ORIGINAL PAPER

Lithium ion conductivity of solid solutions based on $Li₈ZrO₆$

Mariya S. Shchelkanova $1\text{m} \cdot \text{Geor}$ gi Sh. Shekhtman $1 \cdot$ Anastasia V. Kalashnova $1 \cdot$ Olga G. Reznitskikh 2

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Abstract

Lithium ion conductivity of lithium hexaoxozirconate Li₈ZrO₆ doped by Mg²⁺, Sr²⁺, Nb⁵⁺, V⁵⁺, and Ce⁴⁺ cations was studied using impedance spectroscopy. The NMR data indicate that in the low-temperature region, lithium ion migration takes place by exchange of Li⁺ among tetra- and octahedral positions. Data of Raman spectroscopy suggest that the potential reason for the sharp increase of Li_8ZrO_6 and its solid solutions' conductivity near 700 K may be melting of LiOH which occurs as a result of interaction between samples and atmospheric moisture.

Keywords Solid electrolytes . Lithium ion conductivity . Lithium hexaoxozirconate . Lithium batteries . Raman spectroscopy . Nuclear magnetic resonance

Introduction

Among the electrochemical power sources, one of the leading positions belongs to the lithium-ion batteries (LIBs). Compared with most other kinds of batteries, LIBs have a higher output voltage and higher energy storage density. Of special interest is development of all-solid-state lithium-ion batteries using solid lithium-conducting electrolytes instead of traditional liquid organic or polymer electrolytes. The major advantages of allsolid-state lithium-ion batteries are simple design and, the main thing, safety. According to several works [[1\]](#page-5-0), promising solid electrolytes for medium- and low-temperature all-solid-state lithium-ion batteries are LLZ (Li₇La₃Zr₂O₁₂, $\sigma_{300 K}$ = 2.1 × 10^{-4} S × cm⁻¹) [\[2\]](#page-5-0), LISICON (Li₁₄ZnGe₄O₁₆, σ_{300} _K =

Highlights Li₈ZrO₆ doped by M^{2+} , A^{5+} , and Ce⁴⁺ was studied using impedance, Raman study, and NMR. The maximum conductivity $(9 \times$ 10^{-6} S cm⁻¹ at 476 K) is observed for Li_{7.98}Zr_{0.98}V_{0.02}O₆. Below 700 K, Li⁺ ions move via jumps between octahedral and tetrahedral positions. In all cases, conductivity shows a jump-like rise at 680–710 K. One of the reasons for a jump of conductivity near 700 K may be the melting of LiOH.

 \boxtimes Mariya S. Shchelkanova shchelkanova.mariya@mail.ru

¹ Institute of High-Temperature Electrochemistry, Ural Branch of RAS, Akademicheskaya St., 20, Ekaterinburg 620137, Russian Federation

² Institute of Solid State Chemistry, Ural Branch of RAS, GSP, Pervomayskaya St., 91, Yekaterinburg 620990, Russian Federation

 10^{-7} S × cm⁻¹) [[3\]](#page-5-0), NASICON-like Li_{1+x}Al_xTi_{2-x}(PO₄)₃ $(\sigma_{300 \text{ K}} = 7.0 \times 10^{-4} \text{ S} \times \text{cm}^{-1})$ [[4\]](#page-5-0), and glass-ceramic Li_{1.6}Al_{0.6}Ge_{1.4}(PO₄)₃ ($\sigma_{300 \text{ K}} = 4.0 \times 10^{-4} \text{ S} \times \text{cm}^{-1}$) [[5](#page-5-0)]. Unfortunately, most of the above materials have high $Li⁺$ ion conductivity but are unstable in contact with metal lithium. Binary compounds in the system $Li₂O-ZrO₂$ are free from this disadvantage [\[6,](#page-5-0) [7\]](#page-5-0) but their conductivity is too small. Therefore, investigations focused on raising the ionic conductivity of lithium zirconates are of great importance.

Among lithium zirconates, the maximum value of conductivity (10^{-4} – 10^{-5} S × cm⁻¹ at 570 K) is observed for Li₈ZrO₆ [\[8](#page-5-0), [9](#page-5-0)].

