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Strength and deformability of small-scale rock mass models under large confinements

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Abstract Triaxial compression tests were performed to determine the strength and deformability of small-scale rock mass models with multiple joint sets and frequencies under confining stresses up to 12 MPa. The cubical sandstone specimens $(80 \times 80 \times 80 \text{ mm}^3)$ with joint sets simulated by saw-cut surfaces were compressed to failure using a true triaxial load frame. The joint frequencies ranged from 26 to 76 joints per meter. The results indicate that the Hoek–Brown criterion with two material parameters (m and s) can describe rock mass strengths as well as the three parameter criteria of Sheorey, Yudhbir and Ramamurthy-Arora. The parameter s notably decreases with increasing joint frequency, while parameter m is less sensitive to joint frequency. The confining stresses tend to enhance the effects of joint frequencies on rock mass compressive strengths. The deformation moduli in the direction normal to the joints tend to be lower than those parallel to the joints. They decrease with increasing joint frequency. Goodman's equation was modified here to allow calculation of the deformation moduli of the rock mass along the three principal directions. The modified equation can sufficiently describe the deformation moduli normal and parallel to the joints for one-joint set and threejoint set specimens under all confining stresses.

Keywords Triaxial compression - Joint frequency - Deformation modulus · Rock mass strength criterion

Introduction

Several criteria have been proposed to describe rock mass strength based on laboratory testing (Hoek and Brown [1980](#page-12-0); Saroglou and Tsiambaos [2008;](#page-12-0) Singh and Singh [2012;](#page-12-0) Rafiai [2011;](#page-12-0) Kulatilake et al. [2006\)](#page-12-0). Some criteria are developed based on case studies (Sheorey et al. [1989;](#page-12-0) Cai et al. [2004](#page-12-0)) and numerical simulations (Halakatevakis and Sofianos [2010\)](#page-12-0), primarily to determine the effects of joint frequency, joint orientation, and joint set number on rock mass strengths. It has been found that rock mass strength decreases with increasing joint frequency (Ramamurthy and Arora [1994](#page-12-0)) and joint set number (Yang et al. [1998\)](#page-12-0). The joint orientation also affects rock mass strength. The lower strengths are obtained when the joint planes make angles between 30° and 40° with the major principal stress (Ramamurthy and Arora [1994](#page-12-0); Colak and Unlu [2004;](#page-12-0) Goshtasbi et al. [2006\)](#page-12-0). Some strength criteria have been verified by comparison with actual in-situ conditions (Edelbro [2004\)](#page-12-0) and laboratory testing (Sridevi and Sitharam [2000;](#page-12-0) Goshtasbi et al. [2006\)](#page-12-0).

Goodman ([1970\)](#page-12-0) presents an equation to evaluate the elastic constants for an equivalent continuous material representative of a rock mass regularly crossed by a single set of joints using the concept of joint stiffness. Based on testing results, Yoshinaka and Yamabe ([1986\)](#page-12-0) present equations for obtaining equivalent deformation moduli of rock masses with joints in arbitrary orientations and spacings. Ramamurthy ([2001\)](#page-12-0) introduces a joint factor to predict the decrease of deformation modulus with increasing joint frequency. The variation of the deformation moduli due to various joint orientations exhibits U-shaped behavior as reported by Ramamurthy [\(2001](#page-12-0)) and Tiwari and Rao [\(2006](#page-12-0)). Nasseri et al. ([2003\)](#page-12-0), Tiwari and Rao [\(2006](#page-12-0)), and Ebadi et al. [\(2011](#page-12-0)) suggest that the deformation modulus tends to increase with confining pressure. Even though

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several rock mass strength and deformability criteria have been proposed, verification of their predictability under large confinements has rarely been attempted, particularly for rock masses with multiple joint sets.

The objective of this study was to perform triaxial compressive strength tests on sandstone specimens with single and multiple joint sets under confining stresses up to 12 MPa. The joints were simulated by saw-cut surfaces with frequencies ranging from 26 to 76 joints per meter. Some commonly used strength and deformability criteria were applied to the test results to evaluate their validity. From our results, the effects of joint frequency, joint set number, and confining stress on the strengths and deformation moduli of the small-scale rock mass models were determined.

