



G -invariant Szegő kernel asymptotics and CR reduction

Chin-Yu Hsiao¹ · Rung-Tzung Huang²

Received: 20 January 2020 / Accepted: 30 September 2020 / Published online: 28 January 2021
© The Author(s), under exclusive licence to Springer-Verlag GmbH, DE part of Springer Nature 2021

Abstract

Let $(X, T^{1,0}X)$ be a compact connected orientable CR manifold of dimension $2n + 1$ with non-degenerate Levi curvature. Assume that X admits a connected compact Lie group G action. Under certain natural assumptions about the group G action, we show that the G -invariant Szegő kernel for $(0, q)$ forms is a complex Fourier integral operator, smoothing away $\mu^{-1}(0)$ and there is a precise description of the singularity near $\mu^{-1}(0)$, where μ denotes the CR moment map. We apply our result to the case when X admits a transversal CR S^1 action and deduce an asymptotic expansion for the m th Fourier component of the G -invariant Szegő kernel for $(0, q)$ forms as $m \rightarrow +\infty$ and when $q = 0$, we recover Xiaonan Ma and Weiping Zhang's result about the existence of the G -invariant Bergman kernel for ample line bundles. As an application, we show that if m large enough, quantization commutes with reduction.

Mathematics Subject Classification Primary: 58J40 · 32V20 · 53D50; Secondary: 57Q10

1 Introduction and statement of the main results

Let $(X, T^{1,0}X)$ be a CR manifold of dimension $2n + 1$, $n \geq 1$. Let $\square_b^{(q)}$ be the Kohn Laplacian acting on $(0, q)$ forms. The orthogonal projection $S^{(q)} : L^2_{(0,q)}(X) \rightarrow \text{Ker } \square_b^{(q)}$

Communicated by A. Malchiodi.

The first author was partially supported by Taiwan Ministry of Science and Technology Projects 108-2115-M-001-012-MY5, 109-2923-M-001-010-MY4 and Academia Sinica Career Development Award. This work was initiated when the second author was visiting the Institute of Mathematics at Academia Sinica in the summer of 2016. The second author would like to thank the Institute of Mathematics at Academia Sinica for its hospitality and financial support during his stay. The second author was also supported by Taiwan Ministry of Science of Technology Projects 105-2115-M-008-008-MY2, 107-2115-M-008-007-MY2 and 109-2115-M-008-007-MY2.

✉ Chin-Yu Hsiao
chsiao@math.sinica.edu.tw; chinyu.hsiao@gmail.com
Rung-Tzung Huang
rthuang@math.ncu.edu.tw

¹ Institute of Mathematics, Academia Sinica, Astronomy-Mathematics Building, No. 1, Sec. 4, Roosevelt Road, Taipei 10617, Taiwan

² Department of Mathematics, National Central University, Chung-Li 32001, Taoyuan, Taiwan

onto $\text{Ker } \square_b^{(q)}$ is called the Szegő projection, while its distribution kernel $S^{(q)}(x, y)$ is called the Szegő kernel. The study of the Szegő projection and kernel is a classical subject in several complex variables and CR geometry. A very important case is when X is a compact strictly pseudoconvex CR manifold. Assume first that X is the boundary of a strictly pseudoconvex domain. Boutet de Monvel-Sjöstrand [2] showed that $S^{(0)}(x, y)$ is a complex Fourier integral operator.

The Boutet de Monvel-Sjöstrand description of the Szegő kernel had a profound impact in many research areas, especially through [4]: several complex variables, symplectic and contact geometry, geometric quantization, Kähler geometry, semiclassical analysis, quantum chaos, etc. cf. [6,8,11,22,29,32], to quote just a few. These ideas also partly motivated the introduction of the recent direct approaches and their various extensions, see [18,19,21,22].

Now, we consider a connected compact Lie group G acting on X . The study of G -invariant Szegő kernel is closely related to Mathematical physics and geometric quantization of CR manifolds. It is a fundamental problem to establish G -invariant Boutet de Monvel-Sjöstrand type theorems for G -invariant Szegő kernels and study the consequence of the G -invariant Szegő kernel. This is the motivation of this work. In this paper, we consider G -invariant Szegő kernel for $(0, q)$ forms and we show that the G -invariant Szegő kernel for $(0, q)$ forms is a complex Fourier integral operator. In particular, $S^{(q)}(x, y)$ is smoothing outside $\mu^{-1}(0)$ and there is a precise description of the singularity near $\mu^{-1}(0)$, where μ denotes the CR moment map. We apply our result to the case when X admits a transversal CR S^1 action and deduce an asymptotic expansion for the m th Fourier component of the Szegő kernel for $(0, q)$ forms as $m \rightarrow +\infty$. As an application, we show that, if m large enough, quantization commutes with reduction.

In [20], Ma and Zhang have studied the asymptotic expansion of the invariant Bergman kernel of the spin^c Dirac operator associated with high tensor powers of a positive line bundle on a symplectic manifold admitting a Hamiltonian action of a compact connected Lie group and its relation to the asymptotic expansion of Bergman kernel on symplectic reduced space, also the Toeplitz operator aspect in [20, Section 4.5]. Their approach is inspired by the analytic localization techniques developed by Bismut and Lebeau [3]. About the quantization commutes with reduction problem in symplectic geometry, we refer the readers to [22]. In the second part of [22], Ma described how the G -invariant Bergman kernel concentrates on the Bergman kernel of the reduced space. Note that the “quantization commutes with reduction” in the situations in symplectic case is a very active subject. When the action connected Lie group is compact and the symplectic manifold is also compact, this question was solved finally by Meinrenken [24] and Tian-Zhang [31]. When the action connected Lie group is compact and the symplectic manifold is noncompact, this is a famous conjecture of Vergne and was solved by Ma-Zhang in [23].

It should be mentioned that in [7], Charles relates the Toeplitz operators on a compact complex manifold M with the Toeplitz operators on the “reduced” space for torus action, and in [26], Paoletti studied equivariant Szegő kernels on complex manifolds (cf. [20, Remark 0.5]).

We now formulate the main results. We refer to Sect. 2 for some notations and terminology used here. Let $(X, T^{1,0}X)$ be a compact connected orientable CR manifold of dimension $2n + 1, n \geq 1$, where $T^{1,0}X$ denotes the CR structure of X . Fix a global non-vanishing real 1-form $\omega_0 \in C^\infty(X, T^*X)$ such that $\langle \omega_0, u \rangle = 0$, for every $u \in T^{1,0}X \oplus T^{0,1}X$. The Levi form of X at $x \in X$ is the Hermitian quadratic form on $T_x^{1,0}X$ given by $\mathcal{L}_x(U, \bar{V}) = -\frac{1}{2i} \langle d\omega_0(x), U \wedge \bar{V} \rangle$, for all $U, V \in T_x^{1,0}X$. In this work, we assume that

Assumption 1.1 *The Levi form is non-degenerate of constant signature (n_-, n_+) on X . That is, the Levi form has exactly n_- negative and n_+ positive eigenvalues at each point of X , where $n_- + n_+ = n$.*

Let $HX = \{\operatorname{Re} u; u \in T^{1,0}X\}$ and let $J : HX \rightarrow HX$ be the complex structure map given by $J(u + \bar{u}) = iu - i\bar{u}$, for every $u \in T^{1,0}X$. In this work, we assume that X admits a d -dimensional connected compact Lie group G action. We assume throughout that

Assumption 1.2 *The G action preserves ω_0 and J . That is, $g^*\omega_0 = \omega_0$ on X and $g_*J = Jg_*$ on HX , for every $g \in G$, where g^* and g_* denote the pull-back map and push-forward map of G , respectively.*

Let \mathfrak{g} denote the Lie algebra of G . For any $\xi \in \mathfrak{g}$, we write ξ_X to denote the vector field on X induced by ξ . That is, $(\xi_X u)(x) = \frac{\partial}{\partial t}(u(\exp(t\xi) \circ x))|_{t=0}$, for any $u \in C^\infty(X)$.

Definition 1.3 The moment map associated to the form ω_0 is the map $\mu : X \rightarrow \mathfrak{g}^*$ such that, for all $x \in X$ and $\xi \in \mathfrak{g}$, we have

$$\langle \mu(x), \xi \rangle = \omega_0(\xi_X(x)). \tag{1.1}$$

In this work, we assume that

Assumption 1.4 *0 is a regular value of μ , the action G on $\mu^{-1}(0)$ is freely and*

$$\underline{\mathfrak{g}}_x \cap \underline{\mathfrak{g}}_x^{\perp b} = \{0\} \text{ at every point } x \in Y, \tag{1.2}$$

where $\underline{\mathfrak{g}} = \operatorname{Span}(\xi_X; \xi \in \mathfrak{g})$, $\underline{\mathfrak{g}}^{\perp b} = \left\{v \in HX; b(\xi_X, v) = 0, \forall \xi_X \in \underline{\mathfrak{g}}\right\}$, b is the nondegenerate bilinear form on HX given by (2.4).

By Assumption 1.4, $\mu^{-1}(0)$ is a d -codimensional submanifold of X . Let $Y := \mu^{-1}(0)$ and let $HY := HX \cap TY$. Note that if the Levi form is positive at Y , then (1.2) holds. Under the condition (1.2), in Sect. 2.5, we will show that $\dim(HY \cap JHY) = 2n - 2d$ at every point of Y , $\mu^{-1}(0)/G =: Y_G$ is a CR manifold with natural CR structure induced by $T^{1,0}X$ of dimension $2n - 2d + 1$ and we can identify HY_G with $HY \cap JHY$.

Fix a G -invariant smooth Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ so that $T^{1,0}X$ is orthogonal to $T^{0,1}X$, $\underline{\mathfrak{g}}$ is orthogonal to $HY \cap JHY$ at every point of Y , $\langle u | v \rangle$ is real if u, v are real tangent vectors, $\langle T | T \rangle = 1$ and T is orthogonal to $T^{1,0}X \oplus T^{0,1}X$, where T is given by (2.2). The Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ induces, by duality, a Hermitian metric on $\mathbb{C}T^*X$ and also on the bundles of $(0, q)$ forms $T^{*0,q}X$, $q = 0, 1, \dots, n$. We shall also denote all these induced metrics by $\langle \cdot | \cdot \rangle$. Fix $g \in G$. Let $g^* : \Lambda_x^r(\mathbb{C}T^*X) \rightarrow \Lambda_{g^{-1} \circ x}^r(\mathbb{C}T^*X)$ be the pull-back map. Since G preserves J , we have $g^* : T_x^{*0,q}X \rightarrow T_{g^{-1} \circ x}^{*0,q}X$, for all $x \in X$. Thus, for $u \in \Omega^{0,q}(X)$, we have $g^*u \in \Omega^{0,q}(X)$. Put $\Omega^{0,q}(X)^G := \{u \in \Omega^{0,q}(X); g^*u = u, \forall g \in G\}$. Since the Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ is G -invariant, the L^2 inner product $(\cdot | \cdot)$ on $\Omega^{0,q}(X)$ induced by $\langle \cdot | \cdot \rangle$ is G -invariant. Let $u \in L^2_{(0,q)}(X)$ and $g \in G$, we can also define g^*u in the standard way (see the discussion in the beginning of Sect. 3.2). Put $L^2_{(0,q)}(X)^G := \{u \in L^2_{(0,q)}(X); g^*u = u, \forall g \in G\}$. Let $\square_b^{(q)} : \operatorname{Dom} \square_b^{(q)} \rightarrow L^2_{(0,q)}(X)$ be the Gaffney extension of Kohn Laplacian (see (3.1)). Put $(\operatorname{Ker} \square_b^{(q)})^G := \operatorname{Ker} \square_b^{(q)} \cap L^2_{(0,q)}(X)^G$. The G -invariant Szegő projection is the orthogonal projection $S_G^{(q)} : L^2_{(0,q)}(X) \rightarrow (\operatorname{Ker} \square_b^{(q)})^G$ with respect to $(\cdot | \cdot)$. Let $S_G^{(q)}(x, y) \in D'(X \times X, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ be the distribution kernel of $S_G^{(q)}$. The first main result of this work is the following

Theorem 1.5 *With the assumptions and notations above, suppose that $\square_b^{(q)} : \text{Dom } \square_b^{(q)} \rightarrow L^2_{(0,q)}(X)$ has closed range. If $q \notin \{n_-, n_+\}$, then $S_G^{(q)} \equiv 0$ on X .*

Suppose $q \in \{n_-, n_+\}$. Let D be an open set of X with $D \cap \mu^{-1}(0) = \emptyset$. Then, $S_G^{(q)} \equiv 0$ on D .

Let $p \in \mu^{-1}(0)$ and let U be an open set of p and let $x = (x_1, \dots, x_{2n+1})$ be local coordinates defined in U . Then, there exist continuous operators $S_-^G, S_+^G : \Omega^{0,q}_0(U) \rightarrow \Omega^{0,q}(U)$ such that

$$S_G^{(q)} \equiv S_-^G + S_+^G \text{ on } U,$$

and $S_-^G(x, y), S_+^G(x, y)$ satisfy

$$S_{\mp}^G(x, y) \equiv \int_0^\infty e^{i\Phi_{\mp}(x,y)t} a_{\mp}(x, y, t) dt \text{ on } U,$$

with

$$\begin{aligned} a_+(x, y, t), a_-(x, y, t) &\in S_{\text{cl}}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ a_-(x, y, t) &= 0 \text{ if } q \neq n_-, \quad a_+(x, y, t) = 0 \text{ if } q \neq n_+, \\ a_-^0(x, x) &\neq 0, \quad \forall x \in U, \text{ if } q = n_-, \quad a_+^0(x, x) \neq 0, \quad \forall x \in U, \text{ if } q = n_+, \end{aligned} \tag{1.3}$$

$a_-^0(x, x), a_+^0(x, x), x \in \mu^{-1}(0) \cap U$, are given by (1.8) below, $\Phi_-(x, y) \in C^\infty(U \times U)$,

$$\begin{aligned} \text{Im } \Phi_-(x, y) &\geq 0, \\ d_x \Phi_-(x, x) &= -d_y \Phi_-(x, x) = -\omega_0(x), \quad \forall x \in U \cap \mu^{-1}(0), \end{aligned} \tag{1.4}$$

there is a constant $C \geq 1$ such that, for all $(x, y) \in U \times U$,

$$\begin{aligned} |\Phi_-(x, y)| + \text{Im } \Phi_-(x, y) &\leq C (\inf \{d^2(g \circ x, y); g \in G\} + d^2(x, \mu^{-1}(0)) + d^2(y, \mu^{-1}(0))), \\ |\Phi_-(x, y)| + \text{Im } \Phi_-(x, y) &\geq \frac{1}{C} (\inf \{d^2(g \circ x, y); g \in G\} + d^2(x, \mu^{-1}(0)) + d^2(y, \mu^{-1}(0))), \\ Cd^2(x, \mu^{-1}(0)) \geq \text{Im } \Phi_-(x, x) &\geq \frac{1}{C} d^2(x, \mu^{-1}(0)), \quad \forall x \in U, \end{aligned} \tag{1.5}$$

and $\Phi_-(x, y)$ satisfies (1.18) below and (1.19) below, and $\Phi_+(x, y) \in C^\infty(U \times U)$, $-\overline{\Phi}_+(x, y)$ satisfies (1.4), (1.5), (1.18) below and (1.19) below.

We refer the reader to the discussion before (2.1) and Definition 3.1 for the precise meanings of $A \equiv B$ and the symbol space $S_{\text{cl}}^{n-\frac{d}{2}}$, respectively.

Let $\Phi \in C^\infty(U \times U)$. We assume that Φ satisfies (1.4), (1.5), (1.18), (1.19). We will show in Theorem 5.2 that the functions Φ and Φ_- are equivalent on U in the sense of Definition 5.1 if and only if there is a function $f \in C^\infty(U \times U)$ with $f(x, x) = 1$, for every $x \in \mu^{-1}(0)$, such that $\Phi(x, y) - f(x, y)\Phi_-(x, y)$ vanishes to infinite order at $\text{diag}((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U))$. From this observation, we see that the leading term $a_-^0(x, x), x \in \mu^{-1}(0)$, is well-defined. To state the formula for $a_-^0(x, x)$, we introduce some notations. For a given point $x_0 \in X$, let $\{W_j\}_{j=1}^n$ be an orthonormal frame of $(T^{1,0}X, \langle \cdot | \cdot \rangle)$ near x_0 , for which the Levi form is diagonal at x_0 . Put

$$\mathcal{L}_{x_0}(W_j, \overline{W}_\ell) = \mu_j(x_0)\delta_{j\ell}, \quad j, \ell = 1, \dots, n.$$

We will denote by

$$\det \mathcal{L}_{x_0} = \prod_{j=1}^n \mu_j(x_0). \tag{1.6}$$

Let $\{T_j\}_{j=1}^n$ denote the basis of $T^{*0,1}X$, dual to $\{\overline{W}_j\}_{j=1}^n$. We assume that $\mu_j(x_0) < 0$ if $1 \leq j \leq n_-$ and $\mu_j(x_0) > 0$ if $n_- + 1 \leq j \leq n$. Put

$$\begin{aligned} \mathcal{N}(x_0, n_-) &:= \{cT_1(x_0) \wedge \dots \wedge T_{n_-}(x_0); c \in \mathbb{C}\}, \\ \mathcal{N}(x_0, n_+) &:= \{cT_{n_-+1}(x_0) \wedge \dots \wedge T_n(x_0); c \in \mathbb{C}\} \end{aligned}$$

and let

$$\tau_{n_-} = \tau_{x_0, n_-} : T_{x_0}^{*0,q}X \rightarrow \mathcal{N}(x_0, n_-), \quad \tau_{n_+} = \tau_{x_0, n_+} : T_{x_0}^{*0,q}X \rightarrow \mathcal{N}(x_0, n_+), \quad (1.7)$$

be the orthogonal projections onto $\mathcal{N}(x_0, n_-)$ and $\mathcal{N}(x_0, n_+)$ with respect to $\langle \cdot | \cdot \rangle$, respectively.

Fix $x \in \mu^{-1}(0)$, consider the linear map

$$\begin{aligned} R_x : \mathfrak{g}_x &\rightarrow \mathfrak{g}_x, \\ u &\rightarrow R_x u, \quad \langle R_x u | v \rangle = \langle d\omega_0(x), Ju \wedge v \rangle. \end{aligned}$$

Let $\det R_x = \lambda_1(x) \cdots \lambda_d(x)$, where $\lambda_j(x)$, $j = 1, 2, \dots, d$, are the eigenvalues of R_x .

Fix $x \in \mu^{-1}(0)$, put $Y_x = \{g \circ x; g \in G\}$. Y_x is a d -dimensional submanifold of X . The G -invariant Hermitian metric $\langle \cdot | \cdot \rangle$ induces a volume form dv_{Y_x} on Y_x . Put

$$V_{\text{eff}}(x) := \int_{Y_x} dv_{Y_x}.$$

Note that the function $V_{\text{eff}}(x)$ was already appeared in Ma-Zhang [23, (0,10)] as exactly the role in the expansion, cf. [23, (0.14)].

Theorem 1.6 *With the notations used above, for $a_-^0(x, y)$ and $a_+^0(x, y)$ in (1.3), we have*

$$a_{\mp}^0(x, x) = 2^{d-1} \frac{1}{V_{\text{eff}}(x)} \pi^{-n-1+\frac{d}{2}} |\det R_x|^{-\frac{1}{2}} |\det \mathcal{L}_x | \tau_{x, n_{\mp}} |, \quad \forall x \in \mu^{-1}(0). \quad (1.8)$$

We now assume that X admits an S^1 action: $S^1 \times X \rightarrow X$. We write $e^{i\theta}$ to denote the S^1 action. Let $T \in C^\infty(X, TX)$ be the global real vector field induced by the S^1 action given by $(Tu)(x) = \frac{\partial}{\partial \theta} (u(e^{i\theta} \circ x))|_{\theta=0}$, $u \in C^\infty(X)$. We assume that the S^1 action $e^{i\theta}$ is CR and transversal (see Definition 4.1). We take $\omega_0 \in C^\infty(X, T^*X)$ to be the global real one form determined by $\langle \omega_0, u \rangle = 0$, for every $u \in T^{1,0}X \oplus T^{0,1}X$ and $\langle \omega_0, T \rangle = -1$. In this paper, we assume that

Assumption 1.7

$$\begin{aligned} T &\text{ is transversal to the space } \mathfrak{g} \text{ at every point } p \in \mu^{-1}(0), \\ e^{i\theta} \circ g \circ x &= g \circ e^{i\theta} \circ x, \quad \forall x \in X, \quad \forall \theta \in [0, 2\pi[, \quad \forall g \in G, \end{aligned} \quad (1.9)$$

and

$$G \times S^1 \text{ acts freely near } \mu^{-1}(0).$$

Let $u \in \Omega^{0,q}(X)$ be arbitrary. Define

$$Tu := \frac{\partial}{\partial \theta} ((e^{i\theta})^* u)|_{\theta=0} \in \Omega^{0,q}(X).$$

For every $m \in \mathbb{Z}$, let

$$\begin{aligned} \Omega_m^{0,q}(X) &:= \{u \in \Omega^{0,q}(X); Tu = imu\}, \quad q = 0, 1, 2, \dots, n, \\ \Omega_m^{0,q}(X)^G &= \{u \in \Omega^{0,q}(X)^G; Tu = imu\}, \quad q = 0, 1, 2, \dots, n. \end{aligned}$$

We denote $C_m^\infty(X) := \Omega_m^{0,0}(X)$, $C_m^\infty(X)^G := \Omega_m^{0,0}(X)^G$. From the CR property of the S^1 action and (1.9), it is not difficult to see that $Tg^*\bar{\partial}_b = g^*T\bar{\partial}_b = \bar{\partial}_bg^*T = \bar{\partial}_bTg^*$ on $\Omega^{0,q}(X)$, for all $g \in G$. Hence,

$$\bar{\partial}_b : \Omega_m^{0,q}(X)^G \rightarrow \Omega_m^{0,q+1}(X)^G, \quad \forall m \in \mathbb{Z}.$$

We now assume that the Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ is $G \times S^1$ invariant. Then the L^2 inner product $(\cdot | \cdot)$ on $\Omega^{0,q}(X)$ induced by $\langle \cdot | \cdot \rangle$ is $G \times S^1$ -invariant. We then have

$$\begin{aligned} Tg^*\bar{\partial}_b^* &= g^*T\bar{\partial}_b^* = \bar{\partial}_b^*g^*T = \bar{\partial}_b^*Tg^* \text{ on } \Omega^{0,q}(X), \quad \forall g \in G, \\ Tg^*\square_b^{(q)} &= g^*T\square_b^{(q)} = \square_b^{(q)}g^*T = \square_b^{(q)}Tg^* \text{ on } \Omega^{0,q}(X), \quad \forall g \in G, \end{aligned}$$

where $\bar{\partial}_b^*$ is the L^2 adjoint of $\bar{\partial}_b$ with respect to $(\cdot | \cdot)$.

Let $L_{(0,q),m}^2(X)^G$ be the completion of $\Omega_m^{0,q}(X)^G$ with respect to $(\cdot | \cdot)$. We write $L_m^2(X)^G := L_{(0,0),m}^2(X)^G$. Put $H_{b,m}^q(X)^G := (\text{Ker } \square_b^{(q)})_m^G := (\text{Ker } \square_b^{(q)})^G \cap L_{(0,q),m}^2(X)^G$. It is not difficult to see that, for every $m \in \mathbb{Z}$, $(\text{Ker } \square_b^{(q)})_m^G \subset \Omega_m^{0,q}(X)^G$ and $\dim(\text{Ker } \square_b^{(q)})_m^G < \infty$. The m th G -invariant Szegő projection is the orthogonal projection $S_{G,m}^{(q)} : L_{(0,q)}^2(X) \rightarrow (\text{Ker } \square_b^{(q)})_m^G$ with respect to $(\cdot | \cdot)$. Let $S_{G,m}^{(q)}(x, y) \in C^\infty(X \times X, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ be the distribution kernel of $S_{G,m}^{(q)}$. The second main result of this work is the following

Theorem 1.8 *With the assumptions and notations used above, if $q \notin n_-$, then, as $m \rightarrow +\infty$, $S_{G,m}^{(q)} = O(m^{-\infty})$ on X .*

Suppose $q = n_-$. Let D be an open set of X with $D \cap \mu^{-1}(0) = \emptyset$. Then, as $m \rightarrow +\infty$, $S_{G,m}^{(q)} = O(m^{-\infty})$ on D .

Let $p \in \mu^{-1}(0)$ and let U be an open set of p and let $x = (x_1, \dots, x_{2n+1})$ be local coordinates defined in U . Then, as $m \rightarrow +\infty$,

$$\begin{aligned} S_{G,m}^{(q)}(x, y) &= e^{im\Psi(x,y)}b(x, y, m) + O(m^{-\infty}), \\ b(x, y, m) &\in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ b(x, y, m) &\sim \sum_{j=0}^\infty m^{n-\frac{d}{2}-j}b_j(x, y) \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ b_j(x, y) &\in C^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \quad j = 0, 1, 2, \dots, \end{aligned}$$

and

$$b_0(x, x) = 2^{d-1} \frac{1}{\text{Veff}(x)} \pi^{-n-1+\frac{d}{2}} |\det R_x|^{-\frac{1}{2}} |\det \mathcal{L}_x| \tau_{x,n_-}, \quad \forall x \in \mu^{-1}(0), \quad (1.10)$$

where τ_{x,n_-} is given by (1.7), and $\Psi(x, y) \in C^\infty(U \times U)$, $d_x\Psi(x, x) = -d_y\Psi(x, x) = -\omega_0(x)$, for every $x \in \mu^{-1}(0)$, $\Psi(x, y) = 0$ if and only if $x = y \in \mu^{-1}(0)$ and there is a constant $C \geq 1$ such that, for all $(x, y) \in U \times U$,

$$\begin{aligned} \text{Im } \Psi(x, y) &\geq \frac{1}{C} \left(d(x, \mu^{-1}(0))^2 + d(y, \mu^{-1}(0))^2 + \inf_{g \in G, \theta \in S^1} d(e^{i\theta} \circ g \circ x, y)^2 \right), \\ \text{Im } \Psi(x, y) &\leq C \left(d(x, \mu^{-1}(0))^2 + d(y, \mu^{-1}(0))^2 + \inf_{g \in G, \theta \in S^1} d(e^{i\theta} \circ g \circ x, y)^2 \right). \end{aligned}$$

(We refer the reader to Theorem 1.12 for more properties of the phase $\Psi(x, y)$.)

We refer the reader to the discussion in the beginning of Sect. 2.2 and Definition 2.1 for the precise meanings of $A = B + O(m^{-\infty})$ and the symbol space $S_{\text{loc}}^{n-\frac{d}{2}}$, respectively.

It is was proved in Theorem 1.12 in [15]) that when X admits a transversal and CR S^1 action and the Levi form is non-degenerate of constant signature on X , then $\square_b^{(q)}$ has L^2 closed range.

Let $Y_G := \mu^{-1}(0)/G$. In Theorem 2.5, we will show that Y_G is a CR manifold with natural CR structure induced by $T^{1,0}X$ of dimension $2n - 2d + 1$. Let $\mathcal{L}_{Y_G,x}$ be the Levi form on Y_G at $x \in Y_G$ induced naturally from \mathcal{L} . Note that the bilinear form b is non-degenerate on $\mu^{-1}(0)$, where b is given by (2.4). Hence, on $(\underline{\mathfrak{g}}, \underline{\mathfrak{g}})$, b has constant signature on $\mu^{-1}(0)$. Assume that on $(\underline{\mathfrak{g}}, \underline{\mathfrak{g}})$, b has r negative eigenvalues and $d - r$ positive eigenvalues on $\mu^{-1}(0)$. Hence \mathcal{L}_{Y_G} has $q - r$ negative and $n - d - q + r$ positive eigenvalues at each point of Y_G . Let $\square_{b,Y_G}^{(q-r)}$ be the Kohn Laplacian for $(0, q - r)$ forms on Y_G . Fix $m \in \mathbb{N}$. Let $H_{b,m}^{q-r}(Y_G) := \left\{ u \in \Omega^{0,q-r}(Y_G); \square_{b,Y_G}^{(q-r)}u = 0, Tu = imu \right\}$. We will apply Theorem 1.8 to establish an isomorphism between $H_{b,m}^q(X)^G$ and $H_{b,m}^{q-r}(Y_G)$ if m large enough. We introduce some notations.

