividing it by their sum of squares (SOS). In addition, a homogenized SOS image was obtained by dividing the SOS by its Gaussian-blurred ($\sigma = 20$ px) image.

Gradient modulation transfer function (GMTF)

The GMTF was measured using the chirp-based spectroscopy approach by Rahmer et al. [\[17](#page-17-9), [31](#page-17-19)]. In short, in parallel to a thin-slice multi-slice spectroscopy acquisition with in-plane phase-encoding, a chirp pulse was played out in through-slice direction with both positive and negative polarity. The acquisition volume was rotated to obtain measurements of the three cardinal axes. After fltering of low-signal regions, the phase-time course of each voxel was numerically diferentiated and ft by spherical harmonics. By dividing their spectrum by the input chirp spectrum, an estimate of the gradient modulation transfer function was obtained.

Relative microphone amplitude spectra (RMAS)

Concurrent to the GMTF measurement, audio measurements using a microphone (Sennheiser ME 66, Wedemark, Germany), placed into a waveguide of the MR Faraday cage, and a linear audio recorder (Tascam DR-100) were performed on both feld strength confgurations (0.75 T and 3 T). To calibrate microphone levels between ramp-down and product feld strength, a reference sound source at 1 kHz was used. Equivalently to the GMTF, the audio spectrum was divided by the chirp spectrum to obtain the relative microphone amplitude spectra (RMAS), a measure of the frequencydependent acoustic noise emissions allowing direct comparison of the system in its low-feld and product feld strength configuration.

In vivo showcases

Cardiac cine imaging

Cardiac cine scanning was performed at three field strengths, i.e., at 0.75 T, 1.5 T, and 3 T. While measurements at 0.75 T and 3 T were conducted on the same scanner, a separate 3 T Philips Achieva system (Philips Healthcare, Best, the Netherlands), which is permanently ramped down to 1.5 T, was used for the experiments at 1.5 T. Fully sampled 2D balanced steady-state free-precession (bSSFP) cine images were obtained in short-axis, long-axis, and four-chamber views using retrospective vectorcardiogram (VCG) gating.

In addition, the acquired k-space data of the 0.75 T shortaxis cine scans was retrospectively undersampled in k–t space using a variable density sampling mask and reconstructed using a cyclic vectorial total variation regularization [[32\]](#page-17-20) approach to demonstrate achievable image quality and acceleration factors with modern reconstruction techniques [[19\]](#page-17-21). We minimized the following optimization problem

$$
\min_{\rho} \|\widehat{M}\widehat{F}\widehat{S}\rho - d\|_2^2 + \lambda \|\left[\alpha \widehat{\nabla}_x \rho, \alpha \widehat{\nabla}_y \rho, \widehat{\nabla}_t \rho\right]\|_{2,1},
$$

where ρ represents the frames of the reconstructed cardiac cine, \hat{S} is a coil sensitivity operator, \hat{F} the Fourier transformation, \hat{M} the undersampling mask, and \hat{d} the acquired data. For regularization, $\alpha = 0.05$ was used to balance spatial variation $\hat{\nabla}_{xx} \rho$ against temporal variations $\hat{\nabla}_t \rho$. Optimization was performed numerically using the alternating directions method of multipliers (ADMM) [\[33](#page-17-22)]. Inner quadratic optimization was performed using three conjugate gradient steps. The regularization strength λ was chosen manually for best visual reconstruction performance.

Vectorcardiogram

VCGs were obtained at 0.75 T and 1.5 T from the ECGtriggering signal of a cardiac CINE acquisition to investigate the change in the magneto-hydrodynamic efect and potential implications for cardiac trigger accuracy [[34](#page-17-23)]. ECG electrodes were placed on one healthy volunteer and the VCG measurements were performed in direct succession at the 0.75 T and 1.5 T MRIs. As cardiac CINE scans were acquired in parallel, the heart was located in the isocenter position.

Dixon water/fat separation

Abdominal water/fat-separated imaging using multi-acquisition Cartesian and spiral Dixon approaches as well as a multi-echo Cartesian Dixon scan were performed in a coronal, abdominal slice [\[23,](#page-17-11) [31](#page-17-19)–[37\]](#page-17-24).

For the reconstruction of the spiral data, water/fat separation was performed in k-space. A seven-peak fat spectrum was employed to deblur the fat channel [[38,](#page-17-25) [39](#page-17-26)]. Echo and water/fat images were reconstructed and then fed into an iterative reconstruction pipeline to estimate a B0 map from the phase of forward simulated and acquired echo images. Equivalent to the B0 map scans, the B0 map was frst unwrapped and then blurred with the same adaptive Gaussian flter settings. Deblurring of all images was performed using multi-frequency interpolation [[29\]](#page-17-17) and the deblurred echo and water/fat images were then fed back into the B0 estimation code. For homogeneity correction, the frst echo was blurred and used as a fake QBC reference image. Cartesian images were reconstructed equivalently, omitting spiral reconstruction and deblurring steps. Fat spectrum deblurring was performed using the 7-peak model directly in image space.