Sample preparation

The rock specimens used in this study were Phra Wihan sandstone, which is a fine-grained quartz sandstone with highly uniform texture and density (Boonsener and

Fig. 1 Examples of small blocks prepared for three-joint set specimens with 38 joints/m (a) and 76 joints/m (b)

Sonpiron [1997](#page-12-0)). We prepared them to obtain cubical rock mass specimens with nominal dimensions of $80 \times 80 \times 80$ mm³. We also made artificial joints out of saw-cut surfaces using a universal masonry saw (Husqvarna TS 400 F). This saw has a 400-mm diameter blade with a constant rotational speed of 2800 rpm. Water is used as cutting fluid. The saw base has two mutually perpendicular guide rails to provide a precise cutting angle and intervals. For our research, large blocks $(250 \times$ 250×250 mm³) of the sandstone were first cut into thin slabs to obtain a predefined thickness (e.g., 40, 27, 20 and 13 mm) depending on the required joint frequencies. These slabs are then cut across to obtain tabular-shaped specimens, and subsequently, cubical-shaped specimens. These small blocks were assembled to form rock mass models with different joint frequencies and joint set numbers. Figure 1 shows example of rock blocks prepared for three-joint set specimens. Rock blocks with visible crack or chipping at the edges and corners were discarded. In total, we prepared up to 80 rock mass model specimens with two different joint conditions: single-joint set and sets with three mutually perpendicular joints. The joint frequencies and orientations with respect to the applied loads are shown in Table 1. They are briefly described below.

For case I, one-joint set specimens wre prepared to study the effects of joint frequency on the strength of the rock specimens. The joints were parallel to the major principal axis. There were 1, 2, 3, or 4 joints for each set (equivalent to 26–63 joints per meter), and each set had equal spacing.

Case II is similar to case I: one-joint set specimens were prepared to assess the effect of joint frequency when the joints are normal to the major principal axis.

For Case III, specimens with three-mutually perpendicular joint sets were prepared to study the effects of joint

Table 1 Small-scale rock mass specimens prepared for triaxial compression tests with σ_3 up to 12 MPa

set number and joint frequency. This case involved 1, 2, 3, 4, or 5 joints for each set (equivalent to 26–76 joints per meter).

Test method

A true triaxial load frame (Komenthammasopon [2014](#page-12-0)) was used to apply axial stress (σ_1) and constant confining pressures (σ_3) to the intact and jointed rock specimens (Fig. 2). This device was developed to test rock specimens with soft to medium strengths under polyaxial stress states. Three pairs of 100-ton hydraulic pressure cylinders are set in three mutually perpendicular directions. The measurement system comprises pressure transducers, displacement transducers, a switching box, and a data logger. The device can accommodate the cubic or rectangular specimens of different sizes by adjusting the distances between the opposite steel loading platens. For this study the specimens were oven-dried before testing. Neoprene sheets were placed at all interfaces between the loading platens and specimen surfaces to minimize the friction. The constant confining (lateral) stresses ranged from 0, 1, 3, 5, 7, to 12 MPa. First, the specimen was subjected to the pre-defined confining stresses. The axial stress was then increased at a constant rate of 0.1 MPa/s using an electric hydraulic pump. The specimen deformations were monitored along the three loading directions, and were used to calculate the principal strains during loading. The readings were recorded until failure occured. We defined failure by the drop of the applied axial stress. Photographs were taken of the post-test specimens and the modes of failure were identified.

Test results

Tables 2, [3,](#page-3-0) and [4](#page-3-0) give the strength results in terms of the major and minor principal stresses at failure for cases I, II and III. The strengths decrease with increasing joint frequency.

 σ

Figure [3](#page-4-0) shows examples of stress–strain relations in terms of the differential stress as a function of strain for the three mutually perpendicular joint set specimens (Case III). The effect of joint frequency on the rock specimens can be observed by the reduction of failure stresses and the increase of the failure strains. More discussions and analyses of the deformation moduli as affected by the joint frequency are given in the following section.

