ORIGINAL PAPER

# Electrochemical behaviour of hybrid devices based on  $Na<sub>2</sub>SO<sub>4</sub>$ and Rb<sub>2</sub>SO<sub>4</sub> neutral aqueous electrolytes and carbon electrodes within wide cell potential region

Alar Jänes & Jaanus Eskusson & Leonard Mattisen & Enn Lust

Received: 18 December 2013 /Revised: 22 October 2014 /Accepted: 22 October 2014 /Published online: 19 November 2014  $\oslash$  Springer-Verlag Berlin Heidelberg 2014

Abstract Focused ion beam scanning electron microscopy (FIB-SEM), X-ray photoelectron spectroscopy (XPS) and Brunauer–Emmett–Teller gas adsorption methods have been used for the characterisation of physical properties of microporous carbide-derived carbon electrodes, prepared from  $Mo<sub>2</sub>C$  at 600 °C (noted as  $Mo<sub>2</sub>C-CDC$ ) before and after electrochemical tests conducted within a very wide twoelectrode cell potential region. Cyclic voltammetry, constant current charge/discharge and impedance data have been analysed to establish the electrochemical characteristics of the hybrid devices consisting of the 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  and 1 M  $Rb_2SO_4$  aqueous electrolytes and Mo<sub>2</sub>C-CDC electrodes within the very wide cell potential region ( $\Delta E \leq 2.4$  V). The influence of cation chemical composition on the electrochemical characteristics of supercapacitors/electrochemical hybrid devices has been analysed. The complex kinetics behaviour of completed devices (adsorption, blocking adsorption and intercalation of  $Na<sup>+</sup>$  and  $Rb<sup>+</sup>$  ions; faradic and mass transfer; gas adsorption; etc.) has been established at  $\Delta E \ge 1.5$  V. At least three different characteristic time constants dependent on the electrolyte cation composition and cell potential applied have been established.

Keywords Carbide-derived carbon . Aqueous electrolyte . Supercapacitor . Hybrid electrochemical device

A. Jänes  $\cdot$  J. Eskusson  $\cdot$  E. Lust ( $\boxtimes$ ) Institute of Chemistry, University of Tartu, 14a Ravila Str., 50411 Tartu, Estonia e-mail: enn.lust@ut.ee

L. Mattisen

Institute of Physics, University of Tartu, 14c Ravila Str., 50411 Tartu, Estonia

# Introduction

Nowadays, mainly four types of electrochemical energy storage/conversion devices are under development: supercapacitors (SCs), batteries, electrolysers (ELs) and fuel cells (FCs) [\[1](#page-13-0)–[6\]](#page-13-0). Batteries have higher specific energy than supercapacitors but lower specific power, and a very limited number of charging/discharging cycles can be applied [\[1](#page-13-0)–[12\]](#page-14-0). Fuel cells and electrolysers [[13](#page-14-0)–[15](#page-14-0)] are characterised with high energy but moderate power density, and differently from supercapacitors, they are working in the stationary regime. Regardless of low energy density, the high power density supercapacitors are important energy storage systems, which can be used in various areas of modern technology, starting from the pulse energy generation/accumulation systems and finishing with the consumer goods [\[16](#page-14-0)–[21](#page-14-0)]. For the high specific energy–specific power density application, the microporous carbon material [\[22](#page-14-0), [23\]](#page-14-0) and electrolyte characteristics have to be optimised [\[1](#page-13-0)–[6](#page-13-0), [8](#page-13-0)–[11](#page-14-0), [19](#page-14-0)–[22](#page-14-0)]. Based on the results obtained [\[8](#page-13-0)–[11](#page-14-0), [19](#page-14-0)–[23\]](#page-14-0), the electrolyte chemical composition has a big influence on the electrochemical behaviour of aqueous and non-aqueous SCs [[7](#page-13-0)–[12](#page-14-0), [23](#page-14-0)–[26\]](#page-14-0). The potential advantages of  $H<sub>2</sub>O$ -based supercapacitors would be the very cheap electrolyte applicable for SCs, regardless of the narrower region of ideal polarisability compared with that for non-aqueous SCs  $[8-11, 19-27]$  $[8-11, 19-27]$  $[8-11, 19-27]$  $[8-11, 19-27]$  $[8-11, 19-27]$  $[8-11, 19-27]$  $[8-11, 19-27]$  $[8-11, 19-27]$  $[8-11, 19-27]$ . For neutral Na<sub>2</sub>SO<sub>4</sub> solution in  $H_2O$  and microporous and mesoporous carbons, the cell potential operation region from 1.0 up to 1.6 V [[23](#page-14-0)–[25](#page-14-0)], and even 1.9 V, has been demonstrated [\[26\]](#page-14-0). Of course, at  $\Delta E \ge 1.6$  V, these systems are nonideally polarisable and more detailed electrochemical analysis including impedance is inevitable, being the main aim of this paper.

Historically, SCs have been divided into two main groups: electrical double-layer capacitors and so-called hybrid supercapacitors [[1](#page-13-0)–[4](#page-13-0), [8](#page-13-0)–[11](#page-14-0), [16](#page-14-0)–[18\]](#page-14-0). However, based on our

earlier studies [\[8](#page-13-0)–[11](#page-14-0), [19](#page-14-0)–[21](#page-14-0)], there is no rigid border between these types of SCs as some faradic reactions, including gas evolution ( $H_2$  or surface oxidation and  $O_2$  formation and evolution), can be observed outside of the low current density cell potential region (so-called electrical double-layer charging/discharging region, known as the region of ideal polarisability  $\Delta E_{\text{ideal}}$ ) observed for carbon | H<sub>2</sub>O+xM Na<sub>2</sub>SO<sub>4</sub> neutral electrolyte systems usually at  $\Delta E \ge 1.4$  V. Thus, if the cell potential  $\Delta E > \Delta E_{\text{ideal}}$ , the nearly exponential increase of current density with cell potential is usually observed, where the faradic reduction reactions (reduction of  $H_3O^+$  to  $H_2$  and adsorption or intercalation of Li<sup>+</sup>, Na<sup>+</sup>, etc., cations as neutral atoms) at a negatively charged surface are main processes.  $H<sub>2</sub>O$  decomposition or surface oxidation in a neutral medium at a positively charged surface is limited by the activation energy, strongly dependent on the pH of the solution, as well as on the electrode material used. For graphite electrodes, the H<sub>2</sub> evolution overvoltages  $\eta_{H2}$  as well as the surface oxidation/ $O_2$  evolution overpotentials  $\eta_{\text{ox}}$  are usually very high [\[1,](#page-13-0) [8,](#page-13-0) [23\]](#page-14-0).

Using the electrochemical kinetics data for metal electrodes [[28\]](#page-14-0), it was demonstrated that  $H_2$  evolution overpotential  $\eta$  depends strongly on the cation adsorption energy, being highest for Cs<sup>+</sup>-containing electrolytes ( $\eta_{H2}$  ~40 mV, compared with H<sup>+</sup> cations containing salts) due to the shift of  $\psi_1$  potential toward more positive values ( $\psi_1$  potential is the so-called Frumkin or double-layer correction potential, usually fixed/assumed for the plane (position in the Helmholtz layer)), where the electrochemical reaction takes place [[28](#page-14-0)–[32](#page-14-0)]. Thus, usually, it is assumed that the Frumkin correction depends on the chemical composition of the surface active cation used. However, for microporous–mesoporous electrodes, there are some objective complications connected with the establishment of the outer Helmholtz plane position due to the surface roughness and microporosity–mesoporosity, as well as macroscopic non-homogeneity of the electrode surface [\[31](#page-14-0)–[36\]](#page-14-0).