The first crystal structure and transport properties of Li_8ZrO_6 and Li_7LO_6 (L = Nb, Ta), which have close resemblance, were studied in work [[8\]](#page-5-0). All of these compounds have layered structure derivative of CdI₂. Crystal lattice of both Li_8ZrO_6 and Li_7LO_6 consists of layers containing ions of Zr^{4+} (L^{5+}) and 2 of 8 Li^+ cations in octahedral positions, the remaining 6 Li⁺ ions occupy half of tetrahedral sites between layers. Composition of layers may be described as $[(Li₂Zr)O₆]$ _n in case of Li₈ZrO₆ or as $[(Li_□L)O₆]$ _n in case of $Li₇LO₆$. As one can see, $Li₇LO₆$ compounds contain structural lithium vacancies \Box) in octahedral positions; in Li₈ZrO₆ such vacancies are absent. Due to the presence of the lithium vacancies, lithium ion conductivity of $Li₇LO₆$ is considerably higher as compared to $Li₈ZrO₆$, which has ordered crystal structure. Therefore, one can expect that the introduction of additives creating vacancies in the lithium sublattice will increase Li^+ ion conductivity of Li_8ZrO_6 . Actually, the increase of conductivity is observed at substitution of Li^+ by Mg^{2+} or

 Sr^{2+} according to Eq. 1 [\[10\]](#page-5-0) and at substitution of Zr^{4+} by Nb^{5+} or V^{5+} according to Eq. 2 [\[11,](#page-5-0) [12\]](#page-5-0).

$$
2Li^{+}\rightarrow Mg(Sr)^{2+} + V_{Li}
$$
\n⁽¹⁾

$$
Li^{+} + Zr^{4+} \rightarrow Nb(V)^{5+} + V_{Li}
$$
 (2)

Introduction of additives of another type such as substitution of Zr^{4+} by triple-charged cations Al^{3+} , Y^{3+} , and Sc^{3+} with charge compensation by interstitial Li⁺ ions is not followed by conductivity rise [\[13](#page-5-0)–[15\]](#page-5-0). It is of interest that a considerable increase of conductivity takes place at isovalent substitution of Zr^{4+} by Ce⁴⁺ [\[16](#page-5-0)].

In this paper, the values of conductivity in the solid solutions $Li_{8-2x}Me_xZrO_6$ (Me = Mg, Sr), $Li_{8-x}Zr_{1-x}A_xO_6$ $(A = V, Nb)$, and $Li_8Zr_1 - xCe_xO_6$ are compared, and temperature dependences of conductivity of these solid solutions are analyzed. A number of samples were examined by Raman spectroscopy and nuclear magnetic resonance (NMR) spectroscopy. Based on the data obtained, conclusions about factors determining high conductivity of the materials under consideration are made.

Experimental

According to work [[11\]](#page-5-0), the range of solid solutions based on Li₈ZrO₆ in the system Li_{8 – x}Zr_{1 – x}V_xO₆ lies at $x < 0.05$. The data of the present study show that single-phase region boundary in this system corresponds to $x \approx 0.02$. Furthermore, $Li_{7.98}Zr_{0.98}V_{0.02}O_6$ solid solution has the highest conductivity in $Li_{8-x}Zr_1-xV_xO_6$ system. So, to compare the values of conductivity in other systems with heterovalent substitutions, we chose compositions with the same value $x = 0.02$, that is, with the equal concentration of lithium vacancies: $Li_{7.96}\Box_{0.02}Me_{0.02}ZrO_6$ (Me = Mg, Sr) and $Li_{7.98}\Box_{0.02}Zr_{0.98}A_{0.02}O_6$ (A = V, Nb). All these compositions lie in the single-phase regions of the corresponding systems and have conductivity close to maximal value. A sample containing niobium was synthesized and studied in this work. The data on conductivity of Mg- and Srcontaining solid solutions are taken from several works [\[17,](#page-5-0) [18](#page-5-0)]. In the system $Li_8Zr_1 - xCe_xO_6$ as well as in the systems with heterovalent dopants, a sample with $x = 0.02$ was studied.