Water/fat‑separated MR fngerprinting

Water–fat-separated MR Fingerprinting based on the work of Koolstra et al. [[40\]](#page-17-27) was performed by adjusting echo times to the in- and out-of-phase echo times of water/fat at 0.75 T [[24](#page-17-12)]. Each flip angle in a 500 timepoint constant-TR FISP-MRF [[41](#page-17-28)] sequence was acquired twice with 4.6 ms and 9.21 ms echo time in an interleaved fashion, resulting in a total of 1000 timepoints at a TR of 25 ms. A spiral acquisition with seven-fold undersampling was used. K-space data were separated into water and fat, demodulated for multi-frequency interpolation, and the fat channel was fat-spectrum deblurred using a seven-peak fat model [[38](#page-17-25), [39](#page-17-26)]. After transformation to image space, demodulation frequencies were recombined using multi-frequency interpolation [[29\]](#page-17-17) and a separately acquired B0 map (see above's spiral Dixon protocol for details). A dictionary resolving *T*1 ([2:2:100, 100:10:1000, 1000:20:2000, 2000:40:5000] ms) and *T*2 ([2:2:150, 150:10:500, 500:20:1000, 1000:40:2000] ms; $T2 < T1$) for B1 + values between 50 and 80% in 5% steps was obtained using the extended phase graph [[42](#page-17-29), [43](#page-17-30)]. The dictionary was compressed using a sequentially truncated higher-order SVD with singular value thresholding (relative threshold 10^{-4}) [[44\]](#page-17-31). B1 + was prescribed during the matching step $[45]$ $[45]$ $[45]$ by acquiring a co-registered B1+map using DREAM (see Table [1,](#page-4-0) $B1 + M$ apping DREAM (for MRF)) [\[30\]](#page-17-18). The dictionaries were linearly interpolated to intermediate $B1$ + values by interpolation of the $B1$ + basis functions and subsequent multiplication with the core tensor and followed by dictionary normalization.

Abdominal balanced SSFP and dual‑TE GRE

A transverse volume in the abdomen of 15 slices was acquired using Cartesian balanced SSFP in multi-2D acquisition mode.

To assess water/fat separation using a dual-TE Dixon scheme [\[35](#page-17-32)], RF- and gradient-spoiled GRE was performed. The scan duration of 36 s was split over two breath-holds. Equivalent to the three-point Dixon method, water and fat images as well as a B0 map were obtained iteratively based on estimation of a B0 map from the diference of forward simulation and the acquired dual-TE images. In each iteration, the estimated B0 map was frst unwrapped in 3D and then blurred in 3D using the same adaptive in-plane flter strength as in the 2D Dixon and B0 scans and a constant filter of 2 px in through-slice direction. The adaptive blurring was chosen to frst obtain good separation of water and fat by promoting a smooth B0 map before fne-tuning the map to account for local feld inhomogeneity. The updated B0 map was then used to refne the water/fat separation and the resulting forward simulation. Both forward simulation and the water/fat inference used a seven-peak fat-spectrum model [\[38](#page-17-25), [39](#page-17-26)].

Calf muscle spectroscopy

Point-resolved spectroscopy (PRESS) was performed in the calf muscle to compare proton metabolite spectra at 0.75 T, 1.5 T, and 3 T [\[20\]](#page-17-33). Spectra were acquired with 1 kHz, 2 kHz, and 4 kHz bandwidth at 0.75 T, 1.5 T, and 3 T, respectively, with 512 samples each, a TR of 2 s, in an 8 mL $(10 \times 20x40 \text{ mm}^3)$ voxel placed in the soleus muscle (Fig. [10](#page-15-0)a). Sixteen water-unsuppressed and 192 water-suppressed spectra, obtained with chemical-shift-based selective water suppression [\[46](#page-18-1)], were acquired. Power optimization was performed by observing the intensity of the water peak in a series of experiments with proportionally scaled excitation and echo pulses. Shimming was performed using a pencil beam and frst-order shim settings. Spectral data were noise decorrelated, coil channel weights were obtained from water-unsuppressed averages and coil combination was performed using singular value decomposition [[47](#page-18-2)]. Phase

correction of the spectra was performed on the water peak for unsuppressed water spectra and on the main triglyceride peak at 1.3 ppm for water-suppressed spectra. Water-unsuppressed spectra were used to obtain frequency correction information, which was transferred onto water-suppressed spectra by averaging the water frequency shift. To perform eddy current correction, the average phase of the waterunsuppressed signal was subtracted from all signals.

Equal parameters within a group of scans were merged into a single row. *AQ* Acquisition, *2D* single slice, *M-2D* Multiple 2D scans, *MS* Interleaved Multi-Slice, *ME* Multi-Echo, *MA* Multi-Acquisition, *Rect. FOV* Rectangular Fieldof-View, percentage in phase-encode direction, Δ*TE* Echotime increment for multi-echo/multi-acquisition scans, *TE1/TE2* first/second echo time, No. BH \times Dur. Number of breath-holds times duration per breath-hold, *Δ* Slice Gap

Results

System characterization

Figure [2](#page-9-0) shows in vivo B0 and $B1+$ mapping results alongside an anatomical reference. In a transverse slice, a linear feld gradient of approximately 150 Hz between anterior and posterior positions is observed. $B1+$ was found to be inhomogeneous with only 50% transmit efficiency in the center of the transverse slice and higher than nominal tip-angles in the left- and right-posterior regions. In the coronal slice, the B0 map appears more homogenous, however, also shows up to 100 Hz frequency diference between the spine and the liver, spleen, and kidneys. $B1 +$ mapping results are equally more homogenous reaching approximately 80% transmit efficiency in the left kidney and spleen compared with approximately 60% in the liver and right kidney.

