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# **Temperature‑Induced Ductile–Brittle Transition in Porous Carbonates and Change in Compaction Band Growth Revealed by 4‑D X‑Ray Tomography**

**Xiao Chen1 · Klaus Regenauer‑Lieb1,2 · Hamid Roshan1**

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## **Abstract**

Deformation bands featuring localised material failure are ubiquitous in nature. They form important fow barriers and reduce/compartmentalise fuid fow in oil/gas/water reservoir rocks. Moderate temperature changes have been observed to play a fundamental role in the formation and style of these bands, but the mechanisms underpinning these changes are often obscure. Here, we show compaction experiments of highly porous limestone from Mt Gambier, Australia, with chemically non-reacting gaseous (helium) and fuid (kerosene) pore fuids. Gas-flled limestones showed a lower static elastic stifness than fuid-flled specimens. The discrepancy in elastic modulus is particularly noted at high temperature. This indicates the important efect of temperature-sensitive compressibility of gas-flled pores versus nominally incompressible fuid-flled pores. A moderate temperature rise from 25 to 80 °C also led to a sharp transition in compaction band growth from dominantly ductile difuse band growth at low temperatures to prevailing brittle growth at higher temperatures. We attribute this change to a temperature-sensitive change in micro-mechanism from rate-sensitive calcite twinning at room temperature to activation of a near-ideal plastic Peierls mechanism at 80 °C. The inverse-to-normal brittle–ductile transition is documented by time-lapse X-ray CT micro-tomographic images and associated digital image and volume image correlation techniques.

## **Highlights**

- 4D time-lapse triaxial experiments on highly porous carbonates reveal thermomechanical and thermohydromechanical couplings.
- Gas saturated specimens show higher yield stress and lower elastic modulus when compared to fuid-saturated specimens.
- A changeover from ostensibly ductile to a dominantly brittle micro-deformation mechanisms is encountered when raising the temperature from 25 to 80 °C.

**Keywords** Poromechanics · Fluid compressibility efect · Wave propagation · Digital image analysis · Digital volume correlation · Permeability

 $\boxtimes$  Klaus Regenauer-Lieb klaus@curtin.edu.au

# **1 Introduction**

Carbonates are one of the major sedimentary formations widely distributed in the near-surface and upper layers of the Earth's crust. These porous rocks form one of the most considerable hydrocarbon resources accounting for almost 60% of the conventional world's hydrocarbon reserves (Sayers [2008\)](#page-22-0). These hydrocarbon reservoirs are also amongst the best candidate storage facilities for geological sequestration of carbon dioxide when depleted (Lackner [2002](#page-20-0); Paluszny et al. [2020](#page-21-0); Wawersik et al. [2001](#page-22-1)). However, the subsurface

 $1$  School of Mineral and Energy Resources Engineering, UNSW Sydney, Kensington, Sydney, NSW, Australia

Western Australian School of Mines: Minerals, Energy and Chemical Engineering, Curtin University, Perth, WA, Australia

engineering applications of carbonates largely depend on their coupled mechanical, thermal, and chemical behaviours, which can strongly influence the fluid flow properties (Hadgu et al. [1995](#page-20-1); Kümpel [2012](#page-20-2); Zimmerman [2000](#page-23-0)). We target in our investigation the low-temperature regime of up to 80 °C using novel experimentation. Of particular interest are the microstructural response, the possible change of the internal fluid pathways, and the overall permeability effect in both micro- and macro-scale. The analysis will identify the possibility of forming internal microstructural changes such as compaction bands, shear bands, pore collapse, and fractures as a function of temperature and fuid type (gas versus liquid).

### **1.1 Low‑Temperature Efect on Carbonates**

A critical mechanical-related feature of the porous carbonate rocks is the stress/pressure-sensitivity of their skeleton. This enables an early transition from dilation to compaction failure, because the calcite can easily form mechanical twins or causes dislocation glide at relatively low-stress and temperature conditions (Burkhard [1993](#page-19-0); Ferrill et al. [2004](#page-20-3); Turner et al. [1954](#page-22-2)). These localised deformations can occur over a vast range of size scales from the micro-scale grains to large-scale faults extending over tens of kilometres (Bésuelle and Rudnicki [2004](#page-19-1)). The typical deformation bands can be characterised as dilation bands, shear bands, compaction bands, or a combination of several band types, where these strain localisation bands are always associated with instability phenomena (Aydin et al. [2006;](#page-19-2) Fossen et al. [2007\)](#page-20-4). Previous work on temperature sensitivity in rocks focused on thermal pressurisation of the pore fuids and solid phase, resulting in inelastic deformation inside a fault zone (Ghabezloo et al. [2009](#page-20-5); Hu et al. [2020](#page-20-6); Kümpel [1991](#page-20-7); Regenauer-Lieb et al. [2020](#page-21-1); Rice [2006](#page-21-2); Sulem et al. [2007](#page-22-3)). This effect is unlikely to play a significant role in highly porous carbonates due to the inherent high permeability. It was also reported that bulk, pore, and rock-solid compressibility decrease with increasing temperature (Somerton [1992\)](#page-22-4). Temperature can increase and activate chemical reactions such as mineral dehydration, carbonate decomposition, and dissolution or precipitation, which control the deformation bands' initiation and evolution. Heap et al. [\(2009\)](#page-20-8) reported that a moderated temperature increase from 20 to 75 °C could enhance the rate of stress corrosion cracking, thus reducing the short-term strength and time-dependent creep behaviour of three sandstones. However, a few studies focus on low-temperature carbonates' behaviour, especially for highly porous carbonates stable at low to an intermediate depth in the crust, such as the Mount Gambier limestone in South Australia. Most laboratory studies on such carbonate rocks have focused on ambient temperatures, and the experimental data on the changes in transport properties that accompany deformation are scarce. This is especially the case for understanding the micro-scale phenomena driving the macro-scale observations.

# **1.2 Instabilities and Micro‑deformation Mechanisms in Carbonate Reservoir Rocks**

Shear bands are the most common deformation bands described in the geological literature, dominated by shear displacement with or without some additional compaction or dilation processes. Shear bands play a signifcant role in many engineering structures failures, including earthquakes, landslides, and slope instabilities (Fossen et al. [2018;](#page-20-9) Sulem and Stefanou [2016;](#page-22-5) Tondi et al. [2006\)](#page-22-6). On the other hand, the formation of compaction bands corresponds to a tabular structure of localised compressive deformation that forms perpendicular to the maximum compressive principal stress without obvious shear strain (Charalampidou et al. [2011](#page-19-3); Issen and Rudnicki [2000](#page-20-10); Olsson [1999\)](#page-21-3). Compaction bands are usually characterised by a significant pore collapse inducing permeability reduction in the direction perpendicular to the compaction bands, which controls the fuid fow in the reservoir rocks.

Compaction bands were frst reported through feld observation in Jurassic Aztec sandstone of south-eastern Nevada (Hill [1993](#page-20-11)) and later reported in other feld-based studies (Mollema and Antonellini [1996](#page-21-4); Rustichelli et al. [2012](#page-22-7)). Tondi et al. [\(2006\)](#page-22-6) and Rustichelli et al. ([2012\)](#page-22-7) reported narrow compaction bands which oriented oblique and parallel to bedding in naturally deformed carbonates. In addition to feld observations, compaction bands have been reported in laboratory experiments (Baud et al. [2015](#page-19-4); Baxevanis et al. [2006](#page-19-5); Han et al. [2013](#page-20-12); Leuthold et al. [2021](#page-21-5); Oka et al. [2006](#page-21-6); Tembe et al. [2008](#page-22-8); Vajdova and Wong [2003\)](#page-22-9) and have been simulated in numerical–theoretical studies (Croize et al. [2013](#page-19-6); Das et al. [2011,](#page-19-7) [2013;](#page-19-8) Issen and Rudnicki [2001;](#page-20-13) Olsson [1999;](#page-21-3) Shahin et al. [2020](#page-22-10); Sternlof et al. [2005\)](#page-22-11). The controlled laboratory experiments of formation of compaction and shear bands at diferent temperatures, confning pressures, and type and pressure of pore fuids (solely or coupled) have been conducted on siliciclastic rocks to large extent (Baud et al. [2000;](#page-19-9) Heap et al. [2009;](#page-20-8) Jasinski et al. [2015;](#page-20-14) Pons et al. [2011;](#page-21-7) Wong et al. [1997\)](#page-22-12). However, such experiments on carbonates are somewhat limited. Also, the complex mineralogy and microstructure of the carbonate rocks (e.g., mineralogy, deposition, and diagenesis) (Rashid et al. [2017](#page-21-8)) make their coupled behaviour further complex (Abdallah et al. [2021](#page-19-10); Cilona et al. [2014](#page-19-11); Louis et al. [2009](#page-21-9)). Besides the internal characteristics of the carbonates, the coexistence of many external conditions such as stress, type, and pressure of pore fuid and temperature are hypothesized to play a signifcant role in controlling localised deformation. More importantly, it is known that such drivers (stress, temperature, chemical reaction, and pore pressure) are tightly coupled, and if a particular process(es) is of interest, the other driver efects should be isolated.

