# ORIGINAL PAPER



# Carbon-encapsulated nanosphere-assembled  $MoS<sub>2</sub>$  nanosheets with large interlayer distance for flexible lithium-ion batteries

Zuxue Bai<sup>1</sup> • Ya [Yan](http://orcid.org/0000-0002-0234-0477)g<sup>1</sup> • Deyang Zhang<sup>1</sup> • Yangbo Wang<sup>1</sup> • Ying Guo<sup>1</sup> • Hailong Yan<sup>1</sup> • Paul K. Chu<sup>3</sup> • Yongsong Luo<sup>1,2</sup> D

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#### Abstract

2D transition metal dichalcogenides such as  $MoS<sub>2</sub>$  with the unique layered structure have received great attention in the field of lithium-ion batteries (LIBs). However, the low conductivity and poor structural stability adversely affect the rate performance of LIBs. Herein, a flexible and free-standing high-performance lithium-ion battery electrode (MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>) composed of ricecandy-like MoS<sub>2</sub>/C intercalated Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> and PVP-derived carbon with a large interlayer distance of MoS<sub>2</sub> is designed and demonstrated. Lithium-ion batteries have attracted great attention due to their high energy density. Consequently, as an anode material for lithium-ion batteries, MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> provides a high discharge capacity of 538.5 mAh g<sup>-1</sup> at 0.05 A g<sup>-1</sup> and rapid charge/discharge capability of 256.7 mAh g<sup>-1</sup> at 2 A g<sup>-1</sup>, as well as outstanding cycling property (96.7% capacity retention after 150 cycles at 2 A  $g^{-1}$ ). Density-functional theory (DFT) calculation reveals that the rice-candy-like MoS<sub>2</sub>/C structure favors adsorption and diffusion of lithium ions and facilitates the redox reactions. The  $MoS_2/CC@Ti_3C_2T_x$  structure is expected to boost the development of novel 2D materials for high-performance lithium-ion batteries.

Keywords Lithium-ion battery  $\cdot$  MoS<sub>2</sub>  $\cdot$  Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>  $\cdot$  Large interlayer distance  $\cdot$  Flexible electrode

# Introduction

Lithium-ion batteries (LIBs) boasting many advantages such as large energy density, long-term cycle life, and environmental friendliness are widely used in portable electronics and electric vehicles  $[1-3]$  $[1-3]$  $[1-3]$  $[1-3]$ . As a class of 2D materials,

Zuxue Bai and Yangbo Wang contributed equally to this work.

 $\boxtimes$  Deyang Zhang [zdy@xynu.edu.cn](mailto:zdy@xynu.edu.cn)

 $\boxtimes$  Yongsong Luo [ysluo@xynu.edu.cn](mailto:ysluo@xynu.edu.cn)

- <sup>1</sup> Key Laboratory of Microelectronics and Energy of Henan Province, School of Physics and Electronic Engineering, Engineering Research Center for MXene Energy Storage Materials of Henan Province, Xinyang Normal University, Xinyang 464000, People's Republic of China
- <sup>2</sup> College of Physics and Electronic Engineering, Nanyang Normal University, Nanyang 473061, People's Republic of China
- Department of Physics, Department of Materials Science & Engineering, and Department of Biomedical Engineering, City University of Hong Kong, Tat Chee Avenue, Kowloon, Hong Kong, China