Figure [3](#page-10-0) shows single-coil element reconstructions of the B0 scans for both the transverse and the coronal orientation. Due to the relatively small size of the receive coil elements (7.6 cm \times 18.3 cm), their penetration depth and lateral feld of view is limited, which leads to elevated noise in the homogenized sum-of-squares image (left column), especially in the lateral regions of the transverse slice (red circles).

In Fig. [4,](#page-10-1) the gradient modulation transfer function (GMTF) of the *x*-, *y*-, and *z*-axis gradients for the same system operated at 0.75 T (orange) and 3 T (gray) is depicted together with relative microphone amplitude spectra (RMAS) obtained from concurrent audio measurements. At

Fig. 2 In vivo static-feld and transmit-feld inhomogeneity overlayed on an anatomical reference for transverse and coronal angulations. The red lines indicate the anatomical cross-reference of the two slices. NB: For transmitting, a custom-built dual-loop Helmholtz transmitter (previously used for 13 C experiments at 3 T) is used at 0.75 T, not a body coil. In the transverse orientation, a linear gradi-

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ent of approximately 150 Hz in anterior–posterior direction is seen. $B1$ +shows a pronounced drop in transmit efficiency to 50% in the center of the body and only reaches prescribed fip angles in leftand right-posterior regions. In the coronal angulation, both B0 and B1+appear more homogenous. However, a maximum transmit efficiency of 80% is reached in the spleen and kidney

Fig. 3 (Left Column) Homogenized sum-of-squares (SOS) reconstruction. (Column 2 to 5) Single coil element reconstructions divided by their SOS for the coronal and transverse orientations. The red lines indicate the anatomical cross-reference of the two slices. Due to the relatively small size of the receive elements

 $(7.6 \text{ cm} \times 18.3 \text{ cm})$, the penetration depth is reduced leading to amplifcation of noise in the center of the body in the transverse orientation (red circles). In the coronal slice, the anterior elements (coil 3 and coil 4) contribute little signal

Fig. 4 Overlay of the frst-order gradient modulation transfer function (GMTF; line plots) on relative microphone amplitude spectra (RMAS; flled curves) of the system in product confguration (gray) and the lower-feld confguration (orange) for the three gradient axes (**a**, **b**, **c**). Mechanical resonances are marked by dashed vertical lines.

A pronounced reduction in mechanical resonances is observed in the GMTF (e.g., at approximately 1300 Hz on the z-axis), which is accompanied by a reduction in sound pressure. Figure adapted from Ref. [\[17\]](#page-17-9)

lower feld, mechanical resonances are reduced well below the uncertainty threshold of the thin-slice acquisition technique, while acoustic noise is signifcantly reduced over the whole frequency spectrum.

In vivo showcases

Figure [5](#page-11-0) shows a screenshot of the vendor's user interface of a cardiac cine scan in short-axis and four-chamber orientations, a lung bSSFP, and a cardiac T1 mapping scan. All images were reconstructed directly using the vendor's reconstructor. Due to the missing body coil data, contrast equilibration was not possible leading to pronounced hyperintensities, e.g., in the right ventricle of the short-axis view, when compared with the left ventricle. Apart from contrast inhomogeneity, the reconstruction quality is good and allows judging overall scan quality as well as general scan planning.

In Fig. [6a](#page-12-0), a comparison of cardiac short-axis cine scans in peak systole for the same volunteer at 0.75 T, 1.5 T, and 3 T is shown, with Fig. [6](#page-12-0)b showing zoom-ins of the heart.

Fig. 5 Screenshot of the vendor's user interface showing reconstructions performed on the system. Top left: cardiac cine four-chamber view. Top right: cardiac cine short-axis view. Bottom left: free-breathing balanced SSFP of the lung. Bottom right: cardiac T1 mapping

The myocardium, lumen and papillary muscles can be well delineated in all scans. Contrary to 1.5 T and 3 T, the 0.75 T cine image is devoid of banding artifacts as feld homogeneity in relation to the employed repetition time is improved. In Supporting Video 1, a comparison of retrospectively undersampled cardiac cine short-axis, four-chamber, and long-axis views are shown for acceleration factors of 2, 4, and 8 corresponding to breath-hold durations of 8, 4, and 2 s, respectively.

Figure [6c](#page-12-0) shows a comparison of vector cardiogram (VCG) traces at 0.75 and 1.5 T for the same volunteer with unchanged electrode placement. Due to the reduced feld, the magneto-hydrodynamic efect is greatly reduced at 0.75 T, while it can easily surpass the voltage of the R-wave on a 1.5 T system.

Figure [7](#page-13-0) compares water/fat separation results obtained with a three-point Dixon method with multi-acquisition Cartesian, multi-echo Cartesian, and multi-acquisition spiral readouts. The water/fat images as well as the iteratively determined B0 maps are in good agreement. Using spiral imaging, the breath-hold duration could be reduced from 27 s to 6.5 s for a single slice.