This paper reports the results of the systematic studies of the electrochemical devices based on mainly microporous with a small amount of mesopores carbide-derived carbon (CDC) electrodes, prepared from Mo<sub>2</sub>C (−325 mesh powder, Sigma-Aldrich) at  $T<sub>synt</sub>$ = 600 °C, noted as  $Mo_2C-CDC$ , and 1 M  $Rb_2SO_4$  or 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  aqueous electrolyte within the very wide two-electrode cell potential  $\Delta E$ . The very high  $\Delta E$ values have been applied to activate the faradic reactions at Mo2C-CDC electrodes followed by the adsorption/absorption of formed  $H_2$  and  $O_2$  inside of the microporous  $Mo<sub>2</sub>C-CDC$  electrode, as well as carbon electrode surface oxidation at very positive electrode potentials, increasing the faradic capacitance component.

## Experimental

The SC electrodes were prepared from  $Mo<sub>2</sub>C-CDC$  powder and from the mixture of binder (polytetrafluoroethylene, PTFE, 60  $\%$  solution in H<sub>2</sub>O). The mixture received was laminated and roll pressed (HS-160N, Hohsen Corporation) to form a flexible layer of an active electrode material with thickness  $L=100\pm 5$  μm.

The crystallinity and macroscopic structure of the  $Mo<sub>2</sub>C-$ CDC electrodes were evaluated by focused ion beam scanning electron microscopy (FIB-SEM)–energy-dispersive X-ray spectroscopy (EDX) and X-ray diffraction (XRD) (with a Cu anode to generate CuK $\alpha$  radiation ( $\lambda$ =1.5406 Å)). XRD data were collected in the  $2\theta$  range from 15<sup>°</sup> to 70<sup>°</sup> with the step size of 0.05°.

The FIB-SEM data have been obtained using Helios™ NanoLab 600. The X-ray photoelectron spectroscopy (XPS) experiments were carried out with a SCIEHTA SES-100 spectrometer by using an unmonochromated MgKα X-ray source with power of 300 W. The SES-100 system was calibrated using Au 4f photolines. The pressure in the analysis XPS chamber was within the range from  $2\times10^{-10}$  to 5×  $10^{-10}$  mbar.

The two-electrode standard stainless steel test cell (HS Test Cell, Hohsen Corporation) with two identical electrodes (flat cross-sectional area of one electrode ~2.0 cm<sup>2</sup>) was completed with a manometer (Migishita Seiki Co.) for analysis of pressure inside of the cell. All electrochemical experiments were carried out at fixed temperature  $T=20\pm5$  °C. A 25-µm-thick Celgard® 2400 separator sheet was used for mechanical separation of the working  $Mo<sub>2</sub>C-CDC$  electrodes from each other. One molar  $Rb_2SO_4$  and  $Na_2SO_4$  aqueous electrolytes were used.

### Results and discussion

### Adsorption measurement data

Nitrogen adsorption measurements at liquid nitrogen temperature have been conducted using the Micromeritics ASAP 2020 surface area and porosity analysis measurement system. Based on the N<sub>2</sub> adsorption data for Mo<sub>2</sub>C-CDC at  $T_{\text{synt}}$ = 600 °C, the multipoint Brunauer–Emmett–Teller (BET) surface area  $S_{\text{BET}}=1855 \text{ m}^2 \text{ g}^{-1}$ , micropore area (obtained using the *t-plot* method)  $S_{micro} = 1823 \text{ m}^2 \text{ g}^{-1}$ , total pore volume  $V_{\text{tot}}$ =1.139 cm<sup>3</sup> g<sup>-1</sup> and micropore volume  $V_{\text{micro}}$ = 1.077 cm<sup>3</sup> g<sup> $-1$ </sup> were calculated. The pore size distribution function vs. pore width plots, obtained using non-local density functional theory, shows two main peaks at  $d_1$  = 1.04 nm and  $d_2$ =2.64 nm, respectively. Thus, in addition to the micropores  $(d_1)$  lower than 2 nm according to IUPAC classification), there are some amount of mesopores (2 nm  $\leq d_2$  < 50 nm, IUPAC) at/

<span id="page-2-0"></span>inside of the Mo<sub>2</sub>C-CDC, synthesised at  $T=600$  °C. More detailed gas adsorption analysis has been given in our previ-ous paper [\[18](#page-14-0)]. Thus, based on the  $N_2$  sorption data and our previous SC studies  $[16–18]$  $[16–18]$  $[16–18]$  $[16–18]$  $[16–18]$ , Mo<sub>2</sub>C-CDC with hierarchical structure is an interesting material containing micropores for adsorption of charged ions and adsorption/absorption of gaseous  $H_2$  or  $O_2$  into micropores and mesopores. It should be mentioned that mesopores are needed for the quick mass transfer of reactants ( $H_3O^+$  ions (or  $H_2O$  molecules)) into the micropores [\[22](#page-14-0)] as well for desorption of gaseous  $H_2$  and  $O_2$ from micropores into the neutral electrolyte solution.

## FIB-SEM–EDX, Raman and XPS measurements

Focused ion beam scanning electron microscopy combined with the EDX (FIB-SEM–EDX) analysis method was used for the quantitative analysis of SC electrodes (Fig. 1a–d, Table [1\)](#page-4-0). Similarly to the high-resolution transmission electron microscopy (HRTEM) data, discussed in [\[18](#page-14-0)], it was shown that the small amount of graphite-like carbon layers at the top (core level) of mainly amorphous carbon particles was observed. The Raman spectra were recorded using a Renishaw micro-Raman spectrometer equipped with a  $\times$ 50 objective and 488 nm  $Ar^+$  ion laser with maximum radiation power of 15 mW on the sample. XRD and Raman data, in agreement with HRTEM data, show that  $Mo<sub>2</sub>C-CDC$  carbon has mainly an amorphous structure and only very slightly graphitized regions were found at the surface of some particles of Mo<sub>2</sub>C-CDC material [\[11\]](#page-14-0). It was found that before electrochemical tests, the porous  $Mo<sub>2</sub>C-CDC$  structure consists of particles with variable linear dimensions (from 10 to 100 nm). PTFE wires and lumps were observed between CDC particles, connecting the carbon particles into the flexible electrode layer (not shown for shortness). After electrochemical analysis (after 10, 100, 1000 and 10,000 charge/discharge cycles), the SC cells were opened and electrodes were washed many times



Fig. 1 Results of FIB-SEM analysis for Mo<sub>2</sub>C-CDC electrodes in 1 M  $Rb_2SO_4$  (a), 1 M Na<sub>2</sub>SO<sub>4</sub> (b–d) and aqueous solutions for negatively (a– c) and positively (d) charged supercapacitor electrodes. Image b is taken

from the membrane side and  $c$  from the current collector side  $Mo<sub>2</sub>C-CDC$ electrode, tested in  $Na<sub>2</sub>SO<sub>4</sub>+H<sub>2</sub>O$  solution, respectively

with clean Milli- $O<sup>+</sup>$  H<sub>2</sub>O, dried under vacuum and analysed again using the FIB-SEM and XPS methods. It was demonstrated that in agreement with HRTEM data [\[18](#page-14-0)], there is no Cl<sup>−</sup> or Cl<sub>2</sub> and/or other contaminants in electrode materials under testing.

A lot of signals connected with Rb (Fig. [1a,](#page-2-0) Table [1a](#page-4-0)) and Na (Fig. [1b, c](#page-2-0), Table [1b, c](#page-4-0)) absorbed/ deposited/intercalated into the negatively charged  $Mo<sub>2</sub>C-$ CDC were found after 10,000 galvanostatic cycles in  $Rb_2SO_4$  or  $Na_2SO_4$  aqueous electrolyte, differently from the carbon electrodes measured after ten cycles (not shown for shortness) have been observed. Somewhat lower concentrations of  $Na<sup>+</sup>$  and  $Rb<sup>+</sup>$  have been observed in CDC after 1000 charge/discharge cycles (not shown for shortness). It should be noted that for positively charged  $Mo<sub>2</sub>C-CDC$ , the concentration of Na was very low (Fig. [1d,](#page-2-0) Table [1d](#page-4-0)) indicating that the electrochemically activated accumulation of  $Na<sup>+</sup>$  ions (Na intercalation) has been taking place in the negatively charged electrode only.