The effects of the confining stresses on the strength of the specimens can be observed from the σ_1 – σ_3 diagrams, as shown in Fig. [4.](#page-5-0) The relations between σ_1 and σ_3 at failure tend to be non-linear for the intact sandstone and for the specimens with all joint frequencies. The specimens with higher joint frequencies (J_f) show lower strengths than those with lower joint frequencies. The effect of joint frequency on strength tends to be greater for the specimens with joints Fig. 2 True triaxial load frame parallel to σ_1 (Case I, Fig. [4a](#page-5-0)), as compared to those with

Table 2 Deformation moduli for case I

Switching box Displacement transducers Data logger Pressure transducers

Table 3 Deformation moduli for case II

joints normal to σ_1 (Case II, Fig. [4](#page-5-0)b). For example, the strengths of the specimens in case I with 63 joints/m decreases by about 20 % from the intact strength. Under this joint frequency the specimen strengths for case II decrease by about 15 %. The effects of joint frequency are greatest for the three-joint set specimens (Case III, Fig. [4c](#page-5-0)). The strengths of specimens containing three-joint sets with 63 joints per meter drop by nearly 40 % from the intact strength.

The decreases of the specimen strengths with increasing joint frequencies can also be observed from the σ_1 - J_f diagrams in Fig. [5.](#page-6-0) For the one-joint set specimens the differences of the strengths between case I and case II tend to be larger when the confining stresses increase (Fig. [5](#page-6-0)a). The σ_1 -J_f diagrams for case III (Fig. [5](#page-6-0)b) suggest also that the decreases of the threejoint set specimen strengths with increasing joint frequency are more rapid than those of the one-joint set specimens.

Post-test observations

Post-test specimens show two types of failure mode: extensile splitting and compressive shear failure. The specimens with joints parallel to σ_1 (Case I) failed mainly under the extensile splitting mode, particularly under low confining stresses (Fig. [6](#page-6-0)a). Some shear fractures were found in the specimens that failed under high confining stresses. The compressive shear failure was predominant in the specimens with joints normal to σ_1 (Case II, Fig. [6](#page-6-0)b). These observations support the measurement results that the single-joint set specimens with joints normal to σ_1 (Case II) show higher strength than those with joints parallel to σ_1 (Case I). The strength discrepancies are primarily due to the fact that σ_1 direction for case I was parallel to small plates of the sandstone forming the rock mass models. The height-to-width ratios of these plates are relatively high (varying from 2 to 6, depending on the joint frequencies). Under axial loading, each plate can laterally dilate (toward the thinner sides), and hence the extensile fractures can be induced in vertical or nearly vertical directions. The pre-existing joints also help the induced 7 50.0 0.26 1.8 1.8 1.8 12 62.2 0.29 2.2 2.2 2.2

Fig. 3 Examples of differential stresses as a function of strain for three mutually perpendicular joint sets under various confining pressures. a 26 joints/ m, b 38 joints/m, c 51 joints/m, d 63 joints/m and e 76 joints/m

extensile fractures to propagate through the specimen models. The induced fractures can propagate more easily for the specimens with higher joint frequency, as compared to those with lower joint frequencies. This explains why the strengths of case I specimens decreased with increasing joint frequency. Under large confinement the lateral dilation of the sandstone plates toward the pre-existing joints becomes more difficult, and hence some shear fractures are induced across the specimen models.

The sandstone plates assembled for case II specimens were normal to the σ_1 direction. They could not dilate easily under loading, and therefore the compressive shear fractures were predominant. As a result the strengths of case II specimens tended to be greater than those of case I. These shear fractures for case II can propagate more easily in specimens with higher joint frequencies and under lower confining stresses, as compared to those with lower joint frequencies and higher confining stresses.

Fig. 4 Major principal stresses at failure as a function of confining pressure for cases I (a), II (b), and III (c)

Both failure modes were found in the three-joint set specimens (Case III, Fig. [6](#page-6-0)c). This is probably because all case III specimens were formed by small cubical blocks of sandstone (varying from $13 \times 13 \times 13$ to $40 \times 40 \times 40$ mm³, depending on joint frequency). Under the same joint frequency the strengths of the case III specimens were lower than those of cases I and II because each sandstone block had an additional free face to dilate (i.e., there were two mutually perpendicular joint sets parallel to the σ_1 direction).

For all cases the extensile splitting and shear fractures were induced in the intact portion of the specimens. This suggests that if there were invisible cracks or damages along the edges and corners of the small blocks, they may not have had a significant impact on the strength of the specimens.