In Fig. [8,](#page-14-0) proton density, T1 and T2 maps from water/ fat-separated MR Fingerprinting are shown and compared to the water/fat separation from the previous multi-acquisition spiral approach. The water/fat-separated proton density maps are in general agreement with the multi-acquisition Dixon scan and show less pronounced noise.

In Fig. [9](#page-15-1)a, three exemplary slices of a single breathhold, balanced SSFP acquisition in the abdomen are given, showing good soft-tissue and vessel contrast as well as no banding artifacts. Figure [9](#page-15-1)b shows the same slices acquired with a RF- and gradient-spoiled GRE sequence with water/ fat in- and out-of-phase echo times. Using the complex dual-TE images, an iterative water/fat separation with B0 estimation could be performed with results shown on the right of Fig. [9b](#page-15-1). Water and fat could be successfully separated despite the strong feld gradient in anterior–posterior

Fig. 6 a Comparison of short-axis cardiac cine images of the same volunteer obtained at 0.75 T, 1.5 T, and 3 T with **b** zoom-ins of the heart. Red arrows indicate banding artifacts. Compared to 1.5 T and 3 T, the 0.75 T image is nearly banding artifact free due to the improved feld homogeneity relative to the scans repetition time. **c** Vector cardiogram (VCG) traces of the same volunteer with equal

electrode placement obtained at the 0.75 T system (top row) and the 1.5 T (bottom row). The arrow marks the magneto-hydrodynamic efect, which is of equal magnitude as the R-wave on a 1.5 T system for derivation 2 (red curve) and is nearly completely missing on the 0.75 T system. Hence, higher trigger accuracy is expected for cardiac scans on a lower-feld system

direction leading to approximately 150 Hz frequency shift. In comparison, the water/fat frequency shift at 0.75 T is approximately 108 Hz.

In Fig. [10,](#page-15-0) water-suppressed PRESS spectra of the same volunteer are shown for 0.75 T, 1.5 T, and 3 T next to a localizer image with the location of the voxel in the calf muscle. Despite the lower feld strength, the main metabolite peaks of triglyceride (TG-CH₂), creatine (CR-CH₃), and trimethylammonium (TMA) can be resolved even though absolute frequency shifts between the metabolites

are reduced. This is due to the proportional reduction in line width due to longer T_2^* times at lower field.

Discussion

In this work, we have demonstrated that a clinical 3 T system can be temporarily ramped down to 0.75 T and be used in conjunction with existing 13 C transmit and receive hardware to study and compare MRI and MRS on the very same system with limited time and cost overhead. Clinical

Fig. 7 Comparison of coronal water/fat separation using the threepoint Dixon method. **a** A multi-acquisition Cartesian, **b** a multi-echo Cartesian, and **c** multi-acquisition spiral mode is shown. Depicted are

the reconstructed water and fat maps as well as an iteratively determined B0 map

and research protocols could be readily executed at 0.75 T including image and spectroscopy data reconstruction. Additional offline reconstruction allowed us to study data in detail and to improve image and spectral quality, as coil sensitivities could be estimated. This enabled to take phase diferences between the receive elements into account as well as gave the possibility to perform contrast equalization [[48\]](#page-18-3) by emulating a low-resolution body coil image and using it for QBC-regularized reconstruction reducing surface coil fare and leading to similar contrasts and image homogeneity as on clinical systems.

Due to the 3 T factory shims bolted onto the cryostat and the limited available space on the insertable shim rails to compensate them, static-feld homogeneity was lower than expected and did not reach industry standards of \leq 3 ppm for the $50 \times 50 \times 45$ cm³ volume. Together with the limited automatic shimming accuracy in vivo, which was due to the lack of a body coil with homogenous transmit and receive felds, this led to reduced overall shim quality with, e.g.,

a frequency diference of approximately 200 Hz in anterior–posterior direction in abdominal scans. Additionally, the dual-loop Helmholtz transmitter created an inhomogeneous transmit feld which, in combination with the lowpower multi-nuclear amplifer (4 kW), limited achievable fip angles in vivo and further reduced contrast homogeneity. Quantitative protocols, such as MRF, were feasible, but required $B1$ + correction due to the custom-built transmit coil. In addition, the relatively small size of the four-channel receive coils led to a small combined feld of view, which in a transverse orientation in the abdomen led to pronounced noise enhancement in lateral areas half-way between anterior and posterior of the subjects.

