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# **Universal Inequalities for Lower Order Eigenvalues of Self-Adjoint Operators and the Poly-Laplacian**

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**Abstract** In this paper, we first establish an abstract inequality for lower order eigenvalues of a self-adjoint operator on a Hilbert space which generalizes and extends the recent results of Cheng et al. (Calc. Var. Partial Differential Equations, **38**, 409–416 (2010)). Then, making use of it, we obtain some universal inequalities for lower order eigenvalues of the biharmonic operator on manifolds admitting some special functions. Moreover, we derive a universal inequality for lower order eigenvalues of the poly-Laplacian with any order on the Euclidean space.

**Keywords** Eigenvalue, self-adjoint operator, biharmonic operator, poly-Laplacian, Riemannian manifold

**MR(2010) Subject Classification** 35P15, 58C40

## **1 Introduction**

Let  $\Omega$  be a bounded domain in an *n*-dimensional Euclidean space  $\mathbb{R}^n$ . The Dirichlet eigenvalue problem of the poly-Laplacian with any order is described by

$$
\begin{cases}\n(-\Delta)^{l}u = \lambda u, & \text{in } \Omega, \\
u|_{\partial\Omega} = \frac{\partial u}{\partial \nu}\Big|_{\partial\Omega} = \dots = \frac{\partial^{l-1}u}{\partial \nu^{l-1}}\Big|_{\partial\Omega} = 0,\n\end{cases}
$$
\n(1.1)

where  $\Delta$  is the Laplacian and  $\nu$  denotes the outward unit normal vector field of  $\partial\Omega$ . Problem  $(1.1)$  is called the Dirichlet Laplacian problem when  $l = 1$  and the clamped plate problem when  $l = 2$ .

In this paper, we are concerned about inequalities for lower order eigenvalues of self-adjoint operators and problem (1.1). To begin with, we give a brief review of related results. For the Dirichlet Laplacian problem, Payne et al. [22] proved that its lower order eigenvalues satisfy

$$
\lambda_2 + \lambda_3 \le 6\lambda_1 \tag{1.2}
$$

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for  $\Omega \subset \mathbb{R}^2$ . This led to the famous Payne–Pólya–Weinberger conjecture for  $\Omega \subset \mathbb{R}^n$ . In 1993, Ashbaugh and Benguria [2] established the following universal inequality

$$
\frac{1}{n}\sum_{i=1}^{n}\lambda_{i+1}\leq \left(1+\frac{4}{n}\right)\lambda_1\tag{1.3}
$$

for  $\Omega \subset \mathbb{R}^n$ . For more references on the solution of this conjecture, we refer the readers to [3, 4, 10, 15, 21, 24]. In 2008, Sun et al. [25] further derived some universal inequalities for lower order eigenvalues of the Dirichlet Laplacian problem on bounded domains in a complex projective space and a unit sphere. Chen and Cheng [6] proved (1.3) still holds when  $\Omega$  is a bounded domain in a complete Riemannian manifold isometrically minimally immersed in  $\mathbb{R}^n$ .

 $\Gamma_i$  is usually used to denote the *i*-th eigenvalue of the clamped plate problem. In 1998, Ashbaugh [1] announced the following interesting inequalities without proofs

$$
\sum_{i=1}^{n} \left(\Gamma_{i+1}^{\frac{1}{2}} - \Gamma_1^{\frac{1}{2}}\right) \le 4\Gamma_1^{\frac{1}{2}} \tag{1.4}
$$

and

$$
\sum_{i=1}^{n} (\Gamma_{i+1} - \Gamma_1) \le 24\Gamma_1.
$$
 (1.5)

In 2010, for a bounded domain  $\Omega$  in an n-dimensional complete Riemannian manifold M, Cheng et al. [8] proved

$$
\sum_{i=1}^{n} \left(\Gamma_{i+1} - \Gamma_1\right)^{\frac{1}{2}} \le \left[ (2n+4)\Gamma_1^{\frac{1}{2}} + n^2 H_0^2 \right]^{\frac{1}{2}} \left( 4\Gamma_1^{\frac{1}{2}} + n^2 H_0^2 \right)^{\frac{1}{2}},\tag{1.6}
$$

where  $H_0$  is a nonnegative constant which only depends on M and  $\Omega$ . When M is an ndimensional complete minimal submanifold in a Euclidean space, (1.6) implies

$$
\sum_{i=1}^{n} (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}} \le [8(n+2)\Gamma_1]^{\frac{1}{2}}.
$$
 (1.7)