 $Li₂CO₃$ (high purity) and $V₂O₅$, $Nb₂O₅$, $CeO₂$, and $ZrO₂$ (analytical grade) were used as starting materials for the solidstate synthesis of $Li_8Zr_{0.98}Ce_{0.02}O_6$, $L_{7.98}Zr_{0.98}Nb_{0.02}O_6$, and $Li_{7.98}Zr_{0.98}V_{0.02}O_6$. Zirconium, niobium, and cerium oxides were pre-dried at 1270 K to a constant weight. The starting materials were mixed in a stoichiometric ratio, except for $Li₂CO₃$, which was added with 5–8 wt% surplus. The syntheses were carried out in glass carbon crucibles under vacuum in

the hermetic heatproof steel reactor. The synthesis was carried out by stepwise increasing the temperature from 573 to 1253 K. Then, the samples were annealed in the air at 873 K for 6 h to remove traces of carbon.

The crystal structure of the products obtained was determined by the X-ray diffraction (XRD) after each stage of the synthesis. The XRD analysis was carried out using Rigaku D-MAX-2200 V diffractometer in CuKα-radiation, $2\theta = 10^{\circ} - 80^{\circ}$. XRD patterns of compounds were compared with database PDF-2 [\[19](#page-5-0)] and literature data.

Chemical analysis of synthesized samples was performed using inductively coupled plasma atomic emission spectrometer OPTIMA 4300 DV (Perkin Elmer, USA).

Raman spectra of the polycrystalline samples were collected with Raman microscope-spectrometer RENISHAW U-1000 (Ar⁺-radiation, $\lambda = 514$ nm) at 50–4000 cm⁻¹.

Thermal analysis of the materials obtained was performed using synchronous thermal analyzer STA 449 F1 Jupiter over the temperature range 300–870 K in alundum crucibles at the rate of heating 10 K min⁻¹ in the air. The mass of the samples was 10–20 mg. The data obtained were processed by means of NETZSCH Proteus software.

NMR experiments were carried out using an updated Bruker SXP 4-100 pulsed NMR spectrometer at resonance frequency 194.37 MHz in the temperature range 300–670 K. NMR spectra and the spin-lattice relaxation time of $\mathrm{^{7}Li}$ nuclei were measured by the Fourier transformation of the signal of the free induction decay.

The samples for conductivity measuring with diameter 6 mm and thickness 1–2 mm were obtained by pressing the powders of particle size < 50 μm in the steel die at 70 MPa. Then the samples were covered by the powder of the same composition and sintered under helium in the heatproof steel reactor for 5 h at 1179 К. The density of the samples calculated from the geometrical parameters and mass was 90–95% of the theoretical one.

For AC impedance studies, both sides of the sintered samples were painted with silver paste. The impedance measurements were performed in the air at temperatures ranging from 25 to 600 °C, using an immittance meter E7-25 (MNIPI) in the frequency range from 20 Hz to 1 MHz. The results were fitted by using special software program (EQUIVCRT.COM). All of the conductivity measurements were performed on two sets of identical samples to check the reproducibility of the results.

Electron conductivity was measured at direct current using blocking Ag [\[15](#page-5-0)] electrodes in the temperature range 470 to 870 K. For this purpose, direct voltage which was less than solid electrolyte decomposition voltage (20–40 mV) was applied to Ag| solid electrolyte | Ag cell. After a time, the current acquired a fixed value. Under these conditions, ionic current is blocked and the charge is carried by electrons and holes. Knowing the current value and applied voltage, we can calculate the electron conductivity.

Results and discussion

Ionic conductivity and DSC

Figure 1 shows XRD patterns of the samples studied. It is seen that in all cases, diffraction diagram corresponds to the $Li₈ZrO₆$ phase (PDF2 No. 26-0867 card-index data).