Regarding the limited shim quality, it is noted that main feld shimming on more modern clinical scanners is exclusively performed using insertable shim rails (avoiding bolted shim irons altogether), thereby providing more fexibility to shim the same magnet at diferent feld strengths. Also, fullor reduced-rung 13C body resonators are becoming available

a) Multi-Acquisition Dixon

b) Water/Fat Separated FISP-MRF **Water Channel-**

Fig. 8 Water/fat-separated quantitative parameter maps obtained using MR Fingerprinting. **a** Multi-acquisition spiral Dixon scan for reference. **b** Water/fat-separated FISP-MRF images showing the reconstructed water and fat map in the frst column followed by *T*1 and *T*2 matching results. Approximate relaxation parameters are (mean±standard deviation over region of

interest; data from one volunteer only): liver $(T1: 491 \text{ ms} \pm 188 \text{ ms})$; *T*2: 77 ms \pm 97 ms), spleen (*T*1: 911 ms \pm 84 ms; *T*2: 91 ms \pm 24 ms), kidney (*T*1: 958 ms±206 ms; *T*2: 111 ms±50 ms), muscle (*T*1: 744 ms±94 ms; *T*2: 50 ms±13 ms), and fat (*T*1: 195 ms±104 ms; *T*2: 105 ms ± 81 ms)

to allow addressing the limitations of limited transmit-feld homogeneity. In this regard, ramp-down experiments are also considered valuable in assessing 13 C transmit and receive coil performance, which otherwise can only be performed with limitations on enriched and expensive ${}^{13}C$ samples. Using the same rationale, transmit/receive hardware for imaging other nuclei can be repurposed for proton imaging as well. Table [2](#page-16-3) shows a collection of typical MR nuclei and how they translate to equivalent proton feld strengths. E.g., ¹²⁹Xe coils could be used for imaging protons at 0.83 T. ¹³C corresponds to the lowest field strength officially supported by the vendor's multi-nuclear MR product.

Acquisition of the gradient modulation transfer function for both the product and the lower feld strength allowed to conclude that lower feld features improved gradient performance, especially since mechanical resonances are reduced due to a proportional reduction in Lorentz forces. This was also confrmed by a reduction in sound pressure, which benefts image quality and patient comfort.

The magneto-hydrodynamic (MHD) effect in the VCG signal was found to be reduced at 0.75 T compared with

1.5 T. This directly leads to more robust cardiac triggering and hence directly benefts cardiac MR applications.

Feasibility of water/fat separation with both multi-acquisition and multi-echo approaches using spiral and Cartesian acquisitions was demonstrated. In addition, we performed water/fat-separated MR fngerprinting, which showed comparable performance in separating water and fat and in addition delivered separate *T*1/*T*2 maps. Comparing liver *T*1/*T*2 values to 0.55 T, our values (mean \pm std. dev. over region of interest $T1$: 491 ms \pm 188 ms; $T2$: 77 ms \pm 97 ms) show longer T1 and *T*2 values as compared to 0.55 T (*T*1: 359.9 ms, *T*2: 45.2 ms) [\[49\]](#page-18-4). While a prolongation of *T*1 with increasing feld strength is expected, the increase in T2 is unexpected. However, given the relatively high standard deviation over the ROI in both *T*1 and *T*2 estimates, our quantitative fndings should be interpreted with caution. Comparing our spiral MRF results to similar work at 0.55 T [\[49\]](#page-18-4), the beneft of fat-spectrum deblurring should be noted, which lead in our work to well-delineated water/ fat boundaries and little blurring of the visceral fat despite a comparable spiral acquisition duration.

Fig. 9 a Single breath-hold, volumetric balanced SSFP of the abdomen, showing three exemplary slices. **b** Multi-slice dual-TE spoiled GRE with water/fat in- and out-of-phase echo times and iteratively reconstructed water/fat, proton density fat fraction, and off-resonance maps

Fig. 10 a Localizer image of the calf muscle with overlayed location of the 8 mL PRESS voxel in the soleus muscle. **b** Comparison of 0.75 T, 1.5 T, and 3 T water-suppressed PRESS spectra showing

triglyceride (TG-CH₂), creatine (CR-CH₃), and trimethylammonium (TMA). Figure adapted from Ref. [[20](#page-17-33)]

The in vivo calf muscle experiments show that spectroscopy is feasible and utilizable at lower felds and despite the proportional lowering of line separation, peak separation can be maintained as line width is decreasing due to favorable scaling of T2*. Similar to the imaging experiments, an improvement of shim and transmit-feld homogeneity would be beneficial for further reducing line width and hence further improving peak separation and, thus, utility of spectroscopy at 0.75 T [[20\]](#page-17-33).

We claimed the system to be low cost compared to the purchase and installation of a new lower-feld MRI system. Comparing our costs of approximately 100 k CHF (including investments and initial pulse-programming/development costs) to the purchase of a new lower-feld system—likely being in the 1 M CHF range without siting related costs ramp-down of an existing system, especially if transmit/ receive hardware is already available, is a viable, costeffective option.

Table 2 List of nuclei, their reduced gyromagnetic ratio, and their equivalent feld strength, if proton imaging was performed by ramping down a 3 T MRI system to match Larmor frequencies

*Listed as supported by the vendor for multi-nuclear MR at 3 T

The overall experience of ramping down a high-feld MRI was experimental as we faced many issues that have been solved on clinical systems with methods in place that completely shield the high-level user from their existence, e.g., automatic receive gain adjustment to prevent ADC overfow. These issues were reexperienced as measures normally active on a clinical system were not operational due to the lack of tuned body and pickup coils. Despite these challenges, it was a great learning experience especially with regards to the fundamental mechanisms of MRI and the tremendous engineering effort that has been put into state-ofthe-art machines. A body coil tuned to 0.75 T would have been of enormous help for both transmit homogeneity and making preparation phases work. Purchasing a volume transmitter coil to improve transmit homogeneity was also considered, however, as our system has a narrow bore (60 cm), only limited space remains for additional transmit hardware. In hindsight, it would have also been helpful to carefully study other experimental systems that also operate without a body coil, such as the Philips 7 T Achieva, to transfer preparation phases and pulse-programming features. Additionally, there are many user-hidden variables that optimize the scanner for a particular feld strength. It is important to spend exclusive time to ensure these are set properly for low-feld operation.