The proofs of (1.4) and (1.5) were given by Cheng et al. [9]. In fact, they considered problem (1.1) and proved

$$
\sum_{i=1}^{n} (\lambda_{i+1}^{\frac{1}{l}} - \lambda_1^{\frac{1}{l}})^{l-1} \le (2l)^{l-1} \lambda_1^{\frac{l-1}{l}} \tag{1.8}
$$

for  $l \geq 2$ , and

$$
\sum_{i=1}^{n} (\lambda_{i+1} - \lambda_1) \le 4l(2l - 1)\lambda_1.
$$
 (1.9)

It is easy to find that  $(1.8)$  and  $(1.9)$  respectively become  $(1.4)$  and  $(1.5)$  when  $l = 2$ . Moreover,  $(1.9)$  covers  $(1.3)$  when  $l = 1$ . In 2011, Jost et al. [18] derived the inequality

$$
\sum_{i=2}^{n+1} \lambda_i + \sum_{i=1}^{n-1} \frac{2(l-1)i}{2l+i-1} (\lambda_{n+1-i} - \lambda_1) \le (n+4l(2l-1))\lambda_1,\tag{1.10}
$$

which covers (1.3) when  $l = 1$  and improves (1.5) when  $l = 2$ .

It is natural to consider whether these inequalities can be deduced to self-adjoint operators on Hilbert spaces. Harrell and Davies [12] first realized that some results (e.g. the PPW inequality) of Payne et al. [22] for higher order eigenvalues of the Laplcian rely on some facts involving auxiliary operators and their commutators. Some inequalities for higher order eigenvalues of the Laplacian, biharmonic operator and the poly-Laplacian have been deduced to self-adjoint operators on Hilbert spaces (see [5, 13, 14, 16, 17, 19, 23]). To the authors' knowledge, fewer inequalities for lower order eigenvalues of self-adjoint operators on Hilbert spaces have been obtained by using purely algebraic arguments.

The structure of this paper is as follows. In Section 2, we establish an abstract inequality (2.2) for lower order eigenvalues of self-adjoint operators on Hilbert spaces, which extends the recent work of Cheng et al. [8]. The latter sections are devoted to deriving some explicit inequalities for lower order eigenvalues of problem (1.1) on some Riemannian manifolds. In fact, Cheng et al. [8] considered Riemannian manifolds isometrically immersed in a Euclidean space. We study eignvalues of biharmonic operator on manifolds admitting some special functions in Section 3. The first kind is Riemannian manifolds admitting spherical eigenmaps. As we know, any compact homogeneous Riemannian manifold admits eigenmaps for the first eigenvalue of the Laplacian (see [20]). And another kind is Riemannian manifolds admitting some functions  $f_r: M \to \mathbb{R}$  such that

$$
\begin{cases} \langle \nabla f_r, \nabla f_s \rangle = \delta_{rs}, \\ \Delta f_r = 0. \end{cases}
$$
 (1.11)

Product manifolds of Euclidean spaces with any complete manifolds satisfy this condition (see [11, 26]). Utilizing Theorem 2.1, we derive some universal inequalities for lower order eigenvalues of the biharmonic operator on these manifolds in Theorem 3.1. In Section 4, a universal inequality for the poly-Laplacian with any order on a bounded domain in  $\mathbb{R}^n$  is obtained, which covers  $(1.7)$  when  $l = 2$ .

#### **2 Self-Adjoint Operators on Hilbert Spaces**

In this section, we establish an abstract inequality which relates lower order eigenvalues of self-adjoint operators to two collections of auxiliary operators and their commutators.