molybdenum disulfide  $(MoS<sub>2</sub>)$  with a graphene-like layered structure has attracted attention in physics and chemistry [[4,](#page-7-0) [5](#page-7-0)]. The layered structure composed of S–Mo–S units stacked by van der Waals attraction [[6](#page-7-0)–[8](#page-7-0)] has a large theoretical specific capacity of 670 mAh  $g^{-1}$  for Li-ion batteries [[9](#page-7-0)–[11\]](#page-7-0). However, commercial adoption of  $MoS<sub>2</sub>$  is hampered by the low electrical properties, which limit ion and electron transport [\[12](#page-7-0)], large volume variation during adsorption and diffusion of lithium ions [[13](#page-7-0)], and polysulfide dissolution during charging and discharging [[14\]](#page-7-0). To overcome these hurdles, efforts have been made to design nanoscale structures of  $MoS<sub>2</sub>$  and  $MoS<sub>2</sub>/carbon$  composites such as  $MoS<sub>2</sub>$  nanosheets [\[15](#page-7-0), [16](#page-7-0)],  $MoS_2/graphene$  [\[17](#page-7-0), [18](#page-7-0)], and  $MoS_2/CNT$  composites [\[19](#page-7-0), [20](#page-7-0)]. In fact,  $MoS<sub>2</sub>$  with large interlayer spacing and abundant active sites increases the diffusion coefficient of  $Li<sup>+</sup>$  and enhances the battery performance  $[21-23]$  $[21-23]$  $[21-23]$  $[21-23]$ . For example, Jing et al. fabricated  $1-3$  layers  $MoS<sub>2</sub>$  with wide spacing by lithiation-exfoliation process; the LIB shows a rate capacity of 350 mAh  $g^{-1}$  at 0.5 A  $g^{-1}$  and the coulombic efficiency only 52% at the first cycle [\[24](#page-7-0)]. Hence, expansion of the interlayer distance can be realized by introducing carbon materials during the synthesis of  $MoS<sub>2</sub>$  to improve storage of lithium ions.

Recently, MXene (e.g.,  $Ti_3C_2T_x$ ,  $V_2CT_x$ , and  $Ti_2CT_x$ ) as a new candidate of two-dimensional materials has attracted extensive attention in energy storage devices. MXenes are also described as  $M_{n+1}X_nT_x$ , where M stands for early transition metal. X represents a carbon or nitrogen element.  $T_x$  is the surface chemical groups. MXene, a promising material for LIBs with more unique properties than graphene or other 2D materials, has been regarded as a potential material for electrochemical energy storage because of its layered structure and excellent conductivity. For example, Qiu et al. fabricated few-layered MoS<sub>2</sub> nanoplates on Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> nanosheets stabilized by a carbon layer, which exhibited outstanding electrochemical performance for lithium storage [[25\]](#page-7-0). Fabrication of MXene hybrid nanostructures with other materials is a wellknown method, which can be utilized for important applications including supercapacitors [[26](#page-7-0)] and sensors [[27,](#page-7-0) [28](#page-7-0)]. Yola et al. prepared MWCNTs with  $Ti_3C_2T_x$  (mass ratio 3:1) colloidal solution, which exhibited high sensitivity of electrochemistry [\[29](#page-7-0)]. Therefore, these characteristics make  $Ti_3C_2T_x$  a suitable candidate for the base material to construct excellent performance  $MoS<sub>2</sub>$  nanohybrid anode materials for lithium-ion storage.

In the present work, rice-candy-like  $MoS<sub>2</sub>$  is mixed with  $Ti_3C_2T_x$  to form the self-supporting composite film of  $MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  by vacuum filtration. The sandwiched structure of free-standing  $MoS_2/C@Ti_3C_2T_x$  offers many advantages. It solves the stacking problem of  $Ti_3C_2T_x$  and  $MoS<sub>2</sub>$  and dynamically enhances  $Li<sup>+</sup>$  transport. Moreover, the free-standing  $MoS_2/COTi_3C_2T_x$  electrode buffers the volume change and avoids fast capacity decay at high charging/discharging rates. The key for  $MoS_2/COTi_3C_2T_x$ is that PVP-derived carbon is encapsulated on the  $MoS<sub>2</sub>$ layer, so that diffusion of lithium ions is facilitated and the volume variation of  $MoS<sub>2</sub>$  in charging and discharging is buffered. Consequently, the flexible  $M_0S_2/C@Ti_3C_2T_x$ exhibits higher reversible capacity and better cycling stability than  $MoS_2@Ti_3C_2T_x$ . Density-functional theory (DFT) calculation proves that the large interlayer distance enhances the lithium storage capacity.