Conclusion

Clinical 3 T systems can be temporarily ramped down and operated at 0.75 T by exploiting ¹³C transmit/receive hardware allowing to run and reconstruct clinical and research protocols with minimal software modifcations. This approach allows to study and compare MRI and MRS

application on the very same system with limited time and cost overhead.

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Code and data availability The full MATLAB code, raw datasets as well as pre-processed MRI data are openly available for download from:<https://gitlab.ethz.ch/ibt-cmr-public/recon-0.75t-mri>.

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References

- 1. Chandarana H, Bagga B, Huang C, Dane B, Petrocelli R, Bruno M et al (2021) Diagnostic abdominal MR imaging on a prototype low-feld 0.55 T scanner operating at two diferent gradient strengths. Abdom. Radiol. 46:5772–5780. [https://doi.org/](https://doi.org/10.1007/s00261-021-03234-1) [10.1007/s00261-021-03234-1](https://doi.org/10.1007/s00261-021-03234-1)
- 2. Campbell-Washburn AE, Ramasawmy R, Restivo MC, Bhattacharya I, Basar B, Herzka DA et al (2019) Opportunities in interventional and diagnostic imaging by using high-performance low-feld-strength MRI. Radiology 293:384–393. [https://](https://doi.org/10.1148/radiol.2019190452) doi.org/10.1148/radiol.2019190452
- 3. Restivo MC, Ramasawmy R, Bandettini WP, Herzka DA, Campbell-Washburn AE (2020) Efficient spiral in-out and EPI balanced steady-state free precession cine imaging using a highperformance 0.55T MRI. Magn Reson Med 84:2364–2375. <https://doi.org/10.1002/mrm.28278>
- 4. Javed A, Ramasawmy R, O'Brien K, Mancini C, Su P, Majeed W et al (2022) Self-gated 3D stack-of-spirals UTE pulmonary imaging at 0.55T. Magn Reson Med 87:1784–1798. [https://doi.](https://doi.org/10.1002/mrm.29079) [org/10.1002/mrm.29079](https://doi.org/10.1002/mrm.29079)
- 5. Qin C, Murali S, Lee E, Supramaniam V, Hausenloy DJ, Obungoloch J et al (2022) Sustainable low-feld cardiovascular magnetic resonance in changing healthcare systems. Eur Hear J Cardiovasc Imaging 23:e246–e260. [https://doi.org/10.1093/ehjci/](https://doi.org/10.1093/ehjci/jeab286) [jeab286](https://doi.org/10.1093/ehjci/jeab286)
- 6. Varghese J, Jin N, Giese D, Chen C, Liu Y, Pan Y, Nair N, Mahmoud TS, Mahmood K, Matthew ST, Rizwan A, Yuchi H, Orlando PS (2023) Building a comprehensive cardiovascular magnetic resonance exam on a commercial 0.55 T system: A pictorial essay on potential applications Front. Cardiovasc. Med. <https://doi.org/10.3389/fcvm.2023.1120982>
- 7. Juliet VJ, Craft CD, Yingmin C, Ning L, Kelvin J, Rizwan C, Ahmad OP, Simonetti (2020) Assessment of cardiac function blood flow and myocardial tissue relaxation parameters at 0.35 T. NMR Biomed 33(7).<https://doi.org/10.1002/nbm.4317>
- 8. Wald LL, McDaniel PC, Witzel T, Stockmann JP, Cooley CZ (2020) Low-cost and portable MRI. J Magn Reson Imaging 52:686–696. <https://doi.org/10.1002/jmri.26942>
- 9. Hoult DI, Chen C-N, Sank VJ (1986) The feld dependence of NMR imaging. II Arguments concerning an optimal feld strength. Magn Reson Med 3:730–746. [https://doi.org/10.1002/](https://doi.org/10.1002/mrm.1910030509) [mrm.1910030509](https://doi.org/10.1002/mrm.1910030509)
- 10. Marques JP, Simonis FFJJ, Webb AG (2019) Low-feld MRI: An MR physics perspective. J Magn Reson Imaging 49:1528–1542. <https://doi.org/10.1002/jmri.26637>
