#### **ORIGINAL PAPER**



# **Studying the infuence of surface roughness with diferent shapes and quantities on convective heat transfer of fuid within nanochannels using molecular dynamics simulations**

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#### **Abstract**

**Context** In the past decade, rapid advancements in microtechnology and nanotechnology have propelled modern science and technology into the nanoscale domain, where miniaturization and high integration have found extensive applications across various felds, including industry, biotechnology, and information technology. Mechanisms of nanofuid fow and heat transfer properties have received increasing attention. In the fow and convective heat transfer of fuids at the nanoscale, the shape and dimensions of the surfaces play a crucial role. So, the main purpose of our paper is to investigate the infuence of surface roughness with diferent shapes and quantities on fuid fow and convective heat transfer. In this study, we have chosen argon atoms as the fuid and used copper atoms to simulate the nanochannel walls. In order to investigate the infuence of the shape and quantity of roughness on the convective heat transfer of fuids within nanochannels, we computed and analyzed the velocity, temperature, and density distributions of fuids inside channels with triangular, hemispherical, and rectangular roughness. Through simulation results, we found that triangular, hemispherical, and rectangular surface roughness at the same height can result in diferences in temperature and velocity of the fuid within nanochannels. With a nanochannel roughness number of 5, the temperature and velocity of the fuid at the middle position of the nano-channel for the triangular roughness increased by 6% and 25% compared to the rectangular roughness, and by 4% and 10% compared to the hemispherical roughness. The fuid temperature and velocity are highest in channels with triangular surface roughness and lowest in those with rectangular roughness. Furthermore, increasing the quantity of surface roughness decreases the temperature and velocity of the fuid within nanochannels. When the quantity of rectangular surface roughness is 5, the fluid temperature within the nanochannel decreases by 12%, and the velocity decreases by 38% compared to a roughness quantity of 1. We also found, through velocity contours, that the presence of roughness increases the local fuid velocity in the rough regions of nanochannels. Roughness also reduces the density fuctuations of the fuid near the walls within the nanochannel. Roughness signifcantly afects the heat transfer performance of the fuid during its fow, and this infuence should not be overlooked.

**Methods** In this study, molecular dynamics theory was employed, and simulations were conducted using the open-source software LAMMPS to investigate the infuence of diferent shapes and quantities of surface roughness on fuid fow within nanochannels. The model in this paper was constructed using the LAMMPS software, and the surface roughness shapes on the walls were implemented as rectangular, hemispherical, and triangular. The wall surfaces were composed of copper atoms, while the fuid consisted of argon atoms. The copper atoms were arranged in a face-centered cubic (FCC) lattice with a lattice constant of 3.615 A. Similarly, the argon atoms were arranged in a face-centered cubic (FCC) lattice with a lattice constant of 5.62 Å. The interactions between copper atoms were modeled using the EAM (Embedded Atom Method) potential, while the interactions between argon atoms were described using the LJ (Lennard-Jones) potential. The LJ potential was also employed to represent interactions between argon and copper atoms.

**Keywords** Nanochannel · Convective heat transfer · Poiseuille fow · Molecular dynamics simulation

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#### **Introduction**

In recent years, the rapid development of microtechnology and nanotechnology [[1,](#page-13-0) [2\]](#page-13-1) has made nanoscience one of the most promising scientifc felds, with increasing attention to the mechanisms that infuence nanofuid fow and heat transfer performance. The forced convection of fuids within fnite micro and nanochannels serves as an efective method for enhancing the heat dissipation of micro and nanodevices. It has widespread applications in areas such as efficient heat transfer, semiconductor chips, aerospace technology, and more. Therefore, for more efective comprehension and design of nanoscale channels, it is essential to gain a comprehensive understanding of the heat transfer characteristics of fuids within nanochannels. When fluid flow reaches the micro  $[3-5]$  $[3-5]$  $[3-5]$  and nanoscale, microscale efects emerge, and it becomes challenging to elucidate their mechanisms due to the influence of experimental conditions and measurement precision. As a result, numerical simulations remain the primary research method in this context. At the microscale, liquids can be treated as either continuous media or collections of atoms, making Molecular Dynamics Simulation (MDS) [\[6–](#page-13-4)[10\]](#page-13-5) an essential method for studying and predicting the dynamic properties of atomic structures.