Strength criteria

Four commonly used strength criteria were compared against the test results obtained from the three-joint set specimens. These include criteria introduced by Hoek– Brown ([1980\)](#page-12-0), Sheorey [\(1989](#page-12-0)), Yudhbir et al. ([1983\)](#page-12-0), and Ramamurthy and Arora ([1994\)](#page-12-0). Exhaustive reviews of these criteria have been given elsewhere (Hoek and Brown [1980](#page-12-0); Sheorey [1997;](#page-12-0) Hoek et al. [2002;](#page-12-0) Edelbro [2004\)](#page-12-0). They have been widely used to study the strengths of rock mass.

The Hoek–Brown ([1980\)](#page-12-0) criterion defines the relationship between the major and minor principal stresses at failure as:

$$
\sigma_1 = \sigma_3 + \sqrt{m\sigma_c\sigma_3 + s\sigma_c^2} \tag{1}
$$

where *m* and *s* are material constants and σ_c is the uniaxial compressive strength of the intact rock.

Sheorey ([1989\)](#page-12-0) defines the rock mass strength criterion as:

$$
\sigma_1 = \sigma_{\rm cm} \left(1 + \frac{\sigma_3}{\sigma_{\rm tm}} \right)^{b_{\rm m}} \tag{2}
$$

where b_m is a constant, σ_{cm} is the uniaxial compressive strength of rock mass, and σ_{tm} is the uniaxial tensile strength of rock mass.

Yudhbir et al. [\(1983](#page-12-0)) modify the original Bieniawski [\(1974](#page-12-0)) criterion, and propose a more general form as:

$$
\frac{\sigma_1}{\sigma_c} = A + B \left(\frac{\sigma_3}{\sigma_c} \right)^{\alpha} \tag{3}
$$

where α and β are material constants, and β is a dimensionless parameter depending on rock mass quality.

Ramamurthy and Arora ([1994\)](#page-12-0) present a strength criterion that is modified from the Mohr–Coulomb theory:

$$
\frac{\sigma_1 - \sigma_3}{\sigma_3} = \beta \left(\frac{\sigma_{\rm cm}}{\sigma_3} \right)^{\alpha} \tag{4}
$$

where α and β are the material constants. Ramamurthy and Arora [\(1994](#page-12-0)) suggest an alternative formula to determine $\sigma_{\rm cm}$ in terms of joint factor J_F as: $\sigma_{cm} = \sigma_c \exp(-0.008 J_F)$.

Hoek et al. ([2002](#page-12-0)) propose a generalized form of the Hoek–Brown criterion as:

$$
\sigma_1 = \sigma_3 + \left(m_b \sigma_c \sigma_3 + s \sigma_c^2 \right)^a \tag{5}
$$

Fig. 5 Failure stress (σ_1) as a function of joint frequency (J_f) for one-joint set specimens (a) and three-joint set specimens (b)

Fig. 6 Examples of extensile splitting (a), compressive shear failure (b), and combination mode (c)

 (b)

where m_b , s and a are fitting parameters. It is applicable to a heavily jointed rock mass. The parameters m_b and s can be estimated from the geological strength index (GSI) (Hoek [1994;](#page-12-0) Hoek et al. [1995](#page-12-0); Hoek and Brown [1998](#page-12-0); Sonmez and Ulusay [1999\)](#page-12-0). This index is determined from the joint surface conditions, structure rating, and volumetric joint count. The calculated GSI for the rock specimens used here are from 40 to 45 because the tested joints have no infilling and weathering. It is recommended that if GSI is greater than 25, the " a " parameter should be set equal to 0.5 (Sheorey [1997](#page-12-0); Sonmez and Ulusay [1999\)](#page-12-0). In this study the original Hoek–Brown criterion is therefore used instead of the generalized Hoek–Brown criterion.

 (a)

Statistical analyses are performed using the SPSS code (Wendai [2000](#page-12-0)) to fit the test results with the criteria above. The predictive capability of these criteria is determined and compared using the coefficient of correlation (R^2) as an indicator. The material constants and coefficients of correlation calculated for these criteria are summarized in Table [5](#page-7-0). All criteria provide good correlation with the test data, with R^2 greater than 0.9. Figure [7](#page-7-0) compares the test results with the curve fits in the terms of σ_1 as a function of σ_3 at failure.