# **Experimental**

#### Sample preparation

## Synthesis of Ti<sub>3</sub>AlC<sub>2</sub>

#### Preparation of  $Ti_3C_2T_x$

The  $Ti_3C_2T_x$  solution was prepared by selective etching of the  $Ti<sub>3</sub>AIC<sub>2</sub>$  powder. Typically, 1 g of LiF was immersed in 20 mL of 9 M HCl and stirred magnetically until LiF dissolved completely at 35 °C. One gram of  $Ti<sub>3</sub>AIC<sub>2</sub>$  was introduced and stirred vigorously at 35 °C for 12 h. Afterwards, the precipitate was rinsed with deionized water several times until the pH was approximately 6. The sediment was dispersed in 150 mL of DI water and sonicated for 60 min under flowing Ar. After centrifugation at 3500 rpm for 60 min, the dark green supernatant of  $Ti_3C_2T_x$  was collected.

#### Synthesis of MoS<sub>2</sub>/C

A total of 0.1 g of ammonium molybdate, 0.09 g of thiourea, and 0.41 g of polyvinyl pyrrolidone were dissolved in 30 mL of DI water under magnetic stirring and transferred to a Teflon-lined stainless steel autoclave (50 mL) which was heated to 190 °C for 18 h in an electric oven. After natural cooling to room temperature, the  $MoS<sub>2</sub>/C$  composite material was gathered by centrifugation, washed with DI water and anhydrous ethanol several times, and then freeze-dried for 12 h. The MoS<sub>2</sub>/C composite was then heat-treated at 500  $\degree$ C for 3 h under Ar. The  $MoS<sub>2</sub>$  composite was prepared by the same method but without adding PVP.

### Fabrication of the flexible  $MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> film$

The diluted aqueous dispersion (1 mg mL<sup>-1</sup>) of MoS<sub>2</sub>/C was prepared by sonication for 30 min. It was introduced to the  $Ti_3C_2T_x$  solution (0.04 g, 2 mg mL<sup>-1</sup>) and stirred magnetically. The total volume of the solution was 40 mL with the  $Ti_3C_2T_x:MoS_2/C$  ratio of 2:3. After vacuum filtration, the membrane was freeze-dried for 9 h. The flexible  $MoS<sub>2</sub>/$  $C@Ti_3C_2T_x$  film was gathered by peeling from the filter membrane and used directly as an anode for the LIBs. As a reference, the  $MoS_2@Ti_3C_2T_x$  film was also produced by the same procedure.

## Materials characterization

X-Ray diffraction was performed on the D2-Advance (Bruker) automated X-ray diffractometer with Cu  $K_{\alpha}$  radiation ( $\lambda = 1.5418$  Å), and the STA449F5 (NETZSCH) was employed in the thermogravimetric (TG) analysis from 20 to 800 °C at 10 °C min−<sup>1</sup> at an oxygen flow rate of 100 mL min<sup>-1</sup>. Field-emission scanning electron microscopy (FESEM, JEOL S-4800, 10 kV) and transmission electron microscopy (TEM, Tecnai G2 F20) were carried out to examine the structure of the  $MoS_2/C@Ti_3C_2T_x$  film. X-Ray photoelectron spectroscopy was characterized on the K-ALPHA 0.5 eV with monochromatic Al  $K_{\alpha}$  radiation, and Raman

scattering was investigated by the INVIA Raman microprobe (Renishaw Instruments) with a 532-nm laser source.

#### Electrochemical measurements

A lithium foil was utilized for the counter electrode in the 2032 type coin cell, and the electrolyte was  $1 \text{ M } \text{LiPF}_6$  in a mixture of ethylene carbonate (EC):diethylene carbonate (DEC) (1:1 by volume). Assembly was carried out in an argon-filled glovebox. Electrochemical impedance spectroscopy (EIS) was performed in the frequency range from 0.01 Hz to 100 kHz with a potentiostatic signal amplitude of 5 mV. Cyclic voltammetry (CV) was tested on a VMP3 electrochemical workstation at different scanning rates, and galvanostatic charge-discharge experiments were measured using the Neware battery testing system with the voltage range of  $0.01-3$  V (vs Li/Li<sup>+</sup>).