**Theorem 2.1** Let H be a complex Hilbert space with a given inner product  $\langle \cdot, \cdot \rangle$  and cor*responding norm*  $\|\cdot\|$ . Let  $A : \mathcal{D} \subset \mathcal{H} \to \mathcal{H}$  be a self-adjoint operator defined on a dense *domain* D which is semibounded beblow and has a discrete spectrum  $\mu_1 \leq \mu_2 \leq \cdots$ . Let  ${T_i : \mathcal{D} \to \mathcal{H}}_{i=1}^n$  *be a collection of skew-symmetric operators and*  ${B_i : A(\mathcal{D}) \to \mathcal{H}}_{i=1}^n$  *be a collection of symmetric operators which leave*  $\mathcal D$  *invariant. Denote by*  $\{u_i\}_{i=1}^{\infty}$  *the normalized eigenvectors corresponding to the* i*-th eigenvalues* μ<sup>i</sup> *of* A*. This family of eigenvectors are*  $further\ assumed\ to\ be\ an\ orthonormal\ basis\ for\ \mathcal{H}.$  If the operators  $\{B_i\}_{i=1}^n$  satisfy

$$
\langle B_i u_1, u_{j+1} \rangle, \quad \text{for } 1 \le j < i \le n,\tag{2.1}
$$

*we have*

$$
\sum_{i=1}^{n} (\mu_{i+1} - \mu_1)^{\frac{1}{2}} \langle [T_i, B_i]u_1, u_1 \rangle \le 2 \left\{ \sum_{i=1}^{n} \langle [A, B_i]u_1, B_i u_1 \rangle \sum_{i=1}^{n} ||T_i u_1||^2 \right\}^{\frac{1}{2}},\tag{2.2}
$$

*where*  $[A, B] = AB - BA$  *is called the commutator of operators* A *and* B.

*Proof* We consider the vectors  $\phi^i$  given by

$$
\phi^i = B_i u_1 - a^i u_1,\tag{2.3}
$$

where  $a^i = \langle B_i u_1, u_1 \rangle$ . Then, according to (2.1), it is easy to check that

$$
\langle \phi^i, u_{j+1} \rangle = 0, \quad \text{for } 0 \le j < i \le n. \tag{2.4}
$$

Moreover, (2.4) yields

$$
\langle \phi^i, B_i u_1 \rangle = ||\phi^i||^2. \tag{2.5}
$$

Since (2.4) holds, we can take  $\phi^i$  as a trial vector in the Rayleigh–Ritz ratio and get

$$
\mu_{i+1} \le \frac{\langle A\phi^i, \phi^i \rangle}{\langle \phi^i, \phi^i \rangle}.
$$
\n(2.6)

It follows from  $(2.4)$  and  $(2.5)$  that

$$
\langle A\phi^i, \phi^i \rangle = \langle [A, B_i]u_1, \phi^i \rangle + \langle B_i A u_1, \phi^i \rangle = \langle [A, B_i]u_1, \phi^i \rangle + \mu_1 \|\phi^i\|^2. \tag{2.7}
$$

Substituting  $(2.7)$  into  $(2.6)$ , we obtain

$$
(\mu_{i+1} - \mu_1) \|\phi^i\|^2 \le \langle [A, B_i]u_1, \phi^i \rangle. \tag{2.8}
$$

As we know,  $\langle \cdot, \cdot \rangle$  is taken to be linear in its first argument and conjugate linear in its second argument. Since

$$
\langle [A, B_i]u_1, u_1 \rangle = \langle B_i u_1, A u_1 \rangle - \langle B_i A u_1, u_1 \rangle = 0,
$$

we have

$$
\langle [A, B_i]u_1, \phi^i \rangle = \langle [A, B_i]u_1, B_iu_1 \rangle. \tag{2.9}
$$

Hence, substituting (2.9) into (2.8), we get

$$
(\mu_{i+1} - \mu_1) \|\phi^i\|^2 \le \langle [A, B_i]u_1, B_i u_1 \rangle. \tag{2.10}
$$

Since

$$
\overline{a^i}\langle T_i u_1, u_1\rangle = -\langle B_i u_1, u_1\rangle \overline{\langle T_i u_1, u_1\rangle} = -a^i \overline{\langle T_i u_1, u_1\rangle},
$$

we obtain

$$
\operatorname{Re}\overline{a^i}\langle T_i u_1, u_1\rangle = 0. \tag{2.11}
$$

Moreover, we have

$$
2\mathrm{Re}\langle T_i B_i u_1, u_1\rangle = \langle T_i B_i u_1, u_1\rangle + \langle u_1, T_i B_i u_1\rangle = \langle [T_i, B_i] u_1, u_1\rangle.
$$
 (2.12)