Qin and Azimian et al. [\[11](#page-13-6), [12\]](#page-13-7) investigated the infuence of wall surface roughness shape and dimensions on nanofuid fow through non-equilibrium molecular dynamics. They found that changing the roughness shape afects the velocity and temperature distribution of the fuid. At the same height, triangular surface roughness in nanochannels resulted in a 17% increase in fuid velocity compared to cylindrical surfaces. The impact of triangular rough surfaces on flow behavior was greater than that of cylindrical surfaces, and an increase in roughness size led to an augmented fuid fow velocity within nanochannels. Raj et al. [\[13\]](#page-14-0) simulated the Ar–Pt system using MD and found that the hydrodynamic slip length can be adjusted by varying the roughness cycle parameter and the channel width. Li and Liu  $[14]$  $[14]$  used MDS to study the flow of fluids in nanochannels and found that the thermal properties of the fuids are afected by wall interactions and temperature. Sofos et al. [\[15](#page-14-2), [16](#page-14-3)] investigated the effect of rough wall surfaces on the fuid inside the nanochannel based on molecular dynamics and showed that roughness decreases the velocity of the fuid and leads to a decrease in the slip velocity at the wall. Zhang et al. [[17\]](#page-14-4) investigated the efect of wall roughness height on fow rate and found that when solid-liquid interaction is strong, even if the roughness height bump is small, it signifcantly afects the fow rate of fuid through the channel.

Yan et al. [\[18\]](#page-14-5) investigated the effect of external forces on Poisson lobe fow by molecular dynamics and found that increasing the external force increases the average velocity of the fuid. Semiromi and Alipour [\[19,](#page-14-6) [20](#page-14-7)] et al. investigated the velocity-temperature distribution of the fuid for diferent driving forces and found that an increase in driving force leads to an increase in the temperature and velocity of the fuid. Toghraie et al. [\[21\]](#page-14-8) studied the motion characteristics of water and copper nanofuid under Couette and Poiseuille flows. They found that an increase in the height of the nanochannel reduces the density fuctuations of the fuid near the walls.

Yao et al. [\[22–](#page-14-9)[24\]](#page-14-10) investigated the influence of roughness on nanofuid by altering the periodicity of the wall surface structure. They discovered that the combination of surface roughness morphology and the interaction with a weak wall flow has a favorable impact on enhancing heat transfer and reducing resistance. In the case of weak wall flow interaction, the fluid's drag coefficient decreased by  $27.1\%$ , and the Nusselt number decreased by 1.78%. Song et al. [[25](#page-14-11)] investigated the effect of wall roughness of sinusoidal structures with diferent periods on the convective heat transfer of fuids in nanochannels, discussed the intrinsic mechanism of the efect of nanostructures on convective heat transfer, and found that the wall of sinusoidal structures has better heat transfer with increasing amplitude and decreasing period. Marable et al.  $[26]$  $[26]$  concentrated on the effect of solid-liquid interaction, channel height, water velocity, and wall temperature on convective heat transfer. Motlagh [\[27,](#page-14-13) [28](#page-14-14)] studied convective heat transfer in copper-argon nanofuid under nanochannels by increasing and altering the shape of nanoparticles. The research revealed that increasing the number of nanoparticles enhances the heat transfer performance of nanofuids. Additionally, cylindrical-shaped nanoparticles showed a 6% higher Nusselt number compared to sphericalshaped particles.

Convection heat transfer in nanochannels is infuenced by numerous factors, including external forces, channel height, roughness, nanoparticle shape and size, wall wettability, wall material, and temperature. The quantity and shape of roughness are two crucial factors in studying the heat exchange between the wall surface and the fuid. However, in previous studies of convective heat transfer in nanochannels, roughness shapes were typically modeled as a single rectangle. There has been no investigation into whether different shapes of roughness have an efect on fuid fow and convective heat transfer. Moreover, calculating velocity and temperature profles in a single fow direction does not capture the infuence of roughness on the localized fuid within nanochannels. Therefore, we employed triangular, rectangular, and hemispherical roughness with the same height and width to investigate the impact of diferent shapes of roughness on convective heat transfer of the fuid within the channels. The roughness quantities were chosen as 2, 6, and

10. We used the molecular dynamics software LAMMPS to perform the simulation, calculated the temperature velocity and density of the fuid parameters such as a comprehensive understanding of the nano-channel fluid flow and convective heat transfer, and at the same time plotted the velocity and temperature contours to more comprehensive understanding of the channel within the channel rough areas and smooth areas of the channel fuid motion state changes.