Even though the Hoek and Brown criterion uses only two material constants, its predictive capability is as good as those with three material constants. Figure [8](#page-8-0) shows the decrease of parameters m and s of the Hoek–Brown criterion as the joint frequency increases. For intact sandstone specimens, the parameter m is determined as 14.1 where $s = 1.0$. These *m* and *s* values for one-joint set specimens are greater than those for three-joint set specimens. For one-joint set specimens, the effect of joint orientation on m and s is relatively small (comparing cases I and II). The difference of the Hoek–Brown parameters between case I and case II seems to be the same for all joint frequencies (Fig. [8\)](#page-8-0). However, significant decreases of both m and s values are found for the three-joint set specimens when the joint frequencies increase (Case III). The parameter s drops from 1.0 for intact condition to 0.1 for specimens with joint frequency of 76 joints/m. The Ramamurthy and Arora ([1994\)](#page-12-0) criterion can not predict the rock mass strength under unconfined conditions ($\sigma_3 = 0$). As a result, Eq. [\(5](#page-5-0)) was proposed to predict the uniaxial strength of rock mass (σ _{cm}). Both the Sheorey and Ramamurthy-Arora criteria can sufficiently describe the uniaxial strength of the rock specimens as shown in Fig. [9](#page-8-0). The uniaxial strengths

Table 5 Strength criteria and their constants calibrated from test results of case III

Strength criteria	Parameters	Joint frequencies (joints/m)					
		Intact	26	38	51	63	76
		σ_{1}		σ_{1}	σ_1	σ_1	
				翻-03 σ_{3}	σ_3 $\overline{\mathbb{H}}$ σ_3	σ_3 -бз	σ_{3}
Hoek-Brown (1980) $\sigma_1 = \sigma_3 + (m \sigma_c \sigma_3 + s)$ $\sigma_{\rm c}^2$ ^{1/2}	\boldsymbol{m}	14.10	11.30	9.22	7.89	6.03	4.83
	\boldsymbol{S}	1.0	0.7	0.5	0.4	0.3	0.2
	R^2	0.992	0.969	0.984	0.971	0.979	0.986
Sheorey et al. (1989) $\sigma_1 = \sigma_{cm} (1 + \sigma_3/\sigma_{tm})^{bm}$	$\sigma_{\rm cm}$	41.8	35.8	31.0	26.5	22.9	19.6
	$\sigma_{\rm tm}$	2.6	2.1	1.8	1.4	1.2	1.0
	$b_{\rm m}$	0.54	0.51	0.48	0.47	0.46	0.45
	R^2	0.991	0.986	0.988	0.991	0.992	0.997
Yudhbir et al. (1983) $\sigma_1/\sigma_c = A + B(\sigma_3/\sigma_c)^{\alpha}$	\boldsymbol{A}	0.97	0.82	0.70	0.59	0.50	0.43
	\boldsymbol{B}	4.04	3.64	3.31	2.98	2.71	2.44
	α	0.79	0.76	0.72	0.70	0.68	0.65
	R^2	0.986	0.984	0.986	0.987	0.989	0.993
Ramamurthy and Arora (1994) $(\sigma_1 - \sigma_3)$ / $\sigma_3 = \beta (\sigma_{\rm cm}/\sigma_3)^{\alpha}$	$\sigma_{\rm cm}$	43.3	37.4	32.3	27.5	24.0	20.7
	β	3.19	3.15	3.13	3.08	3.00	2.90
	α	0.69	0.69	0.69	0.69	0.69	0.69
	R^2	0.989	0.973	0.979	0.980	0.997	0.995

Fig. 7 Test results (points) and curve fits for four strength criteria

Fig. 8 Hoek–Brown parameters $m(\mathbf{a})$ and $s(\mathbf{b})$ as a function of joint frequency

Fig. 9 Uniaxial compressive strengths of rock mass model ($\sigma_{\rm cm}$) with three joint sets (case III) as a function of joint frequency. The predictions (lines) are compared with the measurements (points)

decrease from 43.3 MPa for intact conditions to about 20 MPa for the joint frequency of 76 joints/m.