## Theoretical calculation

First-principles calculation was performed by the Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional within the general gradient approximation (GGA) using the VASP code and considering all the spin-polarized calculation. The geometric structure was relaxed until the Feynman force on each atom was less than 0.01 eV/Å. The energy was optimized until the convergence condition was less than  $10^{-5}$  eV. Considering the van der Waals interactions, the DFT-D2 approach was adopted in the dispersion correction. The  $3 \times 3$  and  $4 \times 4$  supercells of MoS<sub>2</sub>/C were chosen as the substrate for the bilayer  $MoS<sub>2</sub>$ , and the G sandwiched structure was adopted. The c axis vacuum was set as 15 Å to avoid interactions between slabs. The plane wave cutoff energy was 450 eV. A k-points mesh of  $3 \times 3 \times 1$  in the gamma-center sampling scheme was used for geometry optimization and electronic self-consistency.

# Results and discussion

The fabrication procedures of the flexible  $MoS_2/COTi_3C_2T_x$ film are illustrated in Fig. [1.](#page-3-0)  $Ti_3C_2T_x$  nanosheets are synthesized by selectively etching the Al layer from  $Ti<sub>3</sub>AIC<sub>2</sub>$ . Using ammonium molybdate as the molybdenum source, thiourea as the sulfur source, and PVP as the carbon source, carbonencapsulated nanosphere-assembled  $MoS<sub>2</sub>$  nanosheets are prepared by a hydrothermal reaction and annealing at 500  $\rm{^{\circ}C}$  under flowing Ar. The sandwiched structure of MoS<sub>2</sub>/  $C@Ti_3C_2T_x$  flexible film was synthesized by vacuum filtration and freeze-drying.

The homogeneous  $MoS_2/COTi_3C_2T_x$  solution is the key to the preparation of the flexible film. Figure S1a displays the Tyndall effect of the  $Ti_3C_2T_x$  solution, and after mixing the rice-candy-like  $MoS_2/C$  solution with the  $Ti_3C_2T_x$  solution

under vigorous stirring, a well-dispersed MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> solution is obtained as verified by the Tyndall effect as shown in Fig. S1b [\[30](#page-7-0)]. The X-ray diffraction (XRD) patterns of  $Ti<sub>3</sub>AIC<sub>2</sub>, Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>, MoS<sub>2</sub>/C, and MoS<sub>2</sub> in Fig. 2a and b reveal$  $Ti<sub>3</sub>AIC<sub>2</sub>, Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>, MoS<sub>2</sub>/C, and MoS<sub>2</sub> in Fig. 2a and b reveal$  $Ti<sub>3</sub>AIC<sub>2</sub>, Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>, MoS<sub>2</sub>/C, and MoS<sub>2</sub> in Fig. 2a and b reveal$ the presence of  $Ti<sub>3</sub>AIC<sub>2</sub>$ . After etching and delamination, the sharp diffraction peak at 9.5° disappears but a new peak arising from the (002) plane of  $Ti_3C_2T_x$  is observed at 7° further confirming successful delaminate of the  $Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  nanosheets [\[31](#page-8-0)]. The XRD pattern of  $MoS<sub>2</sub>$  in Fig. [2b](#page-3-0) can be indexed to the JCPDS card (37-1492) with a d (002) spacing of 0.615 nm. The sharp diffraction peak at  $14.378^\circ$  indicates that  $MoS<sub>2</sub>$  is made of a single or a few layers assembled by van der Waals attraction [[32\]](#page-8-0). Although the 14.375° diffraction peaks of MoS<sub>2</sub> cannot be observed from MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>, there are two diffraction peaks at 6.584° and 12.93° corresponding to interlayer distances of 13.41 Å and 6.84 Å, respectively, indicating that the large interlayer distance of 0.67 nm for the PVP-derived carbon layer is combined into the two  $M_0S_2$ interlayers. The calculation shows it is the middle spacing of the MoS<sub>2</sub> interlayer distance [[33\]](#page-8-0). The XRD result of MoS<sub>2</sub>/  $C@Ti_3C_2T_x$  is presented in Fig. S2. As displayed in Fig. S3, The  $E_{2g}^{1}$  (381.68 cm<sup>-1</sup>) and  $A_{1g}$  (404.78 cm<sup>-1</sup>) peaks correspond to the in-plane vibration of Mo and S atoms and the outof-plane vibration of S atoms, respectively. The separation ( $\Delta k$ ) of MoS<sub>2</sub>/C is 23.1 cm<sup>-1</sup>, which smaller than 27.4 cm<sup>-1</sup> of MoS<sub>2</sub> [\[33](#page-8-0), [34](#page-8-0)]. The  $\Delta k$  is positively correlated with the layer number, which increases with the layer number. From the Raman Spectroscopy, it can be concluded the  $M_0S_2/C$  not only has a less number of layers, but also has a large interlayer distance than  $MoS<sub>2</sub>$ . The content of PVP-derived carbon is calculated by Eq. (1) based on the thermogravimetric analysis (TGA) in an oxygen atmosphere:

$$
MoS2(wt\%) = \frac{a}{b} \times c \times 100\%
$$
 (1)

where *a* is the molecular weight of  $MoS<sub>2</sub>$ , *b* is the molecular weight of  $MoO<sub>3</sub>$ , and c is the total residual weight. According to the TG analysis (Fig. [2c\)](#page-3-0), the PVP-derived carbon concentration is 25.69%.

Figure [2d](#page-3-0) exhibits the presence of Ti, C, Mo, S, and O, and the elemental composition is listed in Table S1. Figure [2e](#page-3-0) can be deconvoluted into three peaks: 455.1 and 461.1 eV for Ti-C, 455.6 and 462 eV for  $Ti^{2+}$ , 456.4 eV for  $Ti^{3+}$ , 457.5 eV for Ti-O, 458.9 eV for Ti-F  $[35-38]$  $[35-38]$  $[35-38]$ . Figure [2f](#page-3-0) shows that the peaks at 281.9 eV, 284.6 eV, and 286.2 eV are related to C-Ti, C-C, and C-O, respectively [[39\]](#page-8-0), and Fig. [2g](#page-3-0) shows that the peaks at 529.6 eV, 531.8 eV, 532.2 eV, and 533.1 eV are related to O-Ti, O-Ti/OH, O-C/OH, and  $H<sub>2</sub>O$ , respectively [\[40](#page-8-0)]. As shown in Fig. [2h](#page-3-0) and i, the Mo 3d spectra of  $M_0S_2/$  $C@Ti_3C_2T_x$  exhibiting three main peaks at 229.4 eV, 232.6 eV, and 225.8 eV are assigned to Mo  $3d_{5/2}$ , Mo  $3d_{3/2}$ , and  $Mo^{4+}$ . The S 2p<sub>1/2</sub> and S 2p<sub>3/2</sub> peaks in the S 2p spectrum are at 162.4 and 161.3 eV, respectively [\[41](#page-8-0)].

<span id="page-3-0"></span>

Fig. 2 a XRD patterns of the Ti<sub>3</sub>AlC<sub>2</sub> and Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>. b XRD patterns of the MoS<sub>2</sub>/C and MoS<sub>2</sub>. c TGA curves of MoS<sub>2</sub>/C and MoS<sub>2</sub>. d XPS survey scan spectra. e Ti 2p,  $f C 1s$ ,  $g O 1s$ ,  $h Mo 3d$ , and  $i S 2p of MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> spectra$ 

The morphology of the Ti<sub>3</sub>AlC<sub>2</sub>, Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>, MoS<sub>2</sub>/  $C@Ti_3C_2T_x$ , and  $MoS_2@Ti_3C_2T_x$  is observed by SEM and TEM. The SEM image shows that the pristine  $Ti<sub>3</sub>AIC<sub>2</sub>$  powder has a slightly layered texture but it appears to be seamless (Fig. S4a). The flake-shape folded structure in Fig. S4b corroborates successful synthesis of  $Ti_3C_2T_x$  nanosheets. After vacuum filtration of the  $Ti_3C_2T_x$  suspension, a flexible film is produced as shown in Fig. S4c. The TEM image of  $Ti_3C_2T_x$ in Fig. S4d reveals the flexible characteristic, and Fig. 3a shows that the MoS<sub>2</sub>/C nanosheets and Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> are welldispersed together with the sandwiched structure of  $Ti_3C_2T_x$ and  $MoS<sub>2</sub>/C$ . The inset in Fig. 3a is the digital photograph of the  $MoS_2/COTi_3C_2T_x$  film and the magnified image is in Fig.  $3b$ . The MoS<sub>2</sub> nanosphere/carbon hybrid has a nanosheet morphology resembling rice candies (Fig. 3c).