### **Molecular dynamics simulation**

#### **Physical model and simulation detail**

This study employed the open-source molecular dynamics software known as "LAMMPS" (Large-scale Atomic/ Molecular Massively Parallel Simulator) [[29](#page-14-15)] for conducting molecular dynamics simulations. LAMMPS possesses the capability to simulate interactions and motion among atoms, molecules, and other particles. Developed by Sandia National Laboratories, this software is compatible with both single processors and large-scale parallel computing systems. To visualize the conducted simulations of convective heat transfer, the study utilized "OVITO" (Open Visualization Tool)  $[30]$  $[30]$  $[30]$ . OVITO is an open-source scientific visualization and analysis software designed specifcally for processing, analyzing, and visualizing large-scale atomic, molecular, and particle simulation data. This tool facilitates researchers in achieving a more comprehensive understanding and presentation of simulation data, thereby offering deeper scientifc insights. By integrating simulations using LAMMPS and visualization using OVITO, this study was able to rigorously investigate and present convective heat transfer phenomena in a scientifc manner, providing insights at the molecular level.

In this study, we conducted molecular dynamics simulations of convective heat transfer, utilizing a smooth nanochannel model as depicted in Fig. [1.](#page-2-0) The overall dimensions of the simulation box measured  $345.4 \times 56 \times 270 \text{ Å}^3$ . The upper and lower walls of the nanochannel were composed of copper atoms arranged in a face-centered cubic (FCC) lattice with a lattice constant of 3.615 Å. The fuid within the

nanochannel consisted of argon atoms, likewise arranged in a face-centered cubic (FCC) lattice with a lattice constant of 5.62 Å. The upper and lower walls of the smooth nanochannel have a thickness of 30  $\AA$ , and the height (H) of the channel is 204 Å. Periodic boundary conditions were applied in the X and Y directions. In the Z direction, the fuid was subjected to fxed boundary conditions as it was confned within the upper and lower wall surfaces. The interior of the nanochannel was partitioned into three distinct regions: the force region  $(0 \text{ Å})$  $\langle x \rangle < 10$  Å), the thermostat region (10 Å  $\langle x \rangle < 20$  Å), and the sample collection region (20 Å  $\lt x \lt 314$  Å). The force region was utilized for the application of an external force to the fluid, with a force magnitude of  $0.0006$  eV/Å. The thermostat region employed a Langevin thermostat to control the temperature of the fuid passing through that region at 300 K. The sample collection region was designated for the collection of fuid data. Both the fuid and the wall surfaces were maintained at a temperature of 200 K using Langevin thermostats, ensuring that the fuid within the nanochannel remained in a supercritical state, thereby preventing phase transitions. The simulation employed a timestep of 0.1 fs. The configuration of this model allowed us to conduct an in-depth investigation into the behavior of the fuid within the nanochannel and the phenomenon of convective heat transfer.

As shown in Fig. [2,](#page-3-0) we employed a roughness nanochannel model in which rectangular, half-spherical, and triangular-shaped roughness structures were introduced within the interior of its upper and lower wall surfaces to facilitate investigations under varying roughness conditions. We systematically examined the quantities of roughness, specifcally 2, 6, and 10, with the shapes of these roughness structures being triangular, half-spherical, and rectangular, The triangular and rectangular roughness structures had a width of 31.4  $\AA$  and a height of 28  $\AA$ , while the half-spherical roughness had a height of  $28 \text{ Å}$  as well. All three shapes of roughness shared the same height. In the force region, an external force along the x-direction was applied with a magnitude of 0.0006 eV/Å. After achieving temperature equilibrium within the fuid system and the wall surfaces, an analysis was conducted on the velocity, temperature, and density of the fuid inside both the smooth nanochannel and the rough nanochannel. The aim of this study was to gain

## $200k$ Wall Z Fluid  $200k$  $\mathbf{x}$ Wall

<span id="page-2-0"></span>

#### <span id="page-3-0"></span>**Fig. 2** Nanochannel models with varying roughness



a deeper understanding of the variations in fuid behavior under diferent roughness conditions.