The type of pore fuid, its potential chemical reaction, and the effect of such chemical reaction on specimen's properties have also been shown to afect the deformation bands in porous rocks. Water, a common pore fuid and commonly used as injection fuid for enhanced oil recovery (EOR) objectives, has been studied extensively. The reason behind the waterweakening efect relies on many possible mechanisms. For instance, when water exits as pore fuids, it tends to embrittle rock by pressurisation and lubrication (Anderson [1981;](#page-19-12) Violay et al. [2015](#page-22-13); Violay et al. [2014\)](#page-22-14). It can also decrease rock minerals' surface energy and drop the friction coefficient, promoting subcritical crack growth and stress corrosion (Atkinson [1984](#page-19-13); Baud et al. [2000](#page-19-9); Brantut et al. [2013](#page-19-14); Kodama et al. [2013;](#page-20-15) Sulem and Stefanou [2016\)](#page-22-5). The same mechanism was also hypothesized to explain the water-weakening efect of carbonate rocks (Baud et al. [2009;](#page-19-15) Nicolas et al. [2016](#page-21-10)). Ciantia et al. ([2015\)](#page-19-16) presented a micro-scale mechanism of the water-weakening efect on the decrease of rock strength and introduced three diferent water-weakening processes of carbonates, including (1) short-term depositional bonds destroyed debonding, which is related to the calcite powder that falls into suspension during water saturation; (2) long-term diagenetic bonds dissolved debonding which is afected by long-term dissolution processes; and (3) grain dissolution process which involving the grains, and starting with long-term dissolution simultaneously. More recently, Cai et al. ([2019\)](#page-19-17) reported that the water-weakening efect of rocks depends on the porosity and the mineralogy, especially the proportion of quartz and swelling clays. Besides the chemical reaction efect on the formation of compaction bands, the thermomechanical and thermo-hydromechanical (coupled thermal-pore pressure–stress) effects on compaction bands, especially in carbonate rocks (limestone), is scarce, requiring further research to shed light on the involved processes. Understanding the efect of temperature on the formation of compaction bands is particularly important for geological-scale investigation involving higher temperature conditions.

In this study, we therefore present a systematic micro–macro-scale investigation of thermomechanical and thermo-hydromechanical effects on the formation and evolution of compaction bands in highly porous limestone  $($  ~50% porosity Mount Gambier limestone). To investigate the thermomechanical and thermo-hydromechanical behaviours of the limestone samples, dry (helium saturated) and purifed kerosene were used (e.g. Baxevanis et al. [2006;](#page-19-5) Cerasi and Walle [2016](#page-19-18)), respectively. To shed light on the micro-scale mechanisms, we used our newly developed X-ray transparent hydromechanical triaxial cell (Roshan et al. [2019](#page-21-11)) where the limestone specimens were scanned at confning pressures of 5 MPa with 25 °C, 50 °C and 80 °C temperatures.

Note that the entire analysed temperature range shows sharp stress drops, and the experiment is therefore entirely performed in what is commonly understood as the brittle regime. However, we focus in the present manuscript on the micro-scale mechanisms and classify the brittle and quasiductile response of the material at such microscopic level where brittle creep is here classifed as a micro-scale ductile response. We analysed the time-lapse monitoring of external fuid pressure, mechanical load, and strain in relation to the radiographs and the 3D-CT reconstructions. This allowed us to link the macroscopic constraints to the microscopic response of the investigated limestone.

## **2 Experimental Methodology**

#### **2.1 Rock Material and Sample Preparation**

The triaxial experiments were performed on Mount Gambier limestone retrieved from the Gambier Embayment of Otway, the basin of Paleogene–Neogene age in southern South Australia (Bourman et al. [2016\)](#page-19-19). Mt Gambier limestone is a highly porous (~50% porosity) fossiliferous carbonate that originated from extensive colonies of Bryozoa (or lace coral) that fourished on an open-marine shelf, formed over 30 million years ago. The sample consists of 96% calcite and has a small contribution of quartz, magnesite, and magnetite (accounting for the remaining 4%) from X-ray difraction (XRD) analysis. Chen et al. [\(2020b\)](#page-19-20) reported that the Mt Gambier limestone has been characterised as a dual-porosity profle from the mercury intrusion capillary pressure measurement (MICP): macro-pores ranging from 10 to 200 µm and nano-micro-pores ranging from 0.1 to 5 µm. The average porosity of the sample from MICP measurement and helium porosimetry was obtained as  $\sim$  50% porosity.

Cylindrical cores with 12.7 mm diameter were drilled from the same block. Both sides were then carefully ground to fat and parallel surfaces with 0.02 mm accuracy to a fnal length of 20 mm. This specimen size and geometry are necessary to enable X-ray micro-computed tomography (XRCT) imaging. The X-ray transparent triaxial testing allowed the full-field XRCT measurements with sufficiently high resolution. For limiting the moisture efect, all the specimens were placed in a vacuum oven, dried at 110 °C for 48 h, and the temperature was then reduced to room temperature for another 24 h before conducting the experiments.

For thermo-mechanical testing, helium was used as pore fuid, and for thermo-hydromechanical testing, purifed kerosene was used as pore fuid instead of water. It is noted that no signifcant reaction between helium or kerosene and calcite is reported (Karkush et al. [2013\)](#page-20-16). In dry experiments, the specimens placed in the triaxial system was vacuumed, and helium was injected from the top platen into the specimen; a digital gas pressure regulator was used to control the inlet pressure constant. To ensure that the specimens are saturated with kerosene, dry specimens were frst placed in a desiccator at a vacuum of 1 Torr for 24 h at room temperature. After that, the kerosene was gradually injected into the vessel's bottom until the specimens were flled with kerosene. Complete saturation  $(> 95\%)$  was verified by weighing the specimen for 2 weeks.

## **2.2 Thermo‑hydromechanical Experiments**

Simultaneous triaxial testing and XRCT scanning on the Mount Gambier limestone specimens were conducted using an X-ray transparent triaxial deformation apparatus (Chen et al. [2020b](#page-19-20); Roshan et al. [2019](#page-21-11)). A servo-controlled loading frame was used for all the experiments, and an axial displacement of  $3 \times 10^{-3}$  mm/s was set (corresponding to a nominal strain rate of  $1.5 \times 10^{-4}$  s<sup>-1</sup>). Axial displacement data were logged by the linear variable diferential transducer (LVDT) with a sampling rate of  $0.1 \text{ s}^{-1}$  (10 Hz). Axial load was logged externally by a high precision disk load cell (LPX-1000) with an accuracy of  $\pm$  0.001 kN.

The confining pressure is  $P_c = \sigma_2 = \sigma_3$  and the pore pressure  $P_p$ . We denote the maximum and minimum compressive stress by  $\sigma_1$  and  $\sigma_3$ , respectively. Also, the differential stress is defined by  $q = \sigma_1 - \sigma_3$  and the effective mean stress is defined as  $p = (\sigma_1 + 2\sigma_3)/3 - P_p$ . The compressive stress and strain are considered positive. According to the previous study on Mount Gambier limestone, the confning pressure of 5.0 MPa was used to trigger the pure compaction bands (Chen et al. [2020b](#page-19-20)).