Thus, during electrochemical potential cycling in the wide cell potential region, adsorption/absorption of Rb<sup>+</sup> and  $Na<sup>+</sup>$  ions and the intercalation (or chemisorption) of Rb and Na onto/into the  $Mo<sub>2</sub>C-CDC$  electrode structure have been taking place. It should be mentioned that there is no big difference in amount of Rb or Na accumulated (data in Fig. [1a, b](#page-2-0), Table [1a, b\)](#page-4-0). It is surprising that a small amount of S (probably from  $SO_4^2$ <sup>-</sup> ions), especially from the Rb<sub>2</sub>SO<sub>4</sub> solution, has been observed as well.

The XPS analysis data (Fig. [2\)](#page-5-0) confirm the existence of adsorbed/intercalated Rb and Na (binding energy peak from 235 to 248 eV for Rb and at 498 eV for Na) onto/into the  $Mo_2C-CDC$  structure. Based on these results, the amount of the strongly covalently bonded Rb for Mo2C-CDC electrodes observed after 10,000 galvanostatic cycles is remarkable. For electrodes studied after 100 cycles, the amount of covalently bonded Rb is smaller. For electrodes analysed after ten cycles, it was impossible to verify the existence of the Rb-C surface compound at all. It should be noted that the FIB-SEM–EDX and XPS are ex situ methods and, for more detailed studies, electrochemical in situ synchrotron radiation-based X-ray absorption spectroscopy [[37\]](#page-14-0) must be applied.

#### Cyclic voltammetry data

Cyclic voltammetry (CV) curves, expressed as capacitance for a symmetrical two-electrode system in 1 M  $Na<sub>2</sub>SO<sub>4</sub>$ , as well as in 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  aqueous solutions, are given in panels a and b, respectively, of Fig. [3.](#page-6-0) The current density, j, measured at fixed cell potential scan rate  $v$ , has been used for calculation of the medium capacitance values according to Eq. 1:

$$
C_{\rm CV} = j\nu^{-1} \tag{1}
$$

Equation 1 is correct if the capacitance  $C_{CV}$  is constant  $(C_{\text{CV}} \neq f(E))$  [[8](#page-13-0), [19](#page-14-0), [20\]](#page-14-0) and/or the current density applied is small. In a symmetrical two-electrode system, the gravimetric capacitance  $C_{m;cv}$  (F  $g^{-1}$ ) for one activated carbon electrode can be obtained as follows:

$$
C_{m;\text{CV}} = \frac{2C_{\text{CV}}}{m} \tag{2}
$$

where  $m$  is the weight per one activated carbon electrode assuming that the positively and negatively charged electrodes have the same capacitance at fixed  $\Delta E$  applied. The cyclic voltammetry curves expressed as capacitance vs. cell potential curves up to 1.3 V for 1 M  $Rb_2SO_4$  and 1.2 V for  $Na_2SO_4$ , obtained at small voltage scan rates  $v = d(E)/dt$ , have nearly mirror image symmetry of the current responses about the zero current line  $(t$  is potential scanning time) characteristic of the ideal polarizability of the system under study. However, the discharging current densities, i.e. capacitance values for  $Rb_2SO_4$  as well as for Na<sub>2</sub>SO<sub>4</sub>, depend on the final  $\Delta E$ applied, indicating that the desorption of  $Rb^+$  as well as  $Na^+$ ions is slower than the adsorption step. At  $1.3 \leq \Delta E \leq 1.5$  V, for both electrodes, a small increase in faradic current density has been observed, indicating that the hydrogen evolution and  $H<sub>2</sub>O$  decomposition (or surface oxidation) at  $Mo<sub>2</sub>C-CDC$  is a slow process due to the very high  $H_3O^+$  electroreduction overvoltages [\[33](#page-14-0)–[36\]](#page-14-0).

However, for 1 M  $Rb_2SO_4$  solution, the quick nearly exponential increase of j can be seen at  $\Delta E \ge 1.3$  V. Surprisingly, there is nearly linear dependence of  $\Delta E$  on log *i* (at  $\Delta E \ge$ 1.3 V) (Fig. [4\)](#page-6-0), known as the Tafel plot, if the individual electrode has been analysed. However, for the two-electrode system,  $\Delta E$  vs. log *i* dependence characterises the behaviour of two different (or even more) electrode processes (or the behaviour of one electrode with the rate-limiting reaction kinetics). Therefore, the slope values of  $b$  and Tafel constant value *a* in the Tafel equation  $\Delta E = a + b \log i$  calculated cannot be used for detailed and correct reaction mechanism analysis. The slope  $b$  value, characterising the complex electrochemical behaviour of the system, including the electroreduction step of hydrated protons  $(H_3O^+ + e^- \rightarrow H_2O + H_{ads})$  and interaction of Rb through reaction  $(Rb^+ + e^-)$  and negatively charge electrode as well as the surface oxidation and  $O_2$  evolution (H<sub>2</sub>O) oxidation) reaction at the positively charged  $Mo_2C-CDC$ electrode interface, is very high for the  $Rb_2SO_4$ -based electrolyte.

For  $1 \text{ M Na}_2\text{SO}_4$  solution, the intensive increase of current density starts already only at  $\Delta E \ge 1.5$  V (Fig. [3](#page-6-0)). Differently

<span id="page-4-0"></span>



Processing option: all elements analysed (normalised). All results in weight percent

from the  $Rb_2SO_4$  electrolyte, there is no linear dependence of  $\Delta E$  on log *i* (Fig. [4](#page-6-0)) within the 1.3 < $\Delta E \le 1.5$  V region. Thus,

only at  $\Delta E \ge 1.5$  V the linear dependence of  $\Delta E$  on log *i* has been observed with a noticeably lower slope value.

<span id="page-5-0"></span>Fig. 2 XPS spectra for negatively ( a , b) and positively ( c) charged Mo 2C-CDC electrodes in 1 M  $Rb<sub>2</sub>SO<sub>4</sub>(a)$  and  $1 M Na<sub>2</sub>SO<sub>4</sub>$  (**b**, **c**) aqueous solutions



<span id="page-6-0"></span>

Fig. 3 Cyclic voltammetry curves expressed as capacitance  $C$  (per one electrode) vs. cell potential curves for supercapacitors in 1 M  $\text{Na}_2\text{SO}_4$  (a) and in 1 M Rb<sub>2</sub>SO<sub>4</sub> (b) at potential scan rate 10 mV s<sup>-1</sup> at different maximal cell potentials

Considerably different slope values (Fig. 4) for these two systems indicate that in addition to  $H<sub>2</sub>$  evolution and surface oxidation, the adsorption and reduction (intercalation) of cations play a probably important role. The intensive increase of



the faradic process(es), i.e. hydrogen and  $O<sub>2</sub>$  (surface oxidation) evolution, is shifted ~150±5 mV toward higher  $\Delta E$ values for the  $Rb_2SO_4$  solution, indicating that the hydrogen and oxygen evolution overvoltages at  $Mo_2C-CDC$  from 1 M Na2SO4 solution are somewhat lower than that for 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$ . This is in very good agreement with  $H<sub>2</sub>$  evolution data established for Hg and other metal electrodes [\[28\]](#page-14-0), indicating that the Na<sup>+</sup> or Rb<sup>+</sup> adsorption influences the  $\psi_1$ potential similarly to high hydrogen overvoltages measured for metals (Hg, Pb, Bi, Cd). Thus, the adsorption of  $Rb^+$ (similarly to  $Cs<sup>+</sup>$ ) cations increases the hydrogen as well as the oxygen evolution overvoltages [[28](#page-14-0), [33](#page-14-0)–[36](#page-14-0)]. It is surprising that during the first very wide cell voltage cycles ( $\Delta E = 2.0$  V), the charging (reduction/oxidation) curve is very well reproducible and only a weak increase of discharging currents (reoxidation of  $H<sub>2</sub>$  and reduction of oxygen containing surface functionalities) can be seen after 100 or 1000 cycles but especially after 10,000 cycles (not shown for shortness).