<span id="page-3-1"></span>

#### **Governing equations**

In this study, MDS was conducted based on Newton's second law, employing the Verlet method to determine fundamental dynamical parameters such as the positions, velocities, and interatomic forces of atoms. Newton's second law [\[31\]](#page-14-17) can be expressed by the following equation:

$$
F_i = \mathbf{m}_i a = m_i \frac{d v_i}{dt} = m_i \frac{d^2 r_i}{dt^2}
$$
 (1)

where  $F_i$  represents the external force acting on argon atoms,  $m<sub>i</sub>$  represents the mass of argon atoms,  $v<sub>i</sub>$  represents the velocity of argon atoms, and  $r<sub>i</sub>$  represents the position of argon atoms. Due to the complex interactions within the molecular system, the Velocity-Verlet algorithm was employed to handle the equations of motion. The formula for the Velocity-Verlet algorithm [\[29,](#page-14-15) [32\]](#page-14-18) is expressed as follows:

$$
r_i(t_0 + \Delta t) = r_i(t_0) + \Delta t v_i(t_0) + (\Delta t)^2 f_i(t_0) / 2m_i
$$
  
\n
$$
v_i(t_0 + \Delta t) = v_i(t_0) + \Delta t (F_i(t_0) + F_i(t_0 + \Delta t)) / 2m_i
$$
 (2)

where  $r_i(t_0)$  and  $v_i(t_0)$  represent the initial position and initial force, while  $r_i(t_0 + \Delta t)$ ,  $v_i(t_0 + \Delta t)$ , and  $F_i(t_0 + \Delta t)$  denote the fnal velocity, fnal position, and fnal force, respectively.

The LJ potential/ij  $[6]$  $[6]$  is given by

$$
\phi(r) = 4\varepsilon \left( \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right) \tag{3}
$$

*r* is the distance between two molecules.  $\sigma$  is a molecular length scale, and *ε* is an interaction strength parameter. The cutoff radius interaction is set to 1.0 nm. In addition, the EAM potential [[33](#page-14-19), [34](#page-14-20)] was used between metal atoms. All the LJ potential parameters are listed in Table [1.](#page-3-1) The total energy  $E_i$  of an atom *i* is given by

 $E_i = F_\alpha \left( \sum_{j \neq i} \rho_\beta(r_{ij}) \right) + \frac{1}{2} \sum_{j \neq i} \mathcal{O}_{\alpha\beta}(r_{ij})$  (4) 1  $\frac{1}{2}\sum_{j\neq i}\mathcal{O}_{\alpha\beta}(r_{ij})$ 

where  $r_{ii}$  is the distance between atom *i* and atom *j*,  $\rho_{\beta}$  indicates the contribution of the b-type atom *j* at the position of atom *i* to the electron charge density.  $F_a$  is an embedding function that represents the energy required to place atom *i* in the electron cloud.  $\mathcal{O}_{\alpha\beta}$  refers to the potential energy of the interaction between atom *i* and atom *j*.

The Lorentz-Berthelot mixing law [\[35\]](#page-14-21) is employed for interactions between metal and non-metal atoms, as shown below:

$$
\sigma_{Cu-O} = \frac{\sigma_{Cu} + \sigma_O}{2} \tag{5}
$$

$$
\varepsilon_{Cu-O} = \sqrt{\varepsilon_{Cu}\varepsilon_{O}}\tag{6}
$$

The local average fluid temperature  $T_m(x)$  at different positions along the nanochannel is

$$
T_m(x) = \frac{\int_0^H c \rho u_x(x, z) T(x, z) dz}{\int_0^H c \rho u_x(x, z) dz}
$$
(7)

*C* represents the specifc heat capacity at constant pressure for the fluid, *p* denotes fluid density,  $u<sub>x</sub>(x, z)$  stands for the fluid's velocity in the *x*-direction, and  $T_m(x)$  represents the temperature distribution of the fuid in the x-direction. The local heat transfer coefficient  $h(x)$  can be determined using Eq. ([8\)](#page-4-0), and the local Nusselt number [\[36](#page-14-22)] can be computed using Eq. [\(9](#page-4-1)). Due to solid-liquid interactions, the wall surface signifcantly infuences the near-wall fuid temperature. In order to mitigate the impact of the wall surface on the fuid temperature, the local Nusselt number can be calculated using Eq. ([10](#page-4-2)):

$$
h(x) = \frac{\lambda}{\left(T_m(x) - T_w\right)} \frac{\partial T}{\partial z}\Big|_{z=w} \tag{8}
$$

$$
Nu_x = \frac{h(x)D_h}{\lambda} \tag{9}
$$

$$
Nu_x = \frac{D_h}{\left(T_m(x) - T_w\right)} \frac{\partial T}{\partial z}\Big|_{z=w} \tag{10}
$$

In the equation,  $\lambda$  represents the thermal conductivity of the fuid, *∂T/∂Z* |*z=w* signifes the temperature gradient at the fluid-wall interface, and  $Dh = 2h$ .