Two sets of experiments were conducted to investigate the thermomechanical and thermo-hydromechanical experiments. The thermomechanical experiments involved triaxial testing at 5.0 MPa confnement with helium at 25 °C, 50 °C, and 80 °C temperatures, and thermo-hydromechanical experiments consisted of triaxial testing at 5.0 MPa confnement with kerosene at the same temperatures. At least three repetitions were conducted for each test to capture the variability. The specimens were also CT scanned at several loading stages for each test that are discussed later in the next section.

A soap bubble gas fowmeter was used to measure the downstream gas fow rate (Skoczylas [2015](#page-22-15)) for helium saturated permeability. The inlet pressure of 0.5 psi (3.45 kPa) was chosen because of high Mount Gambier limestone permeability (in order of Darcy).

The temperature of the cell was elevated and controlled by the circulation of high-pressure and high-temperature confning fuid. The temperature-controlled distilled water was circulated with a Vindum pump (VP-3K-HC-T). This pump can work smoothly up to 3500 psi (24.13 MPa) and delivers oscillation-free fow rates ranging up to 97 ml/min. In our experiments, the pump was used as a flow rate-controlled mode to provide enough heat to warm up the specimen to the desired temperature (Fig. [15](#page-16-0) in Appendix 1). A thermostat (Lauda Alpha A6) was used to control the hot water temperature (set as 95 °C) and to refll the Vindum pump. The water was collected from the outlet (back-pressure regulator). This



<span id="page-3-0"></span>**Fig. 1 a** The triaxial deformation apparatus's schematic shows the top and bottom platen pathways, allowing pore fuid injection and outfow. **b** The elevated-temperature experimental set-up

back-pressure regulator was connected to the air gas cylinder, and was used to control the confning pressure. The schematic fgure of the elevated-temperature experimental set-up is presented in Fig. [1](#page-3-0).

#### **2.3 X‑Ray CT Scanning and Image Processing**

X-ray CT (XRCT) is a non-destructive imaging technique that quantifes an object's internal structure in 3D space (Arzilli et al. [2016](#page-19-21); Lock et al. [2002](#page-21-12); McBeck et al. [2020](#page-21-13); Takano et al. [2015](#page-22-16); Viggiani and Tengattini [2019](#page-22-17)). This technique was frst used for medical purposes and later on widely used in engineering applications, including industrial use, geological investigation and material science. In this study, all the XRCT scans were conducted at the Tyree X-ray CT facility at the University of New South Wales, Sydney, Australia. The system has a GE Phoenix Nanofocus X-Ray Tube with a diamond window and a high-quality flatbed detector  $(3072 \times 3072)$  pixels, 3.75 fps readout rate) which was designed to allow easy access and the ability to integrate complex flow experiments with the imaging system. The facility is built in a lead-lined room with excellent temperature and humidity control (∆*T*<0.5 °C). In our study, the specimens were scanned in a circular trajectory with the setting of 100 kV, 120  $\mu$ A, exposure time 0.6 s, three accumulations, 0.5 mm stainless steel flter, and 2880 projection per revolution. A resolution of 15.08 µm/voxel was obtained where the entire specimen was ftted in the feld of view. The raw cone-beam X-ray projection data were reconstructed using the Katsevich algorithm (Kingston et al. [2016](#page-20-17); Sheppard et al. [2014](#page-22-18); Varslot et al. [2010\)](#page-22-19).

In our study, time-lapse X-ray CT imaging was acquired on four specimens during their deformation at diferent conditions: MG-D25: helium saturated at 25 °C; MG-K25: kerosene saturated at 25 °C; MG-D80: helium saturated at 80 °C; MG-K80: kerosene saturated at 80 °C. The rest of the specimens were conducted at the same condition without X-ray CT to verify the repeatability. Before the experiment, each specimen was frst scanned within the cell at unconfned conditions. After that, the specimens were loaded isotropically to 5.0 MPa stress, and triaxial loading was then applied to diferent axial strain levels while acquiring new series of tomographic images. All the experiments were conducted up to  $\sim$  10% axial strain with five scans at different axial strain levels, while the static load (both axial and radial) is maintained during the scan. Each scan took 1.5 h to have suffcient time to obtain a resolution of 15.08 µm/voxel. During the 1.5 h scan, the potential of creep could occur; however, it was negligible in our measurements. This is because any creep would have caused failure of the reconstruction algorithm of CT raw images which relies on recognising features in the immediate vicinity of the previous scan 1.5 h prior.

After X-ray CT scanning, the image processing was conducted using the Mango software developed by the Australian National University. Mango is a powerful tool for parallel segmentation and network generation and the preand post-processing and analysis of associated data. The image processing workflow includes image cropping, edge enhancement, noise reduction, beam hardening correction, intensity calibration, and segmentation. For our time-lapsed X-ray CT images, intensity calibration and image registration were additionally performed. This is an essential part



<span id="page-4-0"></span>**Fig. 2** A simple workfow of X-ray CT image processing and further digital image analysis

of analysing a series of X-ray CT images of the specimen at diferent deformation stages (Fig. [2\)](#page-4-0). Details on image registration and digital image correlation method for analysis can be found in Appendix 2.

# **3 Results**

#### **3.1 Thermo‑hydromechanical Behaviour**

#### **3.1.1 Helium Saturated Specimens**

Figure [3a](#page-5-0)–c shows the stress–strain response of helium saturated specimens deformed at three temperatures (25 °C, 50 °C, and 80 °C). Four stage of stress–strain response is observed from the fgure including Stage 1 where the stress–strain curve shows a short nonlinear response indicating the closure of pre-existing cracks and end-surface effects due to platen engagement (Renner and Rummel [1996](#page-21-14); Walsh [1965](#page-22-20)). Stage 2 exhibits the linear elasticity until the yield point, followed by Stage 3, where the stress–strain relationship diverges from the linear relationship until the peak point. The beginning of the unstable cracking starts in this stage (Xue et al. [2014\)](#page-22-21). At Stage 4, strain-softening,

hardening, and pore collapse are observed (Bieniawski [1967](#page-19-22); Eberhardt et al. [1999;](#page-20-18) Zhou et al. [2018\)](#page-23-1).

Previous studies reported that the pressure drops are linked to nucleation and propagation of compaction bands (Baud et al. [2004;](#page-19-23) Chen et al. [2020b;](#page-19-20) Das et al. [2013](#page-19-8)), meaning that more compaction bands are nucleated and developed with an increase in temperature. The average strength of the rock reduces with increasing temperature, due to failure by a thermal weakening of the structure and thermally induced stresses hence lowering the strength by triggering more collapse of the pores (Mahmutoglu [1998;](#page-21-15) Sengun [2014\)](#page-22-22). This will be further discussed in Sect. [3.2](#page-6-0).

#### **3.1.2 Kerosene Saturated Specimens**

To fnd out the efect of the pore fuid type on the formation of compaction bands, purifed kerosene was used as pore fuid instead of helium. After immersing the specimen in kerosene through the vacuum desiccator, additional kerosene was slowly injected from the bottom port at a low constant flow rate: 3 ml/min to ensure full saturation of pores before the experiments. The same temperatures (25  $\degree$ C, 50  $\degree$ C, and 80 °C) as in the helium experiments were used. As shown in Fig. [3d](#page-5-0)–f, the stress–strain response and its temperature dependency show similar trends to helium saturated



<span id="page-5-0"></span>**Fig. 3** Axial stress versus axial strain for specimens deformed at **a** 25 °C, dry condition, **b** 50 °C at dry condition, **c** 80 °C at dry condition, **d** 25 °C with kerosene, **e** 50 °C with kerosene, and **f** 80 °C with kerosene. In each condition, three experiments were performed

to assess the repeatability. Also, the stages of 1, 2, 3, and 4 on the **a** refer to the closing of pre-existing cracks, linear elasticity, divergence from the linear relationship, and strain softening and hardening, respectively

<span id="page-6-1"></span>

specimens. Figure [4a](#page-6-1) shows the trends of the average stress at the frst peak point and yield point with temperature for the experiments presented in Fig. [3](#page-5-0). The value of axial stress at yield point and frst peak point decreases in both helium and kerosene saturated conditions with a temperature increase from 25 to 80 °C. This temperature-sensitive trend was more pronounced for helium saturated specimens than kerosene saturated specimens. The average elastic modulus of helium and kerosene saturated specimens is plotted as a function of temperature in Fig. [4](#page-6-1)b.