After the polarisation of electrodes at cell  $\Delta E \geq 2.4$  V in the discharging curve at  $\Delta E < 1.0$  V, only a very small and very wide current peak (plateau) (Fig. [5](#page-7-0)) has been observed, indicating the reoxidation of  $H_2$ adsorbed/absorbed in porous  $Mo<sub>2</sub>C-CDC$ , as well as the reduction of the oxygenated functional groups at the electrode surface. Faradic hydrogen reoxidation (ionisation) and surface reduction processes at the  $Mo<sub>2</sub>C-$ CDC electrode surface, within the region of  $\Delta E$  studied, are very slow and irreversible processes as there are no clearly visible current peaks in the discharging curve at lower cell potentials and wide potential scan rates, v, applied.

A comparison of current densities measured in 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  and 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  solutions indicates that both electrodes can be applied within the very wide  $\Delta E$  region; however, for the  $Rb_2SO_4$ -based electrolyte, the faradic current densities are clearly lower (Figs. 3 and [5](#page-7-0)) due to the blocking



<span id="page-7-0"></span>

Fig. 5 Cyclic voltammetry curves expressed as capacitance C vs. cell potential curves in 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  (a) and in 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  (b) at different cell potentials and at different potential scan rates, noted in the figure. Hybrid electrochemical device cell pressure vs. time dependence in 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  (c) and at cell potential 2.4 V

adsorption of  $Rb^+$  at the negatively charged Mo<sub>2</sub>C-CDC electrode (Figs. [1](#page-2-0) and [2\)](#page-5-0).

The activation of microporous and mesoporous  $Mo<sub>2</sub>C-$ CDC electrodes takes place only at  $\Delta E \ge 1.5$  V. Thus,  $C_{m}$ . increases with the decrease of the potential sweep rate applied, indicating a slow rate of reduction and oxidation reactions occurring at negatively and positively charged electrodes. Data in Fig. 5a, b show that the high capacitance values have been calculated in the high-current-density (hydrogen/oxygen evolution/surface oxidation) region, explained by the high



Fig. 6 Surface charge density ( $1-4$ ) and coulombic efficiency  $\eta(5, 6)$  vs. cell potential curves for the Mo<sub>2</sub>C-CDC-based cell in 1 M Na<sub>2</sub>SO<sub>4</sub> (1, 2, 5) and 1 M Rb<sub>2</sub>SO<sub>4</sub> (3, 4, 6) for charging (1, 3) and discharging (2, 4), noted in the figure

faradic process capacitance values for hydrogen and oxygen evolution reactions.

At  $v>5$  mV s<sup>-1</sup> and  $\Delta E>1.7$  V, the C values are lower for the  $Rb_2SO_4$  electrolyte-based SC explained by the higher hydrogen overvoltage for 1 M  $Rb_2SO_4$  solution than that for  $Na<sub>2</sub>SO<sub>4</sub>$ . Thus, the adsorption/reduction/intercalation of  $Rb<sup>+</sup>$  cations onto the Mo<sub>2</sub>C-CDC surface causing the shift of the zero total charge potential (and the  $\psi_1$  potential) toward the less negative electrode potential increases the hydrogen and oxygen evolution overpotential (∼0.04 V).

For the verification of hydrogen and oxygen evolution processes, the pressure inside of the hermetic test cell was monitored. It is surprising that during the first 10–20 min, there was no big increase of pressure inside of the cell during the first cycles or even during short-time holding  $(t<3 \text{ min})$  of the cell at  $\Delta E = 2.4$  V (Fig. 5c). Only during long-lasting holding  $(1-1.5 h)$  of the two-electrode cell at 2.4 V the small overpressure in cells has been observed, as demonstrated in Fig. 5c. As can be seen in Fig. 5c, during 50 to 500 min polarisation at  $\Delta E = 2.4$  V, the pressure inside of the cell increases nearly exponentially and, after polarisation during 1000–1500 min (at  $\Delta E = 2.4$  V), the increase in  $\Delta p$  practically stopped. To the first very rough approximation, this surprising result can be explained by the adsorption/absorption of generated hydrogen/oxygen onto/into the porous  $Mo<sub>2</sub>C-CDC$ electrode [\[28](#page-14-0)–[30,](#page-14-0) [33](#page-14-0)–[35](#page-14-0)].

Surface charge density potential curves (at potential scan rate 10 mV  $s^{-1}$ ) have been integrated to obtain the surface charge density  $Q$ ,  $\Delta E$  curves (Fig. 6). A comparison on  $Q$ ,  $\Delta E$ curves for charging and discharging cycles indicates that the electrochemical coulombic efficiency  $\eta$  noticeably decreases with the increase of  $\Delta E$  applied, if  $\Delta E \ge 1.5$  V. A somewhat lower  $\eta$  has been calculated for the Rb<sub>2</sub>SO<sub>4</sub>-based system, explained by the blocking adsorption of  $Rb<sup>+</sup>$  cations into the negatively charged Mo2C-CDC electrode. The cyclic voltammetry and coulombic efficiency data show that systems under discussion can be used as combined electricity storage, i.e.  $H_2/O_2$  generation/storage systems however with quite low efficiency at  $\Delta E \ge 1.5$  V. However, it is very interesting to mention that even at  $\Delta E = 1.8$  V, the electrochemical reversibility of the Mo<sub>2</sub>C-CDC | 1 M Na<sub>2</sub>SO<sub>4</sub> system is higher ( $\eta \geq$ 85 %) than that for lead accumulators ( $\eta \leq 80$  %) [[1\]](#page-13-0). A very quick decrease of  $\eta$  has been established for the Rb<sub>2</sub>SO<sub>4</sub>-based system at  $\Delta E \geq 1.5$  V.

## Constant current charge/discharge data

The SC cells were tested at constant current (CC) charge/ discharge regimes (from 1 to 50 mA  $cm^{-2}$ ) within the different voltage range from 0 to 1.2, from 0 to 1.5 and from 0 to 2.0 V (Fig. 7a), and from 1.0 to 2.0 V (Fig. 7b). CC curves obtained for the normal  $\Delta E$  region (0–1.2 V) have a very nice

Fig. 7 Galvanostatic charging/ discharging curves for the  $Mo<sub>2</sub>C-$ CDC-based cell in  $1 M Rb<sub>2</sub> SO<sub>4</sub>$  at current density 5 mA cm−<sup>2</sup> at cell potential from 0 to 2 V (a) and from 1 to 2  $V$  (b)

traditional shape [[23\]](#page-14-0) characteristic of ideally polarisable systems (not shown for shortness). However, Fig. 7a, b shows that the very complicated charging/discharging curves have been measured if the wider  $\Delta E$  region has been applied: 0 V < $\Delta E$  < 2.0 V or 1.0 V <  $\Delta E$  < 2.0 V. In principle, the two nearly linear charging regions have been observed: first at lower cell potentials  $(0.5<\Delta E<1.0$  V, double-layer charging) and second from 1.3 V < $\Delta E$  < 1.8 V (probably mainly H<sub>3</sub>O<sup>+</sup> reduction and  $Na<sup>+</sup>$  or  $Rb<sup>+</sup>$  adsorption). In discharging CC curve, the linear regions have been observed at 1.2 V < $\Delta E$  < 1.9 V and 0.1 V < $\Delta E$  < 1.0 V. Interestingly, if these cells have been cycled within the limited cell potential region (1.0  $V < \Delta E < 2.0$  V), the discharging curves are non-linear within the all potential region  $(1.0<\Delta E<2.0$  V) applied, indicating that after the gas (H<sub>2</sub>, O<sub>2</sub>) formation/adsorption step at  $\Delta E$ > 1.2 V, the gas desorption and ionisation is only possible at





lower cell potentials than  $\Delta E < 1.0$  V due to the very high activation energy of  $H_2$  or  $O_2$  ionisation step(s). Thus, the  $H_2$ and/or  $O_2$  ionisation and surface oxide reduction and cation desorption have very high overvoltages at Mo<sub>2</sub>C-CDC electrodes.