Since $\underline{\mathfrak{g}}_x$ is orthogonal to $H_x Y \cap JH_x Y$ and $H_x Y \cap JH_x Y \subset \underline{\mathfrak{g}}_x^{\perp b}$ (see Lemma 2.4 and (2.5) for the meaning of $\underline{\mathfrak{g}}_x^{\perp b}$), for every $x \in Y$, we can find a G -invariant orthonormal basis $\{Z_1, \dots, Z_n\}$ of $T^{1,0}X$ on Y such that

$$\begin{aligned} \mathcal{L}_x(Z_j(x), \bar{Z}_k(x)) &= \delta_{j,k} \lambda_j(x), \quad j, k = 1, \dots, n, \\ Z_j(x) &\in \underline{\mathfrak{g}}_x + iJ\underline{\mathfrak{g}}_x, \quad j = 1, 2, \dots, d, \\ Z_j(x) &\in \mathbb{C}H_x Y \cap J(\mathbb{C}H_x Y), \quad j = d + 1, \dots, n. \end{aligned}$$

Let $\{e_1, \dots, e_n\}$ denote the orthonormal basis of $T^{*0,1}X$ on Y , dual to $\{\bar{Z}_1, \dots, \bar{Z}_n\}$. Fix $s = 0, 1, 2, \dots, n - d$. For $x \in Y$, put

$$B_x^{*0,s}X = \left\{ \sum_{d+1 \leq j_1 < \dots < j_s \leq n} a_{j_1, \dots, j_s} e_{j_1} \wedge \dots \wedge e_{j_s}; a_{j_1, \dots, j_s} \in \mathbb{C}, \forall d + 1 \leq j_1 < \dots < j_s \leq n \right\}$$

and let $B^{*0,s}X$ be the vector bundle of Y with fiber $B_x^{*0,s}X, x \in Y$. Let $C^\infty(Y, B^{*0,s}X)^G$ denote the set of all G -invariant sections of Y with values in $B^{*0,s}X$. Let

$$\iota_G : C^\infty(Y, B^{*0,s}X)^G \rightarrow \Omega^{0,s}(Y_G)$$

be the natural identification.

Assume that $\lambda_1 < 0, \dots, \lambda_r < 0$, and $\lambda_{d+1} < 0, \dots, \lambda_{n-r+d} < 0$. For $x \in Y$, put

$$\hat{N}(x, n_-) = \{ce_{d+1} \wedge \dots \wedge e_{n-r+d}; c \in \mathbb{C}\},$$

and let

$$\begin{aligned} \hat{p} &= \hat{p}_x : \mathcal{N}(x, n_-) \rightarrow \hat{N}(x, n_-), \\ u &= ce_1 \wedge \dots \wedge e_r \wedge e_{d+1} \wedge \dots \wedge e_{n-r+d} \rightarrow ce_{d+1} \wedge \dots \wedge e_{n-r+d}. \end{aligned}$$

Let $\iota : Y \rightarrow X$ be the natural inclusion and let $\iota^* : \Omega^{0,q}(X) \rightarrow \Omega^{0,q}(Y)$ be the pull-back of ι . Let $q = n_-$. Let $S_{Y_G,m}^{(q-r)} : L_{(0,q-r)}^2(Y_G) \rightarrow H_{b,m}^{q-r}(Y_G)$ be the orthogonal projection and let

$$f(x) = \sqrt{V_{\text{eff}}(x)} |\det R_x|^{-\frac{1}{4}} \in C^\infty(Y)^G.$$

Let

$$\begin{aligned} \sigma_m &: H_{b,m}^q(X)^G \rightarrow H_{b,m}^{q-r}(Y_G), \\ u &\rightarrow m^{-\frac{d}{4}} S_{Y_G,m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x,n_-} \circ f \circ \iota^* \circ u. \end{aligned}$$

In Sect. 6.2, we will show that

Theorem 1.9 *With the notations and assumptions above, suppose that $q = n_-$. If m is large, then $\sigma_m : H_{b,m}^q(X)^G \rightarrow H_{b,m}^{q-r}(Y_G)$ is an isomorphism.*

In particular, if m large enough, then $\dim H_{b,m}^q(X)^G = \dim H_{b,m}^{q-r}(Y_G)$.

Remark 1.10 Let's sketch the idea of the proof of Theorem 1.9. We can consider σ_m as a map from $\Omega^{0,q}(X) \rightarrow H_{b,m}^{q-r}(Y_G)$:

$$\begin{aligned} \sigma_m : \Omega^{0,q}(X) &\rightarrow H_{b,m}^{q-r}(Y_G) \subset \Omega^{0,q-r}(Y_G), \\ u &\rightarrow m^{-\frac{d}{4}} S_{Y_G,m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x,n_-} \circ f \circ \iota^* \circ S_{G,m}^{(q)}. \end{aligned}$$

Let $\sigma_m^* : \Omega^{0,q-r}(Y_G) \rightarrow \Omega^{0,q}(X)$ be the adjoint of σ_m . From Theorem 1.8 and some calculation of complex Fourier integral operators, we will show in Sect. 6.2 that $F_m = \sigma_m^* \sigma_m : \Omega^{0,q}(X) \rightarrow \Omega^{0,q}(X)$ is the same type of operator as $S_{G,m}^{(q)}$ and

$$\frac{1}{C_0} F_m = (I + R_m) S_{G,m}^{(q)}, \tag{1.11}$$

where $C_0 > 0$ is a constant and R_m is also the same type of operator as $S_{G,m}^{(q)}$, but the leading symbol of R_m vanishes at $\text{diag}(Y \times Y)$. By using the fact that the leading symbol of R_m vanishes at $\text{diag}(Y \times Y)$, we will show in Lemma 6.8 that $\|R_m u\| \leq \varepsilon_m \|u\|$, for all $u \in \Omega^{0,q}(X)$, for all $m \in \mathbb{N}$, where ε_m is a sequence with $\lim_{m \rightarrow \infty} \varepsilon_m = 0$. In particular, if m is large enough, then the map

$$I + R_m : \Omega^{0,q}(X) \rightarrow \Omega^{0,q}(X) \tag{1.12}$$

is injective. From (1.11) and (1.12), we deduce that, if m is large enough, then $F_m : H_{b,m}^q(X)^G \rightarrow H_{b,m}^q(X)^G$ is injective. Hence $\sigma_m : H_{b,m}^q(X)^G \rightarrow H_{b,m}^{q-r}(Y_G)$ is injective.

Similarly, we can repeat the argument above with minor change and deduce that if m is large enough, then the map $\hat{F}_m = \sigma_m \sigma_m^* : H_{b,m}^{q-r}(Y_G) \rightarrow H_{b,m}^{q-r}(Y_G)$ is injective. Hence, if m is large enough, then the map $\sigma_m^* : H_{b,m}^{q-r}(Y_G) \rightarrow H_{b,m}^q(X)^G$ is injective. Thus, $\dim H_{b,m}^q(X)^G = \dim H_{b,m}^{q-r}(Y_G)$ and σ_m is an isomorphism if m large enough.

Let's apply Theorem 1.9 to complex case. Let (L, h^L) be a holomorphic line bundle over a connected compact complex manifold (M, J) with $\dim_{\mathbb{C}} M = n$, where J denotes the complex structure map of M and h^L is a Hermitian fiber metric of L . Let R^L be the curvature of L induced by h^L . Assume that R^L is non-degenerate of constant signature (n_-, n_+) on M . Let K be a connected compact Lie group with Lie algebra \mathfrak{k} . We assume that $\dim_{\mathbb{R}} K = d$ and K acts holomorphically on (M, J) , and that the action lifts to a holomorphic action on L . We assume further that h^L is preserved by the K -action. Then R^L is a K -invariant form. Let $\omega = \frac{i}{2\pi} R^L$ and let $\tilde{\mu} : M \rightarrow \mathfrak{k}^*$ be the moment map induced by ω . Assume that $0 \in \mathfrak{k}^*$ is regular and the action of K on $\tilde{\mu}^{-1}(0)$ is freely. The analogue of the Marsden-Weinstein reduction holds (see [10]). More precisely, the complex structure J on M induces a complex structure J_K on $M_0 := \tilde{\mu}^{-1}(0)/K$, for which the line bundle $L_0 := L/K$ is a holomorphic line bundle over M_0 .

For any $\xi \in \mathfrak{k}$, we write ξ_M to denote the vector field on M induced by ξ . Let $\mathfrak{k} = \text{Span}(\xi_M; \xi \in \mathfrak{k})$. On $\tilde{\mu}^{-1}(0)$, let b^L be the bilinear form on $\mathfrak{k} \times \mathfrak{k}$ given by $b^L(\cdot, \cdot) = \omega(\cdot, J\cdot)$. Assume that b^L has r negative eigenvalues and $d - r$ positive eigenvalues on $\tilde{\mu}^{-1}(0)$. Let $q = n_-$. For $m \in \mathbb{N}$, let $H^q(M, L^m)^K$ denote the K -invariant q th Dolbeault

cohomology group with values in L^m and let $H^{q-r}(M_0, L_0^m)$ denote the $(q-r)$ th Dolbeault cohomology group with values in L_0^m . Theorem 1.9 implies that, if m is large enough, then there is an isomorphism map: $\tilde{\sigma}_m : H^q(M, L^m)^K \rightarrow H^{q-r}(M_0, L_0^m)$. In particular, if m is large enough, then

$$\dim H^q(M, L^m)^K = \dim H^{q-r}(M_0, L_0^m). \tag{1.13}$$

Note that when $m = 1$ and $q = 0$, the equality (1.13) was first proved in [10, §5]. For $m = 1$, the equality (1.13) was established in [30,33] when L is positive. Zhang [33] combined the methods and results in [31] with Braverman’s idea [5] to construct a suitable quasi-isomorphism to prove the equality (1.13). The proof of the equality (1.13) in [30] is completely algebraic, while the the proof of the equality (1.13) in [33] is purely analytic where different quasi-homomorphisms between Dolbeault complexes under considerations were constructed to prove the equality (1.13). If m is large enough and $q = 0$, an isomorphism map in (1.13) was also constructed in [20, (0.27), Corollary 4.13].

If m large enough and $q = 0$, an isomorphism map in (1.13) was also constructed in [20, (0.27), Corollary 4.13]. The point of [20, (0.27), Corollary 4.13] is to study the isometric aspect of this map, as an consequence of the asymptotic of G -invariant Bergman kernel of Ma-Zhang [20], they gave another proof that it is an isomorphism for m large, and this approaches of the isomorphism for m large is adopted in this paper. It should be mentioned that in this situation, a version of the full asymptotics of $S_{G,m}^{(0)}(x, y)$ including (1.10) was established in [20, Theorem 0.1, 0.2].

1.1 The phase functions $\Phi_-(x, y)$ and $\Psi(x, y)$

In this section, we collect some properties of the phase functions $\Phi_-(x, y)$, $\Psi(x, y)$ in Theorem 1.5 and Theorem 1.8.

Let $v = (v_1, \dots, v_d)$ be local coordinates of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$. From now on, we will identify the element $e \in V$ with $v(e)$. Fix $p \in \mu^{-1}(0)$. In Theorem 3.7, we will show that there exist local coordinates $v = (v_1, \dots, v_d)$ of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$, local coordinates $x = (x_1, \dots, x_{2n+1})$ of X defined in a neighborhood $U = U_1 \times U_2$ of p with $0 \leftrightarrow p$, where $U_1 \subset \mathbb{R}^d$ is an open set of $0 \in \mathbb{R}^d$, $U_2 \subset \mathbb{R}^{2n+1-d}$ is an open set of $0 \in \mathbb{R}^{2n+1-d}$ and a smooth function $\gamma = (\gamma_1, \dots, \gamma_d) \in C^\infty(U_2, U_1)$ with $\gamma(0) = 0 \in \mathbb{R}^d$ such that

$$\begin{aligned} & (v_1, \dots, v_d) \circ (\gamma(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1}) \\ &= (v_1 + \gamma_1(x_{d+1}, \dots, x_{2n+1}), \dots, v_d + \gamma_d(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1}), \\ & \forall (v_1, \dots, v_d) \in V, \quad \forall (x_{d+1}, \dots, x_{2n+1}) \in U_2, \end{aligned} \tag{1.14}$$

$$\begin{aligned} \underline{\mathfrak{g}} &= \text{span} \left\{ \frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_d} \right\}, \\ \mu^{-1}(0) \cap U &= \{x_{d+1} = \dots = x_{2d} = 0\}, \\ \text{On } \mu^{-1}(0) \cap U, \text{ we have } J\left(\frac{\partial}{\partial x_j}\right) &= \frac{\partial}{\partial x_{d+j}} + a_j(x) \frac{\partial}{\partial x_{2n+1}}, \quad j = 1, 2, \dots, d, \end{aligned} \tag{1.15}$$

where $a_j(x)$ is a smooth function on $\mu^{-1}(0) \cap U$, independent of $x_1, \dots, x_{2d}, x_{2n+1}$ and $a_j(0) = 0, j = 1, \dots, d$,

$$\begin{aligned}
 T_p^{1,0} X &= \text{span} \{Z_1, \dots, Z_n\}, \\
 Z_j &= \frac{1}{2} \left(\frac{\partial}{\partial x_j} - i \frac{\partial}{\partial x_{d+j}} \right) (p), \quad j = 1, \dots, d, \\
 Z_j &= \frac{1}{2} \left(\frac{\partial}{\partial x_{2j-1}} - i \frac{\partial}{\partial x_{2j}} \right) (p), \quad j = d + 1, \dots, n, \\
 \langle Z_j | Z_k \rangle &= \delta_{j,k}, \quad j, k = 1, 2, \dots, n, \\
 \mathcal{L}_p(Z_j, \bar{Z}_k) &= \mu_j \delta_{j,k}, \quad j, k = 1, 2, \dots, n
 \end{aligned}
 \tag{1.16}$$

and

$$\begin{aligned}
 \omega_0(x) &= (1 + O(|x|)) dx_{2n+1} + \sum_{j=1}^d 4\mu_j x_{d+j} dx_j \\
 &\quad + \sum_{j=d+1}^n 2\mu_j x_{2j} dx_{2j-1} - \sum_{j=d+1}^n 2\mu_j x_{2j-1} dx_{2j} \\
 &\quad + \sum_{j=d+1}^{2n} b_j x_{2n+1} dx_j + O(|x|^2),
 \end{aligned}
 \tag{1.17}$$

where $b_{d+1} \in \mathbb{R}, \dots, b_{2n} \in \mathbb{R}$. Put $x'' = (x_{d+1}, \dots, x_{2n+1})$, $\hat{x}'' = (x_{d+1}, x_{d+2}, \dots, x_{2d})$, $\hat{y}'' = (x_{d+1}, \dots, x_{2n})$. We have the following (see Theorem 3.11 and Theorem 3.12)

Theorem 1.11 *With the notations above, the phase function $\Phi_-(x, y) \in C^\infty(U \times U)$ is independent of (x_1, \dots, x_d) and (y_1, \dots, y_d) . Hence, $\Phi_-(x, y) = \Phi_-((0, x''), (0, y'')) := \Phi_-(x'', y'')$. Moreover, there is a constant $c > 0$ such that*

$$\text{Im } \Phi_-(x'', y'') \geq c \left(|\hat{x}''|^2 + |\hat{y}''|^2 + |\hat{x}'' - \hat{y}''|^2 \right), \quad \forall ((0, x''), (0, y'')) \in U \times U.
 \tag{1.18}$$

Furthermore,

$$\begin{aligned}
 \Phi_-(x'', y'') &= -x_{2n+1} + y_{2n+1} + 2i \sum_{j=1}^d |\mu_j| y_{d+j}^2 + 2i \sum_{j=1}^d |\mu_j| x_{d+j}^2 \\
 &\quad + i \sum_{j=d+1}^n |\mu_j| |z_j - w_j|^2 + \sum_{j=d+1}^n i \mu_j (\bar{z}_j w_j - z_j \bar{w}_j) \\
 &\quad + \sum_{j=1}^d (-b_{d+j} x_{d+j} x_{2n+1} + b_{d+j} y_{d+j} y_{2n+1}) \\
 &\quad + \sum_{j=d+1}^n \frac{1}{2} (b_{2j-1} - i b_{2j}) (-z_j x_{2n+1} + w_j y_{2n+1}) \\
 &\quad + \sum_{j=d+1}^n \frac{1}{2} (b_{2j-1} + i b_{2j}) (-\bar{z}_j x_{2n+1} + \bar{w}_j y_{2n+1}) \\
 &\quad + (x_{2n+1} - y_{2n+1}) f(x, y) + O(|(x, y)|^3),
 \end{aligned}
 \tag{1.19}$$

where $z_j = x_{2j-1} + i x_{2j}$, $w_j = y_{2j-1} + i y_{2j}$, $j = d + 1, \dots, n$, μ_j , $j = 1, \dots, n$, and $b_{d+1} \in \mathbb{R}, \dots, b_{2n} \in \mathbb{R}$ are as in (1.17) and f is smooth and satisfies $f(0, 0) = 0$, $f(x, y) = \bar{f}(y, x)$.

We now assume that X admits an S^1 action: $S^1 \times X \rightarrow X$. We will use the same notations as in Theorem 1.8. Recall that we work with Assumption 1.7. Let $p \in \mu^{-1}(0)$. We can repeat the proof of Theorem 3.7 with minor change and show that there exist local coordinates $v = (v_1, \dots, v_d)$ of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$, local coordinates $x = (x_1, \dots, x_{2n+1})$ of X defined in a neighborhood $U = U_1 \times (\hat{U}_2 \times]-2\delta, 2\delta[)$ of p with $0 \leftrightarrow p$, where $U_1 \subset \mathbb{R}^d$ is an open set of $0 \in \mathbb{R}^d$, $\hat{U}_2 \subset \mathbb{R}^{2n-d}$ is an open set of $0 \in \mathbb{R}^{2n-d}$, $\delta > 0$, and a smooth function $\gamma = (\gamma_1, \dots, \gamma_d) \in C^\infty(\hat{U}_2 \times]-2\delta, 2\delta[, U_1)$ with $\gamma(0) = 0 \in \mathbb{R}^d$ such that $T = -\frac{\partial}{\partial x_{2n+1}}$ and (1.14), (1.15), (1.16), (1.17) hold. We have the following

Theorem 1.12 *With the notations above, the phase function Ψ satisfies $\Psi(x, y) = -x_{2n+1} + y_{2n+1} + \hat{\Psi}(\hat{x}'', \hat{y}'')$, where $\hat{\Psi}(\hat{x}'', \hat{y}'') \in C^\infty(U \times U)$, $\hat{x}'' = (x_{d+1}, \dots, x_{2n})$, $\hat{y}'' = (y_{d+1}, \dots, y_{2n})$, and Ψ satisfies (1.18) and (1.19).*

2 Preliminaries

2.1 Standard notations

Let M be a C^∞ paracompact manifold. We let TM and T^*M denote the tangent bundle of M and the cotangent bundle of M , respectively. The complexified tangent bundle of M and the complexified cotangent bundle of M will be denoted by $\mathbb{C}TM$ and $\mathbb{C}T^*M$, respectively. Write $\langle \cdot, \cdot \rangle$ to denote the pointwise duality between TM and T^*M . We extend $\langle \cdot, \cdot \rangle$ bilinearly to $\mathbb{C}TM \times \mathbb{C}T^*M$. Let B be a C^∞ vector bundle over M . The fiber of B at $x \in M$ will be denoted by B_x . Let E be a vector bundle over a C^∞ paracompact manifold M_1 . We write $B \boxtimes E^*$ to denote the vector bundle over $M \times M_1$ with fiber over $(x, y) \in M \times M_1$ consisting of the linear maps from E_y to B_x . Let $Y \subset M$ be an open set. From now on, the spaces of distribution sections of B over Y and smooth sections of B over Y will be denoted by $D'(Y, B)$ and $C^\infty(Y, B)$, respectively. Let $E'(Y, B)$ be the subspace of $D'(Y, B)$ whose elements have compact support in Y .

We recall the Schwartz kernel theorem [12, Theorems 5.2.1, 5.2.6], [19, Theorem B.2.7]. Let B and E be C^∞ vector bundles over paracompact orientable C^∞ manifolds M and M_1 , respectively, equipped with smooth densities of integration. If $A : C_0^\infty(M_1, E) \rightarrow D'(M, B)$ is continuous, we write $K_A(x, y)$ or $A(x, y)$ to denote the distribution kernel of A . The following two statements are equivalent

- (1) A is continuous: $E'(M_1, E) \rightarrow C^\infty(M, B)$,
- (2) $K_A \in C^\infty(M \times M_1, B \boxtimes E^*)$.

If A satisfies (1) or (2), we say that A is smoothing on $M \times M_1$. Let $A, \hat{A} : C_0^\infty(M_1, E) \rightarrow D'(M, B)$ be continuous operators. We write

$$A \equiv \hat{A} \text{ (on } M \times M_1) \tag{2.1}$$

if $A - \hat{A}$ is a smoothing operator. If $M = M_1$, we simply write “on M ”.

Let $H(x, y) \in D'(M \times M_1, B \boxtimes E^*)$. We write H to denote the unique continuous operator $C_0^\infty(M_1, E) \rightarrow D'(M, B)$ with distribution kernel $H(x, y)$. In this work, we identify H with $H(x, y)$.

2.2 Some standard notations in semi-classical analysis

Let W_1 be an open set in \mathbb{R}^{N_1} and let W_2 be an open set in \mathbb{R}^{N_2} . Let E and F be vector bundles over W_1 and W_2 , respectively. An m -dependent continuous operator $A_m : C_0^\infty(W_2, F) \rightarrow D'(W_1, E)$ is called m -negligible on $W_1 \times W_2$ if, for m large enough, A_m is smoothing and, for any $K \Subset W_1 \times W_2$, any multi-indices α, β and any $N \in \mathbb{N}$, there exists $C_{K, \alpha, \beta, N} > 0$ such that

$$|\partial_x^\alpha \partial_y^\beta A_m(x, y)| \leq C_{K, \alpha, \beta, N} m^{-N} \text{ on } K, \quad \forall m \gg 1.$$

In that case we write

$$A_m(x, y) = O(m^{-\infty}) \text{ on } W_1 \times W_2, \quad \text{or} \quad A_m = O(m^{-\infty}) \text{ on } W_1 \times W_2.$$

If $A_m, B_m : C_0^\infty(W_2, F) \rightarrow D'(W_1, E)$ are m -dependent continuous operators, we write $A_m = B_m + O(m^{-\infty})$ on $W_1 \times W_2$ or $A_m(x, y) = B_m(x, y) + O(m^{-\infty})$ on $W_1 \times W_2$ if $A_m - B_m = O(m^{-\infty})$ on $W_1 \times W_2$. When $W = W_1 = W_2$, we sometime write “on W ”.

Let X and M be smooth manifolds and let E and F be vector bundles over X and M , respectively. Let $A_m, B_m : C^\infty(M, F) \rightarrow C^\infty(X, E)$ be m -dependent smoothing operators. We write $A_m = B_m + O(m^{-\infty})$ on $X \times M$ if on every local coordinate patch D of X and local coordinate patch D_1 of M , $A_m = B_m + O(m^{-\infty})$ on $D \times D_1$. When $X = M$, we sometime write on X .

We recall the definition of the semi-classical symbol spaces

Definition 2.1 Let W be an open set in \mathbb{R}^N . Let

$$S(1; W) := \left\{ a \in C^\infty(W) \mid \forall \alpha \in \mathbb{N}_0^N : \sup_{x \in W} |\partial^\alpha a(x)| < \infty \right\},$$

$$S_{\text{loc}}^0(1; W) := \left\{ (a(\cdot, m))_{m \in \mathbb{R}} \mid \forall \alpha \in \mathbb{N}_0^N, \forall \chi \in C_0^\infty(W) : \sup_{m \in \mathbb{R}, m \geq 1} \sup_{x \in W} |\partial^\alpha (\chi a(x, m))| < \infty \right\}.$$

For $k \in \mathbb{R}$, let

$$S_{\text{loc}}^k(1) := S_{\text{loc}}^k(1; W) = \left\{ (a(\cdot, m))_{m \in \mathbb{R}} \mid (m^{-k} a(\cdot, m)) \in S_{\text{loc}}^0(1; W) \right\}.$$

Hence $a(\cdot, m) \in S_{\text{loc}}^k(1; W)$ if for every $\alpha \in \mathbb{N}_0^N$ and $\chi \in C_0^\infty(W)$, there exists $C_\alpha > 0$ independent of m , such that $|\partial^\alpha (\chi a(\cdot, m))| \leq C_\alpha m^k$ holds on W .

Consider a sequence $a_j \in S_{\text{loc}}^{k_j}(1)$, $j \in \mathbb{N}_0$, where $k_j \searrow -\infty$, and let $a \in S_{\text{loc}}^{k_0}(1)$. We say

$$a(\cdot, m) \sim \sum_{j=0}^{\infty} a_j(\cdot, m) \text{ in } S_{\text{loc}}^{k_0}(1),$$

if, for every $\ell \in \mathbb{N}_0$, we have $a - \sum_{j=0}^{\ell} a_j \in S_{\text{loc}}^{k_{\ell+1}}(1)$. For a given sequence a_j as above, we can always find such an asymptotic sum a , which is unique up to an element in $S_{\text{loc}}^{-\infty}(1) = S_{\text{loc}}^{-\infty}(1; W) := \cap_k S_{\text{loc}}^k(1)$.

Similarly, we can define $S_{\text{loc}}^k(1; Y, E)$ in the standard way, where Y is a smooth manifold and E is a vector bundle over Y .

2.3 CR manifolds and bundles

Let $(X, T^{1,0}X)$ be a compact, connected and orientable CR manifold of dimension $2n + 1$, $n \geq 1$, where $T^{1,0}X$ is a CR structure of X , that is, $T^{1,0}X$ is a subbundle of rank n of the complexified tangent bundle $\mathbb{C}TX$, satisfying $T^{1,0}X \cap T^{0,1}X = \{0\}$, where $T^{0,1}X = \overline{T^{1,0}X}$, and $[\mathcal{V}, \mathcal{V}] \subset \mathcal{V}$, where $\mathcal{V} = C^\infty(X, T^{1,0}X)$. There is a unique subbundle HX of TX such that $\mathbb{C}HX = T^{1,0}X \oplus T^{0,1}X$, i.e. HX is the real part of $T^{1,0}X \oplus T^{0,1}X$. Let $J : HX \rightarrow HX$ be the complex structure map given by $J(u + \bar{u}) = iu - i\bar{u}$, for every $u \in T^{1,0}X$. By complex linear extension of J to $\mathbb{C}TX$, the i -eigenspace of J is $T^{1,0}X = \{V \in \mathbb{C}HX : JV = \sqrt{-1}V\}$. We shall also write (X, HX, J) to denote a compact CR manifold.

We fix a real non-vanishing 1 form $\omega_0 \in C(X, T^*X)$ so that $\langle \omega_0(x), u \rangle = 0$, for every $u \in H_x X$, for every $x \in X$. For each $x \in X$, we define a quadratic form on HX by

$$\mathcal{L}_x(U, V) = \frac{1}{2} d\omega_0(JU, V), \forall U, V \in H_x X.$$

We extend \mathcal{L} to $\mathbb{C}HX$ by complex linear extension. Then, for $U, V \in T_x^{1,0}X$,

$$\mathcal{L}_x(U, \bar{V}) = \frac{1}{2} d\omega_0(JU, \bar{V}) = -\frac{1}{2i} d\omega_0(U, \bar{V}).$$

The Hermitian quadratic form \mathcal{L}_x on $T_x^{1,0}X$ is called Levi form at x . We recall that in this paper, we always assume that the Levi form \mathcal{L} on $T^{1,0}X$ is non-degenerate of constant signature (n_-, n_+) on X , where n_- denotes the number of negative eigenvalues of the Levi form and n_+ denotes the number of positive eigenvalues of the Levi form. Let $T \in C^\infty(X, TX)$ be the non-vanishing vector field determined by

$$\omega_0(T) = -1, \quad d\omega_0(T, \cdot) \equiv 0 \text{ on } TX. \tag{2.2}$$

Note that X is a contact manifold with contact form ω_0 , contact plane HX and T is the Reeb vector field.