#### **System balance**

In the simulations of convective heat transfer, the fuid temperature was initially controlled at 200 K under the NVT ensemble. Subsequently, the force region and the sample collection region were subjected to the NVE ensemble, while the thermal bath region was regulated using a Langevin <span id="page-4-0"></span>thermostat to maintain a stable temperature of 300 K. The wall surfaces were also temperature-controlled at 200 K using the Langevin thermostat. So the frst step is to make sure that the temperature of each part of our convective heat transfer system is in equilibrium.

<span id="page-4-2"></span><span id="page-4-1"></span>Figure [3](#page-4-3) illustrates the temperature distribution of the fuid and the wall surfaces, both for the smooth wall and when there are ten roughness structures on the wall, over a certain number of timesteps. As shown in Fig. [3](#page-4-3), when both the wall and fuid temperatures reached 200 K, the application of external forces to the fuid resulted in an increase in fuid temperature. After approximately 200,000 time steps, the fuid temperature reached a stable state. Specifcally, the temperature of the smooth plane was 230 K, while the fuid temperatures inside the channels with triangular and hemispherical roughness structures were approximately 235 K, and within the channel with rectangular roughness, the fuid temperature was around 225 K. The temperature of both the roughness wall and the smooth wall remained consistently around 200 K. In the subsequent steps, we will further report information regarding the temperature, velocity, and density of



<span id="page-4-3"></span>**Fig. 3** Nanochannel wall temperature and fuid temperature at diferent timestep

the fuid inside the nanochannel to gain a deeper insight into the nature of the heat transfer process. These results will aid in our comprehensive understanding of the behavior and properties of the fuid under diferent roughness conditions.

### **The temperature distribution of the fuid in the nanochannel**

In our study, we computed the temperature distribution of the fluid in nanochannel models with varying crosssectional shapes and numbers of roughness structures. The temperature distribution is a critical parameter for assessing the influence of different wall shapes and quantities of roughness on temperature under the same external forces. To calculate the temperature distribution of the entire nanofluid, we divided the nanochannel model into 45 segments along the z-directions, with each segment having a length of 6 Å. At each timestep, the average temperature of each atom in different segments was computed. In the force region, an external force of 0.0006 eV/Å was applied to the fluid, and after 4,000,000 timesteps, the temperature distribution along the z-direction was analyzed. As shown in Fig. [4,](#page-6-0) temperature profiles of the fluid under different roughness shapes and quantities indicate that, whether on a smooth wall or with various roughness conditions, the temperature of the fluid at the center of the channel is consistently the highest. Clearly, this suggests that as the fluid moves away from the wall surface, its temperature increases. This is because there is a solid-liquid interaction force between the wall surface and the fluid. As the fluid moves farther away from the wall surface, the interaction force between the solid and liquid weakens. Consequently, the fluid velocity increases, and there is greater viscous friction between fluid atoms, leading to higher temperatures. These observational findings provide crucial insights into the influence of wall shapes and roughness on the temperature distribution of the fluid.

In investigating the influence of different quantities of roughness on fluid temperature in convection heat transfer, we ensured that each roughness shape was identical. As shown in (a), (b), and (c) in Fig. [4,](#page-6-0) this is attributed to the increase in the number of roughness features, which results in an enlarged heat exchange surface area between the wall and the fluid, consequently leading to a further reduction in the temperature of the fluid within the nanochannel. Therefore, we can conclude that the effect of varying quantities of roughness on the temperature of the fluid in convective heat transfer is substantial.