Deformation moduli

The deformation moduli along the three principal axes of rock specimens are determined from the stress–strain curves at about 50 % of the failure stress. It is initially assumed that the Poisson's ratio (v) of the specimens is the same for all principal planes. The deformation moduli along the three principal directions for each case can be derived from the elastic stress– strain relations (Jaeger et al. [2007](#page-12-0)) as shown below. In case I:

$$
\varepsilon_{1,\mathrm{p}} = \sigma_1/E_{\mathrm{m},\mathrm{p}} - v(\sigma_3/E_{\mathrm{m},\mathrm{p}} + \sigma_3/E_{\mathrm{m},\mathrm{n}}) \tag{6}
$$

$$
\varepsilon_{3,\mathrm{p}} = \sigma_3 / E_{\mathrm{m,p}} - \nu (\sigma_1 / E_{\mathrm{m,p}} + \sigma_3 / E_{\mathrm{m,n}}) \tag{7}
$$

$$
\varepsilon_{3,n} = \sigma_3/E_{m,n} - v(\sigma_1/E_{m,p} + \sigma_3/E_{m,p})
$$
 (8)

Case II:

$$
\varepsilon_{1,n} = \sigma_1/E_{m,n} - v(\sigma_3/E_{m,p} + \sigma_3/E_{m,p})
$$
\n(9)

$$
\varepsilon_{3,p} = \sigma_3 / E_{m,p} - \nu (\sigma_1 / E_{m,n} + \sigma_3 / E_{m,p})
$$
 (10)

$$
\varepsilon_{3,p} = \sigma_3/E_{m,p} - \nu(\sigma_1/E_{m,n} + \sigma_3/E_{m,p})
$$
 (11)

where $\varepsilon_{1,p}$ and $\varepsilon_{1,n}$ are the major principal strains parallel and normal to the joints, $\varepsilon_{3,p}$ and $\varepsilon_{3,n}$ are the minor principal strains parallel and normal to the joints, and $E_{\text{m,p}}$ and $E_{\text{m,n}}$, are the deformation moduli parallel and normal to the joints.

In case III:

$$
\varepsilon_{1,m} = \sigma_1/E_{m,1} - v(\sigma_3/E_{m,2} + \sigma_3/E_{m,3})
$$
 (12)

$$
\varepsilon_{3,m} = \sigma_3/E_{m,2} - v(\sigma_1/E_{m,1} + \sigma_3/E_{m,3})
$$
 (13)

$$
\varepsilon_{3,m} = \sigma_3/E_{m,3} - v(\sigma_1/E_{m,1} + \sigma_3/E_{m,2})
$$
 (14)

where $\varepsilon_{1,m}$ is the major principal strain, $\varepsilon_{3,m}$ is the minor principal strains, $E_{m,1}$ is the deformation modulus along the vertical (σ_1) axis, and $E_{\text{m},2}$ and $E_{\text{m},3}$ are the deformation moduli along the two lateral principal axes. Tables [2](#page-2-0), [3,](#page-3-0) and [4](#page-3-0) show the calculation results for all cases. The calculated Poisson's ratios tend to be independent of the joint frequency and loading direction. Their values average between 0.23 and 0.29. For one-joint set specimens the deformation moduli parallel to the joints $(E_{m,p})$ are about 13 % higher than those normal to the joints $(E_{m,n})$ as shown in Fig. [10](#page-9-0). This is true for all joint frequencies regardless of whether the joints are parallel (case I) or normal (case II) to σ_1 . For the three-joint set specimens, the deformation moduli are similar for all principal directions (Fig. [11](#page-9-0)). Figure [12](#page-9-0)a and b plot the average deformation moduli that are parallel $(E_{m,p})$ and normal $(E_{m,n})$ to the joints for cases I and II. For the three-joint set specimens (case III), the deformation moduli averaged from the three principal directions (represented by E_m) are plotted as a function of σ_3 in Fig. [12c](#page-9-0). The deformation moduli ($E_{\rm m,p}$, $E_{\rm m,n}$ and $E_{\rm m}$) decrease with increasing joint frequency.

Deformability criteria

The deformation moduli calculated above are compared against the criteria developed by Goodman [\(1970](#page-12-0)), Yoshinaka and Yamabe ([1986](#page-12-0)), and Ramamurthy [\(2001](#page-12-0)).

Fig. 10 Deformation moduli parallel to joint planes as a function of those normal to joint planes

Fig. 11 Deformation moduli calculated along the two minor principal axes $(E_{m,2}$ and $E_{m,3}$) as a function of those along the major principal axis $(E_{m,1})$ for case III

Goodman ([1970](#page-12-0)) presents an equation to determine the deformation modulus that is normal to the joints as:

$$
\frac{1}{E_{\text{m,n}}} = \frac{1}{k_{\text{n}}s} + \frac{1}{E_{\text{i}}}
$$
\n(15)

where E_i is the elastic modulus of intact rock, k_n is normal joint stiffness and s is the average joint spacing. This equation is only applicable to the results obtained from case II testing.