- 11. Togawa T, Okai O, Oshima M (1967) Observation of blood fow E.M.F. in externally applied strong magnetic feld by surface electrodes. Med Biol Eng 5:169–170. [https://doi.org/10.1007/](https://doi.org/10.1007/BF02474505) [BF02474505](https://doi.org/10.1007/BF02474505)
- 12. Moelker A, Wielopolski PA, Pattynama PMT (2003) Relationship between magnetic feld strength and magnetic-resonance-related acoustic noise levels. MAGMA Magn Reson Mater Physics, Biol Med 16:52–55. [https://doi.org/10.1007/](https://doi.org/10.1007/s10334-003-0005-9) [s10334-003-0005-9](https://doi.org/10.1007/s10334-003-0005-9)
- 13. Sarracanie M, Salameh N (2020) Low-feld MRI: how low can we go? A fresh view on an old debate. Front Phys 8:1–14. <https://doi.org/10.3389/fphy.2020.00172>
- 14. Simonetti OP, Ahmad R (2017) Low-feld cardiac magnetic resonance imaging. Circ Cardiovasc Imaging 10:1–7. [https://](https://doi.org/10.1161/CIRCIMAGING.117.005446) doi.org/10.1161/CIRCIMAGING.117.005446
- 15. Rashid S, Han F, Gao Y, Sung K, Cao M, Yang Y et al (2018) Cardiac balanced steady-state free precession MRI at 0.35 T: a comparison study with 1.5 T. Quant Imaging Med Surg 8:627– 636. <https://doi.org/10.21037/qims.2018.08.09>
- 16. Guenthner C, Peereboom S, Dillinger H, McGrath C, Albannay MM, Vishnevskiy V, et al (2022) Ramp-Down of a High-Performance 3T MRI: Experiences at 0.75T. In: ISMRM Workshop on Low Field MRI. Virtual Conference
- 17. Dillinger H, Kozerke S, Guenthner C (2022) Direct comparison of gradient Fidelity and acoustic noise of the same MRI system at 3 T and 0.75 T. Magn Reson Med 88:1937–1947. [https://doi.](https://doi.org/10.1002/mrm.29312) [org/10.1002/mrm.29312](https://doi.org/10.1002/mrm.29312)
- 18. Dillinger H, Kozerke S, Guenthner C (2021) Direct comparison of gradient modulation transfer functions and acoustic noise spectra of the same MRI at high- (3T) and lower-feld (075T). Intl Soc Mag Reson Med 29:2347
- 19. Vishnevskiy V, Guenthner C, McGrath C, Kozerke S (2021) Highly accelerated cardiac cine imaging on a lower-feld 0.75T MRI. Intl Soc Mag Reson Med 29:3434
- 20. Peereboom SM, Guenthner C, Albannay MM, Kozerke S (2023) Preliminary experience of cardiac proton spectroscopy at 0.75 T. NMR Biomed.<https://doi.org/10.1002/nbm.4892>
- 21. McGrath C, Albannay M, Jafray A, Guenthner C, Kozerke S (2021) First-principle image SNR synthesis depending on feld strength. Intl Soc Mag Reson Med 29:1460
- 22. Peper E, Dillinger H, McGrath C, Guenthner C, Kozerke S. Cine Flow Measurements using Phase-Contrast bSSFP at 0.75 Tesla. In: Proc. Int. Soc. Mag. Reson. Med. 30. London, UK; 2022. p. 0092.
- 23. Guenthner C, Dillinger H, Boernert P, Kozerke S (2021) Threepoint dixon abdominal water/fat separation using a lower-feld 075T MRI. Intl Soc Mag Reson Med 29:1283
- 24. Guenthner C, Koken P, Boernert P, Kozerke S (2021) Abdominal water/fat separated MR fngerprinting on a lower-feld 075T MRI. Intl Soc Mag Reson Med 29:499
- 25. Blumenthal M, Holme C, Roelofs V, Rosenzweig S, Schaten P, Scholand N, Tamir J, Wang X, Uecker M (2022) mrirecon/bart. Zenodo.<https://doi.org/10.5281/zenodo.592960>
- 26. Uecker M, Lai P, Murphy MJ, Virtue P, Elad M, Pauly JM et al (2014) ESPIRiT - An eigenvalue approach to autocalibrating