Usually, the capacitance of the cell ( $C_{\text{CC}}$ ) at  $\Delta E \leq 1.2$  V will be obtained from the slope of the discharge curve according to Eq. 3:

$$
C_{\rm CC} = j \frac{\mathrm{d}t}{\mathrm{d}(E)}\tag{3}
$$

where  $dt/d(E)$  is the slope of the discharge or charging curve with corresponding current density j.

From the constant current charge/discharge curves, the value of internal resistance,  $R_{CC}$ , from the *IR*-drop is observable after the changing of the charging or discharging current direction has been calculated ( $R_{CC} = dE_1/2j$ , where  $dE_1$  is the value of cell potential for 10 ms). The lowest  $R_{\text{CC}}$  value has been calculated for the 1 M  $Rb_2SO_4+H_2O$ -based cell compared with the  $Na<sub>2</sub>SO<sub>4</sub>+H<sub>2</sub>O$ -based cell, explained by the higher molar conductivity of the  $Rb<sub>2</sub>SO<sub>4</sub>$  aqueous solution in agreement with the analysis of impedance data (next chapter) and limiting ionic (molar) conductivity values  $\lambda$  $(S \text{ cm}^2 \text{ mol}^{-1}$ : Na<sup>+</sup> 50.10, Rb<sup>+</sup> 77.8, SO<sub>4</sub><sup>2−</sup> 160). At  $\Delta E$ ≤ 1.5 V, where the gas formation is slow and unimportant, the IR-drop is practically independent of the cycle number applied. However, the  $R_{CC}$  value slightly increases with the cell potential applied at  $\Delta E \ge 1.5$  V, explained by the accumulation of  $H_2$  and  $O_2$  (or surface oxides) inside of the porous carbon matrix and squeezing out of some electrolytes from the microporous matrix, and/or by blocking adsorption/intercalation of Na or Rb atoms into the porous matrix, decreasing the linear dimensions of the pore.

For the  $Mo_2C-CDC$  | 1 M  $Na_2SO_4$  or  $Mo_2C-CDC$  | 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  aqueous interface, there is no remarkable dependence of the relative capacitance,  $C_{rel} = C_{in}/C_x$ , or relative resistance,  $R_{\rm rel} = R_{\rm in}/R_x$ , on the galvanostatic cycle number applied ( $C_{\rm in}$ ) and  $C_x$  are the capacitance values of the third and x number cycle, respectively, and  $R_{\text{in}}$  and  $R_x$  are the corresponding series resistance values, calculated from the *IR*-drop), if  $\Delta E$ <1.3 V. Hence, for 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  and 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  | Mo<sub>2</sub>C-CDC interfaces,  $C_{\text{rel}}$  and  $R_{\text{rel}}$  are quite stable and these salts can be used as aqueous electrolytes for SCs at  $\Delta E \leq 1.3$  V.

The coulombic cycling efficiency, i.e. the so-called round trip efficiency (RTE) has been calculated as a ratio of capacitances measured during discharging and charging of SCs, if  $\Delta E \leq 1.3$  V. The highest RTE  $\geq 98$  % at  $\Delta E \leq 1.3$  V has been calculated for 1 M  $Rb_2SO_4$  aqueous supercapacitors and RTE  $\geq$ 96 % for 1 M Na<sub>2</sub>SO<sub>4</sub>, demonstrating the reasonable dependence of RTE on the electrolyte conductivity discussed before. This effect can be explained by the influence of the *IR*-drop as well as the mass transfer and partial charge transfer resistances on the total charge accumulation efficiency, because the solvation (standard molar hydration enthalpy  $\Delta H$ ) of Rb<sup>+</sup> ions in H<sub>2</sub>O is somewhat lower than that for Na<sup>+</sup> ions ( $\Delta H$  (kJ mol<sup>-1</sup>)  $Na<sup>+</sup>=-321$  and  $Rb<sup>+</sup>=-300$ , respectively). Thus, the adsorption/absorption kinetics of ions (reversibility of charging/discharging of the system) depends on the hydration of cations applied as an electrolyte for SC. Formation of  $H_2$ and  $O<sub>2</sub>$  gases or surface compounds as adsorbed/intercalated reaction products or surface oxides onto/into some places of electrode material at higher  $\Delta E$  is possible, but a weak influence on the RTE of supercapacitors based on the 1 M  $\text{Na}_2\text{SO}_4$ or 1 M  $Rb_2SO_4$  electrolyte indicates that this effect is quite unimportant at lower  $\Delta E \leq 1.3$  V. In addition, the physical adsorption/intercalation of  $Rb^+$  ions at/into the Mo<sub>2</sub>C-CDC electrode shifts the zero charge potential toward more positive values ( $\sim$ 30 mV), thus increasing the  $\psi_1$  potential values at the outer Helmholtz plane and therefore increasing the  $H<sub>2</sub>$  evolution overpotential.

However, for systems tested at  $\Delta E$  > 1.3 V, Eq. 3 cannot be applied as the charging/discharging curves are non-linear. Therefore, for the more detailed analysis, the charging/ discharging curves have been integrated to obtain the charge densities accumulated during charging  $Q<sub>ch</sub>$  and discharging  $Q_{\text{disch}}$  steps. The calculations show that in agreement with CV data, the coulombic efficiency is quite low ( $\eta \le 0.8$ ) at  $\Delta E$ > 2.3 V. The very quick decrease of  $Q_{ch}$  and  $Q_{disch}$  takes place after application of a comparatively small number of charging/ discharging cycles ( $N_{\text{cyc}}$ >1000 at  $\Delta E$ >1.5 V), explained by the quick electrode degradation process after/during gas evolution, adsorption/deionisation steps or by the adsorption/absorption/intercalation of  $Na<sup>+</sup>$  or  $Rb<sup>+</sup>$  ions (thus Na or Rb atoms) into the microporous  $Mo<sub>2</sub>C-CDC$  electrodes.

Analysis of Nyquist and Bode plots

The Nyquist plots [[19](#page-14-0)–[21](#page-14-0), [38](#page-14-0)–[44](#page-14-0)], given in Fig. [8a, b](#page-10-0) for SCs completed from  $Mo<sub>2</sub>C-CDC$  electrodes in 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  or 1 M Na2SO4 aqueous solutions, have been measured within the wide range of ac frequency,  $f$  (from  $1 \times 10^{-3}$  to  $1 \times 10^{5}$  Hz), and at fixed cell potentials from  $\Delta E = 1.2$  to 2.0 V. At the very beginning of the experiments, if  $\Delta E \leq 1.2$  V, the Nyquist plot has a traditional shape (not shown for shortness) [[8](#page-13-0)–[11,](#page-14-0) [19](#page-14-0)–[21,](#page-14-0) [23](#page-14-0)] and they consist mainly of three parts: (1) the very small and noticeably depressed semicircle at higher ac frequency ( $f \ge 3$  Hz) with a characteristic frequency,  $f_{\text{max}}$  (obtained as the frequency at the maximum in the Nyquist plot) (the semicircle shape depends on the adsorption kinetics of ions at the microporous electrode and on the series resistance of a material, mass transfer resistance inside the meso/macroporous carbon structure as well as on the mass transfer resistance in the micropores); (2) the so-called doublelayer capacitance region ("knee" at low frequencies  $f$  < 0.01 Hz), obtained by the finite-length adsorption effect; and

<span id="page-10-0"></span>![](_page_10_Figure_1.jpeg)

Fig. 8 Nyquist ( $-Z''$ , Z') plots for hybrid electrochemical devices at  $\Delta E=$ 2.0 V in 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  (a) and 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  (b) aqueous electrolyte after 1000 (1), 2000 (2), 3000 (3) 4000 (4) and 5000 (5) cycles at constant current charging/discharging cycles from 0 to 2.0 V

(3) the not very-well-expressed so-called "porous" region with a slope of  $\alpha \approx 45^{\circ}$ , characteristic of the mass transfer-limited processes in the microporous–mesoporous matrix of an electrode.