Yoshinaka and Yamabe [\(1986](#page-12-0)) proposed an equation to describe the deformability of samples with more than one joint under various joint orientations.

For case I testing, the joints are parallel to σ_1 and the proposed equation becomes:

$$
\frac{1}{E_{\rm m,p}} = \frac{1}{E_{\rm i}}.\tag{16}
$$

For cases II and III, the joints are parallel and normal to σ_1 and their equations become:

$$
\frac{1}{E_{\text{m,n}}} = \frac{1}{E_{\text{i}}} + \frac{1}{k_{\text{n}}s} \text{(Case II)} \tag{17}
$$

$$
\frac{1}{E_{\rm m}} = \frac{1}{E_{\rm i}} + \frac{1}{k_{\rm n}s} \text{(Case III)}.
$$
 (18)

Fig. 12 Deformation moduli parallel (a) and normal (b) to σ_1 axis as a function of confining stress for one-joint set specimens and for three-joint set specimens (c)

Ramamurthy [\(2001](#page-12-0)) defines the relationship between the ratios of deformation moduli, $E_{\text{m,n}}/E_i$, $E_{\text{m,p}}/E_i$ and joint factor (J_F) as:

$$
E_{\rm m,n} = E_{\rm m,p} = E_{\rm i} \exp(-1.15 \times 10^{-2} J_{\rm F})
$$
 (19)

where J_F is the joint factor, which is defined by:

$$
J_{\rm F} = \frac{J_{\rm f}}{n \cdot r} \tag{20}
$$

where J_f is the joint frequency (number of joints per meter), n is the inclination parameter depending upon the orientation of the joint, and r is a joint strength parameter dependent upon the joint condition. Note that Eqs. (19) and (20) are applicable to cases I and II only.

Fig. 13 Predictions of deformation moduli compared with test data for one-joint set specimens (a, b) and three-joint set specimens (c). Points and error bars represent mean and standard deviation of deformation moduli from all confining pressures

The elastic moduli that are parallel and normal to the joints are plotted as a function of joint frequency for singlejoint set specimens in Fig. 13a and b. Their average value (data point) and standard deviation (shown as error bars) obtained from all confining pressures are shown in the figure. The joint normal stiffness used here is 381.2 GPa/m, which is obtained from Kamonphet ([2012\)](#page-12-0). The elastic modulus of the intact sandstone in the figure is averaged from the measurement results given in Fig. [12.](#page-9-0) The joint factors of Ramamurthy ([2001\)](#page-12-0) used in this study are summarized in Table 6. The Ramamurthy [\(2001](#page-12-0)) equation gives a good estimation for the deformation moduli parallel to the joint planes (Fig. 13a) and normal to the joint planes (Fig. 13b). The Goodman [\(1970](#page-12-0)) and Yoshinaka and

Table 6 Joint factors calculated for this study

Joint orientation	Joint frequency (J_f)	n	r	Joint factor (J_F)
One-joint set parallel to the major principal axis	Intact	1	0.8	
	26	0.85	0.8	38.2
	38	0.85	0.8	55.9
	51	0.85	0.8	75.0
	63	0.85	0.8	92.6
One-joint set normal to the major principal axis	Intact	1	0.8	
	26	0.98	0.8	33.2
	38	0.98	0.8	48.5
	51	0.98	0.8	65.1
	63	0.98	0.8	80.4

Yamabe ([1986\)](#page-12-0) equations can also describe the deformation moduli normal to joint planes (Fig. 13b). Figure 13c shows the elastic moduli averaged from the three principal directions for the three-joint set specimens. Yoshinaka and Yamabe's ([1986\)](#page-12-0) equation adequately describes the deformation moduli of the rock specimens with three-joint sets (Fig. 13c). Note that Yoshinaka and Yamabe ([1986\)](#page-12-0) cannot describe the deformation modulus of the rock mass along the axis that is parallel to the joint plane.