## <span id="page-6-0"></span>**3.2 Temperature‑Dependent Micro‑structural Alteration**

To acquire a detailed map of the internal structure, the microstructural analyses were performed on four specimens, helium saturated condition at 25 and 80 °C, and kerosene saturated condition at 25 and 80 °C, using fulllength 3D X-ray computed tomography. Figure [5](#page-6-2) shows the stress–strain response of four tests with X-ray CT scanning, where each time-lapse experiment had five scans from  $0\%$ (undeformed) up to  $\sim$  10% axial strain marked as S1–S5 in Fig. [5.](#page-6-2)

Digital image correlation analysis was performed on each specimen to track the initiation and evolution of localised compaction bands during the experiments at specifed conditions. In our study, an open-source 2D-DIC code, Ncorr

(a)  $12$ (b)  $12$ **Dry <sup>25</sup> <sup>o</sup> (a) <sup>C</sup> Dry <sup>80</sup> <sup>o</sup> (b) <sup>C</sup>** 10-12 10-12  $\frac{1}{2}$  S<sub>5</sub> S2 S3 10 10 S4 S5 Axial Stress (MPa) Axial Stress (MPa) Axial Stress (MPa) Axial Stress (MPa) Permeability (m<sup>2</sup>)  $\widehat{\phantom{a}}$ Permeability (m<sup>2</sup>) S<sub>2</sub> Permeability (m<sup>2</sup> 8 8 Axial stress 6 6 Permeability 4 4  $\overline{2}$ 2 Axial stress Permeability  $\frac{1}{0.1}$ 10<sup>-13</sup> S1 S1  $10^{-13}$  $0\frac{1}{0}$  $\epsilon$ 0 0.02 0.04 0.06 0.08 0.1 Axial Strain Axial Strain (c)  $12$  $(d)$ <sub>12</sub> **Kerosene 25 <sup>o</sup> (c) <sup>C</sup> Kerosene 80 <sup>o</sup> (d) <sup>C</sup>** S4 10 10 S2 S3 Axial Stress (MPa) Axial Stress (MPa) S4 Axial Stress (MPa) Axial Stress (MPa) S5 8 8  $\overline{\mathbf{s}}$  $\sim$  S<sub>3</sub> 6 6 4 4 2 2 Axial stress Axial stressS1 S1  $0\frac{R}{D}$  $0\frac{V}{0}$ 0 0.02 0.04 0.06 0.08 0.1 0 0.02 0.04 0.06 0.08 0.1 Axial Strain Axial Strain

<span id="page-6-2"></span>**Fig. 5** The axial stress–strain response of **a** specimen MG-D25, **b** specimen-D80, **c** specimen MG-K25, and **d** specimen MG-K80. S1–S5 indicates the number of scans at diferent axial strain levels. The permeability evolution curve in helium saturated samples is plotted in **a** and **b** using the Klinkenberg correction



<span id="page-7-0"></span>**Fig. 6** From top to bottom; the 2D incremental shear strain, 2D incremental compressive strain, 3D volumetric strain, and porosity reduction profle on **a** specimen MG-D25 (dry 25 °C) and **b** specimen MG-D80 (dry 80 °C)

(Blaber et al. [2015\)](#page-19-24), was used to analyse the localisation and propagation of compaction bands over full-length 2D images as previously mentioned. We show two cross-section tomography planes (*XZ* and *YZ* planes) for 2D analysis. To assess the displacements and strain measurements within a sequence of tomographic volume images of a 3D subset, digital volume correlation (DVC) was additionally performed using Avizo 9.0 (Thermo Fisher Scientifc). The contrast patterns were tracked from the reference tomography to the deformed state. Figures [6](#page-7-0) and [7](#page-8-0) represent 2D shear and compressive strain on two representative tomograph slices. The fgures also show a 3D volumetric strain map on a maximum cuboid that was cut from the cylindrical specimen as well as porosity profle reduction between each scan for helium and kerosene saturated specimens at two diferent temperatures (25 and 80  $\degree$ C), respectively. The first X-ray CT scan was acquired at 0% axial strain, followed by the second scan at 2.5% axial strain, which is higher than the yield point  $(1.5-2\%)$ .

Comparing Figs. [6](#page-7-0) and [7](#page-8-0) for helium and kerosene saturated specimens show some distinct diferences. It is seen from Figs. [6](#page-7-0) and [7](#page-8-0) that the specimens saturated by helium at both low and high temperatures experience less compaction band formation than that of kerosene saturated specimens. To confrm this observation, we plot the overall porosity profle of the specimens extracted from segmented tomography images for both helium and kerosene saturated specimens at low and high temperatures, respectively (Figs. [8](#page-8-1) and [9](#page-9-0)). It is noted that porosity extracted from images will have an accuracy limitation to image voxel resolution (15.08  $\mu$ m). The porosity profle across the specimens proves the point observed from DIC–DVC analysis on the formation of compaction bands. These fgures clearly show that kerosene saturated specimens undergo more signifcant pore collapse and associated damage (reduction in porosity) than the helium saturated specimens. For a more detailed investigation of this point, we extracted the radial and axial strains for both helium and kerosene saturated specimens at low and high temperatures from XRCT images at different measurement points and plotted them versus axial stress (Fig. [16](#page-18-0) in Appendix 3). The fgure shows that the helium saturated specimen develops a higher radial strain than the kerosene saturated specimen at both temperatures. We therefore conclude that the compression of the kerosene saturated specimen is accommodated by a higher degree of pore collapse thus reducing the radial strain upon compression. We



<span id="page-8-0"></span>Fig. 7 From top to bottom; the 2D incremental shear strain, 2D incremental compressive strain, 3D volumetric strain, and porosity reduction profle on **a** specimen MG-K25 (kerosene 25 °C) and **b** specimen MG-K80 (kerosene 80 °C)

<span id="page-8-1"></span>**Fig. 8** The porosity profle along *z*-direction of the specimens extracted from segmented images. **a** Specimen MG-D25 (dry 25 °C) and **b** specimen MG-D80 (dry 80 °C)



<span id="page-9-0"></span>**Fig. 9** The porosity profle along *z*-direction of the specimens extracted from segmented images. **a** Specimen MG-K25 (kerosene 25 °C) and **b** specimen MG-K80 (kerosene 80 °C)



suggest that gas saturated matrix micro-pores accommodate a higher degree of elastic deformation prior to macro-pore collapse than the equivalent fuid-saturated matrix micropores. The diference in radial strain is interpreted as indirect evidence of the proposed role of compressibility of the gasflled micro-pores in the matrix surrounding macro-pores.

# **3.3 Efect of Temperature on Permeability Evolution of Compaction Bands**

Compaction bands have been reported to have a negative efect on fuid fow by creating a dramatic decrease in local porosity and permeability within the formed bands (Baxevanis et al. [2006](#page-19-5); Fortin et al. [2005;](#page-20-19) Lenoir et al. [2010](#page-20-20); Sternlof et al. [2006](#page-22-23)). In our study, sample permeabilities were measured on helium saturated specimens during the deformations where the X-ray CT imaging sheds light on the mechanism of permeability reduction due to formation of compaction bands in the presence of thermo-hydromechanical interaction. This technique offers a superior understanding of how permeability changes during the formation of compaction bands.