When investigating the influence of different shapes of roughness on fluid temperature in convective heat transfer, we ensured that the height and quantity of each type of roughness were consistent. As shown in Fig. [4d](#page-6-0), e, f, when the roughness quantity is 2, the temperature distribution of the fluid does not exhibit significant differences. The central fluid temperatures in the smooth plane, triangle, hemisphere, and rectangle nanochannels are 350 K, 343 K, 341 K, and 342 K, respectively. However, when the roughness quantity increases to 6 and 10, the situation begins to differ significantly. In these cases, the central fluid temperatures in the smooth plane, triangle, hemisphere, and rectangle nanochannels are 350 K, 333 K, 327 K, and 322 K, and 350 K, 325 K, 311 K, and 304 K, respectively. Observably, the central fluid temperature in the triangular roughness nanochannel is higher than that in the hemispherical and rectangular roughness channels, and with an increase in the quantity of roughness, the temperature difference becomes increasingly pronounced. This is because there is a solid-liquid interaction force between the wall surface and the fluid. As the fluid moves farther away from the wall surface, the interaction force between the solid and liquid weakens. Consequently, the fluid velocity increases, and there is greater viscous friction between fluid atoms, leading to higher temperatures. Clearly, this suggests that as the fluid moves away from the wall surface, its temperature increases. These observational findings provide crucial insights into the influence of wall shapes and roughness on the temperature distribution of the fluid. Therefore, it can be concluded that the influence of different shapes of roughness on the fluid temperature in convective heat transfer is significant, and this effect becomes more pronounced with an increase in the quantity of roughness.

Temperature contours can offer a more intuitive way to help us comprehend the influence of various wall roughness shapes and quantities on temperature transfer. To calculate the temperature contour of the entire fluid, we divided the nanochannel model into 5762 intervals in the z and x directions, each interval having a length of 4 Å. As shown in Fig. [5,](#page-7-0) this is the temperature distribution of the fluid in the nanochannel with different shapes and quantities of wall surface roughness. When external force is applied to the fluid, it flows along the channel direction. The fluid temperature initially undergoes a thermal development phase before eventually reaching a stable thermal equilibrium state. In the thermal equilibrium state, the fluid temperature initially rises, then decreases, and ultimately stabilizes. When roughness with the same shape on the wall is present, an increase in the number of roughness elements leads to an increased



<span id="page-6-0"></span>**Fig. 4** The temperature distribution of the fuid along the z-direction in the nanochannel with diferent amount and shape of roughness

contact area between the wall and the fluid, resulting in a faster decrease in temperature throughout the nanochannel. When the wall roughness number is the same, the temperature distribution of each part of the fluid in the nanochannel with rectangular roughness is significantly lower than that with hemispherical roughness and triangular roughness. The temperature distribution of the parts of the fluid in the nano-channel with hemispherical roughness is smaller than that with triangular roughness. From this, we can conclude that the shape and quantity of roughness are significant factors affecting heat transfer in nanochannels. Compared to smooth surfaces, increasing



<span id="page-7-0"></span>**Fig. 5** The temperature contours of the fuid within the nanochannel for diferent amount and shape of roughness



<span id="page-7-1"></span>**Fig. 6** The Nusselt number of the fuid in the nanochannel

roughness reduces the temperature of the fluid within nanochannels. Furthermore, the shape of the roughness also significantly influences fluid temperature, which is a factor that cannot be overlooked.

### **The variation of local Nusselt numbers**

This section calculates the local Nusselt numbers for both the smooth plane and roughness with a value of 10, as shown in Fig. [6](#page-7-1). The local Nusselt coefficient is higher at the inlet and gradually decreases along the fow direction. The Nusselt coefficient of the fluid within the roughness wall is greater than the Nusselt coefficient on the smooth wall. This is because the roughness increases the contact area between



<span id="page-8-0"></span>**Fig. 7** The velocity distribution of the fuid along the z-direction in the nanochannel with diferent amount and shape of roughness

the wall and the fuid, and the gaps in the roughness hinder the flow of atoms, thereby enhancing heat transfer. In addition, the reason that the Nusselt number of rectangular roughness is larger than that of hemispherical and triangular is because the gap of rectangular roughness makes it less likely for atoms to spill out of the gap as compared to triangular and hemispherical, so the number of adsorbed fuid atoms is higher and the heat transfer performance is better.

#### **The velocity distribution of the fuid in the nanochannel**

The velocity distribution of the fuid in the nanochannel is an important parameter for assessing fow characteristics. In this part of our study, to evaluate the behavior of argon atoms in nanochannels with roughness, we computed the velocity distribution of the fuid in both smooth-walled nanochannels and nanochannels with varying roughness shapes and quantities. In the force region, an external force of 0.0006 eV/Å was applied to the fuid, and the velocity distribution of the fuid along the z-direction was analyzed after 4,000,000 timesteps.