Modified Goodman equation

The Goodman [\(1970](#page-12-0)) equation is modified here to determine the deformation modulus along three principal directions. It is proposed as:

$$
\frac{1}{E_{\rm m}} = \frac{N}{k_{\rm n}s} + \frac{1}{E_{\rm i}}\tag{21}
$$

Parameter N is introduced as an empirical constant to allow the Goodman equation to be able to predict the deformation moduli that are parallel to the joint plane (Table [7\)](#page-11-0). The N values are defined by the direction of deformation moduli with respect to the joint plane as shown in Table [7.](#page-11-0) Predictions of the deformation moduli for the three cases are given in Fig. [14.](#page-11-0) Good correlations are obtained for all cases ($R^2 > 0.9$). The proposed equation, however, can predict the deformation moduli only in the directions normal and parallel to the joint planes.

Discussions and conclusions

Series of triaxial compression tests were performed to determine the strength and deformability of small-scale rock mass models with single- and multiple-joint sets and joint frequencies under large confinements. It was found that the compressive strengths decrease with increasing

Table 7 Parameter N defined for the modified Goodman equation

Verified by test results in Fig. 14

joint frequency. This agrees with the experimental observations by Ramamurthy and Arora [\(1994](#page-12-0)) on jointed specimens of Plaster of Paris. For one-joint set specimens tested here, the strengths of rock specimens with joints normal to the σ_1 axis are always greater than those with joints parallel to the σ_1 axis. This agrees with experimental observations by Colak and Unlu [\(2004](#page-12-0)), Saroglou and Tsiambaos ([2008\)](#page-12-0), and Goshtasbi et al. ([2006\)](#page-12-0). The decrease in rock mass strengths as the joint frequency increases tends to act equally throughout the range of confining stresses used here (1–12 MPa).

All strength criteria used here can sufficiently predict the strengths of the rock mass specimens under the confining stresses up to 12 MPa. The Hoek–Brown criterion with only two constants $(m \text{ and } s)$ can describe the rock mass strengths as well as the three parameters criteria. The parameter s decreases rapidly with increasing joint frequency while parameter m tends to be insensitive to the joint frequency. The parameters m and s of the one-joint set specimens are higher than those of the three-joint set specimens. This suggests that decreasing joint set numbers will increase the rock mass strength. The measured uniaxial compressive strengths of the rock specimens $(\sigma_{\rm cm})$ decrease with increasing joint frequency. They agree well with the $\sigma_{\rm cm}$ calculated from the Sheorey ([1989](#page-12-0)) and Ramamurthy and Arora ([1994](#page-12-0)) criteria.

It is recognized that the joints studied here were simulated by smooth saw-cut surfaces. The strengths of the rock specimens for all cases, therefore, represent the lower bound of the strengths of actual rock mass where most fractures are rougher. In addition, the major principal stresses applied here are always normal or parallel to the

Fig. 14 Predictions of deformation moduli by modified Goodman equation compared with test data for one-joint set specimens (a, b) and three-joint set specimens (c)

joint planes. It is expected that the rock mass model strengths would be lower if the applied stress makes oblique angles with the joint planes, as evidenced by the test results obtained by Ramamurthy and Arora [\(1994](#page-12-0)), Colak and Unlu ([2004](#page-12-0)) and Goshtasbi et al. [\(2006](#page-12-0)).

For one-joint set specimens the deformation moduli parallel to the joints show the highest values compared to those that are normal to the joints. This is true for all joint frequencies. For three-joint set specimens, the deformation moduli are similar for all principal directions. The deformation moduli decrease with increasing joint frequency. This agrees with the experimental observations by Tiwari and Rao ([2006\)](#page-12-0).

The Goodman (1970) equation can sufficiently predict the deformation moduli for one-joint set specimens with joints normal to the major principal axis (case II). It, however, cannot calculate the deformation modulus for specimens with more than a one-joint set. Yoshinaka and Yamabe's (1986) equation can determine the specimen deformability with more than a one-joint set and orientation but it cannot calculate the deformation modulus parallel to the joints. Ramamurthy's (2001) equation can predict the deformation moduli parallel to the joint planes (case I) and normal to the joint planes (case II). It cannot calculate the deformation moduli of specimens with threejoint sets. A modified Goodman equation is proposed here to determine the deformation moduli along the three principal directions. The parameter N is introduced, whose values depend on joint set directions. The modified equation can adequately describe the deformation moduli parallel and normal to the joints planes for one-joint set and three-joint set specimens.

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