Helium gas was used for dry specimens as a working fluid, and continuous permeability measurements were carried out. The gas permeability is calculated by Darcy's equation for compressible fuids and further converted to a liquid permeability following the Klingenberg correction (Tanikawa and Shimamoto [2006](#page-22-24)). The permeability evolution as a function of axial strain is plotted in Fig. [5](#page-6-2)a, b, along with the stress–strain response of helium saturated specimens at 25 and 80 °C, respectively. The overall trend of permeability reduction has been previously reported, and the processes are discussed (Chen et al. [2020a](#page-19-25)). Both specimens tested under low and high temperatures show similar initial permeability. At 25 °C condition, the permeability reduced from  $8.3 \times 10^{-13}$  to  $1.13 \times 10^{-13}$  m<sup>2</sup>; at 80 °C, the permeability reduced from  $8.2 \times 10^{-13}$  to  $1.77 \times 10^{-13}$  m<sup>2</sup>. However, with the increase of axial load, the reduction in permeability becomes highlighted for the low-temperature specimen compared to the high temperature one. This is especially the case once the 5% axial strain is passed. The permeability result contrasts with the DIC analysis performed on XRCT images which shows that the specimen experiences a much higher degree of compaction and compaction band development at a higher temperature. We will discuss this in the next section.

# **4 Discussion**

We presented our analysis on highly porous carbonate where we systematically investigated the efect of pore fuid type and temperature on the micro-mechanisms of failure, localisation behaviour, and overall strength evolution. Insights into micromechanical aspects were gained indirectly through DIC/DVC analyses of X-ray CT scans and direct radiographs observations.

## **4.1 Pore Fluid Efect**

#### **4.1.1 Efect of Dual Porosity and Drainage Conditions**

Experiments performed with a chemically inert fuid showed only a slight efect of fuid type on yield stress in the hightemperature regime but a more highlighted efect at low temperatures. The results clearly showed that experiments performed with helium have higher yield stress than experiments where specimens were saturated with kerosene as a pore fuid. An obvious candidate for the efect would be chemical weakening. However, the effect of chemical reactions of kerosene with the carbonate matrix can be ruled out as there is no signifcant reaction reported in the literature (Karkush et al. [2013](#page-20-16)).

The limestone sample has a unique dual pore size distribution with distinct clusters of micro-pores ( $\sim$  100 nm–5  $\mu$ m) and macro-pores  $(-10-200 \mu m)$  (Chen et al. [2020b](#page-19-20)). The micro-pores are an integral part of the skeleton around the macro-pores and they therefore have a signifcant efect on mechanical matrix properties. The applied fuid pressure is low and drainage conditions are supported by the percolation network of the macro-pores for both kerosene and helium. We can speak of drained conditions for the macro-pores and the diferent efect of compressibility of fuid or gas does not play a role. The fuid/gas in the micro-pores of the skeleton, however, has to pass through narrow pore throats causing a deviation from the drained response. This near "undrained" response in drained testing condition is well documented for tight rocks such as shale where micro-pores dominate the pore structure (Roshan and Aghighi [2012;](#page-21-16) Roshan and Rahman [2011\)](#page-21-17). Therefore, the restricted permeability in the micro-pores implies that the diference in compressibility of fuid versus gas can manifest itself in the experiment.

This near "undrained" response in specimens saturated with kerosene will have two distinct effects. First, due to partial support of the load by the liquid in micro-pores, one expects to see stifening of the skeleton causing a slight increase in elastic modulus for the specimens saturated with kerosene compared to helium (Levy et al. [2000;](#page-21-18) Mavko et al. [2020](#page-21-19)). On the other hand, the increase in pressure in these micro-pores should lead to a reduction of the yield stress (Han et al. [2019](#page-20-21); Robinson [1959](#page-21-20)). The strength reduction is well observed in the experiments with kerosene and helium at low temperatures, although the increase in elastic modulus is less apparent.

An interesting observation is that the elastic loading path at 80 °C shows a higher elastic modulus for kerosene than for helium, while the difference in yield stress vanishes (Fig. [4](#page-6-1)). The vanishing diference in yield stress is likely linked to a decrease in kerosene viscosity with temperature enabling it to escape the micropore under drained condition, thus leading to a minimal pore-pressure build-up in these micropores. The diference in elastic modulus is attributed to the temperature sensitivity of the elastic deformation that will thus be discussed in the temperature dependence section.

#### **4.1.2 Efect of Reduction in Surface Energy**

As previously mentioned, no chemical weakening such as calcite-pressure solution (Hellmann et al. [2002;](#page-20-22) Newman [1983;](#page-21-21) Plummer and Busenberg [1982](#page-21-22)) or corrosive deterioration (Atkinson [1984;](#page-19-13) Karner et al. [2005;](#page-20-23) Wiederhorn [1967](#page-22-25)) can be triggered by kerosene in our experiment. However, the important physico-chemical mechanism of surface energy reduction remains to be investigated. It is well known that the reduction in surface energy promotes the initiation and growth of micro-cracks (Grifth [1921](#page-20-24)). The surface energy of minerals is formed through interaction between the solid surface and exposed fuids (Siddiqui et al. [2019](#page-22-26)). In our study, we deal with kerosene and helium which should essentially have diferent surface energies with calcite minerals leading to diferent mechanical responses due to resultant crack propagation (Rostom et al. [2013](#page-21-23); Zeng et al. [2020](#page-23-2)). To investigate the change in surface energy and its relation to sample weakening, we use the Griffith crack theory that provides a simple fracture mechanics model stating that fracture energy, *G*, is equal to the energy necessary to produce the two surfaces of a new fracture or twice the surface energy

$$
G - 2\gamma_s^e = 0,\t\t(1)
$$

where *G* is the mechanical energy and  $\gamma_s^e$  is the surface energy of the solid in contact with a given environment. The lowering of the surface energy of the material can enhance crack propagation. Helium is an inert gas and will not induce any surface energy alteration; thus, we focus on kerosene. To investigate the kerosene efect on surface energy and potential weakening of the limestone sample, we choose to compare it with water that is known to have a severe efect on rock mechanical properties including calcite rich rocks (Lisabeth and Zhu [2015;](#page-21-24) Liteanu et al. [2013](#page-21-25); Risnes et al. [2003](#page-21-26), [2005;](#page-21-27) Røyne et al. [2011\)](#page-22-27).

Contact angle measurement is an indirect technique to characterise the surface energy of a solid in presence of fuids (Siddiqui et al. [2019\)](#page-22-26). The results of contact angle measured for water–calcite–air system and kerosene–calcite–air system are presented in Fig. [17](#page-18-1) in Appendix, i.e., 20° and 5° for water and kerosene contact angles, respectively, after stabilisation. Knowing that water–air and kerosene–air surface energies are 0.0728 N/m at 25 °C (Lange and Dean [1967](#page-20-25)) and 0.0267 N/m at 25 °C (Landry et al. [2011\)](#page-20-26), respectively, one can calculate the diference in surface energy of calcite–water and calcite–kerosene as shown in Appendix 4 (Link and Schlünder [1996](#page-21-28)). This calculation shows that the calcite surface energy with water is lower than kerosene. Knowing that the surface energy of minerals are lower when exposed to liquids compared to gases (Brunauer et al. [1956](#page-19-26)), we conclude that kerosene promotes less crack growth and fracture propagation than water, but it still contributes to a reduction in surface energy and promotes cracking higher than helium case. The reduction of the strength of specimens in low temperature (25 °C) is approximately 15% which is similar to Baxevanis et al.'s ([2006\)](#page-19-5) results. Compared with the experimental results using water as pore fuid, e.g., (Liteanu et al. [2013](#page-21-25); Risnes et al. [2003\)](#page-21-26), our results showed less strength reduction when using kerosene saturated specimens consistent with contact angle measurements. Overall, the weaker structure is experienced for specimen saturated by kerosene. This effect is strongly visible in formation of compaction bands and reduction of specimen porosity with kero-sene (Figs. [8](#page-8-1) and [9](#page-9-0)) at any temperature, although such effect is not very pronounced in the values of initial yield strength dominated by the heterogeneity of diferent samples.