At  $\Delta E \ge 1.3$  V, the Nyquist plots for the Na<sub>2</sub>SO<sub>4</sub>-based SC have a very complicated shape and at least five different areas in Z″, Z′ plots have been established (shown in Fig. 8a, b). With the increase of  $\Delta E \geq 1.5$  V applied, the role of faradic processes rises quickly (formation of the semicircle in Nyquist plots at  $f$ <10 mHz). Differently from data measured at  $\Delta E \leq$ 1.3 V, Nyquist plots at a lower region of ac  $f<1$  Hz and at 2.0 Vare non-linear and deviate strongly from a nearly vertical

line explained by the deviation of the  $Na<sub>2</sub>SO<sub>4</sub>$  | Mo<sub>2</sub>C-CDC interface from the simple adsorption-limited step (ideal capacitive) behaviour.

According to the experimental data given in Fig. 8b, the Nyquist plots at  $\Delta E = 2.0$  V for the Rb<sub>2</sub>SO<sub>4</sub> solution have an even more complicated shape, compared with the  $Na<sub>2</sub>SO<sub>4</sub>$ based SC, explained by the strong specific adsorption and intercalation of  $Rb^+$  being possible at  $\Delta E \geq 2.0$  V (Fig. 8b). The shape of the Nyquist plot depends on the number of cycle analysed more pronounced for the  $Rb_2SO_4$ -based system. The relaxation frequency for the kinetically mixed so-called verylow-frequency process  $(f<4 \text{ mHz})$  depends noticeably on the electrolyte studied, and  $f_{\text{low}}$  increases from 1.9 mHz (1 M  $Rb_2SO_4$ ) to 3.8 mHz for 1 M Na<sub>2</sub>SO<sub>4</sub>. Thus, the characteristic relaxation time  $\tau_{\text{max}} = (2\pi f_{\text{max}})^{-1}$  decreases from 0.1 s (1 M  $Rb_2SO_4$ ) to 0.04 s (1 M Na<sub>2</sub>SO<sub>4</sub>), explained by the stronger chemical interaction of  $Rb^+$  cations with the Mo<sub>2</sub>C-CDC surface caused by the lower hydration energy values for  $Rb^+$ ions compared with Na<sup>+</sup> ions [[31,](#page-14-0) [32\]](#page-14-0). In addition, the stronger adsorption bond for Rb<sup>+</sup>-containing systems is caused by the smaller hydration number for Rb<sup>+</sup> ( $\sim$ 4) compared with that for  $Na<sup>+</sup>(~6)$  as well as explained by the shorter effective Debye length values (shorter distance of the closest approach of  $Rb^+$ to the surface for weakly solvated  $Rb<sup>+</sup>$  ions compared with strongly hydrated  $Na<sup>+</sup> ions$  [\[31](#page-14-0), [32\]](#page-14-0).

Thus, probably, dependency of the shape and the slope for the Nyquist plot in the low-frequency region on the SC cell voltage is mainly caused by the rate of the electrical doublelayer formation and  $Rb^+$  adsorption/intercalation processes inside of the mesoporous and microporous carbide-derived carbon electrodes. However, mainly the adsorption-limited step for 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  and 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  electrolyte-based systems has been established only at  $\Delta E \leq 1.3$  V (not shown for shortness). The dependence of the shape of the Nyquist plots on the cell voltage at  $f<0.01$  Hz (so-called finite-length effect region) is mainly caused by the decrease of the effective screening length, well visible for the 1 M  $Rb_2SO_4$  and 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  electrolytes with increasing cell potential. The effective diffuse layer thickness as a function of the electrode rational potential  $E_R$  is given as  $\kappa_{\text{eff}}(E_R) = 2/\kappa \cos(e\beta E_R)$ , where  $\kappa_{\text{eff}}(E_R)$  and  $\kappa$  are the effective and usual Gouy lengths and  $E_R$  is a potential with respect to the zero charge potential,  $\beta = (R_{\rm B}T)^{-1}$ , where T is the absolute temperature and  $R_{\rm B}$  is the Boltzmann constant [\[39](#page-14-0), [40\]](#page-14-0). Thus, the dependence of the shape of the Nyquist plot within the very-low-frequency region on cell voltage, to a first approximation, indicates that the pore dimension is in the same order of magnitude as the effective Debye length for the 1 M electrolyte solutions, i.e. the pore dimension is comparable to the effective diameter of the partially desolvated  $Rb^+$  or Na<sup>+</sup> ions absorbed [[38](#page-14-0)–[40](#page-14-0)]. A similar dependence of the screening length and effective diffuse layer thickness has been established for specifically adsorbed anions [\[38\]](#page-14-0).

<span id="page-11-0"></span>Preliminary non-linear least squares fitting data show that in addition to the slow adsorption and faradic charge transfer steps  $(f \geq 100 \text{ mHz})$  as well as the finite-length Warburg-like mass limited transfer steps  $(5 \text{ mHz} \leq f \leq 50 \text{ mHz})$ , the inductive behaviour  $(f<2.4 \text{ mHz}, \text{Fig. 8b})$ , initiated by the so-called corrosion of electrode material or irreversible blocking adsorption, surface oxide layer formation process, etc., are

![](_page_11_Figure_2.jpeg)

Fig. 9 Bode phase angle, log f and log  $(-Z'')$ , log f plots for hybrid electrochemical devices in 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  (a) and in 1 M  $Rb<sub>2</sub>SO<sub>4</sub>$  (b) at  $\Delta E = 2.0$  V in aqueous electrolyte after 1000 (*rhombs*), 2000 (*triangles*) and 5000 cycles (circles) at current charging/discharging test from 0 to 2.0 V. Open symbols correspond to the Bode phase angle (expressing the magnitude of the frequency response gain), and filled symbols correspond to the log (−Z″), log f plots. Corresponding Bode phase angle, log f plots (1–4) and log (−Z"), log f plots (5–8) in 1 M Na<sub>2</sub>SO<sub>4</sub> (1, 3, 6, 8) and in 1 M Rb<sub>2</sub>SO<sub>4</sub> (2, 4, 5, 7) at  $\Delta E = 2.0 \text{ V } (1, 2, 7, 8)$  at  $\Delta E = 1.0 \text{ V } (3, 4, 5, 6)$ (c), noted in the figure

involved. In addition, accumulation of the gases  $(H<sub>2</sub>$  and  $O<sub>2</sub>)$ into the microporous volume of Mo<sub>2</sub>C-CDC at  $\Delta E$ >1.5 V and reionisation  $\Delta E < 1.0$  V occur in agreement with the rise of high-frequency series resistance of the system under discussion (Fig. [8a, b](#page-10-0)). It should be noted that the exact fitting of Nyquist plots at  $\Delta E$ >1.5 V and  $f$ <100 mHz was impossible due to very many processes involved.

The analysis of the Bode phase angle  $\theta$ , log f plots and log  $(-Z'')$ , log f plots (Fig. 9a–c) shows that differently from the  $Mo_2C-CDC$  | 1 M  $(C_2H_5)_3CH_3NBF_4$ +acetonitrile (nonaqueous) interface ( $\Delta E \leq 3.0$  V) for aqueous Na<sub>2</sub>SO<sub>4</sub> and Rb<sub>2</sub>SO<sub>4</sub> solutions at  $\Delta E$ <1.0 V, at  $\Delta E$ ≥1.3 V, there are noticeable deviations from the ideally polarisable interface model for  $Rb_2SO_4$  and  $Na_2SO_4$  solutions within all ac frequency region studied. For the Mo<sub>2</sub>C-CDC | 1 M Na<sub>2</sub>SO<sub>4</sub> or 1 M  $Rb_2SO_4+H_2O$  system, at least four different regions characterised by different time constants can be calculated. At  $\Delta E = 2.0$  V, the data in Figs. [8](#page-10-0) and 9 show that the veryhigh-frequency region fhigh≥10 kHz is noticeably shorter for the  $Rb_2SO_4$ -based system compared with  $Na_2SO_4$ . The highfrequency minimum (i.e. maximum with negative  $\delta$  values) in δ, log f plot is shifted nearly one order toward lower f values for the  $Na<sub>2</sub>SO<sub>4</sub>$ -based system compared with that for  $Rb<sub>2</sub>SO<sub>4</sub>$ .