Hence, taking the real parts in both sides of the following equality

$$
-2\langle T_i u_1, \phi^i \rangle = 2\langle u_1, T_i B_i u_1 \rangle + 2\overline{a^i} \langle T_i u_1, u_1 \rangle,
$$

and utilizing  $(2.11)$  and  $(2.12)$ , we obtain

$$
\langle [T_i, B_i]u_1, u_1 \rangle = -2\text{Re}\langle \phi^i, T_i u_1 \rangle. \tag{2.13}
$$

Multiplying both sides of (2.13) by  $(\mu_{i+1} - \mu_1)^{\frac{1}{2}}$  and using (2.10), we deduce

$$
(\mu_{i+1} - \mu_1)^{\frac{1}{2}} \langle [T_i, B_i]u_1, u_1 \rangle = -2(\mu_{i+1} - \mu_1)^{\frac{1}{2}} \text{Re} \langle \phi^i, T_i u_1 \rangle
$$
  
\n
$$
\leq \delta(\mu_{i+1} - \mu_1) ||\phi^i||^2 + \frac{1}{\delta} ||T_i u_1||^2
$$
  
\n
$$
\leq \delta \langle [A, B_i]u_1, B_i u_1 \rangle + \frac{1}{\delta} ||T_i u_1||^2,
$$
\n(2.14)

where  $\delta$  is a positive constant. Taking sum on i from 1 to n in (2.14), we have

$$
\delta^2 \sum_{i=1}^n \langle [A, B_i]u_1, B_i u_1 \rangle - \delta \sum_{i=1}^n (\mu_{i+1} - \mu_1)^{\frac{1}{2}} \langle [T_i, B_i]u_1, u_1 \rangle + \sum_{i=1}^n \|T_i u_1\|^2 \ge 0. \tag{2.15}
$$

The left-hand side of  $(2.15)$  is a quadratic polynomial of  $\delta$ . From  $(2.8)$  and  $(2.9)$ , we know that  $\langle [A, B_i]u_1, B_iu_1 \rangle \geq 0$  for  $i = 1, \ldots, n$ . Therefore, we know that its discriminant must be nonpositive. This yields  $(2.2)$ .

# **3 The Biharmonic Operator on Riemannian Manifolds Admitting Some Special Functions**

In this section, we obtain some universal inequalities for lower order eigenvalues of the biharmonic operator on manifolds admitting some special functions.

**Theorem 3.1** *Let* M *be an n-dimensional complete Riemannian manifold. Denote by*  $\Gamma_i$  *the i*-th eigenvalue of the clamped plate problem on a bounded domain  $\Omega$  of M.

(i) *Suppose that* M *admits a spherical eigenmap*  $\varphi$ *. Namely, the components*  $\varphi_1, \ldots, \varphi_{m+1}$ *of map*  $\varphi : M \to \mathbb{S}^m(1)$  *are all eigenfuctions corresponding to the same eigenvalue*  $\lambda$  *of the Laplacian on* M, where  $\mathbb{S}^m(1)$  *is an m-dimensional unite sphere. Then we have* 

$$
\sum_{i=1}^{n} (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}} \le n \left[ \left( \lambda + 6\Gamma_1^{\frac{1}{2}} \right) \left( \lambda + 4\Gamma_1^{\frac{1}{2}} \right) \right]^{\frac{1}{2}}
$$
\n(3.1)

*and*

$$
\sum_{i=1}^{n} \left[ (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}} - 5\Gamma_1^{\frac{1}{2}} \right] \le n\lambda.
$$
 (3.2)

(ii) If there exist m functions  $f_r : M \to \mathbb{R}$  such that (1.11) holds, then

$$
\sum_{r=1}^{m} (\Gamma_{r+1} - \Gamma_1)^{\frac{1}{2}} \le 2(2m+4)^{\frac{1}{2}} \Gamma_1^{\frac{1}{2}}.
$$
 (3.3)

*Proof* (i) Since  $\varphi_1, \ldots, \varphi_{m+1}$  are the components of an eigenmap, it holds

$$
\begin{cases}\n-\Delta \varphi_{\alpha} = \lambda \varphi_{\alpha}, \\
\sum_{\substack{\alpha=1 \\ m+1 \\ \alpha=1}}^{\infty} |\nabla \varphi_{\alpha}|^2 = \lambda, \\
\sum_{\alpha=1}^{\infty} \varphi_{\alpha}^2 = 1.\n\end{cases}
$$
\n(3.4)