#### **4.2 The Role of Temperature**

In the results chapter, we showed a systematic weakening in both helium and purifed kerosene experiments when raising the temperature from 25 to 80 °C. The stress–strain response and its temperature dependency of helium and kerosene saturated experiments show similar trends. The average strength of the rock reduces with an increase in temperature (Fig. [3](#page-5-0)). An interesting observation is the efect of temperature on the near-elastic branch of the load curve. The helium saturated specimen showed a distinct temperature sensitivity with a systematic elastic modulus drop at a higher temperature.

In contrast, the elastic modulus in kerosene saturated specimen appeared to be slightly affected by a temperature change. The combination of temperature sensitivity of the carbonate matrix and near "undrained" response discussed before is the likely cause for this behaviour. While the carbonate matrix weakens with the increase in temperature (Sulem and Stefanou [2016;](#page-22-5) Wong et al. [2020](#page-22-28); Yavuz et al. [2010\)](#page-23-3), the undrained response of kerosene dominates the weakening efect of the matrix by temperature; thus, a lower drop in elastic modulus is observed in kerosene case compared to that of helium. It is, however, difficult to ascertain the exact roots of the weakening of the elastic modulus for the helium saturated specimens as the efect is not directly measurable and is modulated through the dynamics of matrix–pore interactions. In the following section, we will discuss the temperature efect on the yield phenomenon in further detail.

#### **4.2.1 Helium Saturated Experiments**

The comparison of the helium saturated experiments performed at lower and higher temperatures indicates a transition in deformation mode. From macroscopic observation, the strength of the specimens reduces with increasing temperatures from 25 to 80 °C. Similar results have been reported by Liteanu et al. ([2013\)](#page-21-25); a reduction in overall strength (approximately 2 MPa) has been observed when increasing the temperature from 25 to 80 °C at 5 MPa confning pressure using dry Maastrichtian chalk, which has similar composition and porosity to Mt Gambier limestone. More localised deformation features within the specimens can be seen from DIC and DVC analyses at a higher temperature (Fig.  $6$ ).

To fnd out the microstructural changes between lower (25 °C) and higher (80 °C) temperatures, a selected region of interest (ROI) window is shown in Fig.  $10$ . The figure highlights the microstructure evolution in the leading compaction bands and undamaged matrix at Scan 2 (2.5% axial strain), Scan 3 (5% axial strain), and Scan 4 (7.5% axial strain). At lower temperature  $(25 \text{ °C})$ , pore collapse happened within the compaction bands seen by tomographic observations as highlighted by the red dashed box in Fig. [11](#page-13-0)a. There is no pore collapse observed in the undamaged matrix (Fig. [11a](#page-13-0)–c). When increasing the axial strain from 2.5 to 7.5%, the main compaction band that nucleated from the lower boundary propagates towards the middle of the specimen as a wavefront (Fig. [11](#page-13-0)d, e).

We infer that the front-like propagation of damage of the specimen is controlled for the low-temperature experiments by rate dependence of the material. This rate dependence relies on the formation of new surfaces between the calcite grains for which there are many possible deformation mechanisms. We have not performed a microstructural analysis of the sub-micro-scale processes. However, from the literature, one of the possible candidates is calcite twinning. This lowtemperature ductile efect triggers pervasive deformation (Rowe and Rutter [1990\)](#page-22-29). Brittle mechanisms such as pore collapse, crack growth, and grain crushing are competing deformation mechanisms promoting localised deformation features that have a much faster dynamic response than the slow speed of propagation of the wavefront, which is controlled by the twinning reaction. The brittle mechanism presumably operates in conjunction with twinning. However, the macroscopic behaviour appears to be akin to a plastic deformation that afects the propagating front pervasively, while the remainder of the specimen remains intact. The low-temperature deformation mode has been encountered in many similar experiments and is called a "diffuse compaction band" (Louis et al. [2006](#page-21-29); Tembe et al. [2006;](#page-22-30) Wang et al. [2008\)](#page-22-31) which contrasts with the localised style of compaction encountered in the high-temperature experiments.



<span id="page-12-0"></span>**Fig. 10 a**–**c** Microscopic view of the ROI window cropped from the tomography slices (*XZ* plane) of specimen MG-D25 at scan 2, 3, and 4, respectively. The black colour represented the pores, and the grey

colour indicates the grain (matrix). The red dashed box indicated the main compaction band. **d** and **e** are the  $\epsilon_{yy}$  strain map between Scan 2–3 and Scan 3–4, respectively

At a higher temperature (80 $^{\circ}$ C), we observe diffuse and highly localised deformation features (Fig. [11\)](#page-13-0). The new observation at high temperature is the localised compaction band that starts at scan 2–3 in Fig. [5b](#page-6-2) at the top boundary, but does not propagate as a steady wavefront like the difuse band but jumps towards the middle of the specimen. Brittle localised pore collapse or damage can be clearly identifed in the otherwise undamaged matrix from Scan 3 onwards (Fig. [11\)](#page-13-0), indicating that brittle deformation has a more prominent efect at higher temperatures. When zooming into the pore collapse of the smallest pores highlighted in the red box (Fig. [12\)](#page-14-0), we can see that the weaker structures failed in random directions according to the weakest point of the structure itself (Fig. [12\)](#page-14-0), and the broken pieces fall into the open pores as crushed rocks. At scan 4, the pore collapse of the small pores developed as a localised deformation band easily recognised from the strain map showing that the specimen failed in a more complex manner on diferent localised compaction bands. The behaviour of small pores collapse is diferent from the previously reported mechanism of discrete compaction bands in dry specimens (Chen et al. [2020a\)](#page-19-25). The discrepancy between the dominance of difuse compaction band growth versus localised features at higher temperatures experiments is a new observation that has not been reported in the literature before. The importance of the microcrystalline behaviour of the matrix in conjunction with the interaction of the pore filling fluids needs further explanations of the somewhat counterintuitive dominance of a brittle deformation mode at higher temperatures.

In the following, we attempt to fnd explanations for observing the promotion of brittle effects at a higher temperature. Circumstantial evidence for brittle deformation dominance at a higher temperature is that sample heterogeneity appears to be more important at controlling localisation features especially with the reduction in matrix strength. Another indirect evidence is the efect



<span id="page-13-0"></span>**Fig. 11 a**–**c** Microscopic view of the ROI window cropped from the tomography slices (*XZ* plane) of specimen MG-D80 at Scan 2, 3, and 4, respectively. The black colour represents pores, and the grey colour indicates the grain (matrix). The red dashed box indicated the main

compaction band, and the red solid line area indicated deformation features within the undamaged matrix. **d** and **e** are the  $\epsilon_{yy}$  strain map between Scan 2–3 and Scan 3–4, respectively (colour fgure online)

of permeability reduction reported in Fig. [5a](#page-6-2), b, showing that the low-temperature regime shows a higher reduction in permeability with strain. In contrast, the high-temperature experiments have a milder efect on permeability reduction. This is interpreted here through the pervasive destruction of fluid pathways in the low-temperature regime through the difuse deformation bands contrasting to the smaller localised brittle damage zones from the higher temperature experiments. The transition from dominantly ductile at low temperatures (25 °C) to dominantly brittle at higher temperatures (80 °C) is also clearly illustrated in Figs. [11](#page-13-0) and [12](#page-14-0).

#### **4.2.2 Kerosene Saturated Experiments**

A similar transition in deformation mode has been observed in kerosene saturated experiments. The overall reduction of the yield strength of the specimens from 25 to 80 °C is clear through the overall load curve (Fig. [3](#page-5-0)). Interestingly, although there is a strong diference in elastic behaviour, the yield strength of the helium saturated specimen compared to the kerosene saturated one shows similar values at high temperatures. At the low temperature, there is a marked diference in yield strength. Because the diferences in yield strength between high and low temperatures are observed for all specimens, including those where the pores are flled with helium gas, we rule out the dominance of the surface tension efect on temperature. It is instead more likely that the diference is caused by the afore-documented transition in dominant deformation mode from a temperature-sensitive ductile deformation to a temperature-insensitive brittle deformation mode at high temperature.