Fix a smooth Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ so that $T^{1,0}X$ is orthogonal to $T^{0,1}X$, $\langle u | v \rangle$ is real if u, v are real tangent vectors, $\langle T | T \rangle = 1$ and T is orthogonal to $T^{1,0}X \oplus T^{0,1}X$. For $u \in \mathbb{C}TX$, we write $|u|^2 := \langle u | u \rangle$. Denote by $T^{*1,0}X$ and $T^{*0,1}X$ the dual bundles $T^{1,0}X$ and $T^{0,1}X$, respectively. They can be identified with subbundles of the complexified cotangent bundle $\mathbb{C}T^*X$. Define the vector bundle of $(0, q)$ -forms by $T^{*0,q}X := \wedge^q T^{*0,1}X$. The Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ induces, by duality, a Hermitian metric on $\mathbb{C}T^*X$ and also on the bundles of $(0, q)$ forms $T^{*0,q}X$, $q = 0, 1, \dots, n$. We shall also denote all these induced metrics by $\langle \cdot | \cdot \rangle$. Note that we have the pointwise orthogonal decompositions:

$$\begin{aligned} \mathbb{C}T^*X &= T^{*1,0}X \oplus T^{*0,1}X \oplus \{\lambda\omega_0 : \lambda \in \mathbb{C}\}, \\ \mathbb{C}TX &= T^{1,0}X \oplus T^{0,1}X \oplus \{\lambda T : \lambda \in \mathbb{C}\}. \end{aligned}$$

For $x, y \in X$, let $d(x, y)$ denote the distance between x and y induced by the Hermitian metric $\langle \cdot | \cdot \rangle$. Let A be a subset of X . For every $x \in X$, let $d(x, A) := \inf \{d(x, y); y \in A\}$.

Let D be an open set of X . Let $\Omega^{0,q}(D)$ denote the space of smooth sections of $T^{*0,q}X$ over D and let $\Omega_c^{0,q}(D)$ be the subspace of $\Omega^{0,q}(D)$ whose elements have compact support in D .

2.4 Contact reduction

Let G be a connected compact Lie group with Lie algebra \mathfrak{g} such that $\dim_{\mathbb{R}} G = d$. We assume that the Lie group G acts on X preserving ω_0 , i.e. $g^*\omega_0 = \omega_0$, for any $g \in G$. For any $\xi \in \mathfrak{g}$, there is an induced vector field ξ_X on X given by $(\xi_X u)(x) = \frac{\partial}{\partial t} (u(\exp(t\xi) \circ x))|_{t=0}$, for any $u \in C^\infty(X)$.

Definition 2.2 The contact moment map associated to the form ω_0 is the map $\mu : X \rightarrow \mathfrak{g}^*$ such that, for all $x \in X$ and $\xi \in \mathfrak{g}$, we have

$$\langle \mu(x), \xi \rangle = \omega_0(\xi_X(x)). \tag{2.3}$$

We now recall the contact reduction from [1,9]. It was shown in [1,9] that the contact moment map is G -equivariant, so G acts on $Y := \mu^{-1}(0)$, where G acts on \mathfrak{g}^* through co-adjoint represent. Since we assume that the action of G on Y is freely, $Y_G := \mu^{-1}(0)/G$ is a smooth manifold. Let $\pi : Y \rightarrow Y_G$ and $\iota : Y \hookrightarrow X$ be the natural quotient and inclusion, respectively, then there is a unique induced contact form $\tilde{\omega}_0$ on Y_G such $\pi^*\tilde{\omega}_0 = \iota^*\omega_0$. We denote by $HY := \text{Ker } \omega_0 \cap T(\mu^{-1}(0)) = HX \cap TY$, then the induced contact plane on Y_G is $HY_G := \pi_*HY$. In particular, $\dim HY = 2n - d$ and $\dim HY_G = 2n - 2d$.

2.5 CR reduction

In this subsection we study the reduction of CR manifolds with non-degenerate Levi curvature which is a CR analogue of the reduction on complex manifolds considered in [27, §2.1]. For the case of strictly pseudoconvex CR manifolds, the CR reduction was also studied in [17].

Recall that we work with Assumption 1.2. Let b be the nondegenerate bilinear form on HX such that

$$b(\cdot, \cdot) = d\omega_0(\cdot, J\cdot). \tag{2.4}$$

We denote by $\underline{\mathfrak{g}} := \text{Span}(\xi_X, \xi \in \mathfrak{g})$ the tangent bundle of the orbits in X . Let

$$\underline{\mathfrak{g}}^{\perp b} = \left\{ v \in HX; b(\xi_X, v) = 0, \forall \xi_X \in \underline{\mathfrak{g}} \right\}. \tag{2.5}$$

Since we assume that $\underline{\mathfrak{g}}_x \cap \underline{\mathfrak{g}}_x^{\perp b} = \{0\}$, for every $x \in Y$, we immediately get

Lemma 2.3 *When restricted to $\underline{\mathfrak{g}} \times \underline{\mathfrak{g}}$, the bilinear form b is nondegenerate on Y .*

For $x \in Y, V \in H_x X$ and $\xi \in \mathfrak{g}$, by (2.3) and (2.4), we have

$$b_x(\xi_X, JV) = -d\omega_0(x)(\xi_X, V) = -(d\mu(x)(V))(\xi).$$

Therefore,

$$JV \in \underline{\mathfrak{g}}^{\perp b}|_Y \iff d\mu(x)(V) = 0. \tag{2.6}$$

Since $Y = \mu^{-1}(0)$, we have

$$d\mu(x)(V) = 0 \iff V \in T_x Y. \tag{2.7}$$

In particular, for $x \in Y$,

$$\dim \underline{\mathfrak{g}}_x^{\perp b} = \dim(H_x X \cap T_x Y) = \dim H_x Y = 2n - d.$$

By (2.2), (2.7) and the definition of $\underline{\mathfrak{g}}$, we have $\underline{\mathfrak{g}} \subset HX|_Y$. From Lemma 2.3, we can check that $\underline{\mathfrak{g}} + \underline{\mathfrak{g}}^{\perp b} = HX|_Y$. Since $\underline{\mathfrak{g}}_x \cap \underline{\mathfrak{g}}_x^{\perp b} = \{0\}$, for every $x \in Y$, this sum is a direct sum.

Let U be a small open G -invariant neighborhood of Y . Since G acts freely on Y , we can thus also assume that G acts freely on \overline{U} . Since $\underline{\mathfrak{g}}_x \cap \underline{\mathfrak{g}}_x^{\perp b} = \{0\}$, for $x \in Y$, we have, for $x \in Y$,

$$H_x U = \underline{\mathfrak{g}}_x \oplus \underline{\mathfrak{g}}_x^{\perp b}. \tag{2.8}$$

Then, by (2.8), we can choose the horizontal bundles of the fibrations $U \rightarrow U_G := U/G$ and $Y \rightarrow Y_G$ to be

$$H^H U = \underline{\mathfrak{g}}^{\perp b}|_U, \quad H^H Y := H^H U|_Y \cap HY. \tag{2.9}$$

Hence

$$HY = \underline{\mathfrak{g}}|_Y \oplus H^H Y.$$

Lemma 2.4

$$\underline{\mathfrak{g}}^{\perp b}|_Y = JHY. \tag{2.10}$$

$$HU|_Y = J\underline{\mathfrak{g}}|_Y \oplus HY = \underline{\mathfrak{g}}|_Y \oplus J\underline{\mathfrak{g}}|_Y \oplus H^H Y. \tag{2.11}$$

Proof The identity (2.10) follows from (2.6) and (2.7). For $x \in Y$, $V \in H_x Y$ and $\xi \in \mathfrak{g}$,

$$b_x(J\xi_X, V) = d\omega_0(x)(\xi_X, V) = (d\mu(x)(V))(\xi) = 0. \tag{2.12}$$

Using (2.12), $\dim H_x U = \dim H_x Y + \dim J\underline{\mathfrak{g}}_x$, and the fact that b is nondegenerate on JHY , we obtain (2.11). \square

By (2.9), and (2.10), we have $H^H Y = JHY \cap HY$. In particular, $H^H Y$ is preserved by J , so we can define the homomorphism J_G on HY_G in the following way: For $V \in HY_G$, we denote by V^H its lift in $H^H Y$, and we define J_G on Y_G by

$$(J_G V)^H = J(V^H). \tag{2.13}$$

Hence, we have $J_G : HY_G \rightarrow HY_G$ such that $J_G^2 = -\text{id}$, where id denotes the identity map $\text{id} : HY_G \rightarrow HY_G$. By complex linear extension of J_G to $\mathbb{C}TY_G$, we can define the i -eigenspace of J_G is given by $T^{1,0}Y_G = \{V \in \mathbb{C}HY_G : J_G V = \sqrt{-1}V\}$.

Theorem 2.5 *The subbundle $T^{1,0}Y_G$ is a CR structure of Y_G .*

Proof Let $u, v \in C^\infty(Y_G, T^{1,0}Y_G)$, then we can find $U, V \in C^\infty(Y_G, TY_G)$ such that

$$u = U - \sqrt{-1}J_G U, \quad v = V - \sqrt{-1}J_G V.$$

By (2.13), we have

$$u^H = U^H - \sqrt{-1}J U^H, \quad v^H = V^H - \sqrt{-1}J V^H \in T^{1,0}X \cap \mathbb{C}HY.$$

Since $T^{1,0}X$ is a CR structure and it is clearly that $[u^H, v^H] \in \mathbb{C}HY$, we have $[u^H, v^H] \in T^{1,0}X \cap \mathbb{C}HY$. Hence, there is a $W \in C^\infty(X, HX)$ such that

$$[u^H, v^H] = W - \sqrt{-1}J W.$$

In particular, $W, J W \in HY$. Thus, $W \in HY \cap JHY = H^H Y$. Let $X^H \in H^H Y$ be a lift of $X \in TY_G$ such that $X^H = W$. Then we have

$$[u, v] = \pi_*[u^H, v^H] = \pi_*(X^H - \sqrt{-1}J X^H) = X - \sqrt{-1}J_G X \in T^{1,0}Y_G,$$

i.e. we have $[C^\infty(Y_G, T^{1,0}Y_G), C^\infty(Y_G, T^{1,0}Y_G)] \subset C^\infty(Y_G, T^{1,0}Y_G)$. Therefore, $T^{1,0}Y_G$ is a CR structure of Y_G . \square

3 G-invariant Szegő kernel asymptotics

In this section, we will establish asymptotic expansion for the G -invariant Szegő kernel. We first review some known results for Szegő kernel.

3.1 Szegő kernel asymptotics

In this subsection, we don't assume that our CR manifold admits a compact Lie group action but we still assume that the Levi form is non-degenerate of constant signature (n_-, n_+) . The Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ induces, by duality, a Hermitian metric on $\mathbb{C}T^*X$ and also on the bundles of $(0, q)$ forms $T^{*0,q}X$, $q = 0, 1, \dots, n$. We shall also denote all these induced metrics by $\langle \cdot | \cdot \rangle$. For $u \in T^{*0,q}X$, we write $|u|^2 := \langle u | u \rangle$. Let $D \subset X$ be an open

set. Let $\Omega^{0,q}(D)$ denote the space of smooth sections of $T^{*0,q}X$ over D and let $\Omega_0^{0,q}(D)$ be the subspace of $\Omega^{0,q}(D)$ whose elements have compact support in D .

Let

$$\bar{\partial}_b : \Omega^{0,q}(X) \rightarrow \Omega^{0,q+1}(X)$$

be the tangential Cauchy-Riemann operator. Let $dv(x)$ be the volume form induced by the Hermitian metric $\langle \cdot | \cdot \rangle$. The natural global L^2 inner product $(\cdot | \cdot)$ on $\Omega^{0,q}(X)$ induced by $dv(x)$ and $\langle \cdot | \cdot \rangle$ is given by

$$(u | v) := \int_X \langle u(x) | v(x) \rangle dv(x), \quad u, v \in \Omega^{0,q}(X).$$

We denote by $L^2_{(0,q)}(X)$ the completion of $\Omega^{0,q}(X)$ with respect to $(\cdot | \cdot)$. We write $L^2(X) := L^2_{(0,0)}(X)$. We extend $(\cdot | \cdot)$ to $L^2_{(0,q)}(X)$ in the standard way. For $f \in L^2_{(0,q)}(X)$, we denote $\|f\|^2 := (f | f)$. We extend $\bar{\partial}_b$ to $L^2_{(0,r)}(X)$, $r = 0, 1, \dots, n$, by

$$\bar{\partial}_b : \text{Dom } \bar{\partial}_b \subset L^2_{(0,r)}(X) \rightarrow L^2_{(0,r+1)}(X),$$

where $\text{Dom } \bar{\partial}_b := \{u \in L^2_{(0,r)}(X); \bar{\partial}_b u \in L^2_{(0,r+1)}(X)\}$ and, for any $u \in L^2_{(0,r)}(X)$, $\bar{\partial}_b u$ is defined in the sense of distributions. We also write

$$\bar{\partial}_b^* : \text{Dom } \bar{\partial}_b^* \subset L^2_{(0,r+1)}(X) \rightarrow L^2_{(0,r)}(X)$$

to denote the Hilbert space adjoint of $\bar{\partial}_b$ in the L^2 space with respect to $(\cdot | \cdot)$. Let $\square_b^{(q)}$ denote the (Gaffney extension) of the Kohn Laplacian given by

$$\begin{aligned} \text{Dom } \square_b^{(q)} &= \left\{ s \in L^2_{(0,q)}(X); s \in \text{Dom } \bar{\partial}_b \cap \text{Dom } \bar{\partial}_b^*, \bar{\partial}_b s \in \text{Dom } \bar{\partial}_b^*, \bar{\partial}_b^* s \in \text{Dom } \bar{\partial}_b \right\}, \\ \square_b^{(q)} s &= \bar{\partial}_b \bar{\partial}_b^* s + \bar{\partial}_b^* \bar{\partial}_b s \text{ for } s \in \text{Dom } \square_b^{(q)}. \end{aligned} \tag{3.1}$$

By a result of Gaffney, for every $q = 0, 1, \dots, n$, $\square_b^{(q)}$ is a positive self-adjoint operator (see [19, Proposition 3.1.2]). That is, $\square_b^{(q)}$ is self-adjoint and the spectrum of $\square_b^{(q)}$ is contained in $\overline{\mathbb{R}}_+$, $q = 0, 1, \dots, n$. Let

$$S^{(q)} : L^2_{(0,q)}(X) \rightarrow \text{Ker } \square_b^{(q)} \tag{3.2}$$

be the orthogonal projections with respect to the L^2 inner product $(\cdot | \cdot)$ and let

$$S^{(q)}(x, y) \in D'(X \times X, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$$

denote the distribution kernel of $S^{(q)}$.

We recall Hörmander symbol space. Let $D \subset X$ be a local coordinate patch with local coordinates $x = (x_1, \dots, x_{2n+1})$.

Definition 3.1 For $m \in \mathbb{R}$, $S^m_{1,0}(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ is the space of all $a(x, y, t) \in C^\infty(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ such that, for all compact $K \Subset D \times D$ and all $\alpha, \beta \in \mathbb{N}_0^{2n+1}, \gamma \in \mathbb{N}_0$, there is a constant $C_{\alpha,\beta,\gamma} > 0$ such that

$$|\partial_x^\alpha \partial_y^\beta \partial_t^\gamma a(x, y, t)| \leq C_{\alpha,\beta,\gamma} (1 + |t|)^{m-\gamma}, \quad \forall (x, y, t) \in K \times \mathbb{R}_+, t \geq 1.$$

Put

$$S^{-\infty}(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*) := \bigcap_{m \in \mathbb{R}} S^m_{1,0}(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*).$$

Let $a_j \in S_{1,0}^{m_j}(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$, $j = 0, 1, 2, \dots$ with $m_j \rightarrow -\infty$, as $j \rightarrow \infty$. Then there exists $a \in S_{1,0}^{m_0}(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ unique modulo $S^{-\infty}$, such that $a - \sum_{j=0}^{k-1} a_j \in S_{1,0}^{m_k}(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ for $k = 0, 1, 2, \dots$

If a and a_j have the properties above, we write $a \sim \sum_{j=0}^{\infty} a_j$ in $S_{1,0}^{m_0}(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$. We write

$$s(x, y, t) \in S_{\text{cl}}^m(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$$

if $s(x, y, t) \in S_{1,0}^m(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ and

$$s(x, y, t) \sim \sum_{j=0}^{\infty} s^j(x, y) t^{m-j} \text{ in } S_{1,0}^m(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*),$$

$$s^j(x, y) \in C^\infty(D \times D, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \quad j \in \mathbb{N}_0.$$

The following was proved in Theorem 4.8 in [15]

Theorem 3.2 *Given $q = 0, 1, 2, \dots, n$. Assume that $q \notin \{n_-, n_+\}$. Then, $S^{(q)} \equiv 0$ on X .*

We have the following (see Theorem 1.2 in [13], Theorem 4.7 in [15] and see also [2] for $q = 0$)

Theorem 3.3 *We recall that we work with the assumption that the Levi form is non-degenerate of constant signature (n_-, n_+) on X . Let $q = n_-$ or n_+ . Suppose that $\square_b^{(q)}$ has L^2 closed range. Then, $S^{(q)}(x, y) \in C^\infty(X \times X \setminus \text{diag}(X \times X), T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$. Let $D \subset X$ be any local coordinate patch with local coordinates $x = (x_1, \dots, x_{2n+1})$. Then, there exist continuous operators $S_-, S_+ : \Omega_0^{0,q}(D) \rightarrow D'(D, T^{*0,q}X)$ such that*

$$S^{(q)} \equiv S_- + S_+ \text{ on } D,$$

and $S_-(x, y), S_+(x, y)$ satisfy

$$S_{\mp}(x, y) \equiv \int_0^\infty e^{i\varphi_{\mp}(x,y)t} s_{\mp}(x, y, t) dt \text{ on } D,$$

with

$$s_-(x, y, t), s_+(x, y, t) \in S_{\text{cl}}^n(D \times D \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*),$$

$$s_-(x, y, t) = 0 \text{ if } q \neq n_-, \quad s_+(x, y, t) = 0 \text{ if } q \neq n_+,$$

$$s_-^0(x, x) \neq 0, \quad \forall x \in D, \quad s_+^0(x, x) \neq 0, \quad \forall x \in D, \tag{3.3}$$

and the phase functions φ_-, φ_+ satisfy

$$\varphi_+(x, y), \varphi_- \in C^\infty(D \times D), \quad \text{Im } \varphi_-(x, y) \geq 0,$$

$$\varphi_-(x, x) = 0, \quad \varphi_-(x, y) \neq 0 \text{ if } x \neq y,$$

$$d_x \varphi_-(x, y)|_{x=y} = -\omega_0(x), \quad d_y \varphi_-(x, y)|_{x=y} = \omega_0(x),$$

$$\varphi_-(x, y) = -\overline{\varphi_-}(y, x), \quad -\overline{\varphi_+}(x, y) = \varphi_-(x, y).$$

Remark 3.4 It is well-known that for a strictly pseudoconvex CR manifold of dimension 3, $\square_b^{(0)}$ does not have L^2 closed range in general (see [28]). Kohn [16] proved that if $q = n_- = n_+$ or $|n_- - n_+| > 1$ then $\square_b^{(q)}$ has L^2 closed range.

The following result describes the phase function in local coordinates (see chapter 8 of part I in [13])

Theorem 3.5 For a given point $p \in X$, let $\{W_j\}_{j=1}^n$ be an orthonormal frame of $T^{1,0}X$ in a neighborhood of p such that the Levi form is diagonal at p , i.e. $\mathcal{L}_{x_0}(W_j, \bar{W}_s) = \delta_{j,s}\mu_j$, $j, s = 1, \dots, n$. We take local coordinates $x = (x_1, \dots, x_{2n+1})$, $z_j = x_j + ix_{d+j}$, $j = 1, \dots, d$, $z_j = x_{2j-1} + ix_{2j}$, $j = d + 1, \dots, n$, defined on some neighborhood of p such that $\omega_0(p) = dx_{2n+1}$, $x(p) = 0$, and, for some $c_j \in \mathbb{C}$, $j = 1, \dots, n$,

$$W_j = \frac{\partial}{\partial z_j} - i\mu_j \bar{z}_j \frac{\partial}{\partial x_{2n+1}} - c_j x_{2n+1} \frac{\partial}{\partial x_{2n+1}} + \sum_{k=1}^{2n} a_{j,k}(x) \frac{\partial}{\partial x_k} + O(|x|^2), \quad j = 1, \dots, n, \tag{3.4}$$

where $a_{j,k}(x) \in C^\infty$, $a_{j,k}(x) = O(|x|)$, for every $j = 1, \dots, n$, $k = 1, \dots, 2n$. Set $y = (y_1, \dots, y_{2n+1})$, $w_j = y_j + iy_{d+j}$, $j = 1, \dots, d$, $w_j = y_{2j-1} + iy_{2j}$, $j = d + 1, \dots, n$. Then, for φ_- in Theorem 3.3, we have

$$\text{Im } \varphi_-(x, y) \geq c \sum_{j=1}^{2n} |x_j - y_j|^2, \quad c > 0, \tag{3.5}$$

in some neighbourhood of $(0, 0)$ and

$$\begin{aligned} \varphi_-(x, y) &= -x_{2n+1} + y_{2n+1} + i \sum_{j=1}^n |\mu_j| |z_j - w_j|^2 + \sum_{j=1}^n \left(i\mu_j (\bar{z}_j w_j - z_j \bar{w}_j) \right. \\ &\quad \left. + c_j (-z_j x_{2n+1} + w_j y_{2n+1}) + \bar{c}_j (-\bar{z}_j x_{2n+1} + \bar{w}_j y_{2n+1}) \right) \\ &\quad + (x_{2n+1} - y_{2n+1}) f(x, y) + O(|(x, y)|^3), \end{aligned} \tag{3.6}$$

where f is smooth and satisfies $f(0, 0) = 0$, $f(x, y) = \bar{f}(y, x)$.

The following formula for the leading term s_-^0 on the diagonal follows from [13, §9]. The formula for the leading term s_+^0 on the diagonal follows similarly.

Theorem 3.6 We assume that the Levi form is non-degenerate of constant signature (n_-, n_+) at each point of X . Suppose that $\square_b^{(q)}$ has L^2 closed range. If $q = n_\mp$, then, for the leading term $s_\mp^0(x, y)$ of the expansion (3.3) of $s_\mp(x, y, t)$, we have

$$s_\mp^0(x_0, x_0) = \frac{1}{2} \pi^{-n-1} |\det \mathcal{L}_{x_0}| \tau_{x_0, n_\mp}, \quad x_0 \in D,$$

where $\det \mathcal{L}_{x_0}$ is given by (1.6) and τ_{x_0, n_\mp} is given by (1.7).

3.2 G-invariant Szegő kernel

Fix $g \in G$. Let $g^* : \Lambda_x^r(\mathbb{C}T^*X) \rightarrow \Lambda_{g^{-1} \circ x}^r(\mathbb{C}T^*X)$ be the pull-back map. Since G preserves J , we have $g^* : T_x^{*0,q}X \rightarrow T_{g^{-1} \circ x}^{*0,q}X$, $\forall x \in X$. Thus, for $u \in \Omega^{0,q}(X)$, we have $g^*u \in \Omega^{0,q}(X)$ and we write $(g^*u)(x) := u(g \circ x)$. Put $\Omega^{0,q}(X)^G := \{u \in \Omega^{0,q}(X); g^*u = u, \forall g \in G\}$. Now, we assume that the Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ is G -invariant and \mathfrak{g} is orthogonal to $HY \cap JHY$ at every point of Y . The Hermitian metric is G -invariant means that, for any G -invariant vector fields U and V , $\langle U | V \rangle$ is G -invariant. Then the L^2 inner product $(\cdot | \cdot)$ on $\Omega^{0,q}(X)$ induced by $\langle \cdot | \cdot \rangle$ is G -invariant, that is, $(u | v) = (g^*u | g^*v)$, for all $u, v \in \Omega^{0,q}(X)$, $g \in G$. Let $u \in L^2_{(0,q)}(X)$ and let $g \in G$. Take $u_j \in \Omega^{0,q}(X)$, $j = 1, 2, \dots$, with $u_j \rightarrow u$ in $L^2_{(0,q)}(X)$ as $j \rightarrow \infty$. Since $(\cdot | \cdot)$ is G -invariant, there is a $v \in L^2_{(0,q)}(X)$ such that $v = \lim_{j \rightarrow \infty} g^*u_j$. We define $g^*u := v$.

It is clear that the definition is well-defined. We have $g^* : L^2_{(0,q)}(X) \rightarrow L^2_{(0,q)}(X)$. Put $L^2_{(0,q)}(X)^G := \left\{ u \in L^2_{(0,q)}(X); g^*u = u, \forall g \in G \right\}$. It is not difficult to see that $L^2_{(0,q)}(X)^G$ is the completion of $\Omega^{0,q}(X)^G$ with respect to $(\cdot | \cdot)$. We write $L^2(X)^G := L^2_{(0,0)}(X)^G$. Since G preserves J and $(\cdot | \cdot)$ is G -invariant, it is straightforward to see that

$$g^*\bar{\partial}_b = \bar{\partial}_b g^* \text{ on } \text{Dom } \bar{\partial}_b, \quad g^*\bar{\partial}_b^* = \bar{\partial}_b^* g^* \text{ on } \text{Dom } \bar{\partial}_b^*,$$

$$g^*\square_b^{(q)} = \square_b^{(q)} g^* \text{ on } \text{Dom } \square_b^{(q)}.$$

Put $(\text{Ker } \square_b^{(q)})^G := \text{Ker } \square_b^{(q)} \cap L^2_{(0,q)}(X)^G$. The G -invariant Szegő projection is the orthogonal projection $S_G^{(q)} : L^2_{(0,q)}(X) \rightarrow (\text{Ker } \square_b^{(q)})^G$ with respect to $(\cdot | \cdot)$. Let $S_G^{(q)}(x, y) \in D'(X \times X, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ be the distribution kernel of S^G . Let $d\mu$ be a Haar measure on G so that $|G|d\mu := \int_G d\mu = 1$. Then,

$$S_G^{(q)}(x, y) = \int_G S^{(q)}(x, g \circ y) d\mu(g). \tag{3.7}$$

Note that the integral (3.7) is defined in the sense of distribution.

3.3 G-invariant Szegő kernel asymptotics near $\mu^{-1}(0)$

In this section, we will study G -invariant Szegő kernel near $\mu^{-1}(0)$.

Let $e_0 \in G$ be the identity element. Let $v = (v_1, \dots, v_d)$ be the local coordinates of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$. From now on, we will identify the element $e \in V$ with $v(e)$. We first need

Theorem 3.7 *Let $p \in \mu^{-1}(0)$. There exist local coordinates $v = (v_1, \dots, v_d)$ of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$, local coordinates $x = (x_1, \dots, x_{2n+1})$ of X defined in a neighborhood $U = U_1 \times U_2$ of p with $0 \leftrightarrow p$, where $U_1 \subset \mathbb{R}^d$ is an open set of $0 \in \mathbb{R}^d$, $U_2 \subset \mathbb{R}^{2n+1-d}$ is an open set of $0 \in \mathbb{R}^{2n+1-d}$ and a smooth function $\gamma = (\gamma_1, \dots, \gamma_d) \in C^\infty(U_2, U_1)$ with $\gamma(0) = 0 \in \mathbb{R}^d$ such that*

$$(v_1, \dots, v_d) \circ (\gamma(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1})$$

$$= (v_1 + \gamma_1(x_{d+1}, \dots, x_{2n+1}), \dots, v_d + \gamma_d(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1}), \tag{3.8}$$

$$\forall (v_1, \dots, v_d) \in V, \quad \forall (x_{d+1}, \dots, x_{2n+1}) \in U_2,$$

$$\underline{g} = \text{span} \left\{ \frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_d} \right\},$$

$$\mu^{-1}(0) \cap U = \{x_{d+1} = \dots = x_{2d} = 0\}, \tag{3.9}$$

On $\mu^{-1}(0) \cap U$, we have $J\left(\frac{\partial}{\partial x_j}\right) = \frac{\partial}{\partial x_{d+j}} + a_j(x) \frac{\partial}{\partial x_{2n+1}}$, $j = 1, 2, \dots, d$,

where $a_j(x)$ is a smooth function on $\mu^{-1}(0) \cap U$, independent of $x_1, \dots, x_{2d}, x_{2n+1}$ and $a_j(0) = 0$, $j = 1, \dots, d$,

$$T_p^{1,0}X = \text{span} \{Z_1, \dots, Z_n\},$$

$$Z_j = \frac{1}{2} \left(\frac{\partial}{\partial x_j} - i \frac{\partial}{\partial x_{d+j}} \right)(p), \quad j = 1, \dots, d,$$

$$Z_j = \frac{1}{2} \left(\frac{\partial}{\partial x_{2j-1}} - i \frac{\partial}{\partial x_{2j}} \right)(p), \quad j = d + 1, \dots, n, \tag{3.10}$$

$$\langle Z_j | Z_k \rangle = \delta_{j,k}, \quad j, k = 1, 2, \dots, n,$$

$$\mathcal{L}_p(Z_j, \bar{Z}_k) = \mu_j \delta_{j,k}, \quad j, k = 1, 2, \dots, n$$

and

$$\begin{aligned} \omega_0(x) &= (1 + O(|x|))dx_{2n+1} + \sum_{j=1}^d 4\mu_j x_{d+j} dx_j \\ &\quad + \sum_{j=d+1}^n 2\mu_j x_{2j} dx_{2j-1} - \sum_{j=d+1}^n 2\mu_j x_{2j-1} dx_{2j} \\ &\quad + \sum_{j=d+1}^{2n} b_j x_{2n+1} dx_j + O(|x|^2), \end{aligned} \tag{3.11}$$

where $b_{d+1} \in \mathbb{R}, \dots, b_{2n} \in \mathbb{R}$.