Temperature dependences of conductivity for $Li₈ZrO₆$ and solid solutions are similar (Fig. 2). Low- and high-temperature linear regions are separated at 680–710 K by a transition section, which is characterized by a sharp increase in conductivity with increasing temperature. Electron conductivity value of both undoped Li_8ZrO_6 and solid solutions in all cases was by three orders of magnitude less than ionic one. Therefore, further on when discussing the synthesized material conductivity, we mean their ionic conductivity. DSC measurements of Li_8ZrO_6 show the existence of endothermic peak at 680 K (Fig. 3). The enthalpy of effect is -23.6 J g⁻¹. DSC curves of the doped samples also show endothermic peaks in the range 680–700 K. In all cases, the temperature of thermal effect is close to the temperature of the drastic change in conductivity. It gives reason to suggest that at 680–710 K, the materials examined undergo a phase transition followed by lithium sublattice disorder [[10](#page-5-0)]. However, there is another probable interpretation of Li_8ZrO_6 and its solid solutions' conductivity behavior. The fact is that binary compounds of the $Li₂O M_xO_y$ systems with the high content of lithium oxide, such as Li₅MO₄ (M = Al, Ga [\[20](#page-5-0)], Fe [[21](#page-5-0)], Tl [[22](#page-5-0)]), Li₆ZnO₄ [[20\]](#page-5-0), $Li₇LO₆ (L = Bi, Sb, Ta)$ [\[23\]](#page-5-0), and also $Li₈ZrO₆$ [[24](#page-5-0)], are highly hygroscopic and interact with ambient air moisture. For the first time, influence of this effect on conductivity was found for $Li₅AIO₄$. The first conductivity measurements of $Li₅AIO₄$ as well as of Li_5GaO_4 and Li_6ZnO_4 showed [\[20\]](#page-5-0) that at

Fig. 1 XRD patterns of (1) Li_8ZrO_6 PDF2 # 26-0867 and (2) undoped Li_8ZrO_6 phase and (3) $Li_{7.98}Zr_{0.98}V_{0.02}O_6$, (4) $Li_{7.98}Zr_{0.98}Nb_{0.02}O_6$, (5) $Li_8Zr_{0.98}Ce_{0.02}O_6$, (6) $Li_{7.96}Sr_{0.02}ZrO_6$, and (7) $Li_{7.96}Mg_{0.02}ZrO_6$ solid solutions

Fig. 2 Arrhenius plots for (1) Li_8ZrO_6 and (2) $Li_8Zr_{0.98}Ce_{0.02}O_6$, and (3) $Li_{7.96}Sr_{0.02}ZrO_6$, (4) $Li_{7.98}Zr_{0.98}Nb_{0.02}O_6$, (5) $Li_{7.96}Mg_{0.02}ZrO_6$, and (6) $Li_{7.98}Zr_{0.98}V_{0.02}O_6$ solid solutions

heating up to ~ 680 K, the conductivity of these materials increases jump-like achieving 10^{-1} S × cm⁻¹ but activation energy decreases sharply. At first, such a behavior as in the case of Li_8ZrO_6 was explained by lithium sublattice disorder [\[23](#page-5-0)]. However, later, it was shown through the example of $Li₅AIO₄$ that the jump-like increase of conductivity in this case is due to the melting of LiOH, which appears as a result of interaction between $Li₅AIO₄$ and atmospheric moisture [\[25,](#page-5-0) [26\]](#page-5-0). If one compares temperature dependences of conductivity for $Li₅AIO₄$, $Li₅FeO₄$, $Li₅TiO₄$, $Li₈ZrO₆$, and its solid solutions (Fig. 2) and $Li₇LO₆$ (L = Bi, Sb, Ta) [\[23\]](#page-5-0), it can readily be seen that in all cases, $lg\sigma - 1/T$ curves demonstrate the same peculiarities. First, one observes a sharp increase in conductivity and decrease in activation energy, and this at similar temperatures which are close to the melting point of LiOH (680–710 K). Second, in all cases, DTA or DSC measurements indicate an endothermic peak in the same temperature region. Third, the presence of a jump in conductivity and its temperature does not depend on introduction of additives,

Fig. 3 DSC plot of Li_8ZrO_6 : heating rate 10 K min⁻¹

their nature, and concentration. The assumption about the absence of phase transition in $Li₈ZrO₆$ near 700 K is supported by the fact that high-temperature XRD study does not show a change of the crystal structure, which was to be expected at disorder of cation sublattice. The X-ray pattern obtained at 893 K [\[27](#page-5-0)] is identical to that obtained at room temperature [\[28\]](#page-5-0).