The high-resolution TEM (HR-TEM) image in Fig. 3d shows that the  $MoS<sub>2</sub>/C$  nanospheres consist of layered nanoflakes. The interlayer spacing of the  $M_0S_2/C$  (002) plane is 13.41 Å, which is much larger than that of the  $MoS<sub>2</sub>$  (002) plane of 6.2 Å as shown in Fig. 3e and f. This result is in accordance with the above XRD calculations. Figure S5a and  $5b$  show that  $MoS<sub>2</sub>$  are single nanospheres. Figure 3g shows dense and sparse contrast difference suggesting some overlapping of Ti, C, Mo, and S. Furthermore, in the  $MoS<sub>2</sub>/$  $C@Ti_3C_2T_x$  structure, Mo, S, and C are concentrated in the  $MoS_2/C$  nanoflakes and Ti overlaps along the folded  $Ti_3C_2T_x$ ,

providing evidence that the  $MoS<sub>2</sub>/C$  nanoflakes are mixed uniformly with the flake-shaped folded  $Ti_3C_2T_x$  nanosheets. It can be concluded that PVP-derived carbon is intercalated into  $MoS<sub>2</sub>$  to expand the interlayer distance. The abundant active sites and rapid Li<sup>+</sup> diffusion arise from the larger interlayer distance and PVP-derived carbon.

The lithium storage property of the flexible  $M_0S_2/$  $C@Ti_3C_2T_x$  is determined from lithium-ion half-cells. Figure [4a](#page-5-0) presents the first three cycles of the cyclic voltammetry (CV) curves of the  $MoS_2/C@Ti_3C_2T_x$  flexible electrode at  $0.1 \text{ mV s}^{-1}$ . In the first cathodic scan, three reduction peaks at 0.8 V, 0.5 V, and 0.2 V can be seen. The first one (0.8 V) represents intercalation of  $Li^+$  insertion into  $MoS<sub>2</sub>$  to form  $Li<sub>x</sub>MoS<sub>2</sub>$  [[42](#page-8-0)], and the other two peaks (0.5 and 0.2 V) are attributed to reduction of  $Li_xMoS_2$  to metallic Mo and  $Li_2S$  by the conversion reaction and creation of a solid electrolyte interphase (SEI) layer on the electrode [\[43](#page-8-0)]. In addition, anodic peaks appear at 1.3 V and 2.3 V corresponding to partial oxidation of Mo to  $MoS<sub>2</sub>$  and  $Li<sub>2</sub>S$  to S, respectively. In the subsequent cathodic scan, the peak at 1.2 V is related to the conversion of  $MoS<sub>2</sub>$  into Mo and that at 1.9 V is related to the formation of  $Li<sub>2</sub>S$  [[44](#page-8-0)]. As shown in Fig. [4b,](#page-5-0) the galvanostatic discharge/charge curves of the flexible  $M_0S_2/C@Ti_3C_2T_x$ electrode at 0.05 A  $g^{-1}$  over the range between 0.01 V and 3.0 V deliver initial discharge and charge capacities of 733 and 597 mAh  $g^{-1}$ , respectively, in conjunction with an initial



Fig. 3 a, b SEM images of the  $MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  at different magnifications. c The digital photographs of crunchy-rice candy, d TEM images, and e HRTEM image of  $MoS_2/COTi_3C_2T_x$ . **f** HRTEM

image of  $MoS_2@Ti_3C_2T_x$ . g Corresponding mapping images for Ti, C, Mo, and S elements of  $MoS_2/C@Ti_3C_2T_x$ 

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Fig. 4 a CV curves of MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> at 0.1 mV s<sup>-1</sup>. b Galvanostatic discharge/charge curves of the  $MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  electrode at 50 mA  $g^{-1}$ . c Rate capability of the MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>, MoS<sub>2</sub>/Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>, MoS<sub>2</sub>/ C,  $MoS<sub>2</sub>$ , and  $Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  electrode at various current densities. **d** Long-term cycling performance of the MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> and MoS<sub>2</sub>@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> at

0.1 A  $g^{-1}$  and 2 A  $g^{-1}$ . e Nyquist plots of the MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> electrode. f b values determined by the plot of log (i) versus log (v). g Normalized ratio of pseudocapacitive and diffusion-controlled contribution of the  $M_0S_2/C@Ti_3C_2T_x$  at different sweep rates. **h** Capacitive contribution at  $1.0 \text{ mV s}^{-1}$ 

coulombic efficiency of 81.4%. The irreversible capacity loss in the first cycle can be attributed to the formation of the irreversible SEI film. In subsequent cycles, the discharge/ charge capacity is quite similar, illustrating the high reversibility in the electrochemical reactions.