Proof From the standard proof of Frobenius Theorem, it is not difficult to see that there exist local coordinates $v = (v_1, \dots, v_d)$ of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$ and local coordinates $x = (x_1, \dots, x_{2n+1})$ of X defined in a neighborhood U of p with $x(p) = 0$ such that

$$\begin{aligned} &(v_1, \dots, v_d) \circ (0, \dots, 0, x_{d+1}, \dots, x_{2n+1}) \\ &= (v_1, \dots, v_d, x_{d+1}, \dots, x_{2n+1}), \quad \forall (v_1, \dots, v_d) \in V, \quad \forall (0, \dots, 0, x_{d+1}, \dots, x_{2n+1}) \in U, \end{aligned} \tag{3.12}$$

and

$$\underline{\mathfrak{g}} = \text{span} \left\{ \frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_d} \right\}. \tag{3.13}$$

Since $p \in \mu^{-1}(0)$, we have $\omega_0(p)(\frac{\partial}{\partial x_j}(p)) = 0, j = 1, 2, \dots, d$, and hence $\frac{\partial}{\partial x_j}(p) \in H_p X, j = 1, 2, \dots, d$. Consider the linear map

$$\begin{aligned} R : \underline{\mathfrak{g}}_p &\rightarrow \underline{\mathfrak{g}}_p, \\ u &\rightarrow Ru, \quad \langle Ru | v \rangle = \langle d\omega_0, Ju \wedge v \rangle. \end{aligned}$$

Since R is self-adjoint, by using linear transformation in (x_1, \dots, x_d) , we can take (x_1, \dots, x_d) such that, for $j, k = 1, 2, \dots, d$,

$$\left\langle R \frac{\partial}{\partial x_j}(p) \mid \frac{\partial}{\partial x_k}(p) \right\rangle = 4\mu_j \delta_{j,k}, \quad \left\langle \frac{\partial}{\partial x_j}(p) \mid \frac{\partial}{\partial x_k}(p) \right\rangle = 2\delta_{j,k}. \tag{3.14}$$

By taking linear transformation in (v_1, \dots, v_d) , (3.12) still hold.

Let $\omega_0(\frac{\partial}{\partial x_j}) = a_j(x) \in C^\infty(U), j = 1, 2, \dots, d$. Since $a_j(x)$ is G -invariant, we have $\frac{\partial a_j(x)}{\partial x_s} = 0, j, s = 1, 2, \dots, d$. By the definition of the moment map, we have

$$\mu^{-1}(0) \cap U = \{x \in U; a_1(x) = \dots = a_d(x) = 0\}.$$

Since p is a regular value of the moment map μ , the matrix $\left(\frac{\partial a_j}{\partial x_s}(p)\right)_{1 \leq j \leq d, d+1 \leq s \leq 2n+1}$ is of rank d . We may assume that the matrix $\left(\frac{\partial a_j}{\partial x_s}(p)\right)_{1 \leq j \leq d, d+1 \leq s \leq 2d}$ is non-singular. Thus, $(x_1, \dots, x_d, a_1, \dots, a_d, x_{2d+1}, \dots, x_{2n+1})$ are also local coordinates of X . Hence, we can take $v = (v_1, \dots, v_d)$ and $x = (x_1, \dots, x_{2n+1})$ such that (3.12), (3.13), (3.14) hold and

$$\mu^{-1}(0) \cap U = \{x = (x_1, \dots, x_{2n+1}) \in U; x_{d+1} = \dots = x_{2d} = 0\}. \tag{3.15}$$

On $\mu^{-1}(0) \cap U$, let

$$J\left(\frac{\partial}{\partial x_j}\right) = b_{j,1}(x) \frac{\partial}{\partial x_1} + \dots + b_{j,2n+1}(x) \frac{\partial}{\partial x_{2n+1}}, \quad j = 1, 2, \dots, d.$$

Since we only work on $\mu^{-1}(0), b_{j,k}(x)$ is independent of x_{d+1}, \dots, x_{2d} , for all $j = 1, \dots, d, k = 1, \dots, 2n + 1$. Moreover, it is easy to see that $b_{j,k}(x)$ is also independent of x_1, \dots, x_d , for all $j = 1, \dots, d, k = 1, \dots, 2n + 1$. Let $\tilde{x}'' = (x_{2d+1}, \dots, x_{2n+1})$.

Hence, $b_{j,k}(x) = b_{j,k}(\tilde{x}'')$, $j = 1, \dots, d, k = 1, \dots, 2n + 1$. We claim that the matrix $(b_{j,k}(\tilde{x}''))_{1 \leq j \leq d, d+1 \leq k \leq 2d}$ is non-singular near p . If not, it is easy to see that there is a non-zero vector $u \in J\mathfrak{g} \cap HY$, where $Y = \mu^{-1}(0)$. Let $u = Jv, v \in \mathfrak{g}$. Then, $v \in \mathfrak{g} \cap JHY = \mathfrak{g} \cap \mathfrak{g}^{\perp b}$ (see (2.10)). Since $\mathfrak{g} \cap \mathfrak{g}^{\perp b} = \{0\}$ on $\mu^{-1}(0)$, we deduce that $v = 0$ and we get a contradiction. The claim follows. From the claim, we can use linear transformation in (x_{d+1}, \dots, x_{2d}) (the linear transform depends smoothly on \tilde{x}'') and we can take (x_{d+1}, \dots, x_{2d}) such that on $\mu^{-1}(0)$,

$$J\left(\frac{\partial}{\partial x_j}\right) = b_{j,1}(\tilde{x}'')\frac{\partial}{\partial x_1} + \dots + b_{j,d}(\tilde{x}'')\frac{\partial}{\partial x_d} + \frac{\partial}{\partial x_{d+j}} + b_{j,2d+1}(\tilde{x}'')\frac{\partial}{\partial x_{2d+1}} + \dots + b_{j,2n+1}(\tilde{x}'')\frac{\partial}{\partial x_{2n+1}},$$

where $j = 1, 2, \dots, d$. Consider the coordinates change:

$$\begin{aligned} x &= (x_1, \dots, x_{2n+1}) \rightarrow u = (u_1, \dots, u_{2n+1}), \\ (x_1, \dots, x_{2n+1}) &\rightarrow (x_1 - \sum_{j=1}^d b_{j,1}(\tilde{x}'')x_{d+j}, \dots, x_d - \sum_{j=1}^d b_{j,d}(\tilde{x}'')x_{d+j}, x_{d+1}, \dots, x_{2d}, \\ &x_{2d+1} - \sum_{j=1}^d b_{j,2d+1}(\tilde{x}'')x_{d+j}, \dots, x_{2n+1} - \sum_{j=1}^d b_{j,2n+1}(\tilde{x}'')x_{d+j}). \end{aligned}$$

Then,

$$\begin{aligned} \frac{\partial}{\partial x_j} &\rightarrow \frac{\partial}{\partial u_j}, \quad j = 1, \dots, d, 2d + 1, \dots, 2n + 1, \\ \frac{\partial}{\partial x_{d+j}} &\rightarrow -b_{j,1}\frac{\partial}{\partial u_1} - \dots - b_{j,d}\frac{\partial}{\partial u_d} + \frac{\partial}{\partial u_{d+j}} \\ &- b_{j,2d+1}\frac{\partial}{\partial u_{2d+1}} - \dots - b_{j,2n+1}\frac{\partial}{\partial u_{2n+1}}, \quad j = 1, \dots, d. \end{aligned}$$

Hence, on $\mu^{-1}(0) \cap U, J\left(\frac{\partial}{\partial x_j}\right) \rightarrow \frac{\partial}{\partial u_{d+j}}, j = 1, \dots, d$. Thus, we can take $v = (v_1, \dots, v_d)$ and $x = (x_1, \dots, x_{2n+1})$ such that (3.8), (3.13), (3.14), (3.15) hold and on $\mu^{-1}(0) \cap U,$

$$J\left(\frac{\partial}{\partial x_j}\right) = \frac{\partial}{\partial x_{d+j}}, \quad j = 1, 2, \dots, d.$$

Let $Z_j = \frac{1}{2}\left(\frac{\partial}{\partial x_j} - i\frac{\partial}{\partial x_{d+j}}\right)(p) \in T_p^{1,0}X, j = 1, \dots, d$. From (3.14), we can check that

$$\mathcal{L}_p(Z_j, \bar{Z}_k) = \mu_j \delta_{j,k}, \quad \langle Z_j | Z_k \rangle = \delta_{j,k}, \quad j, k = 1, \dots, d.$$

Since \mathfrak{g}_p is orthogonal to $H_pY \cap JH_pY$ and $H_pY \cap JH_pY \subset \mathfrak{g}_p^{\perp b}$, we can find an orthonormal frame $\{Z_1, \dots, Z_d, V_1, \dots, V_{n-d}\}$ for $T_p^{1,0}X$ such that the Levi form \mathcal{L}_p is diagonalized with respect to $Z_1, \dots, Z_d, V_1, \dots, V_{n-d}$, where $V_1 \in \mathbb{C}H_pY \cap J\mathbb{C}H_pY, \dots, V_{n-d} \in \mathbb{C}H_pY \cap J\mathbb{C}H_pY$. Write

$$\operatorname{Re} V_j = \sum_{k=1}^{2n+1} \alpha_{j,k} \frac{\partial}{\partial x_k}, \quad \operatorname{Im} V_j = \sum_{k=1}^{2n+1} \beta_{j,k} \frac{\partial}{\partial x_k}, \quad j = 1, \dots, n - d.$$

We claim that $\alpha_{j,k} = \beta_{j,k} = 0$, for all $k = d + 1, \dots, 2d, j = 1, \dots, n - d$. Fix $j = 1, \dots, n - d$. Since $\operatorname{Re} V_j \in \mathfrak{g}_p^{\perp b}$ and $\operatorname{span} \left\{ \frac{\partial}{\partial x_{d+1}}, \dots, \frac{\partial}{\partial x_{2d}} \right\} \in \mathfrak{g}_p^{\perp b}$, we conclude that

$$\sum_{k=1}^d \alpha_{j,k} \frac{\partial}{\partial x_k} + \sum_{k=2d+1}^{2n+1} \alpha_{j,k} \frac{\partial}{\partial x_k} \in \mathfrak{g}_p^{\perp b} \cap H_pY. \tag{3.16}$$

From (2.10) and (3.16), we deduce that

$$\sum_{k=1}^d \alpha_{j,k} \frac{\partial}{\partial x_k} + \sum_{k=2d+1}^{2n+1} \alpha_{j,k} \frac{\partial}{\partial x_k} \in JH_p Y \cap H_p Y = \underline{\mathfrak{g}}_p^{\perp b} \cap H_p Y$$

and hence

$$J\left(\sum_{k=1}^d \alpha_{j,k} \frac{\partial}{\partial x_k} + \sum_{k=2d+1}^{2n+1} \alpha_{j,k} \frac{\partial}{\partial x_k}\right) \in \underline{\mathfrak{g}}_p^{\perp b} \cap H_p Y. \tag{3.17}$$

From (3.17) and notice that $J(\text{Re } V_j) \in \underline{\mathfrak{g}}_p^{\perp b}$, we deduce that

$$J\left(\sum_{k=d+1}^{2d} \alpha_{j,k} \frac{\partial}{\partial x_k}\right) \in \underline{\mathfrak{g}}_p \cap \underline{\mathfrak{g}}_p^{\perp b} = \{0\}.$$

Thus, $\alpha_{j,k} = 0$, for all $k = d + 1, \dots, 2d, j = 1, \dots, n - d$. Similarly, we can repeat the procedure above and deduce that $\beta_{j,k} = 0$, for all $k = d + 1, \dots, 2d, j = 1, \dots, n - d$.

Since $\text{span}\{\text{Re } V_j, \text{Im } V_j; j = 1, \dots, n - d\}$ is transversal to $\underline{\mathfrak{g}}_p \oplus J\underline{\mathfrak{g}}_p$, we can take linear transformation in $(x_{2d+1}, \dots, x_{2n+1})$ so that

$$\begin{aligned} \text{Re } V_j &= \alpha_{j,1} \frac{\partial}{\partial x_1} + \dots + \alpha_{j,d} \frac{\partial}{\partial x_d} + \frac{\partial}{\partial x_{2j-1+2d}}, \quad j = 1, 2, \dots, n - d, \\ \text{Im } V_j &= \beta_{j,1} \frac{\partial}{\partial x_1} + \dots + \beta_{j,d} \frac{\partial}{\partial x_d} + \frac{\partial}{\partial x_{2j+2d}}, \quad j = 1, 2, \dots, n - d. \end{aligned}$$

Consider the coordinates change:

$$\begin{aligned} x &= (x_1, \dots, x_{2n+1}) \rightarrow u = (u_1, \dots, u_{2n+1}), \\ (x_1, \dots, x_{2n+1}) &\rightarrow (x_1 - \sum_{j=1}^d \alpha_{j,1} x_{2j-1+2d} - \sum_{j=1}^d \beta_{j,1} x_{2j+2d}, \dots, x_d \\ &\quad - \sum_{j=1}^d \alpha_{j,d} x_{2j-1+2d} - \sum_{j=1}^d \beta_{j,d} x_{2j+2d}, x_{d+1}, \dots, x_{2n+1}) \end{aligned}$$

Then,

$$\begin{aligned} \frac{\partial}{\partial x_j} &\rightarrow \frac{\partial}{\partial u_j}, \quad j = 1, \dots, 2d, \\ \frac{\partial}{\partial x_{2j-1+2d}} &\rightarrow -\alpha_{j,1} \frac{\partial}{\partial u_1} - \dots - \alpha_{j,d} \frac{\partial}{\partial u_d} + \frac{\partial}{\partial u_{2j-1+2d}}, \quad j = 1, \dots, n - d, \\ \frac{\partial}{\partial x_{2j+2d}} &\rightarrow -\beta_{j,1} \frac{\partial}{\partial u_1} - \dots - \beta_{j,d} \frac{\partial}{\partial u_d} + \frac{\partial}{\partial u_{2j+2d}}, \quad j = 1, \dots, n - d. \end{aligned}$$

Thus, we can take $v = (v_1, \dots, v_d)$ and $x = (x_1, \dots, x_{2n+1})$ such that (3.8), (3.9) and (3.10) hold.

Now, we can take linear transformation in x_{2n+1} so that $\omega_0(p) = dx_{2n+1}$. Let $W_j, j = 1, \dots, n$ be an orthonormal basis of $T^{1,0}X$ such that $W_j(p) = Z_j, j = 1, \dots, n$, where $Z_j \in T_p^{1,0}X, j = 1, \dots, n$, are as in (3.10). Let $\tilde{x} = (\tilde{x}_1, \dots, \tilde{x}_{2n+1})$ be the coordinates as in Theorem 3.5. It is easy to see that

$$\begin{aligned} \tilde{x}_j &= x_j + a_j x_{2n+1} + h_j(x), \quad h_j(x) = O(|x|^2), \quad a_j \in \mathbb{R}, \quad j = 1, 2, \dots, 2n, \\ \tilde{x}_{2n+1} &= x_{2n+1} + h_{2n+1}(x), \quad h_{2n+1}(x) = O(|x|^2). \end{aligned} \tag{3.18}$$

We may change x_{2n+1} be $x_{2n+1} + h_{2n+1}(0, \dots, 0, x_{d+1}, \dots, x_{2n}, 0)$ and we have

$$\frac{\partial^2 \tilde{x}_{2n+1}}{\partial x_j \partial x_k}(p) = 0, \quad j, k = \{d + 1, \dots, 2n\}. \tag{3.19}$$

Note that when we change x_{2n+1} to $x_{2n+1} + h_{2n+1}(0, \dots, 0, x_{d+1}, \dots, x_{2n}, 0)$, $\frac{\partial}{\partial x_j}$ will change to $\frac{\partial}{\partial x_j} + \alpha_j(x) \frac{\partial}{\partial x_{2n+1}}, j = d + 1, \dots, 2n$, where $\alpha_j(x)$ is a smooth function on

$\mu^{-1}(0) \cap U$, independent of $x_1, \dots, x_d, x_{2n+1}$ and $\alpha_j(0) = 0, j = d + 1, \dots, 2n$. Hence, on $\mu^{-1}(0) \cap U$, we have $J(\frac{\partial}{\partial x_j}) = \frac{\partial}{\partial x_{d+j}} + a_j(x) \frac{\partial}{\partial x_{2n+1}}, j = 1, 2, \dots, d$, where $a_j(x)$ is a smooth function on $\mu^{-1}(0) \cap U$, independent of $x_1, \dots, x_{2d}, x_{2n+1}$ and $a_j(0) = 0, j = 1, \dots, d$.

From (3.4) and (3.18), it is straightforward to see that

$$\begin{aligned} \omega_0(\tilde{x}) &= (1 + O(|\tilde{x}|))d\tilde{x}_{2n+1} + \sum_{j=1}^d 2\mu_j \tilde{x}_{d+j} d\tilde{x}_j + \sum_{j=1}^d (-2\mu_j \tilde{x}_j) d\tilde{x}_{d+j} \\ &\quad + \sum_{j=d+1}^n 2\mu_j x_{2j} dx_{2j-1} + \sum_{j=d+1}^n (-2\mu_j \tilde{x}_{2j-1}) d\tilde{x}_{2j} + \sum_{j=1}^{2n} \hat{b}_j \tilde{x}_{2n+1} d\tilde{x}_j + O(|x|^2) \\ &= (1 + O(|x|))dx_{2n+1} + \sum_{j=1}^d (2\mu_j x_{d+j} + \frac{\partial \tilde{x}_{2n+1}}{\partial x_j}) dx_j + \sum_{j=1}^d (-2\mu_j x_j + \frac{\partial \tilde{x}_{2n+1}}{\partial x_{d+j}}) dx_{d+j} \\ &\quad + \sum_{j=d+1}^n (2\mu_j x_{2j} + \frac{\partial \tilde{x}_{2n+1}}{\partial x_{2j-1}}) dx_{2j-1} + \sum_{j=d+1}^n (-2\mu_j x_{2j-1} + \frac{\partial \tilde{x}_{2n+1}}{\partial x_{2j}}) dx_{2j} \\ &\quad + \sum_{j=1}^{2n} \hat{b}_j x_{2n+1} dx_j + O(|x|^2), \end{aligned} \tag{3.20}$$

where $\tilde{b}_j \in \mathbb{R}, \hat{b}_j \in \mathbb{R}, j = 1, \dots, 2n$. Note that ω_0 is G -invariant. From this observation and (3.20), we deduce that

$$\begin{aligned} \frac{\partial^2 \tilde{x}_{2n+1}}{\partial x_j \partial x_k}(p) &= 0, \quad j \in \{1, \dots, d\}, k \in \{1, \dots, d\} \cup \{2d + 1, \dots, 2n\}, \\ \frac{\partial^2 \tilde{x}_{2n+1}}{\partial x_{d+j} \partial x_k}(p) &= 2\mu_j \delta_{j,k}, \quad j, k \in \{1, \dots, d\}. \end{aligned} \tag{3.21}$$

From (3.21), (3.20) and (3.19), it is straightforward to see that

$$\begin{aligned} \omega_0(x) &= (1 + O(|x|))dx_{2n+1} + \sum_{j=1}^d 4\mu_j x_{d+j} dx_j \\ &\quad + \sum_{j=d+1}^n 2\mu_j x_{2j} dx_{2j-1} - \sum_{j=d+1}^n 2\mu_j x_{2j-1} dx_{2j} + \sum_{j=1}^{2n} b_j x_{2n+1} dx_j + O(|x|^2), \end{aligned} \tag{3.22}$$

where $b_1 \in \mathbb{R}, \dots, b_{2n} \in \mathbb{R}$. Since $\omega_0(p)(\frac{\partial}{\partial x_j}) = 0$ on $x_{d+1} = \dots = x_{2d} = 0, j = 1, 2, \dots, d$, we deduce that $b_1 = \dots = b_d = 0$ and we get (3.11). The theorem follows. \square

We need

Theorem 3.8 *Let $p \in \mu^{-1}(0)$ and take local coordinates $x = (x_1, \dots, x_{2n+1})$ of X defined in an open set U of p with $0 \leftrightarrow p$ such that (3.9), (3.10) and (3.11) hold. Let $\varphi_-(x, y) \in C^\infty(U \times U)$ be as in Theorem 3.3. Then,*

$$\begin{aligned} \varphi_-(x, y) &= -x_{2n+1} + y_{2n+1} - 2 \sum_{j=1}^d \mu_j x_j x_{d+j} + 2 \sum_{j=1}^d \mu_j y_j y_{d+j} + i \sum_{j=1}^n |\mu_j| |z_j - w_j|^2 \\ &\quad + \sum_{j=1}^n i \mu_j (\bar{z}_j w_j - z_j \bar{w}_j) + \sum_{j=1}^d (-\frac{1}{2} b_{d+j}) (-z_j x_{2n+1} + w_j y_{2n+1}) \\ &\quad + \sum_{j=1}^d (\frac{1}{2} b_{d+j}) (-\bar{z}_j x_{2n+1} + \bar{w}_j y_{2n+1}) + \sum_{j=d+1}^n \frac{1}{2} (b_{2j-1} - i b_{2j}) (-z_j x_{2n+1} + w_j y_{2n+1}) \\ &\quad + \sum_{j=d+1}^n \frac{1}{2} (b_{2j-1} + i b_{2j}) (-\bar{z}_j x_{2n+1} + \bar{w}_j y_{2n+1}) + (x_{2n+1} - y_{2n+1}) f(x, y) + O(\|(x, y)\|^3), \end{aligned} \tag{3.23}$$

where $z_j = x_j + ix_{d+j}, j = 1, \dots, d, z_j = x_{2j-1} + ix_{2j}, j = 2d + 1, \dots, 2n, \mu_j, j = 1, \dots, n$, and $b_{d+1} \in \mathbb{R}, \dots, b_{2n} \in \mathbb{R}$ are as in (3.11) and f is smooth and satisfies $f(0, 0) = 0, f(x, y) = \bar{f}(y, x)$.

Proof Let $\tilde{x} = (\tilde{x}_1, \dots, \tilde{x}_{2n+1})$ be the coordinates as in Theorem 3.5. It is easy to see that

$$\begin{aligned} \tilde{x}_j &= x_j + a_j x_{2n+1} + h_j(x), \quad h_j(x) = O(|x|^2), \quad a_j \in \mathbb{R}, \quad j = 1, 2, \dots, 2n, \\ \tilde{x}_{2n+1} &= x_{2n+1} + h_{2n+1}(x), \quad h_{2n+1}(x) = O(|x|^2). \end{aligned} \tag{3.24}$$

From (3.4), it is straightforward to see that

$$\begin{aligned} \omega_0(\tilde{x}) &= (1 + O(|\tilde{x}|))d\tilde{x}_{2n+1} + \sum_{j=1}^d 2\mu_j \tilde{x}_{d+j} d\tilde{x}_j + \sum_{j=1}^d (-2\mu_j \tilde{x}_j) d\tilde{x}_{d+j} \\ &\quad + \sum_{j=d+1}^n 2\mu_j x_{2j} dx_{2j-1} + \sum_{j=d+1}^n (-2\mu_j \tilde{x}_{2j-1}) d\tilde{x}_{2j} + \sum_{j=1}^{2n} \hat{b}_j \tilde{x}_{2n+1} d\tilde{x}_j + O(|x|^2), \end{aligned} \tag{3.25}$$

where

$$\begin{aligned} \hat{b}_j &= c_j + \bar{c}_j, \quad j \in \{1, \dots, d\} \cup \{2d + 1, 2d + 3, \dots, 2n - 1\}, \\ \hat{b}_j &= ic_j - i\bar{c}_j, \quad j \in \{d + 1, \dots, 2d\} \cup \{2d + 2, \dots, 2n\}. \end{aligned}$$

From (3.25) and (3.11), it is not difficult to see that (see also (3.20))

$$\begin{aligned} \frac{\partial^2 \tilde{x}_{2n+1}}{\partial x_j \partial x_k}(p) &= 0, \quad j \in \{1, \dots, d\}, k \in \{1, \dots, d\}, \\ \frac{\partial^2 \tilde{x}_{2n+1}}{\partial x_j \partial x_k}(p) &= 0, \quad j \in \{1, \dots, 2n\}, k \in \{2d + 1, \dots, 2n\}, \\ \frac{\partial^2 \tilde{x}_{2n+1}}{\partial x_{d+j} \partial x_k}(p) &= 2\mu_j \delta_{j,k}, \quad j, k \in \{1, \dots, d\}. \end{aligned} \tag{3.26}$$

From (3.24), (3.26) and (3.6), it is straightforward to check that

$$\begin{aligned} \varphi_-(x, y) &= -x_{2n+1} + y_{2n+1} - 2 \sum_{j=1}^d \mu_j x_j x_{d+j} + 2 \sum_{j=1}^d \mu_j y_j y_{d+j} + i \sum_{j=1}^n |\mu_j| |z_j - w_j|^2 \\ &\quad + \sum_{j=1}^n i \mu_j (\bar{z}_j w_j - z_j \bar{w}_j) + \sum_{j=1}^n \beta_j (-z_j x_{2n+1} + w_j y_{2n+1}) \\ &\quad + \sum_{j=1}^n \bar{\beta}_j (-\bar{z}_j x_{2n+1} + \bar{w}_j y_{2n+1}) + (x_{2n+1} - y_{2n+1}) f(x, y) + O(|(x, y)|^3), \end{aligned} \tag{3.27}$$

where $\beta_j \in \mathbb{C}$, $j = 1, \dots, n$ and f is smooth and satisfies $f(0, 0) = 0$, $f(x, y) = \bar{f}(y, x)$. We now determine β_j , $j = 1, \dots, n$. We can compute that

$$\begin{aligned} \frac{\partial \varphi_-}{\partial x_j}(x, x) &= -4\mu_j x_{d+j} - (\beta_j + \bar{\beta}_j)x_{2n+1} + O(|x|^2), \quad j = 1, \dots, d, \\ \frac{\partial \varphi_-}{\partial x_{d+j}}(x, x) &= -i(\beta_j - \bar{\beta}_j)x_{2n+1} + O(|x|^2), \quad j = 1, \dots, d, \\ \frac{\partial \varphi_-}{\partial x_{2j-1}}(x, x) &= -2\mu_j x_{2j} - (\beta_j + \bar{\beta}_j)x_{2n+1} + O(|x|^2), \quad j = d + 1, \dots, n, \\ \frac{\partial \varphi_-}{\partial x_{2j}}(x, x) &= 2\mu_j x_{2j-1} + (-i\beta_j + i\bar{\beta}_j)x_{2n+1} + O(|x|^2), \quad j = d + 1, \dots, n. \end{aligned} \tag{3.28}$$

Note that $d_x \varphi_-(x, x) = -\omega_0(x)$. From this observation and (3.11), we deduce that

$$\begin{aligned} \frac{\partial \varphi_-}{\partial x_j}(x, x) &= -4\mu_j x_{d+j} + O(|x|^2), \quad j = 1, \dots, d, \\ \frac{\partial \varphi_-}{\partial x_{d+j}}(x, x) &= -b_{d+j} x_{2n+1} + O(|x|^2), \quad j = 1, \dots, d, \\ \frac{\partial \varphi_-}{\partial x_{2j-1}}(x, x) &= -2\mu_j x_{2j} - b_{2j-1} x_{2n+1} + O(|x|^2), \quad j = d + 1, \dots, n, \\ \frac{\partial \varphi_-}{\partial x_{2j}}(x, x) &= 2\mu_j x_{2j-1} - b_{2j} x_{2n+1} + O(|x|^2), \quad j = d + 1, \dots, n. \end{aligned} \tag{3.29}$$

From (3.28) and (3.29), we deduce that

$$\beta_j = -\frac{i}{2} b_{d+j}, \quad j = 1, \dots, d, \quad \text{and} \quad \beta_j = \frac{1}{2} (b_{2j-1} - i b_{2j}), \quad j = d + 1, \dots, n. \tag{3.30}$$

From (3.30) and (3.27), we get (3.23). □

We now work with local coordinates as in Theorem 3.7. From (3.23), we see that near $(p, p) \in U \times U$, we have $\frac{\partial \varphi_-}{\partial y_{2n+1}} \neq 0$. Using the Malgrange preparation theorem [12, Th.7.5.7], we have

$$\varphi_-(x, y) = g(x, y)(y_{2n+1} + \hat{\varphi}_-(x, \hat{y})) \tag{3.31}$$

in some neighborhood of (p, p) , where $\hat{y} = (y_1, \dots, y_{2n})$, $g, \hat{\varphi}_- \in C^\infty$. Since $\text{Im } \varphi_- \geq 0$, it is not difficult to see that $\text{Im } \hat{\varphi}_- \geq 0$ in some neighborhood of (p, p) . We may take U small enough so that (3.31) holds and $\text{Im } \hat{\varphi}_- \geq 0$ on $U \times U$. From [25, Th.4.2], we see that since $\varphi_-(x, y)$ and $\hat{\varphi}_-(x, y)$ satisfy (3.31), $\varphi_-(x, y)t$ and $(y_{2n+1} + \hat{\varphi}_-(x, \hat{y}))t$ are equivalent in the sense of Melin–Sjöstrand. More precisely, for any $k \in \mathbb{R}$ and any $b_1(x, y, t) \in S_{\text{cl}}^k(U \times U \times$

$\mathbb{R}_+, T^{*0,q} X \boxtimes (T^{*0,q} X)^*$, we can find $b_2(x, y, t) \in S_{cl}^k(U \times U \times \mathbb{R}_+, T^{*0,q} X \boxtimes (T^{*0,q} X)^*)$ such that

$$\int_0^\infty e^{i\varphi_-(x,y)t} b_1(x, y, t) dt \equiv \int_0^\infty e^{i\hat{\varphi}_-(x,y)t} b_2(x, y, t) dt \text{ on } U$$

and vice versa. We can replace the phase φ_- by $y_{2n+1} + \hat{\varphi}_-(x, \hat{y})$. From now on, we assume that $\varphi_-(x, y)$ has the form

$$\varphi_-(x, y) = y_{2n+1} + \hat{\varphi}_-(x, \hat{y}). \tag{3.32}$$

It is easy to check that $\varphi_-(x, y)$ satisfies (3.5) and (3.23) with $f(x, y) = 0$.