Figure [7](#page-8-0) illustrates the velocity distribution of the fluid along the flow direction in both smooth and rough wall surfaces. As shown in Fig. [7a](#page-8-0), b, c, when the roughness shapes are the same, increasing the roughness quantity leads to a decrease in the fluid velocity compared to the smooth wall surface. This is because the increase in roughness leads to an increase in the contact area between the wall and the fluid, resulting in more solid-liquid interactions near the wall surface. Additionally, the gaps in the roughness accommodate fluid atoms, impeding fluid flow. Due to the presence of fluid viscosity, the decrease in fluid velocity near the wall results in a reduction in fluid velocity at the center of the channel.

As shown in Fig. [7](#page-8-0)d, e, f, diferent shapes of wall roughness also lead to diferences in the velocity of the fuid, and these diferences become more and more pronounced as the amount of roughness increases. The velocity of the fuid inside triangular roughness channels is greater than that in the spherical and rectangular roughness channels. This is because the contact area between the triangular roughness and the fuid is smaller than that of rectangular and



<span id="page-9-0"></span>**Fig. 8** The velocity contours of the fuid within the nanochannel for diferent amount and shape of roughness

hemispherical roughness. The solid-liquid interaction force on the fuid is reduced, and the triangular roughness has a larger opening at the top of the gap compared to the rectangular gap, making it easier for atoms to flow out of the gap. Therefore, it can be concluded that the shape and quantity of roughness signifcantly infuence the fuid fow velocity within nanochannels, which is a factor that cannot be overlooked.

Velocity contours provide a more intuitive representation of the infuence of diferent wall shapes and quantities on the flow velocity. Figure [8](#page-9-0) illustrates the velocity distribution of the fuid in both smooth and roughened nanochannels. In a smooth channel, the velocity distribution of the fuid is in a state of full development and the velocity remains constant along the fow direction. However, unlike the smooth wall, in the presence of roughness under constant external force, the existence of roughness reduces the cross-sectional area in the middle of the nanochannel, increasing pressure, resulting in higher fuid velocity in the central portion of the nanochannel compared to other regions. Compared to smooth walls, the presence of roughness in nanochannels in varying shapes and quantities leads to changes in fuid velocity at diferent positions within the channel. From Fig. [8](#page-9-0), it can be observed that as the quantity of roughness increases, the fuid velocity within the channel decreases. Under the same quantity of roughness, triangular and hemispherical roughness result in higher fuid velocities compared to rectangular roughness. Furthermore, as the quantity of roughness increases, this trend becomes more pronounced. This is because at the same height, rectangular roughness has a larger contact area with the fuid compared to triangular and hemispherical shapes.

For a more intuitive understanding of the impact of roughness on the fuid's velocity in the x-direction, as shown in Fig. [9,](#page-10-0) we computed the fuid velocity along the x-direction. When the fuid passes through the roughness region, the velocity increases. To gain a more intuitive understanding of the infuence of roughness on the fluid's velocity in the x-direction, as depicted in Fig. [9,](#page-10-0) we conducted calculations of the fuid's velocity along the x-direction. When the fuid passes through the roughness region, the velocity increases.

In summary, the shape and quantity of roughness are signifcant factors infuencing heat transfer in nanochannels. Compared to smooth surfaces, roughness can reduce the velocity of the fuid within nanochannels. The velocity distribution of the fuid within nanochannels with rectangular roughness is lower than that of triangular and hemispherical roughness. Furthermore, the quantity of roughness signifcantly infuences the fuid velocity, with an increase in roughness quantity leading to a reduction in fuid velocity. Under the same external force conditions, the presence of roughness also increases the local fuid velocity.



(c). The velocity distribution for ten quantities of triangular, half-spherical, and rectangular roughn

<span id="page-10-0"></span>**Fig. 9** The velocity distribution of the fuid along the x-direction in the nanochannel with diferent amount and shape of roughness

## **The density distribution of the fuid in the nanochannel**

In the fnal step of this study, we analyzed the density distribution of the fuid within the nanochannel. The fuid density distribution is a crucial parameter that aids in assessing the impact of diferent surface roughness on the fuid fow and heat transfer within nanochannels. Figure [10](#page-11-0) illustrates the density distribution of fuid along the fow direction within nanochannels with varying shapes and quantities of surface roughness. Due to the interaction forces between solid and fuid, and the relatively weaker interaction forces among