The observation supports the interpretation of a transition in dominant micromechanical deformation mode that at lower temperatures, kerosene saturated specimens exhibit the same pervasive deformation mode as helium saturated ones. Similarly, both gaseous and fuid-saturated specimens

<span id="page-14-0"></span>

<span id="page-14-1"></span>**Fig. 13 a**–**c** Microscopic view of the ROI window cropped from the tomography slices (*XZ* plane) of specimen MG-K25 at scan 2, 3, and 4, respectively. The black colour represented the pores, and the grey

colour indicates the grain (matrix). The red dashed box indicated the main compaction bands. **d** and **e** are the  $\epsilon_{yy}$  strain map between Scan 2–3 and Scan 3–4, respectively

show the same localised deformation features within the specimens, as seen from DIC and DVC analysis at a higher temperature (Fig. [7\)](#page-8-0). A representative example of the prominent propagating wavefront of the difuse compaction bands in both helium and kerosene cases is shown in Fig. [13.](#page-14-1)

This deformation mode is in stark contrast to the DIC map recorded at higher temperatures. Localised pore collapse becomes predominant, and individual pores collapse and crush to form coherent localised deformation features at higher axial strain (Fig. [14](#page-15-0)). A narrow compaction localisation is nucleated between Scan 2 to Scan 3, both seen in the tomography images and the strain map (Fig. [14b](#page-15-0), d). The deformation bands are more evident in particular selected tomographic slices (Fig. [14c](#page-15-0)) than in the strain map



<span id="page-15-0"></span>**Fig. 14 a–c** Microscopic view of the ROI window cropped from the tomography slices (*XZ* plane) of specimen MG-K80 at scan 2, 3, and 4, respectively. The black colour represented the pores, and the grey

colour indicates the grain (matrix). The red solid line area indicated the main deformation bands. **d** and **e** are the  $\epsilon_{yy}$  strain map between Scan 2–3 and Scan 3–4, respectively

(Fig. [14e](#page-15-0)). This is due to the high degree of localisation and the inability of the DIC algorithm to track the crushed calcite powder that flls the larger pores as reference points are vanishing.

In summary, our systematic comparison of the results from gaseous and fuid-saturated experiments shows that a pervasive ductile deformation mode dominates at a lower temperature and a localised brittle deformation mode at a higher temperature. We attribute this effect to two main reasons. The low-temperature ductile efect is rate controlled by calcite twinning, which is known to occur at room temperature (Friedman and Heard [1974](#page-20-27); Guéguen and Boutéca [2004;](#page-20-28) Parlangeau et al. [2019;](#page-21-30) Turner et al. [1954\)](#page-22-2). This efect controls the slow propagation of the main compaction bands from the boundary towards the middle of the specimen. At higher temperatures, the Peierls mechanism can be triggered, which shows little rate dependency and is called low-temperature plasticity (Sly et al. [2020](#page-22-32)). Although this mechanism is a ductile micro-mechanism, its viscous damping is very low due to the near-ideal plastic response. It, therefore, can assist or trigger the macroscopic brittle failure of the skeleton, assisted by the surface energy efect of the fuid phase.

Our new observations may alternatively be interpreted by a proposed temperature-sensitive mechanism proposed by earlier analyses. These report that moderate temperature changes can signifcantly afect the mechanical behaviour of deformation bands. Explanation in the literature (Ghabezloo et al. [2009;](#page-20-5) Kuempel et al. [2017;](#page-20-29) Lemos [2003;](#page-20-30) Renard et al. [2009](#page-21-31)) attributed the efect to changes in the stress-state afected by thermally induced pore-pressure changes, thus lowering the threshold for induced brittle damage. Yavuz et al. ([2010\)](#page-23-3) examined three diferent limestones: monomineralic rock (almost calcite), similar to Mt Gambier limestone used in our study. Dilation of calcite and compaction of grains has been observed when temperature increased to 100 °C; this compaction of grains resulted in an efective porosity decrease. Nicolas et al. ([2016\)](#page-21-10) reported that the temperature diference between 20 and 70 °C can result in shear-enhanced compaction triggered at lower confning pressure in dry Tavel limestone. However, we consider the temperature-sensitive pore-pressure efect to be negligible in our specimens. All mechanisms can co-occur, making it difficult to differentiate between thermally induced weakening through pore-pressure changes and alternative transitions in deformation behaviour of the solid matrix as suggested by us.

# **5 Conclusion**

We have conducted a series of triaxial compression experiments on a highly porous Mt Gambier limestone and investigated the efect of pore fuids and temperature on mechanical behaviour. Concurrent compression tests with X-ray CT scans were used to offer a time-lapse evolution of the deformation localisations and digital image analysis, offering a micro-structural observation under in-situ conditions. This study used an ideal thermo-hydro-poro-mechanical system to eliminate the chemical reaction, using purifed kerosene and helium gas as pore fuids. Two novel fundamental observations were made:

1. We have observed that the yield stress of nominally inert gas (helium) saturated specimens is higher than the nominally inert liquid (purifed kerosene) saturated specimens. Conversely, the elastic modulus was lower for the inert gas saturated specimens than inert fuid-saturated specimens. We attribute the effects to the near-undrained response of the sub-micro-scale pore system of the specimen.

2. A second important observation is a changeover from ostensibly ductile to a dominantly brittle deformation when raising the temperature. Regarding the core scale mechanical behaviour afected by moderate temperature change, our time-lapse experiment with X-ray CT tomography provided a microscopic view of structural alteration during compaction band formation. A transition in deformation mode was observed in both helium and kerosene saturated conditions when increasing the temperature from 25 to 80 °C. At lower temperature, we have observed wave-like propagation of compaction bands propagating from the boundary towards the specimen's middle. At high-temperature, pore collapse, shear or shear-enhanced compaction localisation and Mode 1 initial crack orientation indicate that brittle deformation has a more prominent effect at higher temperatures. In the low-temperature ductile regime, compaction bands' formation shows pervasive deformation with no visible damage or pore collapse in the undamaged matrix. The undamaged matrix remained the stress supporting structure, and all the pore collapse happened within the difuse compaction regions (Figs. [6a](#page-7-0) and [7a](#page-8-0)).

The observation of the inverse-to-normal material behaviour for brittle–ductile transition can be characterised as follows. Specimens deformed at low temperature (25 °C) in a pervasive ductile-like manner with propagating difuse compaction bands and specimens deforming at a higher temperature (80 °C) by the localised brittle collapse of pores in stationary highly localised compaction bands. Although direct insight into the micro-mechanisms of deformation was not possible at the investigation scale, we attribute this perplexing behaviour to a dominant deformation mode transition. We postulate that a transition from calcite twinning dominated rate-sensitive deformation at low temperatures to a near-rate insensitive Peierls mechanism ideal plastic yield with subsequent brittle failure at high temperature. An alternative explanation to the change in micro-mechanism proposed in the literature (Ghabezloo et al. [2009](#page-20-5); Lemos [2003](#page-20-30);



<span id="page-16-0"></span>**Fig. 15** The relationship between desired elevated temperature versus time at diferent fow rates

Renard et al. [2009](#page-21-31)) cannot be ruled out. Thermal expansion of pore fuid as a trigger for brittle deformation would also explain the observation of temperature-induced weakening by changes in pore pressure facilitating brittle deformation.

When fuids rather than gas coexist within the matrix, the problem becomes more complicated. As carbonate rocks are primary sedimentary materials widely distributed in the surface and upper layers of the Earth's surface, the efect of water as pore fuid has been widely studied in the past. Water has been reported to have a signifcant weakening efect on carbonate rocks' mechanical behaviour, which is the result of several mechanisms (Cai et al. [2019](#page-19-17); Ciantia et al. [2015](#page-19-16); Liteanu et al. [2013\)](#page-21-25). This water-weakening effect can be accelerated when the temperature is raised (Korsnes et al. [2008](#page-20-31); Lisabeth and Zhu [2015](#page-21-24); Nicolas et al. [2016](#page-21-10)). The surface energy reduction by kerosene was also apparent through our experimentation; thus, similar behaviour to that of water in weakening the structure is expected although to a lesser degree. To gain further insight into the micro-mechanism of deformation, future work needs to focus on direct subnanoscale time-lapse observations to identify the counterintuitive ductile to brittle transition in Mt Gambier limestone.