In order to make use of Theorem 2.1, we construct some functions satisfying (2.1) by using  $\{\varphi_{\alpha}\}_{\alpha=1}^{m+1}$ . We consider an  $(m+1)\times(m+1)$  matrix  $Q = (\int_{\Omega} \varphi_{\alpha} u_1 u_{\beta+1})_{(m+1)\times(m+1)}$ . According to the QR-factorization theorem, we know that there exists an orthogonal  $(m + 1) \times (m + 1)$ matrix  $P = (p_{\alpha\beta})_{(m+1)\times(m+1)}$  such that  $U = PQ$  is an upper triangle matrix. Namely, we have

$$
\sum_{\gamma=1}^{m+1} p_{\alpha\gamma} \int_{\Omega} \varphi_{\gamma} u_1 u_{\beta+1} = 0, \quad \text{for } 1 \le \beta < \alpha \le m+1.
$$

Define functions  $\psi_{\alpha}$  by

$$
\psi_{\alpha} = \sum_{\gamma=1}^{m+1} p_{\alpha\gamma} \varphi_{\gamma}.
$$

Thus we infer

$$
\int_{\Omega} \psi_{\alpha} u_1 u_{\beta+1} = 0, \quad \text{for } 1 \le \beta < \alpha \le m+1.
$$
\n(3.5)

Moreover, because  $P$  is an orthogonal matrix, it follows from  $(3.4)$  that

$$
\begin{cases}\n-\Delta \psi_{\alpha} = \lambda \psi_{\alpha}, \\
\sum_{\substack{\alpha=1 \\ \alpha+1}}^{\alpha+1} |\nabla \psi_{\alpha}|^2 = \lambda, \\
\sum_{\alpha=1}^{\alpha+1} \psi_{\alpha}^2 = 1.\n\end{cases}
$$
\n(3.6)

According to (3.5), taking  $A = \Delta^2$ ,  $B_{\alpha} = \psi_{\alpha}$  and  $T_{\alpha} = [\Delta, \psi_{\alpha}]$  in Theorem 2.1, we derive

$$
\sum_{\alpha=1}^{m+1} (\Gamma_{\alpha+1} - \Gamma_1)^{\frac{1}{2}} \langle [[\Delta, \psi_{\alpha}], \psi_{\alpha}] u_1, u_1 \rangle
$$
  
 
$$
\leq 2 \Bigg\{ \sum_{\alpha=1}^{m+1} \langle [\Delta^2, \psi_{\alpha}] u_1, \psi_{\alpha} u_1 \rangle \sum_{\alpha=1}^{m+1} \| [\Delta, \psi_{\alpha}] u_1 \|^2 \Bigg\}^{\frac{1}{2}}.
$$
 (3.7)

Now we need to calculate and estimate both sides of (3.7). A straightforward calculation gives

$$
\langle [\Delta^2, \psi_\alpha] u_1, \psi_\alpha u_1 \rangle
$$
  
=  $\int_{\Omega} u_1 \psi_\alpha [u_1 \Delta^2 \psi_\alpha + 2 \Delta \psi_\alpha \Delta u_1 + 2 \langle \nabla \psi_\alpha, \nabla \Delta u_1 \rangle + 2 \langle \nabla \Delta \psi_\alpha, \nabla u_1 \rangle$   
+  $2 \Delta \langle \nabla \psi_\alpha, \nabla u_1 \rangle ]$   
=  $\int_{\Omega} [u_1^2 (\Delta \psi_\alpha)^2 + 4u_1 \Delta \psi_\alpha \langle \nabla \psi_\alpha, \nabla u_1 \rangle + 4 \langle \nabla \psi_\alpha, \nabla u_1 \rangle^2 - 2 |\nabla \psi_\alpha|^2 u_1 \Delta u_1].$  (3.8)

Moreover, using the Cauchy–Schwarz inequality, one gets

$$
\int_{\Omega} |\nabla u_1|^2 = -\int_{\Omega} u_1 \Delta u_1 \le \left(\int_{\Omega} u_1^2\right)^{\frac{1}{2}} \left[\int_{\Omega} (-\Delta u_1)^2\right]^{\frac{1}{2}} = \Gamma_1^{\frac{1}{2}}.
$$
\n(3.9)