<span id="page-11-0"></span>**Fig. 10** The density distribution of the fuid along the x-direction in the nanochannel with diferent amount and shape of roughness

argon atoms compared to the interaction forces between argon atoms and the wall surface, the fuid density is higher at smooth and rough wall surfaces, resulting in lower fuid velocity. In the case where the roughness shapes are the same but differ in quantity, as observed from Fig. [10,](#page-11-0) due to the uneven contact area between roughness and the fuid, the fuid density at the smooth wall position is higher than at the bottom of the roughness. As shown in Fig. [10a](#page-11-0), b, c, with an increase in the quantity of roughness, the surface area at the top of the roughness increases, leading to an increase in fuid density at the top of the roughness. Furthermore, in the simulation, the increase in roughness also results in an increase



<span id="page-12-0"></span>**Fig. 11** The density contours of the fuid inside the nanochannels for diferent amount and shape of roughness

in fuid density in the middle of the channel. From the e, d, and f sections of Fig. [10](#page-11-0), it can be observed that due to the larger top surface area of rectangular roughness compared to triangular and hemispherical roughness, the fuid density at the top of rectangular roughness is greater than that of triangular and hemispherical roughness. In the simulations, due to the larger volume and smaller gaps of rectangular and hemispherical roughness, it results in higher fuid density within the channel compared to triangular roughness and smooth walls.

Density contour provide a more intuitive representation of the infuence of diferent wall shapes and quantities on fuid density distribution (Fig. [11\)](#page-12-0). When the scale reaches the micro-nanoscale, fuids can be considered to be composed of atoms. Since the distribution of atoms in a fuid is disordered, it is susceptible to external forces and solid-liquid interaction forces, resulting in an inhomogeneous density distribution. The atoms at diferent positions in the nanochannel experience varying degrees of external forces due to their diferent distances from the force region. As a result, the density distribution of the fuid exhibits a parabolic layered structure, with higher fuid density near the walls. In the presence of roughness on the wall, due to the resistance posed by the rough surface to fuid fow, fuid atoms tend to accumulate at the bottom of the roughness and within the gaps, making it difficult

for them to overflow. The forces exerted by rectangular and hemispherical roughness on the fuid are greater than those exerted by triangular roughness and a smooth surface. With increasing wall roughness, the fuid density near the wall increases, while the velocity of the fuid at the wall decreases.

## **Conclusion**

In our study, the fow and convective heat transfer characteristics of fuids in nanochannels with diferent roughness were investigated using the molecular dynamics software LAMMPS. Compared to studying only a single roughness and statistically analyzing velocity-temperature curves along a single fow direction, which cannot refect the localized impact of diferent roughness within nanochannels, we generated velocity and temperature contours to gain a more comprehensive understanding of the changes in fuid motion within the rough and smooth areas of channels. By analyzing these contours and curve plots, the following conclusions were drawn:

- 1. In the study of fuid fow and convective heat transfer, the temperature and velocity of the fuid in the triangular channel were reduced by 5% and 36%, hemispherical by 10% and 54%, and rectangular by 10% and 62%, respectively, in the case of a wall roughness number of 10 as compared to the case of roughness 2. Therefore, the efect of the shape of the roughness within the nanochannel on the fuid motion should be fully considered.
- 2. For roughness quantities of 10, 6, and 2, triangular roughness in nanochannels exhibited higher fuid temperatures and velocities compared to hemispherical roughness by 4%, 8%, and 1.4%, 2%, and 0.2%, 3%, and compared to rectangular roughness by 6%, 25%, and 8%, 12%, and 0.5%, 5%. In channels with triangular roughness, the fuid velocity and temperature are the highest, while in channels with rectangular roughness, they are the lowest. Therefore, it is crucial to thoroughly consider the infuence of diferent shapes of roughness within nanochannels on fuid fow and heat transfer.
- 3. When the external force applied to the fluid at the entrance of the nanochannel remains constant, roughness causes the area through which the fuid passes inside the channel to decrease, resulting in an increase in pressure. This increase in pressure leads to an increase in fuid velocity as it passes through the rough region.
- 4. In the case of a roughness quantity of 10, the Nusselt numbers for rough nanochannels are superior to those for smooth nanochannels. Among them, the Nusselt number for the rectangular shape is the highest, demonstrating the most outstanding heat transfer performance.

5. Compared to smooth channels, surface roughness reduces the density of the fuid near the wall surface because there are fewer fuid particles near the roughness surface.

**Author contributions** YL: writing—review and editing, investigation. CC: methodology, software.

**Data Availability** The data and materials used in this study are available upon request.

#### **Declarations**

**Ethical approval** In this study, no ethical approval was required as it did not involve human participants or animals.

**Competing interests** The authors declare no competing interests.

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