We now study $S_G^{(q)}(x, y)$. From Theorem 3.2, we get

Theorem 3.9 *Assume that $q \notin \{n_-, n_+\}$. Then, $S_G^{(q)} \equiv 0$ on X .*

Assume that $q = n_-$ and $\square_b^{(q)}$ has L^2 closed range. Fix $p \in \mu^{-1}(0)$ and let $v = (v_1, \dots, v_d)$ and $x = (x_1, \dots, x_{2n+1})$ be the local coordinates of G and X as in Theorem 3.7. Assume that $d\mu = m(v)dv = m(v_1, \dots, v_d)dv_1 \cdots dv_d$ on V , where V is an open neighborhood of $e_0 \in G$ as in Theorem 3.7. From (3.7), we have

$$S_G^{(q)}(x, y) = \int_G \chi(g)S^{(q)}(x, g \circ y)d\mu(g) + \int_G (1 - \chi(g))S^{(q)}(x, g \circ y)d\mu(g),$$

where $\chi \in C_0^\infty(V)$, $\chi = 1$ near e_0 . Since G is freely on $\mu^{-1}(0)$, if U and V are small, there is a constant $c > 0$ such that

$$d(x, g \circ y) \geq c, \quad \forall x, y \in U, g \in \text{Supp}(1 - \chi), \tag{3.33}$$

where U is an open set of $p \in \mu^{-1}(0)$ as in Theorem 3.7. From now on, we take U and V small enough so that (3.33) holds. In view of Theorem 3.3, we see that $S^{(q)}(x, y)$ is smoothing away from diagonal. From this observation and (3.33), we conclude that $\int_G (1 - \chi(g))S^{(q)}(x, g \circ y)d\mu(g) \equiv 0$ on U and hence

$$S_G^{(q)}(x, y) \equiv \int_G \chi(g)S^{(q)}(x, g \circ y)d\mu(g) \text{ on } U. \tag{3.34}$$

From Theorem 3.3 and (3.34), we have

$$\begin{aligned} S_G^{(q)}(x, y) &\equiv \hat{S}_-(x, y) + \hat{S}_+(x, y) \text{ on } U, \\ \hat{S}_\mp(x, y) &= \int_G \chi(g)S_\mp(x, g \circ y)d\mu(g), \end{aligned} \tag{3.35}$$

Write $x = (x', x'') = (x', \hat{x}'', \tilde{x}'')$, $y = (y', y'') = (y', \hat{y}'', \tilde{y}'')$, where $\hat{x}'' = (x_{d+1}, \dots, x_{2d})$, $\hat{y}'' = (y_{d+1}, \dots, y_{2d})$, $\tilde{x}'' = (x_{2d+1}, \dots, x_{2n+1})$, $\tilde{y}'' = (y_{2d+1}, \dots, y_{2n+1})$. Since $S_G^{(q)}(x, y)$ is G -invariant, we have $S_G^{(q)}(x, y) = S_G^{(q)}((0, x''), (\gamma(y''), y''))$, where $\gamma \in C^\infty(U_2, U_1)$ is as in Theorem 3.7. From this observation and (3.35), we have

$$S_G^{(q)}(x, y) \equiv \hat{S}_-((0, x''), (\gamma(y''), y'')) + \hat{S}_+((0, x''), (\gamma(y''), y'')) \text{ on } U. \tag{3.36}$$

Write $\hat{x}'' = (x_{d+1}, \dots, x_{2n})$, $\hat{y}'' = (y_{d+1}, \dots, y_{2n})$. From (3.32), (3.36), Theorem 3.7 and Theorem 3.3, we have

$$\begin{aligned} &\hat{S}_-((0, x''), (\gamma(y''), y'')) \\ &\equiv \int e^{i(y_{2n+1} + \hat{\varphi}_-((0, x''), (v + \gamma(y''), \hat{y}''))))t} s_-((0, x''), (v + \gamma(y''), y''), t) m(v) dv dt. \end{aligned} \tag{3.37}$$

From (3.23), it is straightforward to see that

$$\det \left(\left(\frac{\partial^2 \hat{\varphi}_-}{\partial v_k \partial v_j}(p, p) \right)_{j,k=1}^d \right) = (2i)^d |\mu_1| \cdots |\mu_d| \neq 0. \tag{3.38}$$

We pause and introduce some notations. Let W be an open set of \mathbb{R}^N , $N \in \mathbb{N}$. From now on, we write $W^{\mathbb{C}}$ to denote an open set in \mathbb{C}^N with $W^{\mathbb{C}} \cap \mathbb{R}^N = W$ and for $f \in C^\infty(W)$, from now on, we write $\tilde{f} \in C^\infty(W^{\mathbb{C}})$ to denote an almost analytic extension of f (see Section 2 in [25]). Let $h(x'', y'') \in C^\infty(U \times U, \mathbb{C}^d)$ be the solution of the system

$$\frac{\partial \tilde{\varphi}_-}{\partial \tilde{y}_j}((0, x''), (h(x'', y'') + \gamma(y''), \hat{y}'')) = 0, \quad j = 1, 2, \dots, d, \tag{3.39}$$

and let

$$\Phi_-(x'', y'') := y_{2n+1} + \tilde{\varphi}_-((0, x''), (h(x'', y'') + \gamma(y''), \hat{y}'')). \tag{3.40}$$

It is known that (see page 147 in [25]) $\text{Im } \Phi_-(x'', y'') \geq 0$. Note that

$$\frac{\partial \hat{\varphi}_-}{\partial v_j} |_{\hat{x}''=\hat{y}''=0, \tilde{x}''=\tilde{y}''=0, x'=v+\gamma(y'')=0} = -\langle \omega_0(x), \frac{\partial}{\partial x_j} \rangle = 0,$$

where $x = (0, (0, \tilde{x}''))$. We deduce that for $\hat{x}'' = \hat{y}'' = 0, \tilde{x}'' = \tilde{y}'' = 0, v = -\gamma(y'')$ are real critical points. From this observation, we can calculate that

$$d_x \Phi_- |_{x''=y'', \hat{x}''=0} = -f(x'')\omega_0(x), \quad d_y \Phi_- |_{x''=y'', \hat{x}''=0} = f(x'')\omega_0(x), \tag{3.41}$$

where $x = (0, \tilde{x}'')$ and $f \in C^\infty$ is a positive function with $f(p) = 1$. By using stationary phase formula of Melin–Sjöstrand [25], we can carry out the v integral in (3.37) and get

$$\hat{S}_-((0, x''), (\gamma(y''), y'')) \equiv \int e^{i\Phi_-(x'', y'')t} a_-(x'', y'', t) dt \quad \text{on } U,$$

where $a_-(x'', y'', t) \sim \sum_{j=0}^\infty t^{n-\frac{d}{2}-j} a_-^j(x'', y'')$ in $S_{1,0}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$,

$$\begin{aligned} a_-^j(x'', y'') &\in C^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \quad j = 0, 1, \dots, \\ a_-^0(p, p) &= \frac{1}{2}m(0)\pi^{-n-1+\frac{d}{2}}|\mu_1|^{\frac{1}{2}} \cdots |\mu_d|^{\frac{1}{2}}|\mu_{d+1}| \cdots |\mu_n|\tau_{p,n}. \end{aligned} \tag{3.42}$$

We now study the property of the phase $\Phi_-(x'', y'')$. We need the following which is known (see Section 2 in [25])

Theorem 3.10 *There exist a constant $c > 0$ and an open set $\Omega \in \mathbb{R}^d$ such that*

$$\text{Im } \Phi_-(x'', y'') \geq c \inf_{v \in \Omega} \{ |\text{Im } \hat{\varphi}_-((0, x''), (v + \gamma(y''), \hat{y}'')) + |d_v \hat{\varphi}_-((0, x''), (v + \gamma(y''), \hat{y}''))|^2 \}, \tag{3.43}$$

for all $((0, x''), (0, y'')) \in U \times U$.

We can now prove

Theorem 3.11 *If U is small enough, then there is a constant $c > 0$ such that*

$$\text{Im } \Phi_-(x'', y'') \geq c \left(|\hat{x}''|^2 + |\hat{y}''|^2 + |\hat{x}'' - \hat{y}''|^2 \right), \quad \forall ((0, x''), (0, y'')) \in U \times U. \tag{3.44}$$

Proof From (3.5), we see that there is a constant $c_1 > 0$ such that

$$\text{Im } \hat{\varphi}_-((0, x''), (v + \gamma(y''), \hat{y}'')) \geq c_1(|v + \gamma(y'')|^2 + |\hat{x}'' - \hat{y}''|^2), \quad \forall v \in \Omega, \quad (3.45)$$

where Ω is any open set of $0 \in \mathbb{R}^d$. From (3.45) and (3.43), we conclude that there is a constant $c_2 > 0$ such that

$$\text{Im } \Phi_-(x'', y'') \geq c_2(|\hat{x}'' - \hat{y}''|^2 + |d_{y'} \hat{\varphi}_-((0, x''), (0, \hat{x}''))|^2). \quad (3.46)$$

From (3.23), we see that the matrix

$$\left(\frac{\partial^2 \hat{\varphi}_-}{\partial x_j \partial x_k}(p, p) + \frac{\partial^2 \hat{\varphi}_-}{\partial y_j \partial y_k}(p, p) \right)_{1 \leq k \leq d, d+1 \leq j \leq 2d}$$

is non-singular. From this observation and notice that $d_{y'} \hat{\varphi}_-((0, x''), (0, \hat{x}''))|_{\hat{x}''} = 0$, we deduce that if U is small enough then there is a constant $c_3 > 0$ such that

$$|d_{y'} \hat{\varphi}_-((0, x''), (0, x''))| \geq c_3 |\hat{x}''|. \quad (3.47)$$

From (3.47) and (3.46), the theorem follows. □

From now on, we assume that U is small enough so that (3.44) holds.

We now determine the Hessian of $\Phi_-(x'', y'')$ at (p, p) . Let $\hat{h}(x'', y'') := h(x'', y'') + \gamma(y'')$. From (3.39), we have

$$\frac{\partial^2 \hat{\varphi}_-}{\partial x_{d+1} \partial y_1}(p, p) + \sum_{j=1}^d \frac{\partial^2 \hat{\varphi}_-}{\partial y_1 \partial y_j}(p, p) \frac{\partial \hat{h}_j}{\partial x_{d+1}}(p, p) = 0. \quad (3.48)$$

From (3.23), we can calculate that

$$\frac{\partial^2 \hat{\varphi}_-}{\partial x_{d+1} \partial y_1}(p, p) = 2\mu_1, \quad \frac{\partial^2 \hat{\varphi}_-}{\partial y_1 \partial y_j}(p, p) = 2i|\mu_1| \delta_{1,j}, \quad j = 1, 2, \dots, d. \quad (3.49)$$

From (3.49) and (3.48), we obtain $\frac{\partial \hat{h}_1}{\partial x_{d+1}}(p, p) = i \frac{\mu_1}{|\mu_1|}$. We can repeat the procedure above several times and deduce that

$$\frac{\partial \hat{h}_j}{\partial x_{d+k}}(p, p) = \frac{\partial \hat{h}_j}{\partial y_{d+k}}(p, p) = i \frac{\mu_j}{|\mu_j|} \delta_{j,k}, \quad j, k = 1, 2, \dots, d. \quad (3.50)$$

From (3.50), (3.23), (3.40) and by some straightforward computation (we omit the details), we get

Theorem 3.12 *With the notations used above, we have*

$$\begin{aligned} \Phi_-(x'', y'') = & -x_{2n+1} + y_{2n+1} + 2i \sum_{j=1}^d |\mu_j| y_{d+j}^2 + 2i \sum_{j=1}^d |\mu_j| x_{d+j}^2 \\ & + i \sum_{j=d+1}^n |\mu_j| |z_j - w_j|^2 + \sum_{j=d+1}^n i \mu_j (\bar{z}_j w_j - z_j \bar{w}_j) \\ & + \sum_{j=1}^d (-b_{d+j} x_{d+j} x_{2n+1} + b_{d+j} y_{d+j} y_{2n+1}) \\ & + \sum_{j=d+1}^n \frac{1}{2} (b_{2j-1} - i b_{2j}) (-z_j x_{2n+1} + w_j y_{2n+1}) \\ & + \sum_{j=d+1}^n \frac{1}{2} (b_{2j-1} + i b_{2j}) (-\bar{z}_j x_{2n+1} + \bar{w}_j y_{2n+1}) \\ & + (x_{2n+1} - y_{2n+1}) f(x, y) + O(|(x, y)|^3), \end{aligned} \quad (3.51)$$

where $z_j = x_{2j-1} + i x_{2j}$, $j = 2d+1, \dots, 2n$, μ_j , $j = 1, \dots, n$, and $b_{d+1} \in \mathbb{R}, \dots, b_{2n} \in \mathbb{R}$ are as in (3.11) and f is smooth and satisfies $f(0, 0) = 0$, $f(x, y) = \bar{f}(y, x)$.

We can change $\Phi_-(x'', y'')$ be $\Phi_-(x'', y'') \frac{1}{f(x'')}$, where $f(x'')$ is as in (3.41). Thus,

$$d_x \Phi_-|_{x''=y'', \hat{x}''=0} = -\omega_0(x), \quad d_y \Phi_-|_{x''=y'', \hat{x}''=0} = \omega_0(x), \tag{3.52}$$

where $x = (0, \tilde{x}'')$. It is clear that $\Phi_-(x'', y'')$ still satisfies (3.44) and (3.51).

We now determine the leading term $a_-^0(p, p)$. In view of (3.42), we only need to calculate $m(0)$. Put $Y_p = \{g \circ p; g \in G\}$. Y_p is a d -dimensional submanifold of X . The G -invariant Hermitian metric $\langle \cdot | \cdot \rangle$ induces a volume form dv_{Y_p} on Y_p . Put

$$V_{\text{eff}}(p) := \int_{Y_p} dv_{Y_p}.$$

For $f(g) \in C^\infty(G)$, let $\hat{f}(g \circ p) := f(g), \forall g \in G$. Then, $\hat{f} \in C^\infty(Y_p)$. Let $d\hat{\mu}$ be the measure on G given by $\int_G f d\hat{\mu} := \int_{Y_p} \hat{f} dv_{Y_p}$, for all $f \in C^\infty(G)$. It is not difficult to see that $d\hat{\mu}$ is a Haar measure and

$$\int_G d\hat{\mu} = V_{\text{eff}}(p). \tag{3.53}$$

Recall that we work with the local coordinates in Theorem 3.7. In view of (3.10), we see that $\left\{ \frac{1}{\sqrt{2}} \frac{\partial}{\partial x_1}, \dots, \frac{1}{\sqrt{2}} \frac{\partial}{\partial x_d} \right\}$ is an orthonormal basis for \mathfrak{g}_p . Hence $m(0) = 2^{\frac{d}{2}} \frac{1}{V_{\text{eff}}(p)}$. From this observation, (3.53) and (3.42), we get

$$a_-^0(p, p) = 2^{\frac{d}{2}-1} \frac{1}{V_{\text{eff}}(p)} \pi^{-n-1+\frac{d}{2}} |\mu_1|^{\frac{1}{2}} \cdots |\mu_d|^{\frac{1}{2}} |\mu_{d+1}| \cdots |\mu_n| \tau_{p,n_-}. \tag{3.54}$$

Similarly, we can repeat the procedure above and deduce that

$$\hat{S}_+((0, x''), (\gamma(y''), y'')) \equiv \int e^{i\Phi_+(x'', y'')t} a_+(x'', y'', t) dt \quad \text{on } U,$$

where $a_+(x'', y'', t) \sim \sum_{j=0}^\infty t^{n-\frac{d}{2}-j} a_+^j(x'', y'')$ in $S_{1,0}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q} X \boxtimes (T^{*0,q} X)^*)$,

$$a_+^j(x'', y'') \in C^\infty(U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \quad j = 0, 1, \dots,$$

$$a_+^0(p, p) = 2^{\frac{d}{2}-1} \frac{1}{V_{\text{eff}}(p)} \pi^{-n-1+\frac{d}{2}} |\mu_1|^{\frac{1}{2}} \cdots |\mu_d|^{\frac{1}{2}} |\mu_{d+1}| \cdots |\mu_n| \tau_{p,n_+}, \tag{3.55}$$

and $\Phi_+(x'', y'') \in C^\infty(U \times U)$, $\text{Im } \Phi_+(x'', y'') \geq 0$, $-\overline{\Phi}_+(x'', y'')$ satisfies (3.44), (3.51) and (3.52).

Summing up, we get one of the main result of this work

Theorem 3.13 *We recall that we work with the assumption that the Levi form is non-degenerate of constant signature (n_-, n_+) on X . Let $q = n_-$ or n_+ . Suppose that $\square_b^{(q)}$ has L^2 closed range. Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates defined in an open set U of p such that $x(p) = 0$ and (3.8), (3.9), (3.10), (3.11) hold. Write $x'' = (x_{d+1}, \dots, x_{2n+1})$. Then, there exist continuous operators $S_-^G, S_+^G : \Omega_0^{0,q}(U) \rightarrow \Omega^{0,q}(U)$ such that*

$$S_G^{(q)} \equiv S_-^G + S_+^G \quad \text{on } U,$$

and $S_-^G(x, y), S_+^G(x, y)$ satisfy

$$S_\mp^G(x, y) \equiv \int_0^\infty e^{i\Phi_\mp(x'', y'')t} a_\mp(x'', y'', t) dt \quad \text{on } U,$$

with

$$a_-(x, y, t), a_+(x, y, t) \in S_{\text{cl}}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*),$$

$$a_-(x, x) \neq 0, \forall x \in U, \quad a_+^0(x, x) \neq 0, \forall x \in U,$$

$a_-^0(p, p)$ and $a_+^0(p, p)$ are given by (3.54) and (3.55) respectively, $\Phi_-(x'', y'') \in C^\infty(U \times U)$ satisfies (3.52), (3.44) and (3.51), $\Phi_+(x'', y'') \in C^\infty(U \times U)$, $-\Phi_+(x'', y'')$ satisfies (3.52), (3.44) and (3.51).

3.4 G-invariant Szegő kernel asymptotics away $\mu^{-1}(0)$

The goal of this section is to prove the following

Theorem 3.14 *Let D be an open set of X with $D \cap \mu^{-1}(0) = \emptyset$. Then, $S_G^{(q)} \equiv 0$ on D .*

We first need

Lemma 3.15 *Let $p \notin \mu^{-1}(0)$. Then, there are open sets U of p and V of $e \in G$ such that for any $\chi \in C_0^\infty(V)$, we have*

$$\int_G S^{(q)}(x, g \circ y) \chi(g) d\mu(g) \equiv 0 \text{ on } U. \tag{3.56}$$

Proof If $q \notin \{n_-, n_+\}$. By Theorem 3.2, we get (3.56). We may assume that $q = n_-$. Take local coordinates $v = (v_1, \dots, v_d)$ of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$, local coordinates $x = (x_1, \dots, x_{2n+1})$ of X defined in a neighborhood $U = U_1 \times U_2$ of p with $0 \leftrightarrow p$, where $U_1 \subset \mathbb{R}^d$ is an open set of $0 \in \mathbb{R}^d$, $U_2 \subset \mathbb{R}^{2n+1-d}$ is an open set of $0 \in \mathbb{R}^{2n+1-d}$, such that

$$(v_1, \dots, v_d) \circ (\gamma(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1})$$

$$= (v_1 + \gamma_1(x_{d+1}, \dots, x_{2n+1}), \dots, v_d + \gamma_d(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1}),$$

$$\forall (v_1, \dots, v_d) \in V, \quad \forall (x_{d+1}, \dots, x_{2n+1}) \in U_2,$$

and

$$\underline{\mathfrak{g}} = \text{span} \left\{ \frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_d} \right\},$$

where $\gamma = (\gamma_1, \dots, \gamma_d) \in C^\infty(U_2, U_1)$ with $\gamma(0) = 0 \in \mathbb{R}^d$. Note that we don't use the local coordinates in Theorem 3.7. It should be notice that G needs not act locally freely on near p , (3.33) need not be true. We can not use off-diagonal expansion for the Szegő kernel to get this lemma. We will use some properties of the phase φ_- and integrations by parts to obtain this lemma. From Theorem 3.3, we have

$$\int_G S^{(q)}(x, g \circ y) \chi(g) d\mu(g) \equiv \int_G S_-(x, g \circ y) \chi(g) d\mu(g) + \int_G S_+(x, g \circ y) \chi(g) d\mu(g) \text{ on } U. \tag{3.57}$$

From Theorem 3.3, we have

$$\int_G S_-(x, g \circ y) \chi(g) d\mu(g)$$

$$\equiv \int e^{i(\varphi_-(x, (v+\gamma(y''), y''))t)} S_-(x, (v + \gamma(y''), y''), t) \chi(v) m(v) dv dt,$$

where $y'' = (y_{d+1}, \dots, y_{2n+1})$, $m(v)dv = d\mu|_V$. Since $p \notin \mu^{-1}(0)$ and notice that $d_y \varphi_-(x, x) = \omega_0(x, x)$, we deduce that if V and U are small then $d_v(\varphi_-(x, (v + \gamma(y''), y''))) \neq 0$, for every $v \in V$, $(x, y) \in U \times U$. Hence, by using integration by parts with respect to v , we get

$$\int_G S_-(x, g \circ y)\chi(g)d\mu(g) \equiv 0 \text{ on } U. \tag{3.58}$$

Similarly, we have

$$\int_G S_+(x, g \circ y)\chi(g)d\mu(g) \equiv 0 \text{ on } U. \tag{3.59}$$

From (3.57), (3.58) and (3.59), the lemma follows. □

Lemma 3.16 *Let $p \notin \mu^{-1}(0)$ and let $h \in G$. We can find open sets U of p and V of h such that for every $\chi \in C_0^\infty(V)$, we have $\int_G S^{(q)}(x, g \circ y)\chi(g)d\mu(g) \equiv 0$ on U .*

Proof Let U and V be open sets as in Lemma 3.15. Let $\hat{V} = hV$. Then, \hat{V} is an open set of G . Let $\hat{\chi} \in C_0^\infty(\hat{V})$. We have

$$\begin{aligned} \int_G S^{(q)}(x, g \circ y)\hat{\chi}(g)d\mu(g) &= \int_G S^{(q)}(x, h \circ g \circ y)\hat{\chi}(h \circ g)d\mu(g) \\ &= \int_G S^{(q)}(x, g \circ y)\chi(g)d\mu(g), \end{aligned} \tag{3.60}$$

where $\chi(g) := \hat{\chi}(h \circ g) \in C_0^\infty(V)$. From (3.60) and Lemma 3.15, we deduce that

$$\int_G S^{(q)}(x, g \circ y)\hat{\chi}(g)d\mu(g) \equiv 0 \text{ on } U.$$

The lemma follows. □

Proof of Theorem 3.14 Fix $p \in D$. We need to show that $S_G^{(q)}$ is smoothing near p . Let $h \in G$. By Lemma 3.16, we can find open sets U_h of p and V_h of h such that for every $\chi \in C_0^\infty(V_h)$, we have

$$\int_G S^{(q)}(x, g \circ y)\chi(g)d\mu(g) \equiv 0 \text{ on } U_h. \tag{3.61}$$

Since G is compact, we can find open sets U_{h_j} and V_{h_j} , $j = 1, \dots, N$, such that $G = \bigcup_{j=1}^N V_{h_j}$. Let $U = D \cap \left(\bigcap_{j=1}^N U_{h_j}\right)$ and let $\chi_j \in C_0^\infty(V_{h_j})$, $j = 1, \dots, N$, with $\sum_{j=1}^N \chi_j = 1$ on G . From (3.61), we have

$$S_G^{(q)}(x, y) = \int_G S^{(q)}(x, g \circ y)d\mu(g) = \sum_{j=1}^N \int_G S^{(q)}(x, g \circ y)\chi_j(g)d\mu(g) \equiv 0 \text{ on } U.$$

The theorem follows. □

From Theorems 3.9, 3.13 and 3.14, we get Theorem 1.5.

4 G-invariant Szegő kernel asymptotics on CR manifolds wit S^1 action

Let $(X, T^{1,0}X)$ be a compact CR manifold of dimension $2n + 1, n \geq 1$. We assume that X admits an S^1 action: $S^1 \times X \rightarrow X$. We write $e^{i\theta}$ to denote the S^1 action. Let $T \in C^\infty(X, TX)$ be the global real vector field induced by the S^1 action given by $(Tu)(x) = \frac{\partial}{\partial \theta} (u(e^{i\theta} \circ x))|_{\theta=0}, u \in C^\infty(X)$. We recall

Definition 4.1 We say that the S^1 action $e^{i\theta}$ is CR if $[T, C^\infty(X, T^{1,0}X)] \subset C^\infty(X, T^{1,0}X)$ and the S^1 action is transversal if for each $x \in X, \mathbb{C}T(x) \oplus T_x^{1,0}X \oplus T_x^{0,1}X = \mathbb{C}T_xX$. Moreover, we say that the S^1 action is locally free if $T \neq 0$ everywhere. It should be mentioned that transversality implies locally free.