Then it follows from  $(3.6)$ ,  $(3.8)$  and  $(3.9)$  that

$$
\sum_{\alpha=1}^{m+1} \langle [\Delta^2, \psi_{\alpha}] u_1, \psi_{\alpha} u_1 \rangle = \lambda^2 \int_{\Omega} u_1^2 + 4 \int_{\Omega} \sum_{\alpha=1}^{m+1} \langle \nabla \psi_{\alpha}, \nabla u_1 \rangle^2 - 2\lambda \int_{\Omega} u_1 \Delta u_1
$$
  

$$
\leq \lambda^2 + 4 \int_{\Omega} \sum_{\alpha=1}^{m+1} |\nabla \psi_{\alpha}|^2 |\nabla u_1|^2 - 2\lambda \int_{\Omega} u_1 \Delta u_1
$$
  

$$
\leq \lambda^2 + 6\lambda \Gamma_1^{\frac{1}{2}}.
$$
 (3.10)

At the same time, it follows from (3.6) and (3.9) that

$$
\sum_{\alpha=1}^{m+1} \| [\Delta, \psi_{\alpha}] u_1 \|^2 = \int_{\Omega} \sum_{\alpha=1}^{m+1} \left[ u_1^2 (\Delta \psi_{\alpha})^2 + 4 \langle \nabla \psi_{\alpha}, \nabla u_1 \rangle^2 + 4u_1 \Delta \psi_{\alpha} \langle \nabla \psi_{\alpha}, \nabla u_1 \rangle \right]
$$
  
\$\leq \lambda^2 + 4\lambda \Gamma\_1^{\frac{1}{2}}\$. (3.11)

On the other hand, according to the definition of the commutator, it holds

$$
\langle [[\Delta, \psi_{\alpha}], \psi_{\alpha}]u_1, u_1 \rangle
$$
  
=  $\int_{\Omega} u_1 [\Delta(\psi_{\alpha}^2 u_1) - 2\psi_{\alpha} \Delta(\psi_{\alpha} u_1) + \psi_{\alpha}^2 \Delta u_1]$   
=  $\int_{\Omega} \{2u_1 \Delta(\psi_{\alpha}^2 u_1) - 2u_1 \psi_{\alpha} [\psi_{\alpha} \Delta u_1 + 2 \langle \nabla \psi_{\alpha}, \nabla u_1 \rangle + u_1 \Delta \psi_{\alpha}] \}$   
=  $-2 \int_{\Omega} [u_1 \psi_{\alpha} \langle \nabla \psi_{\alpha}, \nabla u_1 \rangle + u_1^2 \psi_{\alpha} \Delta \psi_{\alpha}]$   
=  $2 \int_{\Omega} u_1^2 |\nabla \psi_{\alpha}|^2$ . (3.12)

Moreover, since

$$
\sum_{j=1}^{n} |\nabla \psi_j|^2 \le \sum_{\alpha=1}^{m+1} |\nabla \psi_\alpha|^2 = \lambda,
$$

we observe that

$$
|\nabla \psi_j|^2 \le \frac{\lambda}{n}, \quad \text{for } j = 1, \dots, n. \tag{3.13}
$$

Therefore, making use of (3.13), we can deduce

$$
\sum_{\alpha=1}^{m+1} (\Gamma_{\alpha+1} - \Gamma_1)^{\frac{1}{2}} |\nabla \psi_{\alpha}|^2
$$
\n
$$
\geq \sum_{i=1}^n (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}} |\nabla \psi_i|^2 + (\Gamma_{n+1} - \Gamma_1)^{\frac{1}{2}} \sum_{k=n+1}^{m+1} |\nabla \psi_k|^2
$$
\n
$$
= \sum_{i=1}^n (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}} |\nabla \psi_i|^2 + (\Gamma_{n+1} - \Gamma_1)^{\frac{1}{2}} \sum_{j=1}^n \left(\frac{\lambda}{n} - |\nabla \psi_j|^2\right)
$$
\n
$$
\geq \sum_{i=1}^n (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}} |\nabla \psi_i|^2 + \sum_{j=1}^n (\Gamma_{j+1} - \Gamma_1)^{\frac{1}{2}} \left(\frac{\lambda}{n} - |\nabla \psi_j|^2\right)
$$
\n
$$
= \frac{\lambda}{n} \sum_{i=1}^n (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}}.
$$
\n(3.14)

