## **A clarification note about hitting times densities for Ornstein-Uhlenbeck processes**

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**Abstract.** In this note, we point out that the formula given in the correction note by Leblanc et al. [5] for the distribution of the first hitting time of  $b$  by an Ornstein– Uhlenbeck process starting from a is only true in case  $b = 0$ .

**Key words:** Hitting time, Ornstein-Uhlenbeck process

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Let  $(U_t, t \ge 0)$  be an Ornstein-Uhlenbeck process with parameter  $\lambda > 0$ , starting from  $a \in \mathbb{R}$ , that is the solution of:

$$
U_t = a + B_t - \lambda \int_0^t U_s \, ds = e^{-\lambda t} \left( a + \int_0^t e^{\lambda s} \, dB_s \right), \ t \ge 0,
$$
 (1)

where  $(B_t, t \ge 0)$  denotes a Brownian motion starting from 0.

Below, we give an expression of the density of  $T_b^U = \inf\{t : U_t = b\}$ , for  $b \in \mathbb{R}$ , in terms of some integrals involving the BES<sup>(3)</sup> bridges of lengths  $t \ge 0$ , starting at  $(b - a)$ , for  $b > a$ , and conditioned to end at 0 at time t. It is seen on this expression that, in the discussion in Leblanc et al. [5], one term has been omitted, which explains why Formula (3) in Leblanc et al. [5] is incorrect, and how to correct it, at least in terms of  $BES^{(3)}$  bridges.

For general discussions of first hitting times of Ornstein-Uhlenbeck processes we refer to Breiman [2], Horowitz [4] and Borodin-Salminen [1].

Denote by  $Q_a^{(\lambda)}$  the law of U, solution of (1), and by  $W_a = Q_a^{(0)}$  the law of  $\{(a + B_t), t \ge 0\}$ , both laws being defined on the canonical space  $C(\mathbb{R}_+, \mathbb{R})$ , where  $X_t(\omega) = \omega(t)$ , and  $\mathcal{F}_t = \sigma\{X_s, s \leq t\}$ . We denote  $T_b = \inf\{t : X_t = b\}$ .

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The absolute continuity relationship:

$$
Q_a^{(\lambda)}\Big|_{\mathcal{F}_t} = \exp\left[-\frac{\lambda}{2}(X_t^2 - a^2 - t) - \frac{\lambda^2}{2} \int_0^t X_s^2 ds\right] \cdot W_a|_{\mathcal{F}_t} \tag{2}
$$

is well-known (see, e.g. Yor [7], Chapter 2). It also holds with  $t$  being replaced by T, any stopping time being assumed to be finite both under  $Q_a^{(\lambda)}$  and  $W_a$ . Consequently, the following holds:

$$
Q_a^{(\lambda)}(T_b \in dt) = \exp\left[-\frac{\lambda}{2} \left(b^2 - a^2 - t\right)\right] W_a \left[\exp\left(-\frac{\lambda^2}{2} \int_0^t X_s^2 ds\right); T_b \in dt\right].
$$
\n(3)

In Leblanc et al. [5], the authors use the (obvious, but important) fact that  $(X_s, s \geq 1)$ 0) under  $W_a$  is distributed as  $(b - X_s, s \ge 0)$  under  $W_{(b-a)}$ . Thus, (3) may be written as:

$$
Q_a^{(\lambda)}(T_b \in dt) =
$$
  
exp $\left[-\frac{\lambda}{2} (b^2 - a^2 - t)\right]$   $W_{b-a} \left[\exp\left(-\frac{\lambda^2}{2} \int_0^t (b - X_s)^2 ds\right); T_0 \in dt\right]$ . (4)

Now, recall that for  $c > 0$ , under  $W_c$ , the process  $(X_s, s \leq T_0)$  conditioned by  $(T_0 = t)$ , is a BES<sup>(3)</sup> bridge of length t, starting at c, and ending at 0, whose law we denote by  $P_c^{(3)}(\cdot | X_t = 0)$ . Thus, denoting  $c = |b - a|$ , we obtain:

$$
W_{b-a} \left[ \exp\left( -\frac{\lambda^2}{2} \int_0^t (b - X_s)^2 ds \right); T_0 \in dt \right]
$$
  

$$
\int E_c^{(3)} \left[ \exp\left( -\frac{\lambda^2}{2} \int_0^t (b - X_s)^2 ds \right) | X_t = 0 \right] W_c(T_0 \in dt), \text{ if } b > a,
$$
  
(5)

$$
= \begin{cases} E_c^{(3)} \left[ \exp\left( -\frac{\lambda^2}{2} \int_0^t (b - X_s)^2 ds \right) | X_t = 0 \right] W_c(T_0 \in dt), & \text{if } b > a, \\ E_c^{(3)} \left[ \exp\left( -\frac{\lambda^2}{2} \int_0^t (b + X_s)^2 ds \right) | X_t = 0 \right] W_c(T_0 \in dt), & \text{if } b < a \end{cases}
$$
(6)

Now, we remark that, for  $b \neq 0$ , the two conditional expectations  $E_c^{(3)}$   $\left[\exp\left(-\frac{\lambda^2}{2}\int_0^t (b \mp X_s)^2 ds\right) |X_t = 0\right]$  differ from that written in Leblanc et al. [5] on top of p.111, where we find instead:

$$
E_c^{(3)}\left[\exp\left(-\frac{\lambda^2}{2}\int_0^t X_s^2 ds\right)|X_t=0\right],\tag{7}
$$

which is known to be equal to (see Pitman and Yor [6], or Yor ([7], Formula (2.5) for  $\delta = 3$ :

$$
\left(\frac{\lambda t}{\sinh(\lambda t)}\right)^{\frac{3}{2}} \exp\left[-\frac{c^2}{2t}(\lambda t \coth(\lambda t) - 1)\right]
$$

Thus, to summarize, we first obtain, by continuing (3) and (7)

$$
Q_a^{(\lambda)}(T_0 \in dt) = |a| \frac{\exp(\lambda a^2/2)}{\sqrt{2\pi}} \exp\left[\frac{\lambda}{2}(t - a^2 \coth(\lambda t))\right] \left(\frac{\lambda}{\sinh(\lambda t)}\right)^{\frac{3}{2}} dt
$$
\n(8)

as found in a number of papers, e.g., Yor ([8], Exercise p. 56), Elworthy et al. [3]. Secondly, the knowledge of the density of  $Q_a^{(\lambda)}(T_b \in dt)/dt$ , for all a and b is equivalent to the knowledge of the joint Laplace transform of  $(\int_0^t X_s ds, \int_0^t X_s^2 ds)$ under the BES<sup>(3)</sup> bridges laws  $P_c^{(3)}(\cdot|X_t=0)$ ; we do not solve this problem here, but simply point out that the Formula (7) above, found in Leblanc et al. [5], should be replaced by the correct Formula (6) in order to make (4) somewhat more explicit.

The error in Leblanc et al. [5] is to identify

$$
W_a \left[ \exp \left( -\frac{\lambda^2}{2} \int_0^t X_s^2 ds \right) | T_b = t \right]
$$

with (7), which is only true in the case  $b = 0$ .

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