Raman spectroscopy

Table 1 shows the observed bands of the Raman spectra for $Li₈ZrO₆$ and some of the solid solutions synthesized in this work as compared to the literature data. One can see that the spectra of Li_8ZrO_6 and its solid solutions contain a number of extra bands, which were not observed for $Li₈ZrO₆$ in work [\[23\]](#page-5-0). However, most of these extra bands may be referred to $Li₂CO₃$, LiOH, and LiOH·H₂O [\[29](#page-5-0)–[31\]](#page-5-0). The formation of LiOH as a result of interaction between $Li₈ZrO₆$ and water vapor according to the reaction

$$
Li_8ZrO_6 + 2H_2O \rightarrow Li_4ZrO_4 + 4LiOH \tag{3}
$$

was experimentally proved in work [\[24](#page-5-0)]. It seems that the formation of LiOH·H₂O and Li₂CO₃ is a result of the further interaction of LiOH and atmosphere moisture and $CO₂$. Therefore, it is reasonable to assume that the anomalous behavior of conductivity of Li_8ZrO_6 and its solid solutions may be due to the formation of LiOH, which melts at elevated temperatures as in the case of $Li₅AlO₄$. As concerns $Li₇LO₆$ $(L = Nb, Ta, Sb, Bi)$, the authors $[23]$ $[23]$ $[23]$ do not find in the samples studied any interaction with atmosphere products, but do not register any change in crystal lattice at temperatures above 700 K either. Therefore, the point of sharp increase in conductivity of the materials under discussion when heated to \sim 700 K remains unclear.

NMR spectroscopy

⁷Li NMR experiments were performed for $Li_{7.98}Zr_{0.98}V_{0.02}O_6$ solid solution in the temperature range 300–750 K (Fig. 4).

Fig. 4 7 Li NMR spectra as a function of temperature for $Li_{7.98}Zr_{0.98}V_{0.02}O_6$

Modeling of the spectrum shape shows that spectral line can be described as superposition of two lines which correspond to different structural positions of lithium Li1 and Li2. Such positions in solid solutions with $Li₈ZrO₆$ structure are represented by octahedral positions in layers and tetrahedral ones between layers. ⁷Li NMR spectra for undoped Li_8ZrO_6 are similar [[9](#page-5-0)].

Figure [5](#page-4-0) shows ⁷Li NMR line width of the central transition as function of temperature. As one can see, at temperature increase, the width of ${}^{7}Li$ NMR line is gradually reduced. Such behavior of NMR spectra is typical for systems with fast ion transport. Sharp narrowing of NMR line takes place when characteristic frequency of diffusion τ_d^{-1} approaches the width of the NMR line of rigid lattice. For $Li_{7.98}Zr_{0.98}V_{0.02}O_6$ solid solution, this temperature area is \sim 450–600 K (Fig. [5\)](#page-4-0). Earlier, close temperature interval of NMR line narrowing was observed for undoped Li_8ZrO_6 [\[9](#page-5-0)]. It should be pointed out that for $Li₇LO₆$, which has a similar crystal structure but contains

temperature for $Li_{7.98}Zr_{0.98}V_{0.02}O_6$

structural lithium vacancies close to \sqrt{L} LiNMR, the line is ob-served at essentially lower temperatures 200–400 K [\[23](#page-5-0)], which correlates with higher Li^+ conductivity of $Li₇LO₆$ as compared to $Li₈ZrO₆$.

Activation energy for jumps of Li⁺ ions in solid solution $Li_{7.98}Zr_{0.98}V_{0.02}O_6$ calculated from NMR data at low temperatures is equal to 44 kJ × mol⁻¹, whereas the value obtained from the temperature dependence of conductivity is more than twice as large (99 kJ × mol⁻¹). It suggests that the total mobility of lithium is to a great degree due to the local jumps of Li⁺.

The results of the transport property study of materials under discussion are summarized in Table 2. One can see that introduction of dopants in all cases results in conductivity increase and activation energy lowering as compared to undoped Li_8ZrO_6 . In case of systems with two- and fifthcharged dopants, it is due to formation of lithium vacancies. Unlike $Li₇LO₆$, charge carriers in $Li₈ZrO₆$ are thermal

Table 2 Ionic conductivities and the activation energies for $Li₈ZrO₆$ and solid solutions $Li_{7.98}Zr_{0.98}A_{0.02}O_6$ (A = V, Nb), $Li_{7.96}Me_{0.02}ZrO_6$ (Me = Mg, Sr), and $Li_8Zr_{1-x}Ce_xO_6$