A comparison of the rate and cycling performance of the  $MoS_2/COTi_3C_2T_x$ ,  $MoS_2@Ti_3C_2T_x$ ,  $MoS_2/C$ ,  $MoS_2$ , and  $Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  samples at different various current densities is shown in Fig. 4c. As expected, the flexible  $M_0S_2/C@Ti_3C_2T_x$  electrode exhibits excellent rate capabilities of 538.5, 510.5, 461.8, 405.6, 337.7, and 256.7 mAh  $g^{-1}$  at 0.05, 0.1, 0.2, 0.5, 1, and 2 A  $g^{-1}$ , respectively. When the current density rises again to 0.05 A  $g^{-1}$ , the discharge capacity of 538 mAh g−<sup>1</sup> can be restored rapidly implying remarkable reversibility. In addition, the long cycling stability of the flexible  $MoS_2/C@Ti_3C_2T_x$  electrode is assessed. Figure 4d shows the cycling performances at current densities of 0.1 A  $g^{-1}$  and 2 A  $g^{-1}$ ; the retain capacities of 489.3 and 248.4 mA h  $g^{-1}$  are observed after 150 cycles, respectively. The excellent rate performance of  $MoS_2/C@Ti_3C_2T_x$  mainly ascribed to PVP-derived carbon is intercalated into  $MoS<sub>2</sub>$  to expand the interlayer distance, and dynamically facilitates electron/ion transport. Apparently, the result indicates that the ultra-wide interlayer spacing  $MoS<sub>2</sub>$  and  $Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  flexible film could dramatically enhance lithium storage capacity as expected. Compared with other 2D materials for LIBs, the as-prepared  $M_0S_2/C@Ti_3C_2T_x$  shows a high rate performance (Table S2). Figure S6 depicts the cross-sectional SEM image of  $MoS<sub>2</sub>/$  $C@Ti_3C_2T_x$  after 150 cycles at 0.1 A g<sup>-1</sup>, and MoS<sub>2</sub>/  $C@Ti_3C_2T_x$  retains the sandwiched morphology confirming the good structural stability. Therefore, the superior rate performance of  $MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  is mainly attributed to the larger interlayer distance of MoS<sub>2</sub> and the high electrical conductivity of  $Ti_3C_2T_x$ .

To further confirm the impact of PVP-derived carbon, the Nyquist plots of the flexible electrodes are shown in Fig. 4e. It has been shown that the diameter of the semicircle is related to the charge transfer resistance  $(R_{\rm ct})$  and the slope in the low-frequency region is related to the Warburg impedance  $(Z_w)$ . In the equivalent circuit,  $R_s$  corresponds to the contact resistance,  $R_{\rm ct}$  is the charge transfer resistance, CPE corresponds to the constant phase element, and  $Z_w$  is the Warburg impedance (insets in Fig. 4e and Fig. S7). The flexible  $MoS_2/C@Ti_3C_2T_x$  electrode has a smaller  $R_{\text{ct}}$  (38.9  $\Omega$ ) than MoS<sub>2</sub>@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> (74.65  $\Omega$ ),

suggesting that addition of PVP-derived carbon enhances the charge transfer ability  $[45]$  $[45]$ . In addition, the large slope indicates large lithium-ion diffusion rates as a result of the larger interlayer distance.