We assume now that $(X, T^{1,0}X)$ is a compact connected CR manifold with a transversal CR locally free S^1 action $e^{i\theta}$ and we let T be the global vector field induced by the S^1 action. Let $\omega_0 \in C^\infty(X, T^*X)$ be the global real one form determined by $\langle \omega_0, u \rangle = 0$, for every $u \in T^{1,0}X \oplus T^{0,1}X$, and $\langle \omega_0, T \rangle = -1$. Note that ω_0 and T satisfy (2.2). Assume that X admits a compact connected Lie group G action and the Lie group G acts on X preserving ω_0 and J . We recall that we work with Assumption 1.7.

We now assume that the Hermitian metric $\langle \cdot | \cdot \rangle$ on $\mathbb{C}TX$ is $G \times S^1$ invariant. Then the L^2 inner product $(\cdot | \cdot)$ on $\Omega^{0,q}(X)$ induced by $\langle \cdot | \cdot \rangle$ is $G \times S^1$ -invariant. We then have

$$Tg^* \bar{\partial}_b^* = g^* T \bar{\partial}_b^* = \bar{\partial}_b^* g^* T = \bar{\partial}_b^* T g^* \text{ on } \Omega^{0,q}(X), \forall g \in G,$$

$$Tg^* \square_b^{(q)} = g^* T \square_b^{(q)} = \square_b^{(q)} g^* T = \square_b^{(q)} T g^* \text{ on } \Omega^{0,q}(X), \forall g \in G.$$

Let $L^2_{(0,q),m}(X)^G$ be the completion of $\Omega^{0,q}(X)^G$ with respect to $(\cdot | \cdot)$. We write $L^2_m(X)^G := L^2_{(0,0),m}(X)^G$. Put $(\text{Ker } \square_b^{(q)})^G_m := (\text{Ker } \square_b^{(q)})^G \cap L^2_{(0,q),m}(X)^G$. It is not difficult to see that, for every $m \in \mathbb{Z}, (\text{Ker } \square_b^{(q)})^G_m \subset \Omega^{0,q}(X)^G$ and $\dim (\text{Ker } \square_b^{(q)})^G_m < \infty$. The m th G -invariant Szegő projection is the orthogonal projection $S_{G,m}^{(q)} : L^2_{(0,q)}(X) \rightarrow (\text{Ker } \square_b^{(q)})^G_m$ with respect to $(\cdot | \cdot)$. Let $S_{G,m}^{(q)}(x, y) \in C^\infty(X \times X, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ be the distribution kernel of $S_{G,m}^{(q)}$. We can check that

$$S_{G,m}^{(q)}(x, y) = \frac{1}{2\pi} \int_{-\pi}^\pi S_G^{(q)}(x, e^{i\theta} \circ y) e^{im\theta} d\theta. \tag{4.1}$$

The goal of this section is to study the asymptotics of $S_{G,m}^{(q)}$ as $m \rightarrow +\infty$.

From Theorem 3.14, (4.1) and by using integration by parts several times, we get

Theorem 4.2 *Let $D \subset X$ be an open set with $D \cap \mu^{-1}(0) = \emptyset$. Then, $S_{G,m}^{(q)} = O(m^{-\infty})$ on D .*

We now study $S_{G,m}^{(q)}$ near $\mu^{-1}(0)$. We can repeat the proof of Theorem 3.7 with minor change and get

Theorem 4.3 *Let $p \in \mu^{-1}(0)$. There exist local coordinates $v = (v_1, \dots, v_d)$ of G defined in a neighborhood V of e_0 with $v(e_0) = (0, \dots, 0)$, local coordinates $x = (x_1, \dots, x_{2n+1})$ of X defined in a neighborhood $U = U_1 \times (\hat{U}_2 \times]-2\delta, 2\delta[)$ of p with $0 \leftrightarrow p$, where $U_1 \subset \mathbb{R}^d$ is an open set of $0 \in \mathbb{R}^d, \hat{U}_2 \subset \mathbb{R}^{2n-d}$ is an open set of $0 \in \mathbb{R}^{2n-d}, \delta > 0$, and a smooth function $\gamma = (\gamma_1, \dots, \gamma_d) \in C^\infty(\hat{U}_2 \times]-2\delta, 2\delta[, U_1)$ with $\gamma(0) = 0 \in \mathbb{R}^d$ such that*

$$(v_1, \dots, v_d) \circ (\gamma(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1})$$

$$= (v_1 + \gamma_1(x_{d+1}, \dots, x_{2n+1}), \dots, v_d + \gamma_d(x_{d+1}, \dots, x_{2n+1}), x_{d+1}, \dots, x_{2n+1}),$$

$$\forall (v_1, \dots, v_d) \in V, \forall (x_{d+1}, \dots, x_{2n+1}) \in \hat{U}_2 \times]-2\delta, 2\delta[,$$

$$\begin{aligned}
 T &= -\frac{\partial}{\partial x_{2n+1}}, \quad \mathfrak{g} = \text{span} \left\{ \frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_d} \right\}, \\
 \mu^{-1}(0) \cap U &= \{x_{d+1} = \dots = x_{2d} = 0\}, \\
 \text{On } \mu^{-1}(0) \cap U, \text{ we have } J\left(\frac{\partial}{\partial x_j}\right) &= \frac{\partial}{\partial x_{d+j}} + a_j(x) \frac{\partial}{\partial x_{2n+1}}, \quad j = 1, 2, \dots, d,
 \end{aligned}
 \tag{4.2}$$

where $a_j(x)$ is a smooth function on $\mu^{-1}(0) \cap U$, independent of $x_1, \dots, x_{2d}, x_{2n+1}$ and $a_j(0) = 0, j = 1, \dots, d$,

$$\begin{aligned}
 T_p^{1,0} X &= \text{span} \{Z_1, \dots, Z_n\}, \\
 Z_j &= \frac{1}{2} \left(\frac{\partial}{\partial x_j} - i \frac{\partial}{\partial x_{d+j}} \right) (p), \quad j = 1, \dots, d, \\
 Z_j &= \frac{1}{2} \left(\frac{\partial}{\partial x_{2j-1}} - i \frac{\partial}{\partial x_{2j}} \right) (p), \quad j = d + 1, \dots, n, \\
 \langle Z_j | Z_k \rangle &= \delta_{j,k}, \quad j, k = 1, 2, \dots, n, \\
 \mathcal{L}_p(Z_j, \bar{Z}_k) &= \mu_j \delta_{j,k}, \quad j, k = 1, 2, \dots, n
 \end{aligned}$$

and

$$\begin{aligned}
 \omega_0(x) &= (1 + O(|x|)) dx_{2n+1} + \sum_{j=1}^d 4\mu_j x_{d+j} dx_j \\
 &\quad + \sum_{j=d+1}^n 2\mu_j x_{2j} dx_{2j-1} - \sum_{j=d+1}^n 2\mu_j x_{2j-1} dx_{2j} + O(|x|^2).
 \end{aligned}$$

Remark 4.4 Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Theorem 4.3. We can change x_{2n+1} be $x_{2n+1} - \sum_{j=1}^d a_j(x)x_{d+j}$, where $a_j(x), j = 1, \dots, d$, are as in (4.2). With this new local coordinates $x = (x_1, \dots, x_{2n+1})$, on $\mu^{-1}(0) \cap U$, we have $J\left(\frac{\partial}{\partial x_j}\right) = \frac{\partial}{\partial x_{d+j}}, j = 1, 2, \dots, d$. Moreover, it is clear that $\Phi_-(x, y) + \sum_{j=1}^d a_j(x)x_{d+j} - \sum_{j=1}^{d-1} a_j(y)y_{d+j}$ satisfies (1.19). Note that $a_j(x)$ is a smooth function on $\mu^{-1}(0) \cap U$, independent of $x_1, \dots, x_{2d}, x_{2n+1}$ and $a_j(0) = 0, j = 1, \dots, d$.

We now work with local coordinates as in Theorem 4.3. From (3.51), we see that near $(p, p) \in U \times U$, we have $\frac{\partial \Phi_-}{\partial y_{2n+1}} \neq 0$. Using the Malgrange preparation theorem [12, Th. 7.5.7], we have

$$\Phi_-(x, y) = g(x, y)(y_{2n+1} + \hat{\Phi}_-(x'', \hat{y}'')) \tag{4.3}$$

in some neighborhood of (p, p) , where $\hat{y}'' = (y_{d+1}, \dots, y_{2n}), g, \hat{\Phi}_- \in C^\infty$. Since $\text{Im } \Phi_- \geq 0$, it is not difficult to see that $\text{Im } \hat{\Phi}_- \geq 0$ in some neighborhood of (p, p) . We may take U small enough so that (4.3) holds and $\text{Im } \hat{\Phi}_- \geq 0$ on $U \times U$. From [25, Th. 4.2], we see that since $\Phi_-(x, y)$ and $y_{2n+1} + \hat{\Phi}_-(x'', \hat{y}'')$ satisfy (4.3), $\Phi_-(x, y)t$ and $(y_{2n+1} + \hat{\Phi}_-(x'', \hat{y}''))t$ are equivalent in the sense of Melin–Sjöstrand (see the discussion after (3.31), for the meaning of equivalent in the sense of Melin–Sjöstrand). We can replace the phase Φ_- by $y_{2n+1} + \hat{\Phi}_-(x, \hat{y}'')$. From now on, we assume that

$$\Phi_-(x, y) = y_{2n+1} + \hat{\Phi}_-(x'', \hat{y}''). \tag{4.4}$$

It is easy to check that $\Phi_-(x, y)$ satisfies (3.44) and (3.51) with $f(x, y) = 0$. Similarly, from now on, we assume that

$$\Phi_+(x, y) = -y_{2n+1} + \hat{\Phi}_+(x'', \hat{y}''). \tag{4.5}$$

We now study $S_{G,m}^{(q)}(x, y)$. From Theorem 3.9, we get

Theorem 4.5 Assume that $q \notin \{n_-, n_+\}$. Then, $S_{G,m}^{(q)} = O(m^{-\infty})$ on X .

Assume that $q = n_-$. It was proved in Theorem 1.12 in [15] that when X admits a transversal S^1 action, then $\square_b^{(q)}$ has L^2 closed range. Fix $p \in \mu^{-1}(0)$ and let $v = (v_1, \dots, v_d)$

and $x = (x_1, \dots, x_{2n+1})$ be the local coordinates of G and X as in Theorem 4.3 and let U and V be open sets as in Theorem 4.3. We take U small enough so that there is a constant $c > 0$ such that

$$d(e^{i\theta} \circ g \circ x, y) \geq c, \quad \forall(x, y) \in U \times U, \quad \forall g \in G, \theta \in [-\pi, -\delta] \cup [\delta, \pi], \tag{4.6}$$

where $\delta > 0$ is as in Theorem 4.3. We will study $S_{G,m}^{(q)}(x, y)$ in U . From (4.1), we have

$$\begin{aligned} S_{G,m}^{(q)}(x, y) &= \frac{1}{2\pi} \int_{-\pi}^{\pi} S_G^{(q)}(x, e^{i\theta} \circ y) e^{im\theta} d\theta \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-imx_{2n+1} + imy_{2n+1}} S_G^{(q)}(\hat{x}, e^{i\theta} \circ \hat{y}) e^{im\theta} d\theta \\ &= I + II, \\ I &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-imx_{2n+1} + imy_{2n+1}} \chi(\theta) S_G^{(q)}(\hat{x}, e^{i\theta} \circ \hat{y}) e^{im\theta} d\theta, \\ II &= \frac{1}{2\pi} \int_{-\pi}^{\pi} e^{-imx_{2n+1} + imy_{2n+1}} (1 - \chi(\theta)) S_G^{(q)}(\hat{x}, e^{i\theta} \circ \hat{y}) e^{im\theta} d\theta, \end{aligned}$$

where $\hat{x} = (x_1, \dots, x_{2n}, 0) \in U, \hat{y} = (y_1, \dots, y_{2n}, 0) \in U, \chi \in C_0^\infty(1 - 2\delta, 2\delta], \chi = 1$ on $[-\delta, \delta]$. We first study II . We have

$$II = \frac{1}{2\pi} \int_{-\pi}^{\pi} \int_G e^{-imx_{2n+1} + imy_{2n+1}} (1 - \chi(\theta)) S^{(q)}(\hat{x}, e^{i\theta} \circ g \circ \hat{y}) e^{im\theta} d\mu(g) d\theta. \tag{4.7}$$

From (4.7), (4.6) and notice that $S^{(q)}$ is smoothing away from diagonal, we deduce that

$$II = O(m^{-\infty}).$$

We now study I . From Theorem 3.13, (4.1), (4.4) and (4.5), we have

$$\begin{aligned} I &= I_0 + I_1, \\ I_0 &= \frac{1}{2\pi} \int_0^\infty \int_{-\pi}^{\pi} e^{-imx_{2n+1} + imy_{2n+1}} \chi(\theta) e^{i(-\theta + \hat{\Phi}_-(\hat{x}'', \hat{y}''))t + im\theta} a_-(\hat{x}'', (\hat{y}'', -\theta), t) dt d\theta, \\ I_1 &= \frac{1}{2\pi} \int_0^\infty \int_{-\pi}^{\pi} e^{-imx_{2n+1} + imy_{2n+1}} \chi(\theta) e^{i(\theta + \hat{\Phi}_+(\hat{x}'', \hat{y}''))t + im\theta} a_+(\hat{x}'', (\hat{y}'', -\theta), t) dt d\theta. \end{aligned}$$

We first study I_1 . From $\frac{\partial}{\partial \theta} (i(\theta + \hat{\Phi}_+(\hat{x}'', \hat{y}''))t + im\theta) \neq 0$, we can integrate by parts with respect to θ several times and deduce that

$$I_1 = O(m^{-\infty}).$$

We now study I_0 . We have

$$I_0 = \frac{1}{2\pi} \int_0^\infty \int_{-\pi}^{\pi} e^{-imx_{2n+1} + imy_{2n+1}} \chi(\theta) e^{im(-\theta t + \hat{\Phi}_-(\hat{x}'', \hat{y}'')t + \theta)} m a_-(\hat{x}'', (\hat{y}'', -\theta), mt) dt d\theta. \tag{4.8}$$

We can use the complex stationary phase formula of Melin–Sjöstrand [25, Theorem 2.3] to carry the $dt d\theta$ integration in (4.8) and get (the calculation is similar as in the proof of Theorem 3.17 in [14], we omit the details)

$$\begin{aligned} I_0 &= e^{im\Psi(x,y)} b(\hat{x}'', \hat{y}'', m) + O(m^{-\infty}), \\ \Psi(x, y) &= \hat{\Phi}_-(\hat{x}'', \hat{y}'') - x_{2n+1} + y_{2n+1}, \\ b(\hat{x}'', \hat{y}'', m) &\in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\ b(\hat{x}'', \hat{y}'', m) &\sim \sum_{j=0}^\infty m^{n-\frac{d}{2}-j} b_j(\hat{x}'', \hat{y}'') \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\ b_j(\hat{x}'', \hat{y}'') &\in C^\infty(U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \quad j = 0, 1, 2, \dots, \\ b_0(p, p) &= a_-^0(p, p) = 2^{\frac{d}{2}-1} \frac{1}{\text{Veff}(p)} \pi^{-n-1+\frac{d}{2}} |\mu_1|^{\frac{1}{2}} \cdots |\mu_d|^{\frac{1}{2}} |\mu_{d+1}| \cdots |\mu_n| \tau_{p,n-}. \end{aligned}$$

Assume that $q = n_+ \neq n_-$. If $m \rightarrow -\infty$, then the expansion for $S_{G,m}^{(q)}(x, y)$ as $m \rightarrow -\infty$ is similar to $q = n_-$ case. When $m \rightarrow +\infty$, we can repeat the method above with minor change and deduce that $S_{G,m}^{(q)}(x, y) = O(m^{-\infty})$ on X . Summing up, we get Theorem 1.8.

5 Equivalent of the phase function $\Phi_-(x, y)$

Let $p \in \mu^{-1}(0)$ and let U be a small open set of p . We need

Definition 5.1 With the assumptions and notations used in Theorem 3.13, let $\Phi_1, \Phi_2 \in C^\infty(U \times U)$. We assume that Φ_1 and Φ_2 satisfy (3.52), (3.51) and (3.44). We say that Φ_1 and Φ_2 are equivalent on U if for any $b_1(x, y, t) \in S_{cl}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ we can find $b_2(x, y, t) \in S_{cl}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$ such that

$$\int_0^\infty e^{i\Phi_1(x,y)t} b_1(x, y, t) dt \equiv \int_0^\infty e^{i\Phi_2(x,y)t} b_2(x, y, t) dt \text{ on } U$$

and vice versa.

We characterize now the phase Φ_- .

Theorem 5.2 Let $\Phi_-(x, y) \in C^\infty(U \times U)$ be as in Theorem 3.13. Let $\Phi \in C^\infty(U \times U)$. We assume that Φ satisfies (3.52), (3.51) and (3.44). The functions Φ and Φ_- are equivalent on U in the sense of Definition 5.1 if and only if there is a function $f \in C^\infty(U \times U)$ with $f(x, x) = 1$ such that $\Phi(x, y) - f(x, y)\Phi_-(x, y)$ vanishes to infinite order at $\text{diag}((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U))$.

Proof The “ \Leftarrow ” part follows from global theory of complex Fourier integral operator of Melin–Sjöstrand [25]. We only need to prove the “ \Rightarrow ” part. Take local coordinates $x = (x_1, \dots, x_{2n+1})$ defined in some small neighbourhood of p such that $x(p) = 0$ and $\omega_0(p) = dx_{2n+1}$. Since $d_y \Phi(x, y)|_{x=y \in \mu^{-1}(0)} = d_y \Phi_-(x, y)|_{x=y \in \mu^{-1}(0)} = \omega_0(x)$, we have $\frac{\partial \Phi}{\partial y_{2n+1}}(p, p) = \frac{\partial \Phi_-}{\partial y_{2n+1}}(p, p) = 1$. From this observation and the Malgrange preparation theorem [12, Theorem 7.5.7], we conclude that in some small neighborhood of (p, p) , we can find $f(x, y), f_1(x, y) \in C^\infty$ such that

$$\Phi_-(x, y) = f(x, y)(y_{2n+1} + h(x, \hat{y})), \quad \Phi(x, y) = f_1(x, y)(y_{2n+1} + h_1(x, \hat{y})) \tag{5.1}$$

in some small neighborhood of (p, p) , where $\hat{y} = (y_1, \dots, y_{2n})$. For simplicity, we assume that (5.1) hold on $U \times U$. It is clear that $\Phi_-(x, y)$ and $y_{2n+1} + h(x, \hat{y})$ are equivalent in the sense of Definition 5.1, $\Phi(x, y)$ and $y_{2n+1} + h_1(x, \hat{y})$ are equivalent in the sense of Definition 5.1, we may assume that $\Phi_-(x, y) = y_{2n+1} + h(x, \hat{y})$ and $\Phi(x, y) = y_{2n+1} + h_1(x, \hat{y})$. Fix $x_0 \in \mu^{-1}(0) \cap U$. We are going to prove that $h(x, \hat{y}) - h_1(x, \hat{y})$ vanishes to infinite order at $(x_0, x_0) \in (\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U)$. Take

$$b(x, y, t) \sim \sum_{j=0}^\infty b_j(x, y) t^{n-\frac{d}{2}-j} \in S_{cl}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$$

with $b_0(x, x) \neq 0$ at each $x \in U \cap \mu^{-1}(0)$. Since Φ and Φ_- are equivalent on U in the sense of Definition 5.1, we can find $a(x, y, t) \in S_{cl}^{n-\frac{d}{2}}(U \times U \times \mathbb{R}_+, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$

such that

$$\int_0^\infty e^{i\Phi_{-(x,y)t}} b(x, y, t) dt \equiv \int_0^\infty e^{i\Phi_{(x,y)t}} a(x, y, t) dt \text{ on } U.$$

Put $x_0 = (x_0^1, x_0^2, \dots, x_0^{2n+1})$ and $\hat{x}_0 = (x_0^1, \dots, x_0^{2n})$. Take $\tau \in C_0^\infty(\mathbb{R}^{2n+1})$, $\tau_1 \in C_0^\infty(\mathbb{R}^{2n})$, $\chi \in C_0^\infty(\mathbb{R})$ so that $\tau = 1$ near x_0 , $\tau_1 = 1$ near \hat{x}_0 , $\chi = 1$ near x_0^{2n+1} and $\text{Supp } \tau \Subset U$, $\text{Supp } \tau_1 \times \text{Supp } \chi \Subset U' \times \text{Supp } \chi \Subset U$, where U' is an open neighborhood of \hat{x}_0 in \mathbb{R}^{2n} . For each $k > 0$, we consider the distributions

$$\begin{aligned} A_k : u &\mapsto \int_0^\infty e^{i(y_{2n-1}+h(x,\hat{y}))t-iky_{2n+1}} \tau(x)b(x, y, t)\tau_1(\hat{y})\chi(y_{2n+1})u(\hat{y})dydt, \\ B_k : u &\mapsto \int_0^\infty e^{i(y_{2n+1}+h_1(x,\hat{y}))t-iky_{2n+1}} \tau(x)a(x, y, t)\tau_1(\hat{y})\chi(y_{2n+1})u(\hat{y})dydt, \end{aligned}$$

for $u \in C_0^\infty(U', T^{*0,q}X)$. By using the stationary phase formula of Melin–Sjöstrand [25], we can show that (cf. the proof of [14, Theorem 3.12]) A_k and B_k are smoothing operators and

$$\begin{aligned} A_k(x, \hat{y}) &\equiv e^{ikh(x,\hat{y})} g(x, \hat{y}, k) + O(k^{-\infty}), \\ B_k(x, \hat{y}) &\equiv e^{ikh_1(x,\hat{y})} p(x, \hat{y}, k) + O(k^{-\infty}), \\ g(x, \hat{y}, k), p(x, \hat{y}, k) &\in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U', T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ g(x, \hat{y}, k) &\sim \sum_{j=0}^\infty g_j(x, \hat{y})k^{n-\frac{d}{2}-j} \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U', T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ p(x, \hat{y}, k) &\sim \sum_{j=0}^\infty p_j(x, \hat{y})k^{n-\frac{d}{2}-j} \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U', T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ g_j(x, \hat{y}), p_j(x, \hat{y}) &\in C^\infty(U \times U', T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \quad j = 0, 1, \dots, \\ g_0(x_0, \hat{x}_0) &\neq 0. \end{aligned}$$

Since

$$\int_0^\infty e^{i(y_{2n+1}+h(x,\hat{y}))t} b(x, y, t) dt - \int_0^\infty e^{i(y_{2n+1}+h_1(x,\hat{y}))t} a(x, y, t) dt$$

is smoothing, by using integration by parts with respect to y_{2n+1} , it is easy to see that $A_k - B_k = O(k^{-\infty})$ (see [14, Section 3]). Thus,

$$\begin{aligned} e^{ikh(x,\hat{y})} g(x, \hat{y}, k) &= e^{ikh_1(x,\hat{y})} p(x, \hat{y}, k) + F_k(x, \hat{y}), \\ F_k(x, \hat{y}') &= O(k^{-\infty}). \end{aligned} \tag{5.2}$$

Now, we are ready to prove that $h(x, \hat{y}) - h_1(x, \hat{y})$ vanishes to infinite order at (x_0, \hat{x}_0) . We assume that there exist $\alpha_0 \in \mathbb{N}_0^{2n+1}$, $\beta_0 \in \mathbb{N}_0^{2n}$, $|\alpha_0| + |\beta_0| \geq 1$ such that

$$|\partial_x^{\alpha_0} \partial_{\hat{y}}^{\beta_0} (ih(x, \hat{y}) - ih_1(x, \hat{y}))|_{(x_0, \hat{x}_0)} = C_{\alpha_0, \beta_0} \neq 0$$

and

$$|\partial_x^\alpha \partial_{\hat{y}}^\beta (ih(x, \hat{y}) - ih_1(x, \hat{y}))|_{(x_0, \hat{x}_0)} = 0 \text{ if } |\alpha| + |\beta| < |\alpha_0| + |\beta_0|.$$

From (5.2), we have

$$\begin{aligned} &|\partial_x^{\alpha_0} \partial_{\hat{y}}^{\beta_0} (e^{ikh(x,\hat{y})-ikh_1(x,\hat{y})} g(x, \hat{y}, k) - p(x, \hat{y}, k))|_{(x_0, \hat{x}_0)} \\ &= -|\partial_x^{\alpha_0} \partial_{\hat{y}}^{\beta_0} (e^{-ikh_1(x,\hat{y})} F_k(x, \hat{y}))|_{(x_0, \hat{x}_0)}. \end{aligned} \tag{5.3}$$

Since $h_1(x_0, \hat{x}_0) = -x_0^{2n+1}$ and $F_k(x, \hat{y}) = O(k^{-\infty})$, we have

$$\lim_{k \rightarrow \infty} k^{-n+\frac{d}{2}-1} |\partial_x^{\alpha_0} \partial_{\hat{y}}^{\beta_0} (e^{-ikh_1(x,\hat{y})} F_k(x, \hat{y}))|_{(x_0, \hat{x}_0)} = 0. \tag{5.4}$$

On the other hand, we can check that

$$\begin{aligned} & \lim_{k \rightarrow \infty} k^{-n+\frac{d}{2}-1} |\partial_x^{\alpha_0} \partial_y^{\beta_0} (e^{ikh(x,\hat{y})-ikh_1(x,\hat{y})} g(x, \hat{y}, k) - p(x, \hat{y}, k))|_{(x_0, \hat{x}_0)} \\ & = C_{\alpha_0, \beta_0} g_0(x_0, \hat{x}_0) \neq 0 \end{aligned} \tag{5.5}$$

since $g_0(x_0, \hat{x}_0) \neq 0$. From (5.3), (5.4) and (5.5), we get a contradiction. Thus, $h(x, \hat{y}) - h_1(x, \hat{y})$ vanishes to infinite order at (x_0, \hat{x}_0) . Since x_0 is arbitrary, the theorem follows. \square

6 The proof of Theorem 1.9

6.1 Preparation

Fix $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4 defined in an open set U of p . We may assume that $U = \Omega_1 \times \Omega_2 \times \Omega_3 \times \Omega_4$, where $\Omega_1 \subset \mathbb{R}^d$, $\Omega_2 \subset \mathbb{R}^d$ are open sets of $0 \in \mathbb{R}^d$, $\Omega_3 \subset \mathbb{R}^{2n-2d}$ is an open set of $0 \in \mathbb{R}^{2n-2d}$ and Ω_4 is an open set of $0 \in \mathbb{R}$. From now on, we identify Ω_2 with

$$\{(0, \dots, 0, x_{d+1}, \dots, x_{2d}, 0, \dots, 0) \in U; (x_{d+1}, \dots, x_{2d}) \in \Omega_2\},$$

Ω_3 with $\{(0, \dots, 0, x_{2d+1}, \dots, x_{2n}, 0) \in U; (x_{d+1}, \dots, x_{2n}) \in \Omega_3\}$, $\Omega_2 \times \Omega_3$ with

$$\{(0, \dots, 0, x_{d+1}, \dots, x_{2n}, 0) \in U; (x_{d+1}, \dots, x_{2n}) \in \Omega_2 \times \Omega_3\}.$$

For $x = (x_1, \dots, x_{2n+1})$, we write $x'' = (x_{d+1}, \dots, x_{2n+1})$, $\hat{x}'' = (x_{d+1}, \dots, x_{2n})$, $\tilde{x}'' = (x_{d+1}, \dots, x_{2d})$,

$$\tilde{\tilde{x}}'' = (x_{2d+1}, \dots, x_{2n+1}), \quad \tilde{\tilde{\tilde{x}}}' = (x_{2d+1}, \dots, x_{2n}).$$