Compound	σ (S \times cm ⁻¹)		E_a (kJ \times mol ⁻¹)	
	833 K	454 K	Temperature region	
			High	Low
Li_8ZrO_6	8.7×10^{-3}	4.8×10^{-8}	76	99
$Li_{7.98}Zr_{0.98}Nb_{0.02}O_6$	1.0×10^{-1}	1.5×10^{-6}	53	80
$Li_{7.98}Zr_{0.98}V_{0.02}O_6$	2.5×10^{-1}	4.6×10^{-6}	50	75
$Li_8Zr_{0.98}Ce_{0.02}O_6$	1.3×10^{-2}	5.2×10^{-7}	74	88
$Li_{7.96}Mg_{0.02}ZrO_6$	2.3×10^{-1}	1.2×10^{-6}	65	85
$Li_{7.96}Sr_{0.02}ZrO_6$	6.3×10^{-2}	1.2×10^{-7}	70	95

activated defects rather than structural vacancies. At substitution of Li^+ or Zr^{4+} ions by M^{2+} or A^{5+} ions respectively, the concentration of lithium vacancies drastically increases but activation energy decreases by a value of defect formation. Both factors lead to the increase of conductivity. The growth of conductivity at isovalent substitution of Zr^{4+} by Ce^{4+} seems to be a result of geometrical changes. As ion radius of $Ce⁴⁺$ is essentially larger compared to Zr^{4+} (0.87 and 0.76 Å respec-tively) [[32](#page-5-0)], the substitution of Ce^{4+} for Zr^{4+} should lead to broadening of lithium migration channels and lowering of the average value of potential barrier which separates an occupied $Li⁺$ position in the crystal lattice from a vacant one.

It should be noted, however, that conductivity increase of $Li₈ZrO₆$ at introduction of dopants in low-temperature region in all cases is not very large. Maximal value of conductivity is observed for $Li_{7.98}Zr_{0.98}V_{0.02}O_6$ and is equal to $8.9 \cdot 10^{-6}$ S × cm⁻¹ at 476 K (Table 2). That is only six times as large as compared to undoped $Li₈ZrO₆$. Rather low value of conductivity is the result of low solubility of the dopants and, as a consequence, low concentration of lithium vacancies. In $Li₇LO₆$, half of the octahedral $Li⁺$ positions are vacant, so its lithium ion conductivity is essentially higher and for Li₇BiO₆ achieves 3×10^{-3} S \times cm⁻¹ at 476 K [\[23\]](#page-5-0).

In the high-temperature region, the conductivity of Li_8ZrO_6 and solid solutions based on it drastically increases and at 870 K achieves 10^{-1} S \times cm⁻¹ (Table 2). However, the reasons for such a behavior are not well established. A possible solution to this problem could be a precise neutron diffraction study of crystal structure in the temperature region of existence of supposed disordered form of $Li₈ZrO₆$ or an investigation into the dependence of conductivity as function of time at \sim 800 K at a change of dry atmosphere into moisture and vice versa as it was done in [\[33\]](#page-5-0).

Conclusion

Temperature dependences of conductivity and Raman spectra of solid solutions based on lithium hexaoxozirconate $Li_{7.96}Me_{0.02}ZrO_6$ (Me = Mg, Sr), $Li_{7.98}Zr_{0.98}A_{0.02}O_6$ (A = V, Nb), and $Li_{7.98}Zr_{0.98}Ce_{0.02}O_6$ are investigated. Ionic conductivity at doping increases as compared to $Li₈ZrO₆$; however, below \sim 680 K, it remains rather low and significantly smaller as compared to $Li₇LO₆$ (L = Nb, Ta, Sb, Bi), which have the similar crystal structure but contain structural lithium vacancies.

Conductivity of both undoped $Li₈ZrO₆$ and solid solutions shows a jump-like rise with increasing temperature at 680– 710 K. Earlier, such a behavior was explained by a structural phase transition with disordering of the lithium sublattice. The data of Raman spectroscopy, however, shows that the melting of lithium hydroxide formed due to interaction of samples with atmospheric moisture may be another reason for a sharp growth of conductivity near 700 K.

The results of NMR experiments show that at low temperatures, the movement of lithium cations in the crystal lattice takes place via exchange between octahedral and tetrahedral positions.

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