![](_page_11_Figure_8.jpeg)

Fig. 10 Gravimetric series  $(C_s)$  (filled triangles) and parallel  $(C_p)$  (open *triangles*) capacitances vs. frequency plots at  $\Delta E = 2.0$  V for hybrid electrochemical devices in 1 M  $Na<sub>2</sub>SO<sub>4</sub>(a)$  and in 1 M  $Rb<sub>2</sub>SO<sub>4</sub>(b)$ 

<span id="page-12-0"></span>The  $|\delta|$  values are higher  $(|\delta|=40^{\circ})$  for the Na<sub>2</sub>SO<sub>4</sub>-based system than for the  $Rb_2SO_4$ -based SC ( $|\delta|=20^\circ$ ), indicating the more pronounced mixed kinetics for  $Rb_2SO_4$  and lower ratio adsorption but more pronounced mass transfer processes for the  $Na<sub>2</sub>SO<sub>4</sub>$ -based SC.

Phase angle absolute values are nearly comparable with the so-called medium-frequency ac (from 0.2 to 10 Hz for  $Na<sub>2</sub>SO<sub>4</sub>$  and from 1.0 to 500 Hz or  $Rb<sub>2</sub>SO<sub>4</sub>$ ) region, indicating that the faradic processes are mainly occurring at the  $Mo<sub>2</sub>C-$ CDC electrode surfaces. Within the low-frequency ac region  $f_{\text{low}}$  < 1 Hz for  $\text{Rb}_2\text{SO}_4$  and  $f_{\text{low}}$  < 2 Hz for  $\text{Na}_2\text{SO}_4$ , the second minimum in  $\delta$ , log f plots can be seen indicating the slow mixed kinetics processes (faradic mass transfer and adsorption step-limited processes) at the  $Mo<sub>2</sub>C-CDC$  | electrolyte interface. Within the very-low-frequency ac region,  $f < 2.4$  mHz, the phase angle values increase to zero, indicating that the very slow faradic reaction(s) determine(s) the  $Na<sub>2</sub>SO<sub>4</sub>$  or  $Rb<sub>2</sub>SO<sub>4</sub>$  aqueous solution | Mo<sub>2</sub>C-CDC interface behaviour.

The relaxation frequency for the high-frequency process  $f_{\text{high}}$  from 100 to 500 Hz for the Na<sub>2</sub>SO<sub>4</sub>-based system depends noticeably on the cycle number applied. With the increase of cycle number applied from 1 to 1000, there is no big shift of  $f_{\text{high}}$  values, but  $f_{\text{high}}$  decreases very remarkably after 2000 or 5000 cycles  $(f_{\text{high}}=120 \text{ Hz})$ . Surprisingly,  $f_{\text{high}}=$ 8000 Hz is noticeably higher for the  $Rb_2SO_4$ -based SC and there is no remarkable shift of  $f_{\text{high}}$  with the cycle number (material degradation step) studied.

Interestingly, the low-frequency processes  $f_{\text{med}}$ =0.3 Hz have systematically lower values for the  $Na<sub>2</sub>SO<sub>4</sub>$ -based SC compared with  $f_{\text{med}} = 5$  Hz for the  $Rb_2SO_4$ -based supercapacitor. At very low frequency, there are no clearly visible maxima in phase angle absolute value, log f plots, and the  $|\delta|$  are approaching zero, indicating that the very slow faradic processes are rate-determining processes for the  $Na<sub>2</sub>SO<sub>4</sub>$ - and  $Rb<sub>2</sub>SO<sub>4</sub>$ -based SC (Fig. [9a, b\)](#page-11-0). At  $f<sub>med</sub>$  and  $f<sub>high</sub>$ frequency regions, the mixed kinetics processes take place at

![](_page_12_Figure_5.jpeg)

![](_page_12_Figure_6.jpeg)

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<span id="page-13-0"></span>mesoporous–macroporous regions of  $Mo<sub>2</sub>C-CDC$  electrodes, respectively. Thus, a comparison of data with our previous work [[23\]](#page-14-0) and data in Fig. [9c](#page-11-0) shows that the  $\Delta E$  has very noticeable influence on the medium- and low-frequency process parameters.

Based on the Orazem et al. model [\[42\]](#page-14-0) within the highfrequency region, the slope of log  $\left|-\frac{Z''}{N}\right|$ , log f plot is nearly −0.85, indicating the mixed kinetics adsorption, mass transfer step and faradic processes at the  $Rb_2SO_4$  | Mo<sub>2</sub>C-CDC or Na<sub>2</sub>SO<sub>4</sub> | Mo<sub>2</sub>C-CDC interfaces. More detailed analysis of log  $\left| \nabla \right|$ , log f plots demonstrates that the slope value of log  $|-Z''|$ , log f plot for the Rb<sub>2</sub>SO<sub>4</sub>-based system is somewhat less negative  $(-0.92)$  compared with the Na<sub>2</sub>SO<sub>4</sub>-based system, indicating the more pronounced deviation of the  $Rb_2SO_4+$  $H_2O$  | Mo<sub>2</sub>C-CDC supercapacitor system from the ideally polarisable electrode model. Within the medium-frequency region 1.0<f <100 Hz, there are linear areas in log  $\left|-\frac{Z''}{N}\right|$ , log  $f$  plots with the slope nearly zero, indicating that within this frequency region, the slow faradic processes determine the kinetics behaviour of the SC under study.

The absolute values of the imaginary part of impedance strongly depend on  $\Delta E$  applied and at fixed  $\Delta E$ =2.0 V are lower for the  $Rb_2SO_4$ -based SC compared with the Na<sub>2</sub>SO<sub>4</sub>based system. With the increase of  $\Delta E$  from 1.0 to 2.0 V,  $|-Z''|$ values increase, indicating that capacitance values increase with  $\Delta E$  applied. At very low ac frequency and very high  $\Delta E = 2.0$  V, the limiting rate-determining process is the very slow faradic charge transfer step (the slope of log  $\left|-\mathbb{Z}^n\right|$ , log f plots is nearly zero for both systems under study [\[42](#page-14-0)–[44\]](#page-14-0), but at  $\Delta E = 1.0$  V, nearly ideal capacitive behaviour has been established (slope −0.95) [\[23](#page-14-0)].

The values of gravimetric series  $(C_s)$  and parallel  $(C_p)$ capacitances (Figs. [10a, b](#page-11-0)), calculated from the Nyquist plots  $(Z''=-1/(j2\pi fC_s; C_p=C_s/(1+\tan^2\delta))$ ; tan  $\delta=Z'/Z''$ ), are in a good agreement with the values of  $C_{m}$ ; cy and  $C_{CC}$  only at very low ac frequency ( $f \le 1$  mHz).  $C_s$  weakly increases in the order 1 M Na<sub>2</sub>SO<sub>4</sub>≤1 M Rb<sub>2</sub>SO<sub>4</sub> aqueous electrolyte-based supercapacitors. At higher frequency, there is a well-expressed dependence of  $C_s$  or  $C_p$  on f for the all electrolyte systems studied, caused by the faradic reactions as well as by the small ac penetration depth compared with the pore length.