From now on, we identify x'' with $(0, \dots, 0, x_{d+1}, \dots, x_{2n+1}) \in U$, $\hat{x}'' = (x_{d+1}, \dots, x_{2n})$ with $(0, \dots, 0, x_{d+1}, \dots, x_{2n}, 0) \in U$, \tilde{x}'' with $(0, \dots, 0, x_{d+1}, \dots, x_{2d}, 0, \dots, 0) \in U$, $\tilde{\tilde{x}}''$ with $(0, \dots, 0, x_{2d+1}, \dots, x_{2n+1}) \in U$, $\tilde{\tilde{\tilde{x}}}'$ with $(0, \dots, 0, x_{2d+1}, \dots, x_{2n}, 0)$. Since $G \times S^1$ acts freely on $\mu^{-1}(0)$, we take Ω_2 and Ω_3 small enough so that if $x, x_1 \in \Omega_2 \times \Omega_3$ and $x \neq x_1$, then

$$g \circ e^{i\theta} \circ x \neq g_1 \circ e^{i\theta_1} \circ x_1, \quad \forall (g, e^{i\theta}) \in G \times S^1, \quad \forall (g_1, e^{i\theta_1}) \in G \times S^1. \tag{6.1}$$

We now assume that $q = n_-$ and let $\Psi(x, y) \in C^\infty(U \times U)$ be as in Theorem 1.8. From $S_{G,m}^{(q)} = (S_{G,m}^{(q)})^*$, we get

$$e^{im\Psi(x,y)} b(x, y, m) = e^{-im\bar{\Psi}(y,x)} b^*(x, y, m) + O(m^{-\infty}), \tag{6.2}$$

where $(S_{G,m}^{(q)})^* : L^2_{(0,q)}(X) \rightarrow L^2_{(0,q)}(X)$ is the adjoint of $S_{G,m}^{(q)} : L^2_{(0,q)}(X) \rightarrow L^2_{(0,q)}(X)$ with respect to $\langle \cdot | \cdot \rangle$ and $b^*(x, y, m) : T_x^{*0,q} X \rightarrow T_y^{*0,q} X$ is the adjoint of $b(x, y, m) : T_y^{*0,q} X \rightarrow T_x^{*0,q} X$ with respect to $\langle \cdot | \cdot \rangle$. From (6.2), we can repeat the proof of Theorem 5.2 with minor change and deduce that

$$\Psi(x, y) + \bar{\Psi}(y, x) \text{ vanishes to infinite order at } \text{diag} \left((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U) \right). \tag{6.3}$$

From $\bar{\partial}_b S_{G,m}^{(q)} = 0$, we can check that

$$\bar{\partial}_b \Psi(x, y) \text{ vanishes to infinite order at } \text{diag} \left((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U) \right). \tag{6.4}$$

From (6.3), (6.4) and notice that $\frac{\partial}{\partial x_j} - i \frac{\partial}{\partial x_{d+j}} \in T_x^{0,1} X, j = 1, \dots, d$, where $x \in \mu^{-1}(0)$ (see Remark 4.4), and $\frac{\partial}{\partial x_j} \Psi(x, y) = \frac{\partial}{\partial y_j} \Psi(x, y) = 0, j = 1, \dots, d$, we conclude that

$$\frac{\partial}{\partial x_{d+j}} \Psi(x, y)|_{x_{d+1}=\dots=x_{2d}=0} \text{ and } \frac{\partial}{\partial y_{d+j}} \Psi(x, y)|_{y_{d+1}=\dots=y_{2d}=0} \text{ vanish to infinite order at } \text{diag} \left((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U) \right).$$

Let $G_j(x, y) := \frac{\partial}{\partial y_{d+j}} \Psi(x, y)|_{y_{d+1}=\dots=y_{2d}=0}, H_j(x, y) := \frac{\partial}{\partial x_{d+j}} \Psi(x, y)|_{x_{d+1}=\dots=x_{2d}=0}$. Put

$$\Psi_1(x, y) := \Psi(x, y) - \sum_{j=1}^d y_{d+j} G_j(x, y), \quad \Psi_2(x, y) := \Psi(x, y) - \sum_{j=1}^d x_{d+j} H_j(x, y).$$

Then, for $j = 1, 2, \dots, d$,

$$\frac{\partial}{\partial y_{d+j}} \Psi_1(x, y)|_{y_{d+1}=\dots=y_{2d}=0} = 0 \quad \text{and} \quad \frac{\partial}{\partial x_{d+j}} \Psi_2(x, y)|_{x_{d+1}=\dots=x_{2d}=0} = 0, \tag{6.5}$$

and, for $j = 1, 2$,

$$\Psi(x, y) - \Psi_j(x, y) \text{ vanishes to infinite order at } \text{diag} \left((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U) \right). \tag{6.6}$$

We also write $u = (u_1, \dots, u_{2n+1})$ to denote the local coordinates of U . Recall that for any smooth function $f \in C^\infty(U)$, we write $\tilde{f} \in C^\infty(U^{\mathbb{C}})$ to denote an almost analytic extension of f (see the discussion after (3.38)). We consider the following two systems

$$\frac{\partial \tilde{\Psi}_1}{\partial \tilde{u}_{2d+j}}(\tilde{x}, \tilde{u}'') + \frac{\partial \tilde{\Psi}_2}{\partial \tilde{x}_{2d+j}}(\tilde{u}'', \tilde{y}) = 0, \quad j = 1, 2, \dots, 2n - 2d, \tag{6.7}$$

and

$$\frac{\partial \tilde{\Psi}_1}{\partial \tilde{u}_{d+j}}(\tilde{x}, \tilde{u}'') + \frac{\partial \tilde{\Psi}_2}{\partial \tilde{x}_{d+j}}(\tilde{u}'', \tilde{y}) = 0, \quad j = 1, 2, \dots, 2n - d, \tag{6.8}$$

where $\tilde{u}'' = (0, \dots, 0, \tilde{u}_{2d+1}, \dots, \tilde{u}_{2n+1}), \tilde{u}'' = (0, \dots, 0, \tilde{u}_{d+1}, \dots, \tilde{u}_{2n+1})$. From (6.5) and Theorem 1.12, we can take $\tilde{\Psi}_1$ and $\tilde{\Psi}_2$ so that for every $j = 1, 2, \dots, d$,

$$\frac{\partial \tilde{\Psi}_1}{\partial \tilde{u}_{d+j}}(\tilde{x}, \tilde{u}'') = 0 \quad \text{and} \quad \frac{\partial \tilde{\Psi}_2}{\partial \tilde{x}_{d+j}}(\tilde{u}'', \tilde{y}) = 0, \quad \text{if } \tilde{u}_{d+1} = \dots = \tilde{u}_{2d} = 0, \tag{6.9}$$

and, for $j = 1, 2$,

$$\tilde{\Psi}_j(\tilde{x}, \tilde{y}) = -\tilde{x}_{2n+1} + \tilde{y}_{2n+1} + \tilde{\Psi}_j(\tilde{x}'', \tilde{y}''), \quad \tilde{\Psi}_j \in C^\infty(U^{\mathbb{C}} \times U^{\mathbb{C}}), \tag{6.10}$$

where $\tilde{x}'' = (0, \dots, 0, \tilde{x}_{d+1}, \dots, \tilde{x}_{2n}, 0), \tilde{y}'' = (0, \dots, 0, \tilde{y}_{d+1}, \dots, \tilde{y}_{2n}, 0)$.

From Theorem 1.12, (1.19) and $d_x \Psi(x, x) = -d_y \Psi(x, x) = -\omega_0(x), \forall x \in \mu^{-1}(0)$, it is not difficult to see that

$$\frac{\partial \tilde{\Psi}_1}{\partial \tilde{u}_{d+j}}(\tilde{x}'', \tilde{x}'') + \frac{\partial \tilde{\Psi}_2}{\partial \tilde{x}_{d+j}}(\tilde{x}'', \tilde{x}'') = 0, \quad j = 1, 2, \dots, 2n - d,$$

and the matrices

$$\left(\frac{\partial^2 \Psi}{\partial u_{2d+j} \partial u_{2d+k}}(p, p) + \frac{\partial^2 \Psi}{\partial x_{2d+j} \partial x_{2d+k}}(p, p) \right)_{j,k=1}^{2n-2d},$$

$$\left(\frac{\partial^2 \Psi}{\partial u_{d+j} \partial u_{d+k}}(p, p) + \frac{\partial^2 \Psi}{\partial x_{d+j} \partial x_{d+k}}(p, p) \right)_{j,k=1}^{2n-d}$$

are non-singular. Moreover,

$$\det \left(\frac{\partial^2 \Psi}{\partial u_{2d+j} \partial u_{2d+k}}(p, p) + \frac{\partial^2 \Psi}{\partial x_{2d+j} \partial x_{2d+k}}(p, p) \right)_{j,k=1}^{2n-2d} = (4i|\mu_{d+1}| \cdots 4i|\mu_n|)^2,$$

$$\det \left(\frac{\partial^2 \Psi}{\partial u_{d+j} \partial u_{d+k}}(p, p) + \frac{\partial^2 \Psi}{\partial x_{d+j} \partial x_{d+k}}(p, p) \right)_{j,k=1}^{2n-d} = (8i|\mu_1| \cdots 8i|\mu_d|)(4i|\mu_{d+1}| \cdots 4i|\mu_n|)^2.$$

Hence, near (p, p) , we can solve (6.7) and (6.8) and the solutions are unique. Let $\alpha(x, y) = (\alpha_{2d+1}(x, y), \dots, \alpha_{2n}(x, y)) \in C^\infty(U \times U, \mathbb{C}^{2n-2d})$ and $\beta(x, y) = (\beta_{d+1}(x, y), \dots, \beta_{2n}(x, y)) \in C^\infty(U \times U, \mathbb{C}^{2n-d})$ be the solutions of (6.7) and (6.8), respectively. From (6.9), it is easy to see that

$$\beta(x, y) = (\beta_{d+1}(x, y), \dots, \beta_{2n}(x, y)) = (0, \dots, 0, \alpha_{2d+1}(x, y), \dots, \alpha_{2n}(x, y)). \tag{6.11}$$

From (6.11), we see that the value of $\tilde{\Psi}_1(x, \tilde{u}'') + \tilde{\Psi}_2(\tilde{u}'', y)$ at critical point $\tilde{u}'' = \alpha(x, y)$ is equal to the value of $\tilde{\Psi}_1(x, \tilde{u}'') + \tilde{\Psi}_2(\tilde{u}'', y)$ at critical point $\tilde{u}'' = \beta(x, y)$. Put

$$\Psi_3(x, y) := \tilde{\Psi}_1(x, \alpha(x, y)) + \tilde{\Psi}_2(\alpha(x, y), y) = \tilde{\Psi}_1(x, \beta(x, y)) + \tilde{\Psi}_2(\beta(x, y), y). \tag{6.12}$$

$\Psi_3(x, y)$ is a complex phase function. From (6.10), we have

$$\Psi_3(x, y) = -x_{2n+1} + y_{2n+1} + \hat{\Psi}_3(\hat{x}'', \hat{y}''), \quad \hat{\Psi}_3(\hat{x}'', \hat{y}'') \in C^\infty(U \times U).$$

Moreover, we have the following

Theorem 6.1 *The function $\Psi_3(x, y) - \Psi(x, y)$ vanishes to infinite order at $\text{diag} \left((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U) \right)$.*

Proof We consider the kernel $S_{G,m}^{(q)} \circ S_{G,m}^{(q)}$ on U . Let $V \Subset U$ be an open set of p . Let $\chi(\hat{x}'') \in C_0^\infty(\Omega_2 \times \Omega_3)$. From (6.1), we can extend $\chi(\hat{x}'')$ to $W := \{g \circ e^{i\theta} \circ x; (g, e^{i\theta}) \in G \times S^1, x \in \Omega_2 \times \Omega_3\}$ by $\chi(g \circ e^{i\theta} \circ \hat{x}'') := \chi(\hat{x}'')$, for every $(g, e^{i\theta}) \in G \times S^1$. Assume that $\chi = 1$ on some neighborhood of V . Let $\chi_1 \in C_0^\infty(U)$ with $\chi_1 = 1$ on some neighborhood of V and $\text{Supp } \chi_1 \subset \{x \in X; \chi(x) = 1\}$. We have

$$\chi_1 S_{G,m}^{(q)} \circ S_{G,m}^{(q)} = \chi_1 S_{G,m}^{(q)} \chi \circ S_{G,m}^{(q)} + \chi_1 S_{G,m}^{(q)} (1 - \chi) \circ S_{G,m}^{(q)}. \tag{6.13}$$

Let's first consider $\chi_1 S_{G,m}^{(q)} (1 - \chi) \circ S_{G,m}^{(q)}$. We have

$$(\chi_1 S_{G,m}^{(q)} (1 - \chi))(x, u) = \frac{1}{2\pi} \chi_1(x) \int_{-\pi}^\pi \int_G S^{(q)}(x, g \circ e^{i\theta} \circ u) (1 - \chi(u)) e^{im\theta} d\mu(g) d\theta. \tag{6.14}$$

If $u \notin \{x \in X; \chi(x) = 1\}$. Since $\text{Supp } \chi_1 \subset \{x \in X; \chi(x) = 1\}$ and $\chi(x) = \chi(g \circ e^{i\theta} \circ x)$, for every $(g, e^{i\theta}) \in G \times S^1$, for every $x \in X$, we conclude that $g \circ e^{i\theta} \circ u \notin \text{Supp } \chi_1$, for every $(g, e^{i\theta}) \in G \times S^1$. From this observation and notice that $S^{(q)}$ is smoothing away from diagonal, we can integrate by parts with respect to θ in (6.14) and deduce that $\chi_1 S_{G,m}^{(q)} \circ (1 - \chi) = O(m^{-\infty})$ and hence

$$\chi_1 S_{G,m}^{(q)} (1 - \chi) \circ S_{G,m}^{(q)} = O(m^{-\infty}). \tag{6.15}$$

From (6.13) and (6.15), we get

$$\chi_1 S_{G,m}^{(q)} \circ S_{G,m}^{(q)} = \chi_1 S_{G,m}^{(q)} \chi \circ S_{G,m}^{(q)} + O(m^{-\infty}). \tag{6.16}$$

We can check that on U ,

$$\begin{aligned}
 & (\chi_1 S_{G,m}^{(q)} \chi \circ S_{G,m}^{(q)})(x, y) \\
 &= (2\pi) \int e^{im\Psi(x,u'')+im\Psi(u'',y)} \chi_1(x)b(x, \hat{u}'', m)\chi(\hat{u}'')b(\hat{u}'', y, m)dv(\hat{u}'') + O(m^{-\infty}) \\
 &= (2\pi) \int e^{im\Psi_1(x,u'')+im\Psi_2(u'',y)} \chi_1(x)b(x, \hat{u}'', m)\chi(\hat{u}'')b(\hat{u}'', y, m)dv(\hat{u}'') + O(m^{-\infty}) \\
 & \text{(here we use (6.6)),}
 \end{aligned}
 \tag{6.17}$$

where $d\mu(g)d\theta dv(\hat{u}'') = dv(x)$ on U . We use complex stationary phase formula of Melin–Sjöstrand [25] to carry out the integral (6.17) and get

$$\begin{aligned}
 & (\chi_1 S_{G,m}^{(q)} \chi \circ S_{G,m}^{(q)})(x, y) = e^{im\Psi_3(x,y)}a(x, y, m) + O(m^{-\infty}) \text{ on } U, \\
 & a(x, y, m) \in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\
 & a(x, y, m) \sim \sum_{j=0}^{\infty} m^{n-\frac{d}{2}-j} a_j(x, y) \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\
 & a_j(x, y) \in C^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \quad j = 0, 1, 2, \dots, \\
 & a_0(p, p) \neq 0.
 \end{aligned}
 \tag{6.18}$$

From (6.16), (6.18) and notice that $(\chi_1 S_{G,m}^{(q)} \circ S_{G,m}^{(q)})(x, y) = (\chi_1 S_{G,m}^{(q)})(x, y)$, we deduce that

$$e^{im\Psi_3(x,y)}a(x, y, m) = e^{im\Psi(x,y)}\chi_1(x)b(x, y, m) + O(m^{-\infty}) \text{ on } U.
 \tag{6.19}$$

From (6.19), we can repeat the proof of Theorem 5.2 with minor change and deduce that $\Psi_3(x, y) - \Psi(x, y)$ vanishes to infinite order at $\text{diag}((\mu^{-1}(0) \cap U) \times (\mu^{-1}(0) \cap U))$. \square

The following two theorems follow from (6.6), (6.12), Theorem 6.1, complex stationary phase formula of Melin–Sjöstrand [25] and some straightforward computation. We omit the details.

Theorem 6.2 *With the notations used above, let*

$$\begin{aligned}
 & A_m(x, y) = e^{im\Psi(x,y)}a(x, y, m), \quad B_m(x, y) = e^{im\Psi(x,y)}b(x, y, m), \\
 & a(x, y, m) \in S_{\text{loc}}^k(1; U \times U, H \boxtimes F^*), \quad b(x, y, m) \in S_{\text{loc}}^\ell(1; U \times U, F \boxtimes E^*), \\
 & a(x, y, m) \sim \sum_{j=0}^{\infty} m^{k-j} a_j(x, y) \text{ in } S_{\text{loc}}^k(1; U \times U, H \boxtimes F^*), \\
 & b(x, y, m) \sim \sum_{j=0}^{\infty} m^{\ell-j} b_j(x, y) \text{ in } S_{\text{loc}}^\ell(1; U \times U, F \boxtimes E^*), \\
 & a_j(x, y) \in C^\infty(U \times U, H \boxtimes F^*), \quad b_j(x, y) \in C^\infty(U \times U, F \boxtimes E^*), \quad j = 0, 1, 2, \dots,
 \end{aligned}$$

where E, F and H are vector bundles over X . Let $\chi(\hat{x}'') \in C_0^\infty(\Omega_2 \times \Omega_3)$. Then, we have

$$\begin{aligned}
 & \int A_m(x, u)\chi(\hat{u}'')B_m(u, y)dv(\hat{u}'') = e^{im\Psi(x,y)}c(x, y, m) + O(m^{-\infty}), \\
 & c(x, y, m) \in S_{\text{loc}}^{k+\ell-(n-\frac{d}{2})}(1; U \times U, H \boxtimes E^*), \\
 & c(x, y, m) \sim \sum_{j=0}^{\infty} m^{k+\ell-(n-\frac{d}{2})-j} c_j(x, y) \text{ in } S_{\text{loc}}^{k+\ell-(n-\frac{d}{2})}(1; U \times U, H \boxtimes E^*), \\
 & c_0(x, x) = 2^{-n-\frac{d}{2}}\pi^{n-\frac{d}{2}}|\det \mathcal{L}_x|^{-1}|\det R_x|^{\frac{1}{2}}a_0(x, x)b_0(x, x)\chi(\hat{x}''), \quad \forall x \in \mu^{-1}(0) \cap U,
 \end{aligned}$$

where $|\det R_x|$ is in the discussion before Theorem 1.6.

Moreover, if there are $N_1, N_2 \in \mathbb{N}$, such that $|a_0(x, y)| \leq C|(x, y) - (x_0, x_0)|^{N_1}$, $|b_0(x, y)| \leq C|(x, y) - (x_0, x_0)|^{N_2}$, for all $x_0 \in \mu^{-1}(0) \cap U$, where $C > 0$ is a constant, then,

$$|c_0(x, y)| \leq \hat{C} |(x, y) - (x_0, x_0)|^{N_1+N_2},$$

for all $x_0 \in \mu^{-1}(0) \cap U$, where $\hat{C} > 0$ is a constant.

Theorem 6.3 *With the notations used above, let*

$$\begin{aligned} & \mathcal{A}_m(x, \tilde{y}'') = e^{im\Psi(x, \tilde{y}'')} \alpha(x, \tilde{y}'', m), \quad \mathcal{B}_m(\tilde{x}'', y) = e^{im\Psi(\tilde{x}'', y)} \beta(\tilde{x}'', y, m), \\ & \alpha(x, \tilde{y}'', m) \in S_{\text{loc}}^k(1; U \times (\Omega_3 \times \Omega_4), H \boxtimes F^*), \quad \beta(\tilde{x}'', y, m) \\ & \in S_{\text{loc}}^\ell(1; (\Omega_3 \times \Omega_4) \times U, F \boxtimes E^*), \\ & \alpha(x, \tilde{y}'', m) \sim \sum_{j=0}^\infty m^{k-j} \alpha_j(x, \tilde{y}'') \text{ in } S_{\text{loc}}^k(1; U \times (\Omega_3 \times \Omega_4), H \boxtimes F^*), \\ & \beta(\tilde{x}'', y, m) \sim \sum_{j=0}^\infty m^{\ell-j} \beta_j(\tilde{x}'', y) \text{ in } S_{\text{loc}}^\ell(1; (\Omega_3 \times \Omega_4) \times U, F \boxtimes E^*), \\ & \alpha_j(x, \tilde{y}'') \in C^\infty(U \times (\Omega_3 \times \Omega_4), H \boxtimes F^*), \quad \beta_j(\tilde{x}'', y) \\ & \in C^\infty((\Omega_3 \times \Omega_4) \times U, F \boxtimes E^*), \quad j = 0, 1, 2, \dots, \end{aligned}$$

where E, F and H are vector bundles over X . Let $\chi_1(\tilde{x}'') \in C_0^\infty(\Omega_3)$. Then, we have

$$\begin{aligned} & \int \mathcal{A}_m(x, \tilde{u}'') \chi_1(\tilde{u}'') \mathcal{B}_m(\tilde{u}'', y) dv(\tilde{u}'') = e^{im\Psi(x, y)} \gamma(x, y, m) + O(m^{-\infty}), \\ & \gamma(x, y, m) \in S_{\text{loc}}^{k+\ell-(n-d)}(1; U \times U, H \boxtimes E^*), \\ & \gamma(x, y, m) \sim \sum_{j=0}^\infty m^{k+\ell-(n-d)-j} \gamma_j(x, y) \text{ in } S_{\text{loc}}^{k+\ell-(n-d)}(1; U \times U, H \boxtimes E^*), \\ & \gamma_0(x, x) = 2^{-n} \pi^{n-d} |\det \mathcal{L}_x|^{-1} |\det R_x| \alpha_0(x, \tilde{x}'') \beta_0(\tilde{x}'', x) \chi_1(\tilde{x}''), \quad \forall x \in \mu^{-1}(0) \cap U, \end{aligned}$$

where $|\det R_x|$ is in the discussion before Theorem 1.6.

Moreover, if there are $N_1, N_2 \in \mathbb{N}$, such that $|\alpha_0(x, \tilde{y}'')| \leq C|(x, \tilde{y}'') - (x_0, x_0)|^{N_1}$, $|\beta_0(x, \tilde{y}'')| \leq C|(x, \tilde{y}'') - (x_0, x_0)|^{N_2}$, for all $x_0 \in \mu^{-1}(0) \cap U$, where $C > 0$ is a constant, then,

$$|\gamma_0(x, y)| \leq \hat{C} |(x, y) - (x_0, x_0)|^{N_1+N_2},$$

for all $x_0 \in \mu^{-1}(0) \cap U$, where $\hat{C} > 0$ is a constant.

6.2 The proof of Theorem 1.9

Since \underline{g}_x is orthogonal to $H_x Y \cap JH_x Y$ and $H_x Y \cap JH_x Y \subset \underline{g}_x^{\perp b}$, for every $x \in Y$, we can find a G -invariant orthonormal basis $\{Z_1, \dots, Z_n\}$ of $T^{1,0}X$ on Y such that

$$\begin{aligned} & \mathcal{L}_x(Z_j(x), \bar{Z}_k(x)) = \delta_{j,k} \lambda_j(x), \quad j, k = 1, \dots, n, \quad x \in Y, \\ & Z_j(x) \in \underline{g}_x + iJ\underline{g}_x, \quad \forall x \in Y, \quad j = 1, 2, \dots, d, \\ & Z_j(x) \in \mathbb{C}H_x Y \cap J(\mathbb{C}H_x Y), \quad \forall x \in Y, \quad j = d + 1, \dots, n. \end{aligned}$$

Let $\{e_1, \dots, e_n\}$ denote the orthonormal basis of $T^{*0,1}X$ on Y , dual to $\{\bar{Z}_1, \dots, \bar{Z}_n\}$. Fix $s = 0, 1, 2, \dots, n - d$. For $x \in Y$, put

$$B_x^{*0,s} X = \left\{ \sum_{d+1 \leq j_1 < \dots < j_s \leq n} a_{j_1, \dots, j_s} e_{j_1} \wedge \dots \wedge e_{j_s}; a_{j_1, \dots, j_s} \in \mathbb{C}, \forall d + 1 \leq j_1 < \dots < j_s \leq n \right\}$$

and let $B^{*0,s} X$ be the vector bundle of Y with fiber $B_x^{*0,s} X, x \in Y$. Let $C^\infty(Y, B^{*0,s} X)^G$ denote the set of all G -invariant sections of Y with values in $B^{*0,s} X$. Let

$${}_{tG} : C^\infty(Y, B^{*0,s} X)^G \rightarrow \Omega^{0,s}(Y_G)$$

be the natural identification.

Assume that $\lambda_1 < 0, \dots, \lambda_r < 0$, and $\lambda_{d+1} < 0, \dots, \lambda_{n-r+d} < 0$. For $x \in Y$, put

$$\hat{\mathcal{N}}(x, n_-) = \{ce_{d+1} \wedge \dots \wedge e_{n-r+d}; c \in \mathbb{C}\},$$

and let

$$\begin{aligned} \hat{p} &= \hat{p}_x : \mathcal{N}(x, n_-) \rightarrow \hat{\mathcal{N}}(x, n_-), \\ u &= ce_1 \wedge \cdots \wedge e_r \wedge e_{d+1} \wedge \cdots \wedge e_{n-r+d} \rightarrow ce_{d+1} \wedge \cdots \wedge e_{n-r+d}. \end{aligned}$$

Let $\iota : Y \rightarrow X$ be the natural inclusion and let $\iota^* : \Omega^{0,q}(X) \rightarrow \Omega^{0,q}(Y)$ be the pull-back of ι . Recall that we work with the assumption that $q = n_-$. Let $\square_{b,Y_G}^{(q-r)}$ be the Kohn Laplacian for $(0, q - r)$ forms on Y_G . Fix $m \in \mathbb{N}$. Let $H_{b,m}^{q-r}(Y_G) := \{u \in \Omega^{0,q-r}(Y_G); \square_{b,Y_G}^{(q-r)} u = 0, Tu = imu\}$. Let $S_{Y_G,m}^{(q-r)} : L^2_{(0,q-r)}(Y_G) \rightarrow H_{b,m}^{q-r}(Y_G)$ be the orthogonal projection and let $S_{Y_G,m}^{(q-r)}(x, y)$ be the distribution kernel of $S_{Y_G,m}^{(q-r)}$. Let

$$f(x) = \sqrt{V_{\text{eff}}(x)} |\det R_x|^{-\frac{1}{4}} \in C^\infty(Y)^G.$$

Let

$$\begin{aligned} \sigma_m : \Omega^{0,q}(X) &\rightarrow H_{b,m}^{q-r}(Y_G), \\ u &\rightarrow m^{-\frac{d}{4}} S_{Y_G,m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x,n_-} \circ f \circ \iota^* \circ S_{G,m}^{(q)}. \end{aligned}$$

Recall that τ_{x,n_-} is given by (1.7). Let $\sigma_m^* : \Omega^{0,q-r}(Y_G) \rightarrow \Omega^{0,q}(X)$ be the adjoints of σ_m . It is easy to see that $\sigma_m^* u \in H_{b,m}^q(X)^G := (\text{Ker } \square_b^{(q)})_m^G, \forall u \in \Omega^{0,q-r}(Y_G)$. Let $\sigma_m(x, y)$ and $\sigma_m^*(x, y)$ denote the distribution kernels of σ_m and σ_m^* , respectively.

Let's pause and recall some results for $S_{Y_G,m}^{(q-r)}$. We first introduce some notations. Let $\mathcal{L}_{Y_G,x}$ be the Levi form on Y_G at $x \in Y_G$ induced naturally from \mathcal{L} . The Hermitian metric $\langle \cdot | \cdot \rangle$ on $T^{1,0}X$ induces a Hermitian metric $\langle \cdot | \cdot \rangle$ on $T^{1,0}Y_G$. Let $\det \mathcal{L}_{Y_G,x} = \lambda_1 \dots \lambda_{n-d}$, where $\lambda_j, j = 1, \dots, n - d$, are the eigenvalues of $\mathcal{L}_{Y_G,x}$ with respect to the Hermitian metric $\langle \cdot | \cdot \rangle$. For $x \in Y_G$, let $\hat{\tau}_x : T_x^{*0,q-r} Y_G \rightarrow \hat{\mathcal{N}}(x, n_-)$ be the orthogonal projection.