To detect the electrochemical kinetics of the flexible  $MoS<sub>2</sub>/$  $C@Ti_3C_2T_x$  electrode, CV is performed at scanning rates from 0.1 to 1 mV  $s^{-1}$  to determine the redox pseudocapacitance contributions (Fig. S8). The relationship between the peak current (i) and scanning rate (v) follows the following power law [[46\]](#page-8-0):

$$
i = av^b \tag{2}
$$

where  $b = 0.5$  or 1 corresponds to the diffusion-controlled or behavior pseudocapacitive effect, respectively. Figure [4f](#page-5-0) shows that the b values of the anodic and cathodic peaks are 0.80 and 0.87, respectively, suggesting the surface faradaic redox reaction is predominant in the electrochemical reaction. The ratio of the capacitive contribution is further quantified by Eq. (3) [\[47\]](#page-8-0):

$$
i = k_1 v + k_2 v^{1/2}
$$
 (3)

where  $k_1v$  is the non-diffusion contribution and  $k_2v^{1/2}$  is the diffusion-controlled contribution [\[48\]](#page-8-0). As observed in Fig. [4g,](#page-5-0) with increasing scanning rates, the charge storage rate increases in the surface-controlled process. The pseudocapacitive contribution exhibits an increasing tendency with scanning rates and is 81% for a scanning rate of 1 mV s<sup>-1</sup> (Fig. [4h\)](#page-5-0). The high pseudocapacitive contribution of the MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> composite mainly originates from the PVP-derived carbon, which can provide more active redox sites for the pseudocapacitive behavior [[49](#page-8-0)]. The results indicate that a large high pseudocapacitive contribution is crucial to the electron transport kinetics.



Fig. 5 Lithium adsorption in a  $MoS<sub>2</sub>/MoS<sub>2</sub>$  and b  $MoS<sub>2</sub>/C/MoS<sub>2</sub>$ interlayer. c Charge density of Li adsorption in  $MoS<sub>2</sub>/C$  interface. Yellow and turquoise regions indicate the accumulation and depletion, respectively

Since the PVP-derived carbon increases the interlayer distance of  $MoS<sub>2</sub>$ , the Li<sup>+</sup> storage mechanism is determined by DFT calculation of MoS<sub>2</sub>/C [\[50,](#page-8-0) [51](#page-8-0)] in which the carbon layer is graphene. As shown in Fig.  $5a$  and b, MoS<sub>2</sub>/C exhibits stronger chemical adsorption with lower binding energies  $(E_a = -0.245 \text{ eV})$  than MoS<sub>2</sub> which shows weaker physical adsorption with a larger binding energy ( $E_a = 1.727$  eV) (Table S3). Therefore, the PVP-derived carbon enhances adsorption of  $Li<sup>+</sup>$  due to the larger negative adsorption energy and better charge transfer [\[34](#page-8-0), [52](#page-8-0)]. To further study the electronic structure of  $MoS<sub>2</sub>/C$ , the charge density is analyzed as shown in Fig. 5c. Charge exchange occurs between the PVPderived carbon and  $MoS<sub>2</sub>$ , and charges are mainly lost from Li to accumulate in the Li-S and Li-C bonds. In this way, the intercalated PVP-derived carbon (graphene in this model) creates a stable channel for Li<sup>+</sup> transport consequently improving charge transport.

# Conclusions

The flexible sandwiched structure consisting of rice-candylike MoS<sub>2</sub>/C nanosheets intercalated in Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub> is designed and demonstrated as high-performance electrode materials in Li-ion batteries. The large interlayer spacing provides abundant and rapid diffusion channels for the electrolyte and there are more adsorption sites to enhance the electrochemical characteristics. The PVP-derived carbon layer enhances the electrical conductivity and structural robustness. As a demonstration of the commercial viability in Li-ion batteries, the flexible  $MoS<sub>2</sub>/C@Ti<sub>3</sub>C<sub>2</sub>T<sub>x</sub>$  anode exhibits remarkable reversible capacities and coulombic efficiency, improved rate performance, and outstanding cycling stability. DFT calculation reveals that lithium storage on the  $MoS<sub>2</sub>/C$  nanosheets leads to more intercalation to enhance adsorption/diffusion of lithium ions. The flexible  $MoS_2/COTi_3C_2T_x$  electrode has large commercial potential, and the design and fabrication strategy provide insights into the development of superior electrochemical properties for flexible energy storage equipment.

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#### <span id="page-7-0"></span>**Declarations**

Conflict of interest The authors declare no competing interests.

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