It should be mentioned that the essential increase of the IRdrop inside of the microporous carbide-derived carbon electrode matrix, expressed as  $R_s$  and  $R_p$  ( $R_s = Z'$ ) and  $R_p = R_s (1 + 1/\sqrt{1 + \frac{Z_0}{Z_0}})$  $\tan^2 \delta$ , with the increase of cycle number analysed, explained by the accumulation of gaseous products and electrodeposited non-conducting intermediates onto/into the electrode surface, has been observed for both electrolytes. However, the lesspronounced dependence of  $C_s$  and  $R_s$  on  $f$  (as well as  $C_p$  and  $R_p$ ) and  $C_p/C_s \approx 1.0$  at  $f \rightarrow 0$  Hz for 1 M Rb<sub>2</sub>SO<sub>4</sub> at  $\Delta E = 1.0$  V (Fig. [11a, b](#page-12-0)) indicates that for this electrolyte, the distances of the maximal approach of cations onto the  $Mo<sub>2</sub>C-CDC$  porous system are shorter and the quicker establishment of the

adsorption equilibrium has been established at relatively moderate ac frequency  $(f_{low}=10 \text{ mHz})$ . At  $f\leq 1 \text{ mHz}$ , the ratio of  $C_p/C_s$  is very close to unity which indicates that the blocking adsorption step-limited processes at  $Mo_2C-CDC$  | 1 M  $Na<sub>2</sub>SO<sub>4</sub>$  or 1 M  $Rb<sub>2</sub>SO<sub>4</sub>+H<sub>2</sub>O$  interface are taking place at  $\Delta E$ =1.0 V. At  $\Delta E$ =2.0 V and  $f \le 1$  mHz,  $C_p/C_s$  for the  $Na<sub>2</sub>SO<sub>4</sub>$ -based system is nearly 0.8 and, for  $Rb<sub>2</sub>SO<sub>4</sub>$ , lower  $R_p/R_s$  values at  $\Delta E$ =2.0 V differently from  $\Delta E$ =1.0 V have been calculated at  $f \rightarrow 0$ , indicating that adsorption-limited faradic reduction processes are prevailing at the  $Rb_2SO_4$ based supercapacitor. At  $f \rightarrow 0$ , for the Rb<sub>2</sub>SO<sub>4</sub>-based system, somewhat higher increase of  $R_p/R_s$  has been observed explained by more effective surface blocking compared with the  $Na<sub>2</sub>SO<sub>4</sub>$ -based SC. However, more detailed time stability tests are inevitable for SC applications under development.

#### **Conclusions**

The analysis of impedance data demonstrates very complex kinetics behaviour of completed devices (adsorption and blocking adsorption, faradic, mass transfer, reduction of Na<sup>+</sup> and  $Rb<sup>+</sup>$  ions and intercalation, surface oxidation, gas adsorption, ionisation, etc.) at  $\Delta E \ge 1.5$  V. At least four different characteristic time constants dependent on the electrolyte cation composition and cell potential applied have been calculated and discussed. Preliminary non-linear least squared fitting data show that low and high  $Rb_2SO_4$ -based system frequency inductive effect (characteristic of the decomposition of cell components) must be taken into account to have a reasonable fitting of calculated Nyquist plots with experimental data.

Acknowledgments This work was supported by the Estonian Science Foundation under Project No. 9184, Estonian Ministry of Education and Research project SF0180002s08, European Regional Development Fund Project SLOKT10209T, Estonian Centre of Excellence in Research Project TK117T "High-technology Materials for Sustainable Development" and Project IUT20-13.

#### References

- 1. Conway BE (1999) Electrochemical supercapacitors: scientific fundamentals and technological applications. Kluwer Academic/Plenum Publishers, New York
- 2. Burke A (2000) J Power Sources 91:37–50
- 3. Kötz R, Carlen M (2000) Electrochim Acta 45:2483–2498
- 4. Chmiola J, Yushin G, Gogotsi Y, Portet C, Simon P, Taberna PL (2006) Science 313:1760–1763
- 5. Miller JR, Simon P (2008) Science 321:651–654
- 6. Simon P, Gogotsi Y (2008) Nat Mater 7:845–854
- 7. Salitra G, Soffer A, Eliad L, Cohen Y, Aurbach D (2000) J Electrochem Soc 147:2486–2493
- 8. Jänes A, Permann L, Arulepp M, Lust E (2004) Electrochem Commun 6:313–318
- <span id="page-14-0"></span>9. Jänes A, Kurig H, Lust E (2007) Carbon 45:1226–1233
- 10. Jänes A, Thomberg T, Lust E (2007) Carbon 45:2717–2722
- 11. Jänes A, Thomberg T, Kurig H, Lust E (2009) Carbon 47:23–29
- 12. Eliad, Salitra G, Soffer A, Aurbach D (2005) Langmuir 21:3198– 3202
- 13. Wang Y, Chen KS, Mishler J, Cho SC, Adroher XC (2011) Appl Energy 88:981–1007
- 14. Carmo M, Fritz DL, Mergel J, Stolten D (2013) Int J Hydrogen Energy 38:4901–4934
- 15. Härk E, Nerut J, Vaarmets K, Tallo I, Kurig H, Eskusson J, Kontturi K, Lust E (2013) J Electroanal Chem 689:176–184
- 16. Lee HY, Goodenough JB (1999) J Solid State Chem 144:220–223
- 17. Portet C, Lillo-Ródenas MÁ, Linares-Solano A, Gogotsi Y (2009) Phys Chem Chem Phys 11:4943–4945
- 18. Roldan S, Villar I, Ruiz V, Blanco C, Granda M, Menendez R, Santamaria R (2010) Energy Fuels 24:3422–3428
- 19. Jänes A, Lust E (2006) J Electroanal Chem 588:285–295
- 20. Jänes A, Lust E (2006) J Electrochem Soc 153:A113–A116
- 21. Thomberg T, Jänes A, Lust E (2009) J Electroanal Chem 630:55–62
- 22. Lota G, Frackowiak E (2009) Electrochem Commun 11:87–90
- 23. Eskusson J, Jänes A, Kikas A, Matisen L, Lust E (2011) J Power Sources 196:4109–4116
- 24. Khomenko V, Raymundo-Piñero E, Béguin F (2010) J Power Sources 195:4234–4241
- 25. Bichat MP, Raymundo-Piñero E, Béguin F (2010) Carbon 48:4351– 4361
- 26. Demarconnay L, Raymundo-Piñero E, Béguin F (2010) Electrochem Commun 12:1275–1278
- 27. Gao Q, Demarconnay L, Raymundo-Piñero E, Béguin F (2012) Energy Environ Sci 5:9611–9617
- 28. Damaskin BB, Petrii OA, Tsirlina G (2006) Electrochemistry. Khimia, Moscow
- 29. Frumkin AN (1987) Elektrodnye protsessy (The electrode processes). Nauka, Moscow
- 30. Conway BE (1965) Theory and principles of electrode processes. Ronald, New York
- 31. Trasatti S, Lust E, in White RE, Bockris J O'M, Conway BE (eds) (1999) Modern aspects of electrochemistry. vol 33 Kluwer Academic/Plenum Publishers, New York, p 1
- 32. Lust E, in: Bard JA, Stratman M (eds) (2002) Encyclopedia of electrochemistry, Wiley–VCH, Weinheim, pp 188–224
- 33. Frumkin AN, in: Yeager E (ed) (1961) Advances in electrochemistry and electrochemical engineering. New York, Wiley, pp 1-15
- 34. Frumkin AN (1965) J Electroanal Chem 9:173–183
- 35. Bian X, Scanlon MD, Wang S, Liao L, Tang Y, Liu B, Girault HH (2013) Chem Sci 4:1432–1441
- 36. Chizmadzhev YA, Chirkov YG, in: (Yaeger E, Bockris JOM, Conway BE . Sarangapani S (eds) (1983) Comprehensive treatise of electrochemistry. Plenum Press, New York, pp 317-91
- 37. Tõnisoo A, Kruusma J, Pärna R, Kikas A, Hirsimäki M, Nõmmiste E, Lust E (2013) J Electrochem Soc 160:A1084–A1093
- 38. Lust E, Jänes A, Sammelselg V, Miidla P (2000) Electrochim Acta 46:185–191
- 39. Lust E, Jänes A, Arulepp M (2004) J Solid State Electrochem 8:488– 496
- 40. Jänes A, Permann L, Arulepp M, Lust E (2004) J Electroanal Chem 569:257–269
- 41. Nurk G, Jänes A, Lust K, Lust E (2001) J Electroanal Chem 515:17–32
- 42. Orazem ME, Pébère N, Tribollet B (2006) J Electrochem Soc 153: B129–B136
- 43. Lasia A (2006) J Electroanal Chem 593:159–166
- 44. Macdonald JR, Johnson WB (1987) Fundamentals of impedance spectroscopy. In: Macdonald JR (ed) Impedance spectroscopy. Wiley, New York, pp 1–26