Combining  $(3.12)$  and  $(3.14)$ , we obtain

$$
\sum_{\alpha=1}^{m+1} (\Gamma_{\alpha+1} - \Gamma_1)^{\frac{1}{2}} \langle [[\Delta, \psi_{\alpha}], \psi_{\alpha}] u_1, u_1 \rangle = 2 \sum_{\alpha=1}^{m+1} (\Gamma_{\alpha+1} - \Gamma_1)^{\frac{1}{2}} \int_{\Omega} u_1^2 |\nabla \psi_{\alpha}|^2
$$
  

$$
\geq \frac{2\lambda}{n} \sum_{i=1}^{n} (\Gamma_{i+1} - \Gamma_1)^{\frac{1}{2}}.
$$
 (3.15)

Substituting (3.10), (3.11) and (3.15) into (3.7), we can infer (3.1).

From  $[(\lambda + 6\Gamma_1^{\frac{1}{2}})(\lambda + 4\Gamma_1^{\frac{1}{2}})]^{\frac{1}{2}} \leq \lambda + 5\Gamma_1^{\frac{1}{2}}$ , we know that (3.2) is true.

(ii) Similar to the proof of (i), using the QR-factorization theorem, we can construct functions  $\{h_r\}_{r=1}^m$  by using  $\{f_r\}_{r=1}^m$ , such that  $h_r$  satisfy

$$
\begin{cases}\n\int_{\Omega} h_r u_1 u_{t+1} = 0, \quad 1 \le t < r \le m, \\
\langle \nabla h_r, \nabla h_s \rangle = \delta_{rs}, \\
\Delta h_r = 0.\n\end{cases} \tag{3.16}
$$

 $\Box$ 

Therefore, taking  $A = \Delta^2$ ,  $B_r = h_r$  and  $T_r = [\Delta, h_r]$  in Theorem 2.1, we get

$$
\sum_{r=1}^{m} (\Gamma_{r+1} - \Gamma_1)^{\frac{1}{2}} \langle [[\Delta, h_r], h_r] u_1, u_1 \rangle
$$
  
 
$$
\leq 2 \bigg\{ \sum_{r=1}^{m} \langle [\Delta^2, h_r] u_1, h_r u_1 \rangle \sum_{r=1}^{m} ||[\Delta, h_r] u_1||^2 \bigg\}^{\frac{1}{2}}.
$$
 (3.17)

According to (3.16),  $\{\nabla h_r\}_{r=1}^m$  is a set of orthonormal vector fields. Hence, we have

$$
\sum_{r=1}^{m} \langle \nabla h_r, \nabla u_1 \rangle^2 \le |\nabla u_1|^2. \tag{3.18}
$$

Using (3.9), (3.16) and (3.18), we obtain

$$
\sum_{r=1}^{m} \langle [\Delta^2, h_r] u_1, h_r u_1 \rangle
$$
\n
$$
= \int_{\Omega} \sum_{r=1}^{m} \left[ u_1^2 (\Delta h_r)^2 + 4u_1 \Delta h_r \langle \nabla h_r, \nabla u_1 \rangle + 4 \langle \nabla h_r, \nabla u_1 \rangle^2 - 2 |\nabla h_r|^2 u_1 \Delta u_1 \right]
$$
\n
$$
= \int_{\Omega} \sum_{r=1}^{m} \left[ 4 \langle \nabla h_r, \nabla u_1 \rangle^2 - 2 |\nabla h_r|^2 u_1 \Delta u_1 \right]
$$
\n
$$
\leq (2m+4) \int_{\Omega} |\nabla u_1|^2
$$
\n
$$
\leq (2m+4) \Gamma_1^{\frac{1}{2}}
$$
\n(3.19)

and

$$
\sum_{r=1}^{m} \| [\Delta, h_r] u_1 \|^2 = \int_{\Omega} \sum_{r=1}^{m} \left[ u_1^2 (\Delta h_r)^2 + 4 \langle \nabla h_r, \nabla u_1 \rangle^2 + 4 u_1 \Delta h_r \langle \nabla h_r, \nabla u_1 \rangle \right] \le 4 \Gamma_1^{\frac{1}{2}}.
$$
\n(3.20)

Moreover, according to  $(3.12)$  and  $(3.16)$ , the term on the left-hand side of  $(3.17)$  is

$$
\langle [[\Delta, h_r], h_r]u_1, u_1 \rangle = 2 \int_{\Omega} u_1^2 |\nabla h_r|^2 = 2. \tag{3.21}
$$

Substituting  $(3.19)$ – $(3.21)$  into  $(3.17)$ , we derive  $(3.3)$ .