parallel MRI: Where SENSE meets GRAPPA. Magn Reson Med 71:990–1001. <https://doi.org/10.1002/mrm.24751>

- 27. Roemer PB, Edelstein WA, Hayes CE, Souza SP, Mueller OM (1990) The NMR phased array. Magnet Resonance Med 16(2):192–225.<https://doi.org/10.1002/mrm.1910160203>
- 28. Fuderer M, van den Brink J, Jurrissen M (2004) SENSE reconstruction using feed forward regularization. In: Proc. Intl. Soc. Mag. Res. Med. vol 11. p. 2130
- 29. Man L-C, Pauly JM, Macovski A (1997) Multifrequency interpolation for fast off-resonance correction. Magn Reson Med 37:785–792.<https://doi.org/10.1002/mrm.1910370523>
- 30. Nehrke K, Börnert P (2012) DREAM-a novel approach for robust, ultrafast, multislice B1 mapping. Magn Reson Med 68:1517– 1526. <https://doi.org/10.1002/mrm.24158>
- 31. Rahmer J, Mazurkewitz P, Börnert P, Nielsen T (2019) Rapid acquisition of the 3D MRI gradient impulse response function using a simple phantom measurement. Magn Reson Med 82:2146–2159.<https://doi.org/10.1002/mrm.27902>
- 32. Blomgren P, Chan TF (1998) Color TV: Total variation methods for restoration of vector-valued images. IEEE Trans Image Process 7:304–309.<https://doi.org/10.1109/83.661180>
- 33. Boyd S (2010) Distributed optimization and statistical learning via the alternating direction method of multipliers. found. trends® mach. Learn 3:1–122.<https://doi.org/10.1561/2200000016>
- 34. Fischer SE, Wickline SA, Lorenz CH (1999) Novel real-time R-wave detection algorithm based on the vectorcardiogram for accurate gated magnetic resonance acquisitions. Magn Reson Med 42:361–370. [https://doi.org/10.1002/\(SICI\)1522-2594\(199908\)](https://doi.org/10.1002/(SICI)1522-2594(199908)42:2%3c361::AID-MRM18%3e3.0.CO;2-9) [42:2%3c361::AID-MRM18%3e3.0.CO;2-9](https://doi.org/10.1002/(SICI)1522-2594(199908)42:2%3c361::AID-MRM18%3e3.0.CO;2-9)
- 35. Dixon WT (1984) Simple proton spectroscopic imaging. Radiology 153:189–194. [https://doi.org/10.1148/radiology.153.1.60892](https://doi.org/10.1148/radiology.153.1.6089263) [63](https://doi.org/10.1148/radiology.153.1.6089263)
- 36. Glover GH (1991) Multipoint dixon technique for water and fat proton and susceptibility imaging. J Magn Reson Imaging 1:521– 530.<https://doi.org/10.1002/jmri.1880010504>
- 37. Glover GH, Schneider E (1991) Three-point dixon technique for true water/fat decomposition withB0 inhomogeneity correction. Magn Reson Med 18:371–383. [https://doi.org/10.1002/mrm.](https://doi.org/10.1002/mrm.1910180211) [1910180211](https://doi.org/10.1002/mrm.1910180211)
- 38. Brodsky EK, Holmes JH, Yu H, Reeder SB (2008) Generalized K-space decomposition with chemical shift correction for non-Cartesian water-fat imaging. Magn Reson Med 59:1151–1164. <https://doi.org/10.1002/mrm.21580>
- 39. Yu H, Shimakawa A, McKenzie CA, Brodsky E, Brittain JH, Reeder SB (2008) Multiecho water-fat separation and simultaneous R 2* estimation with multifrequency fat spectrum modeling. Magn Reson Med 60:1122–1134. [https://doi.org/10.1002/mrm.](https://doi.org/10.1002/mrm.21737) [21737](https://doi.org/10.1002/mrm.21737)
- 40. Koolstra K, Webb AG, Veeger TTJ, Kan HE, Koken P, Börnert P (2020) Water–fat separation in spiral magnetic resonance fngerprinting for high temporal resolution tissue relaxation time quantifcation in muscle. Magn Reson Med 84:646–662. [https://](https://doi.org/10.1002/mrm.28143) doi.org/10.1002/mrm.28143
- 41. Jiang Y, Ma D, Seiberlich N, Gulani V, Griswold M (2015) MR fngerprinting using fast imaging with steady state precession (FISP) with spiral readout. Magn Reson Med 74:1621–1631. <https://doi.org/10.1002/mrm.25559>
- 42. Weigel M (2015) Extended phase graphs: dephasing, RF pulses, and echoes—pure and simple. J Magn Reson Imaging 41:266– 295.<https://doi.org/10.1002/jmri.24619>
- 43. Guenthner C, Amthor T, Doneva M, Kozerke S (2021) A unifying view on extended phase graphs and Bloch simulations for quantitative MRI. Sci Rep 11:21289. [https://doi.org/10.1038/](https://doi.org/10.1038/s41598-021-00233-6) [s41598-021-00233-6](https://doi.org/10.1038/s41598-021-00233-6)
- 44. Vannieuwenhoven N, Vandebril R, Meerbergen K (2012) A new truncation strategy for the higher-order singular value

decomposition. SIAM J Sci Comput 34:A1027–A1052. [https://](https://doi.org/10.1137/110836067) doi.org/10.1137/110836067

- 45. Ma D, Coppo S, Chen Y, McGivney DF, Jiang Y, Pahwa S et al (2017) Slice profle and B1 corrections in 2D magnetic resonance fngerprinting. Magn Reson Med 78:1781–1789. [https://doi.org/](https://doi.org/10.1002/mrm.26580) [10.1002/mrm.26580](https://doi.org/10.1002/mrm.26580)
- 46. Haase A, Frahm J, Hanicke W, Matthaei D (1985) 1H NMR chemical shift selective (CHESS) imaging. Phys Med Biol 30:341–344. <https://doi.org/10.1088/0031-9155/30/4/008>
- 47. Weiss K, Martini N, Boesiger P, Kozerke S (2013) Cardiac proton spectroscopy using large coil arrays. NMR Biomed 26:276–284. <https://doi.org/10.1002/nbm.2845>
- 48. Harvey PR, Rozijn TH, Ijperen GH, Prins WM, Mens WRM, Benschop FJM (2015) Image intensity correction for magnetic resonance imaging. US 9036884 B2, 2015.
- 49. Liu Y, Hamilton J, Jiang Y, Seiberlich N (2022) Assessment of MRF for simultaneous T1 and T2 quantifcation and water–fat separation in the liver at 0.55 T. Magn Reson Mater Physics, Biol Med.<https://doi.org/10.1007/s10334-022-01057-9>

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