## **Appendix 1**

A series of pre-tests were conducted at a constant water flow rate to warm up the specimen to the desired temperature. The temperature of the specimen was measured with a thermocouple probe made of Cu–CuNi wires which was placed inside the specimen. The thermocouple was connected to ALMEMO 2590 data logger, which has a resolution of 0.1 °C. Figure [15](#page-16-0) shows the desired temperature of 80 °C requiring circulating hot water (95 °C) at 90 ml/min

for 60 min and temperature of 50 °C requiring circulating hot water (95 °C) at 90 ml/min for 8 min and then reduced to 24 ml/min for another 30 min until the temperature stabilises. The temperature fluctuation is  $\pm 0.5$  °C throughout.

## **Appendix 2**

## **Digital Image Processing**

Each scan might have slightly diferent greyscale intensity because of the flament lifetime changes. Diferences in greyscale intensity can result in bias when obtaining the accurate registration and thresholding segmentation process across all the images. The greyscale intensity was carefully calibrated according to selected homogeneous regions: air, kerosene, triaxial cell body, top and bottom platens, and rubber sleeve. All diferences were corrected using an obtained linear function.

When the X-ray beam passes through a cylindrical specimen, the outer regions absorb and scatter the lower energies in the X-ray spectrum, which results in the exterior regions being brighter than the inner regions. Beam-hardening correlation thus applied a Gaussian smoothing kernel on specifed regions of the data to reduce the beam hardening efect (Ketcham and Hanna [2014](#page-20-32)). Noise is another artifact that makes it challenging to diferentiate low-density areas, thereby reducing the ability to segment efectively (Nagarajappa et al. [2015\)](#page-21-32). A nonlinear anisotropic difusion flter denoising the images and preserving the edges using a similar method by Frangakis and Hegerl ([2001\)](#page-20-33) was used in our study.

Image registration is a crucial step in time-lapse experiments. A 3D registration algorithm developed by Latham et al. ([2008\)](#page-20-34) was used for image registration. This technique brings two or more images into geometric alignment for further digital image correlation analysis. In our study, each scan was carefully registered to a previous scan to ensure that undeformed parts are overlapped.

The histogram of all the image slices can be exported before image segmentation (Fig. [1\)](#page-3-0). The greyscale values of the tomographs correspond to the X-ray attenuation produced by material, i.e., for example, lower for air than kerosene and the calcite grains. The histogram shows two peaks corresponding to the pores (air, low X-ray attenuation) and matrix (mainly calcite, high X-ray attenuation). With compression of the specimen, the volume of the pore phase reduces, and the matrix phase increases. Two dashed lines indicate the greyscale value of the pore phase and the matrix

phase. The converging active contours' (CAC) method was used for the segmentation of these images (Sheppard et al. [2014](#page-22-18); Sheppard et al. [2004\)](#page-22-33). This method uses a combination of the watershed and active contour methods to segment the greyscale data. The real boundary of the pore and matrix can be determined by this method. Later on, the layer-bylayer porosity profle can be extracted from the segmented images along with *Z*-axis.

## **Digital Image Correlation (DIC) and Digital Volume Correlation (DVC) Analysis**

Since the 1980s, the 2D digital image correlation (DIC) technique has been widely used to measure real-time fullfeld data of displacements and strains (Chu et al. [1985](#page-19-27); Vendroux and Knauss [1998](#page-22-34)). This non-destructive testing method can measure the deformation behaviour of a material over a wide area in exceptional detail. DIC uses image registration algorithms to track the relative displacement of material points between a reference image and a deformed image. Our study uses an open-source 2D subset-based DIC software package: Ncorr (Blaber et al. [2015\)](#page-19-24). Ncorr uses the reliability guided (RG-DIC) method to obtain displacement values for a subset. The shear and compressive strain feld can be extracted based on the displacement feld using the Green–Lagrangian strain tensor, which can be used to determine the nucleation and propagation of the compaction bands. This DIC code has been shown to work well for tracking the materials displacements (Caselle et al. [2019](#page-19-28); Lv et al. [2019;](#page-21-33) Siddiqui et al. [2021;](#page-22-35) Stanier et al. [2016\)](#page-22-36). With the wide use of X-ray CT and synchrotron 3D images, the digital volume correlation (DVC) technique has become popular (Bay et al. [1999\)](#page-19-29). This technique can be considered as the extended version of the 2D-DIC method to the 3D-DVC domain in conjunction with 3D images, which efectively determines the internal volumetric deformation behaviours of solid materials. Avizo (Thermofsher Scientifc) digital volume correlation (DVC) was used for 3D-DVC analysis in our study. A subset-based (local) approach is used to capture the large displacements on a coarse, regular grid.

## **Appendix 3**

The radial deformations for both helium and kerosene saturated specimens at low and high temperatures were extracted from every X-ray CT scan. The radial and axial strain versus axial stress is plotted in Fig. [16.](#page-18-0)



<span id="page-18-0"></span>**Fig. 16** The axial and radial strain versus axial stress of **a** specimen MG-D25 (dry 25 °C) and MG-D80 (dry 80 °C) and **b** specimen MG-K25 (kerosene 25 °C) and MG-K80 (kerosene 80 °C)

# **Appendix 4**

The contact angles of kerosene–calcite–air and water–calcite–air were measured by viewing the drop profle (Link and Schlünder [1996](#page-21-28); Siddiqui et al. [2019\)](#page-22-26) and are shown in Fig. [17.](#page-18-1)

The equilibrium contact angle  $(\theta_C)$  is determined from Young equation (Butt et al. [2013\)](#page-19-30)

$$
\gamma_{SG} + \gamma_{SL} + \gamma_{LG} \cos(\theta_C) = 0,
$$

where,  $\gamma$  is the surface energy, and SG, SL, and LG represent the solid–gas, solid–liquid, and liquid–gas, respectively. We write the contact angle equations for water–calcite–air and kerosene–calcite–air system as

$$
\gamma_{\text{Sa}} + \gamma_{\text{Sw}} + \gamma_{\text{wa}} \cos(\theta_{\text{w}}) = 0,
$$



<span id="page-18-1"></span>**Fig. 17** Contact angle measurement using distilled water and purifed kerosene with crushed powder of Mt Gambier limestone. Because of the highly permeable sample powder, the contact angle of two different fuids was measured using a high-speed camera. The results were plotted by contact angle versus time. Two dashed lines show the power-law ft of the kerosene and water contact angles, respectively

$$
\gamma_{\text{Sa}} + \gamma_{\text{Sk}} + \gamma_{\text{ka}} \cos(\theta_k) = 0,
$$

where Sa, Sw, wa, and w represent calcite–air, calcite–water, water–air, and water, and Sk, ka, and k represent calcite–kerosene, kerosene–air, and kerosene, respectively. The diference between  $\gamma_{\text{Sw}}$  and  $\gamma_{\text{Sk}}$  is

$$
\Delta \gamma = (\gamma_{\rm Sk} - \gamma_{\rm Sw}) = \gamma_{\rm ka} \cos(\theta_{\rm k}) - \gamma_{\rm wa} \cos(\theta_{\rm w}),
$$

where  $\gamma_{wa}$  is 0.0728 N/m at 25 °C (Lange and Dean [1967\)](#page-20-25) and  $\gamma_{ka}$  is 0.0267 N/m at 25 °C (Landry et al. [2011\)](#page-20-26). Also,  $\theta_k$ and  $\theta_w$  are 5° and 20°, respectively, leading to

Δ*𝛾*0.042N∕m.

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**Availability of Data and Materials** The datasets used and/or analysed during the current study are available from the corresponding author on reasonable request.

**Code Availability** Not applicable.

#### **Declarations**

**Conflict of Interest** The authors declare that they have no competing interests.

**Ethics Approval** Not applicable.

**Consent to Participate** Not applicable.

**Consent for Publication** Not applicable.

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