## **4 The Poly-Laplacian with Any Order on the Euclidean Space**

In this section, we obtain a universal inequality for problem (1.1). Moreover, it covers (1.7) of [8] when  $l = 2$ .

**Theorem 4.1** *Let*  $\Omega$  *be a bounded domain in*  $\mathbb{R}^n$ *. Denote by*  $\lambda_i$  *the i-th eigenvalue of problem* (1.1)*. Then we have*

$$
\sum_{i=1}^{n} (\lambda_{i+1} - \lambda_1)^{\frac{1}{2}} \le [4l(n+2l-2)]^{\frac{1}{2}} \lambda_1^{\frac{1}{2}}.
$$
 (4.1)

*Proof* Similar to the proof of Theorem 3.1, by utilizing the QR-factorization theorem, one can prove that there exists a set of Cartesian coordinate system  $(x_1,...,x_n)$  of  $\mathbb{R}^n$  such that the

following orthogonality conditions are satisfied:

$$
\int_{\Omega} x_i u_1 u_{j+1} = 0 \quad \text{for } 1 \le j < i \le n. \tag{4.2}
$$

Namely,  $B_i = x_i$   $(i = 1, ..., n)$  satisfy (2.1). Therefore, taking  $A = (-\Delta)^{l}$ ,  $B_i = x_i$  and  $T_i = \frac{\partial}{\partial x_i}$  in Theorem 2.1, we have

$$
\sum_{i=1}^{n} (\mu_{i+1} - \mu_1)^{\frac{1}{2}} \left\langle \left[ \frac{\partial}{\partial x_i}, x_i \right] u_1, u_1 \right\rangle
$$
\n
$$
\leq 2 \left\{ \sum_{i=1}^{n} \left\langle [(-\Delta)^l, x_i] u_1, x_i u_1 \right\rangle \sum_{i=1}^{n} \left\| \frac{\partial}{\partial x_i} u_1 \right\|^2 \right\}^{\frac{1}{2}}.
$$
\n(4.3)

Utilizing

$$
\int_{\Omega} |\nabla^k u_1|^2 \le \lambda_1^{\frac{k}{l}} \tag{4.4}
$$

(cf. [7]) and

$$
(-\Delta)^{l}(x_{i}u_{1})=x_{i}(-\Delta)^{l}u_{1}-2l\nabla(-\Delta)^{l}u_{1}\cdot\nabla x_{i},
$$

we have

$$
\sum_{i=1}^{n} \langle [(-\Delta)^{l}, x_{i}]u_{1}, x_{i}u_{1} \rangle
$$
\n
$$
= -2l \sum_{i=1}^{n} \int_{\Omega} x_{i}u_{1} \nabla (-\Delta)^{l} u_{1} \cdot \nabla x_{i}
$$
\n
$$
= l \sum_{i=1}^{n} \int_{\Omega} u_{1}(-\Delta)^{l-1} u_{1} - 2l(l-1) \sum_{i=1}^{n} \int_{\Omega} (-\Delta)^{l-2} u_{1} \nabla x_{i} \cdot \nabla (\nabla u_{1} \cdot \nabla x_{i})
$$
\n
$$
= l \sum_{i=1}^{n} \int_{\Omega} |\nabla^{l-1} u_{1}|^{2} - 2l(l-1) \int_{\Omega} (-\Delta)^{l-2} u_{1} \Delta u_{1}
$$
\n
$$
= l(n+2l-2) \int_{\Omega} |\nabla^{l-1} u_{1}|^{2}
$$
\n
$$
\leq l(n+2l-2) \lambda_{1}^{\frac{l-1}{2}}.
$$
\n(4.5)

At the same time, it is not difficult to get

$$
\left\langle \left[ \frac{\partial}{\partial x_i}, x_i \right] u_1, u_1 \right\rangle = 1 \tag{4.6}
$$

and

$$
\sum_{i=1}^{n} \left\| \frac{\partial}{\partial x_i} u_1 \right\|^2 = \int_{\Omega} |\nabla u_1|^2 \le \lambda_1^{\frac{1}{l}}.
$$
\n(4.7)

Substituting  $(4.5)$ – $(4.7)$  into  $(4.3)$ , we obtain  $(4.1)$ .

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