Let $\pi : Y \rightarrow Y_G$ be the natural quotient. Let $S_{Y_G}^{(q-r)} : L^2_{(0,q-r)}(Y_G) \rightarrow \text{Ker } \square_{b,Y_G}^{(q-r)}$ be the Szegő projection as (3.2). Since $S_{Y_G}^{(q-r)}$ is smoothing away from diagonal (see Theorem 3.3), it is easy to see that for any $x, y \in Y$, if $\pi(e^{i\theta} \circ x) \neq \pi(e^{i\theta} \circ y)$, for every $\theta \in [0, 2\pi[$, then there are open sets U of $\pi(x)$ in Y_G and V of $\pi(y)$ in Y_G such that for all $\hat{\chi} \in C_0^\infty(U), \tilde{\chi} \in C_0^\infty(V)$, we have

$$\hat{\chi} S_{Y_G,m}^{(q-r)} \tilde{\chi} = O(m^{-\infty}) \text{ on } Y_G. \tag{6.20}$$

Fix $p \in Y$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4. We will use the same notations as in the beginning of Sect. 6.1. From now on, we identify \tilde{x}'' as local coordinates of Y_G near $\pi(p) \in Y_G$ and we identify $W := \Omega_3 \times \Omega_4$ with an open set of $\pi(p)$ in Y_G . It was proved in Theorem 4.11 in [14] that as $m \rightarrow +\infty$,

$$\begin{aligned} S_{Y_G,m}^{(q-r)}(\tilde{x}'', \tilde{y}'') &= e^{im\phi(\tilde{x}'', \tilde{y}'')} b(\tilde{x}'', \tilde{y}'', m) + O(m^{-\infty}) \text{ on } W, \\ \beta(\tilde{x}'', \tilde{y}'', m) &\in S_{\text{loc}}^{n-d}(1; W \times W, T^{*0,q-r} Y_G \boxtimes (T^{*0,q-r} Y_G)^*), \\ \beta(\tilde{x}'', \tilde{y}'', m) &\sim \sum_{j=0}^\infty m^{n-d-j} b_j(\tilde{x}'', \tilde{y}'') \text{ in } S_{\text{loc}}^{n-d}(1; W \times W, T^{*0,q-r} Y_G \boxtimes (T^{*0,q-r} Y_G)^*), \\ \beta_j(\tilde{x}'', \tilde{y}'') &\in C^\infty(W \times W, T^{*0,q-r} Y_G \boxtimes (T^{*0,q-r} Y_G)^*), \quad j = 0, 1, 2, \dots, \\ \beta_0(\tilde{x}'', \tilde{x}'') &= \frac{1}{2} \pi^{-(n-d)-1} |\det \mathcal{L}_{Y_G, \tilde{x}''}| \hat{\tau}_{\tilde{x}''}, \quad \forall \tilde{x}'' \in W, \end{aligned}$$

and

$$\begin{aligned}
 \phi(\tilde{x}'', \tilde{y}'') &= -x_{2n+1} + y_{2n+1} + \hat{\phi}(\tilde{x}'', \tilde{y}'') \in C^\infty(W \times W), \\
 d_{\tilde{y}''} \phi(\tilde{x}'', \tilde{y}'') &= -d_{\tilde{x}''} \phi(\tilde{x}'', \tilde{y}'') = -\omega_0(\tilde{x}''), \\
 \text{Im } \hat{\phi}(\tilde{x}'', \tilde{y}'') &\geq c|\tilde{x}'' - \tilde{y}''|^2, \text{ where } c > 0 \text{ is a constant,} \\
 p_0(\tilde{x}'', d_{\tilde{x}''} \phi(\tilde{x}'', \tilde{y}'')) &\text{ vanishes to infinite order at } \tilde{x}'' = \tilde{y}'', \\
 \phi(\tilde{x}'', \tilde{y}'') &= -x_{2n+1} + y_{2n+1} + i \sum_{j=d+1}^n |\mu_j| |z_j - w_j|^2 \\
 &\quad + \sum_{j=d+1}^n i \mu_j (\bar{z}_j w_j - z_j \bar{w}_j) + O(|(\tilde{x}'', \tilde{y}'')|^3),
 \end{aligned} \tag{6.21}$$

where p_0 denotes the principal symbol of $\square_{b, Y_G}^{(q-r)}$, $z_j = x_{2j-1} + ix_{2j}$, $j = d + 1, \dots, n$, and μ_{d+1}, \dots, μ_n are the eigenvalues of $\mathcal{L}_{Y_G, p}$.

Note that for any $\phi_1(\tilde{x}'', \tilde{y}'') \in C^\infty(W \times W)$, if ϕ_1 satisfies (6.21), then $\phi_1 - \phi$ vanishes to infinite order at $\tilde{x}'' = \tilde{y}''$ (see Remark 3.6 in [14]). It is not difficult to see that the phase function $\Psi(\tilde{x}'', \tilde{y}'')$ satisfies (6.21). Hence, we can replace the phase $\phi(\tilde{x}'', \tilde{y}'')$ by $\Psi(\tilde{x}'', \tilde{y}'')$ and we have

$$S_{Y_G, m}^{(q-r)}(\tilde{x}'', \tilde{y}'') = e^{im\Psi(\tilde{x}'', \tilde{y}'')} \beta(\tilde{x}'', \tilde{y}'', m) + O(m^{-\infty}) \text{ on } W. \tag{6.22}$$

We can now prove

Theorem 6.4 *With the notations used above, if $y \notin Y$, then for any open set D of y with $\overline{D} \cap Y = \emptyset$, we have*

$$\sigma_m = O(m^{-\infty}) \text{ on } Y_G \times D. \tag{6.23}$$

Let $x, y \in Y$. If $\pi(x) \neq \pi(e^{i\theta} \circ y)$, for every $\theta \in [0, 2\pi[$, then there are open sets U_G of $\pi(x)$ in Y_G and V of y in X such that

$$\sigma_m = O(m^{-\infty}) \text{ on } U_G \times V. \tag{6.24}$$

Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4. Then,

$$\begin{aligned}
 \sigma_m(\tilde{x}'', y) &= e^{im\Psi(\tilde{x}'', y'')} \alpha(\tilde{x}'', y'', m) + O(m^{-\infty}) \text{ on } W \times U, \\
 \alpha(\tilde{x}'', y'', m) &\in S_{\text{loc}}^{n-\frac{3}{4}d} (1; W \times U, T^{*0, q-r} Y_G \boxtimes (T^{*0, q} X)^*), \\
 \alpha(\tilde{x}'', y'', m) &\sim \sum_{j=0}^\infty m^{n-\frac{3}{4}d-j} \alpha_j(\tilde{x}'', y'') \text{ in } S_{\text{loc}}^{n-\frac{3}{4}d} (1; W \times U, T^{*0, q-r} Y_G \boxtimes (T^{*0, q} X)^*), \\
 \alpha_j(\tilde{x}'', y'') &\in C^\infty(W \times U, T^{*0, q-r} Y_G \boxtimes (T^{*0, q} X)^*), \quad j = 0, 1, 2, \dots,
 \end{aligned} \tag{6.25}$$

$$\alpha_0(\tilde{x}'', \tilde{x}'') = 2^{-n+2d-1} \pi^{\frac{d}{2}-n-1} \frac{1}{\sqrt{V_{\text{eff}}(\tilde{x}'')}} |\det \mathcal{L}_{\tilde{x}''}| |\det R_x|^{-\frac{3}{4}} \hat{\tau}_{\tilde{x}'', n_-}, \quad \forall \tilde{x}'' \in W, \tag{6.26}$$

where U is an open set of p , $W = \Omega_3 \times \Omega_4$, Ω_3 and Ω_4 are open sets as in the beginning of Sect. 6.1.

Proof Note that $S_{G, m}^{(q)} = O(m^{-\infty})$ away Y . From this observation, we get (6.23). Let $x, y \in Y$. Assume that $\pi(x) \neq \pi(e^{i\theta} \circ y)$, for every $\theta \in [0, 2\pi[$. Since

$$S_{G, m}^{(q)}(x, y) = \frac{1}{2\pi |G|_{d\mu}} \int_{-\pi}^\pi \int_G S^{(q)}(x, e^{i\theta} \circ g \circ y) e^{im\theta} d\mu(g) d\theta$$

and $S^{(q)}$ is smoothing away from diagonal, we can integrate by parts with respect to θ and deduce that there are open sets U_1 of x in X and V_1 of y in X such that

$$S_{G, m}^{(q)} = O(m^{-\infty}) \text{ on } U_1 \times V_1. \tag{6.27}$$

From (6.20), we see that there are open sets \hat{U}_G of $\pi(x)$ in Y_G and \hat{V}_G of $\pi(y)$ in Y_G such that

$$S_{Y_G, m}^{(q-r)} = O(m^{-\infty}) \text{ on } \hat{U}_G \times \hat{V}_G. \tag{6.28}$$

From (6.27) and (6.28), we get (6.24).

Fix $u = (u_1, \dots, u_{2n+1}) \in Y \cap U$. From (6.23) and (6.24), we only need to show that (6.25) and (6.26) hold near u and we may assume that $u = (0, \dots, 0, u_{2d+1}, \dots, u_{2n}, 0) = \tilde{u}''$. Let V be a small neighborhood of u . Let $\chi(\tilde{x}'') \in C_0^\infty(\Omega_3)$. From (6.1), we can extend $\chi(\tilde{x}'')$ to

$$Q = \left\{ g \circ e^{i\theta} \circ x; (g, e^{i\theta}) \in G \times S^1, x \in \Omega_3 \right\}$$

by $\chi(g \circ e^{i\theta} \circ \tilde{x}'') := \chi(\tilde{x}'')$, for every $(g, e^{i\theta}) \in G \times S^1$. Assume that $\chi = 1$ on some neighborhood of V . Let $V_G = \{\pi(x); x \in V\}$. Let $\chi_1 \in C_0^\infty(Y_G)$ with $\chi_1 = 1$ on some neighborhood of V_G and $\text{Supp } \chi_1 \subset \{\pi(x) \in Y_G; x \in Y, \chi(x) = 1\}$. We have

$$\begin{aligned} \chi_1 \sigma_m &= m^{-\frac{d}{4}} \chi_1 S_{Y_G, m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x, n_-} \circ f \circ \iota^* \circ S_{G, m}^{(q)} \\ &= m^{-\frac{d}{4}} \chi_1 S_{Y_G, m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x, n_-} \circ f \circ \iota^* \circ \chi S_{G, m}^{(q)} \\ &\quad + m^{-\frac{d}{4}} \chi_1 S_{Y_G, m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x, n_-} \circ f \circ \iota^* \circ (1 - \chi) S_{G, m}^{(q)}. \end{aligned} \tag{6.29}$$

If $u \in Y$ but $u \notin \{x \in X; \chi(x) = 1\}$. Since $\text{Supp } \chi_1 \subset \{\pi(x) \in X; x \in Y, \chi(x) = 1\}$ and $\chi(x) = \chi(g \circ e^{i\theta} \circ x)$, for every $(g, e^{i\theta}) \in G \times S^1$, for every $x \in X$, we conclude that $\pi(e^{i\theta} \circ u) \notin \text{Supp } \chi_1$, for every $e^{i\theta} \in S^1$. From this observation and (6.20), we get

$$m^{-\frac{d}{4}} \chi_1 S_{Y_G, m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x, n_-} \circ f \circ \iota^* \circ (1 - \chi) S_{G, m}^{(q)} = O(m^{-\infty}) \text{ on } Y_G \times X. \tag{6.30}$$

From (6.29) and (6.30), we get

$$\chi_1 \sigma_m = m^{-\frac{d}{4}} \chi_1 S_{Y_G, m}^{(q-r)} \circ \iota_G \circ \hat{p} \circ \tau_{x, n_-} \circ f \circ \iota^* \circ \chi S_{G, m}^{(q)} + O(m^\infty) \text{ on } Y_G \times X.$$

From (6.22) and Theorem 1.8, we can check that on U ,

$$\chi_1 \sigma_m(\tilde{x}'', y) = (2\pi) \int e^{im\Psi(\tilde{x}'', \tilde{v}'') + im\Psi(v'', y)} \chi_1(\tilde{x}) \beta(\tilde{x}'', \tilde{v}'', m) \hat{b}(\tilde{v}'', y, m) dv(\tilde{v}'') + O(m^{-\infty}), \tag{6.31}$$

where $\hat{b}(\tilde{v}'', y, m) = (\iota_G \circ \hat{p} \circ \tau_{x, n_-} \circ f \circ \iota^* \circ \chi(\tilde{v}'') \circ b)(\tilde{v}'', y, m)$. From (6.31) and Theorem 6.3, we see that (6.25) and (6.26) hold near u . The theorem follows. \square

Let

$$F_m := \sigma_m^* \sigma_m : \Omega^{0, q}(X) \rightarrow H_{b, m}^q(X)^G, \quad \hat{F}_m := \sigma_m \sigma_m^* : \Omega^{0, q-r}(Y_G) \rightarrow H_{b, m}^{q-r}(Y_G).$$

Let $F_m(x, y)$ and $\hat{F}_m(x, y)$ be the distribution kernels of F_m and \hat{F}_m respectively. From Theorems 6.2 and 6.3, we can repeat the proof of Theorem 6.4 with minor change and deduce the following two theorems

Theorem 6.5 *With the notations used above, if $y \notin Y$, then for any open set D of y with $\overline{D} \cap Y = \emptyset$, we have $F_m = O(m^{-\infty})$ on $X \times D$.*

Let $x, y \in Y$. If $\pi(x) \neq \pi(e^{i\theta} \circ y)$, for every $\theta \in [0, 2\pi[$, then there are open sets D_1 of x in X and D_2 of y in X such that $F_m = O(m^{-\infty})$ on $D_1 \times D_2$.

Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4. Then,

$$\begin{aligned}
 F_m(x, y) &= e^{im\Psi(x'', y'')} a(x'', y'', m) + O(m^{-\infty}) \text{ on } U \times U, \\
 a(x'', y'', m) &\in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\
 a(x'', y'', m) &\sim \sum_{j=0}^{\infty} m^{n-\frac{d}{2}-j} a_j(x'', y'') \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\
 a_j(x'', y'') &\in C^\infty(U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \quad j = 0, 1, 2, \dots,
 \end{aligned}$$

and

$$a_0(\tilde{x}'', \tilde{x}'') = 2^{-3n+4d-1} \pi^{-n-1} \frac{1}{\text{Veff}(\tilde{x}'')} |\det \mathcal{L}_{\tilde{x}''}| |\det R_x|^{-\frac{1}{2}} \tau_{\tilde{x}'', n-}, \quad \forall \tilde{x}'' \in U \cap Y, \tag{6.32}$$

where U is an open set of p .

Theorem 6.6 Let $x, y \in Y$. If $\pi(x) \neq \pi(e^{i\theta} \circ y)$, for every $\theta \in [0, 2\pi[$, then there are open sets D_G of $\pi(x)$ in Y_G and V_G of $\pi(y)$ in Y_G such that $\hat{F}_m = O(m^{-\infty})$ on $D_G \times V_G$.

Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4. Then,

$$\begin{aligned}
 \hat{F}_m(x, y) &= e^{im\Psi(\tilde{x}'', \tilde{y}'')} \hat{a}(\tilde{x}'', \tilde{y}'', m) + O(m^{-\infty}) \text{ on } W \times W, \\
 \hat{a}(\tilde{x}'', \tilde{y}'', m) &\in S_{\text{loc}}^{n-d}(1; W \times W, T^{*0,q-r} Y_G \boxtimes (T^{*0,q-r} Y_G)^*), \\
 \hat{a}(\tilde{x}'', \tilde{y}'', m) &\sim \sum_{j=0}^{\infty} m^{n-d-j} \hat{a}_j(\tilde{x}'', \tilde{y}'') \text{ in } S_{\text{loc}}^{n-d}(1; W \times W, T^{*0,q-r} Y_G \boxtimes (T^{*0,q-r} Y_G)^*), \\
 \hat{a}_j(\tilde{x}'', \tilde{y}'') &\in C^\infty(W \times W, T^{*0,q-r} Y_G \boxtimes (T^{*0,q-r} Y_G)^*), \quad j = 0, 1, 2, \dots, \\
 \hat{a}_0(\tilde{x}'', \tilde{x}'') &= 2^{-3n+\frac{5}{2}d-1} \pi^{-n+\frac{d}{2}-1} |\det \mathcal{L}_{Y_G, \tilde{x}''}| \hat{\tau}_{\tilde{x}''}, \quad \forall \tilde{x}'' \in W
 \end{aligned}$$

where $W = \Omega_3 \times \Omega_4$, Ω_3 and Ω_4 are open sets as in the beginning of Sect. 6.1.

Let $R_m := \frac{1}{C_0} F_m - S_{G,m}^{(q)} : \Omega^{0,q}(X) \rightarrow H_{b,m}^q(X)^G$, where $C_0 = 2^{-3d+3n} \pi^{\frac{d}{2}}$. Since $F_m = F_m S_{G,m}^{(q)}$, it is clear that

$$\frac{1}{C_0} F_m = S_{G,m}^{(q)} + R_m = S_{G,m}^{(q)} + R_m S_{G,m}^{(q)} = (I + R_m) S_{G,m}^{(q)}. \tag{6.33}$$

Our next goal is to show that for m large, $I + R_m : \Omega^{0,q}(X) \rightarrow \Omega^{0,q}(X)$ is injective. From Theorem 6.5 and Theorem 1.8, we see that if $y \notin Y$, then for any open set D of y with $\overline{D} \cap Y = \emptyset$, we have

$$R_m = O(m^{-\infty}) \text{ on } X \times D. \tag{6.34}$$

Let $x, y \in Y$. If $\pi(x) \neq \pi(e^{i\theta} \circ y)$, for every $\theta \in [0, 2\pi[$, then there are open sets D_1 of x in X and D_2 of y in X such that

$$R_m = O(m^{-\infty}) \text{ on } D_1 \times D_2. \tag{6.35}$$

Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4. Then,

$$\begin{aligned}
 R_m(x, y) &= e^{im\Psi(x'', y'')} r(x'', y'', m) + O(m^{-\infty}) \text{ on } U \times U, \\
 r(x'', y'', m) &\in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\
 r(x'', y'', m) &\sim \sum_{j=0}^{\infty} m^{n-\frac{d}{2}-j} r_j(x'', y'') \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\
 r_j(x'', y'') &\in C^\infty(U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \quad j = 0, 1, 2, \dots.
 \end{aligned} \tag{6.36}$$

Moreover, from (6.32) and (1.10), it is easy to see

$$|r_0(x, y)| \leq C|(x, y) - (x_0, x_0)|, \tag{6.37}$$

for all $x_0 \in \mu^{-1}(0) \cap U$, where $C > 0$ is a constant. We need

Lemma 6.7 *Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4 defined in an open set U of p . Let*

$$\begin{aligned} H_m(x, y) &= e^{im\Psi(x'', y'')}h(x, y, m) \text{ on } U \times U, \\ h(x, y, m) &\in S_{\text{loc}}^{n-1-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ h(x, y, m) &\sim \sum_{j=0}^{\infty} m^{n-1-\frac{d}{2}-j}h_j(x, y) \text{ in } S_{\text{loc}}^{n-1-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ h_j(x, y) &\in C_0^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \quad j = 0, 1, 2, \dots \end{aligned}$$

Assume that $h(x, y, m) \in C_0^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*)$. Then, there is a constant $\hat{C} > 0$ independent of m such that

$$\|H_m u\| \leq \delta_m \|u\|, \quad \forall u \in \Omega^{0,q}(X), \quad \forall m \in \mathbb{N}, \tag{6.38}$$

where δ_m is a sequence with $\lim_{m \rightarrow \infty} \delta_m = 0$.

Proof Fix $N \in \mathbb{N}$. It is not difficult to see that

$$\|H_m u\| \leq \left\| (H_m^* H_m)^{2^N} u \right\|^{\frac{1}{2^{N+1}}} \|u\|^{1-\frac{1}{2^{N+1}}}, \quad \forall u \in \Omega^{0,q}(X), \tag{6.39}$$

where H_m^* denotes the adjoint of H_m . From Theorem 6.2, we can repeat the proof of Theorem 6.4 with minor change and deduce that

$$\begin{aligned} (H_m^* H_m)^{2^N}(x, y) &= e^{im\Psi(x'', y'')}p(x, y, m) + O(m^{-\infty}) \text{ on } U \times U, \\ p(x, y, m) &\in S_{\text{loc}}^{n-2^{N+1}-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ p(x, y, m) &\in C_0^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*). \end{aligned}$$

Hence,

$$|(H_m^* H_m)^{2^N}(x, y)| \leq \hat{C}m^{n-2^{N+1}-\frac{d}{2}}, \quad \forall (x, y) \in U \times U, \tag{6.40}$$

where $\hat{C} > 0$ is a constant independent of m . Take N large enough so that $n - 2^{N+1} - \frac{d}{2} < 0$. From (6.40) and (6.39), we get (6.38). \square

We also need

Lemma 6.8 *Let $p \in \mu^{-1}(0)$ and let $x = (x_1, \dots, x_{2n+1})$ be the local coordinates as in Remark 4.4 defined in an open set U of p . Let*

$$\begin{aligned} B_m(x, y) &= e^{im\Psi(x'', y'')}g(x, y, m) \text{ on } U \times U, \\ g(x, y, m) &\in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ g(x, y, m) &\sim \sum_{j=0}^{\infty} m^{n-\frac{d}{2}-j}g_j(x, y) \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \\ g_j(x, y) &\in C_0^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*), \quad j = 0, 1, 2, \dots, \\ g(x, y) &\in C_0^\infty(U \times U, T^{*0,q}X \boxtimes (T^{*0,q}X)^*). \end{aligned}$$

Suppose that

$$|g_0(x, y)| \leq C|(x, y) - (x_0, x_0)|,$$

for all $x_0 \in \mu^{-1}(0) \cap U$, where $C > 0$ is a constant. Then,

$$\|B_m u\| \leq \varepsilon_m \|u\|, \quad \forall u \in \Omega^{0,q}(X), \quad \forall m \in \mathbb{N}, \tag{6.41}$$

where ε_m is a sequence with $\lim_{m \rightarrow \infty} \varepsilon_m = 0$.

Proof Fix $N \in \mathbb{N}$. It is not difficult to see that

$$\|B_m u\| \leq \left\| (B_m^* B_m)^{2^N} u \right\|^{\frac{1}{2^{N+1}}} \|u\|^{1 - \frac{1}{2^{N+1}}}, \quad \forall u \in \Omega^{0,q}(X), \tag{6.42}$$

where B_m^* denotes the adjoint of B_m . From Theorem 6.2, we can repeat the proof of Theorem 6.4 with minor change and deduce that

$$\begin{aligned} (B_m^* B_m)^{2^N}(x, y) &= e^{im\Psi(x'', y'')} \hat{g}(x, y, m) + O(m^{-\infty}) \text{ on } U \times U, \\ \hat{g}(x, y, m) &\in S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\ \hat{g}(x, y, m) &\sim \sum_{j=0}^{\infty} m^{n-\frac{d}{2}-j} \hat{g}_j(x, y) \text{ in } S_{\text{loc}}^{n-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \\ \hat{g}_j(x, y) &\in C_0^\infty(U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \quad j = 0, 1, 2, \dots, \\ \hat{g}(x, y, m) &\in C_0^\infty(U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*), \end{aligned}$$

and

$$|\hat{g}_0(x, y)| \leq C |(x, y) - (x_0, x_0)|^{2^{N+1}}, \tag{6.43}$$

for all $x_0 \in \mu^{-1}(0) \cap U$, where $C > 0$ is a constant. Let

$$(B_m^* B_m)_0^{2^N}(x, y) = e^{im\Psi(x'', y'')} \hat{g}_0(x, y, m), \quad (B_m^* B_m)_1^{2^N}(x, y) = e^{im\Psi(x'', y'')} h(x, y, m),$$

where $h(x, y, m) = \hat{g}(x, y, m) - \hat{g}_0(x, y, m)$. It is clear that $h(x, y, m) \in S_{\text{loc}}^{n-1-\frac{d}{2}}(1; U \times U, T^{*0,q} X \boxtimes (T^{*0,q} X)^*)$. From Lemma 6.7, we see that

$$\left\| (B_m^* B_m)_1^{2^N} u \right\| \leq \delta_m \|u\|, \quad \forall u \in \Omega^{0,q}(X), \quad \forall m \in \mathbb{N}, \tag{6.44}$$

where δ_m is a sequence with $\lim_{m \rightarrow \infty} \delta_m = 0$.

From (6.43), we see that

$$|\hat{g}_0(x, y)| \leq C_1 \left(|\hat{x}''| + |\hat{y}''| + |\tilde{x}'' - \tilde{y}''| \right)^{2^{N+1}}, \tag{6.45}$$

where $C_1 > 0$ is a constant. From (3.44), we see that

$$|\text{Im } \Psi(x, y)| \geq c \left(|\hat{x}''|^2 + |\hat{y}''|^2 + |\tilde{x}'' - \tilde{y}''|^2 \right), \tag{6.46}$$

where $c > 0$ is a constant. From (6.45) and (6.46), we conclude that

$$|(B_m^* B_m)_0^{2^N}(x, y)| \leq \hat{C} m^{-2^N + n - \frac{d}{2}}, \quad \forall (x, y) \in U \times U, \tag{6.47}$$

where $\hat{C} > 0$ is a constant independent of m . From (6.47), we see that if N large enough, then

$$\left\| (B_m^* B_m)_0^{2^N} u \right\| \leq \hat{\delta}_m \|u\|, \quad \forall u \in \Omega^{0,q}(X), \quad \forall m \in \mathbb{N}, \tag{6.48}$$

where $\hat{\delta}_m$ is a sequence with $\lim_{m \rightarrow \infty} \hat{\delta}_m = 0$.

From (6.42), (6.44) and (6.48), we get (6.41). □

From (6.34), (6.35), (6.36), (6.37) and Lemma 6.8, we get

Theorem 6.9 *With the notations above, we have $\|R_m u\| \leq \varepsilon_m \|u\|$, $\forall u \in \Omega^{0,q}(X)$, $\forall m \in \mathbb{N}$, where ε_m is a sequence with $\lim_{m \rightarrow \infty} \varepsilon_m = 0$.*

In particular, if m is large enough, then the map $I + R_m : \Omega^{0,q}(X) \rightarrow \Omega^{0,q}(X)$ is injective.

Proof of Theorem 1.9 From (6.33) and Theorem 6.9, we see that if m is large enough, then the map $F_m = \sigma_m^* \sigma_m : H_{b,m}^q(X)^G \rightarrow H_{b,m}^q(X)^G$ is injective. Hence, if m is large enough, then the map $\sigma_m : H_{b,m}^q(X)^G \rightarrow H_{b,m}^{q-r}(Y_G)$ is injective and $\dim H_{b,m}^q(X)^G \leq \dim H_{b,m}^{q-r}(Y_G)$.

Similarly, we can repeat the proof of Theorem 6.9 with minor change and deduce that, if m is large enough, then the map $\hat{F}_m = \sigma_m \sigma_m^* : H_{b,m}^{q-r}(Y_G) \rightarrow H_{b,m}^{q-r}(Y_G)$ is injective. Hence, if m is large enough, then the map $\sigma_m^* : H_{b,m}^{q-r}(Y_G) \rightarrow H_{b,m}^q(X)^G$ is injective. Thus, $\dim H_{b,m}^q(X)^G = \dim H_{b,m}^{q-r}(Y_G)$ and σ_m is an isomorphism if m large enough. \square

References

1. Albert, C.: Le théorème de réduction de Marsden-Weinstein en géométrie cosymplectique et de contact. *J. Geom. Phys.* **6**, 627–649 (1989)
2. Boutet de Monvel, L., Sjöstrand, J.: Sur la singularité des noyaux de Bergman et de Szegő. *Astérisque* **34–35**, 123–164 (1976)
3. Bismut, J.-M., Lebeau, G.: Complex immersions and Quillen metrics. *Inst. Hautes Études Sci. Publ. Math.* **74**(1991), ii+298 (1992)
4. Boutet de Monvel, L., Guillemin, V.: The Spectral Theory of Toeplitz Operators. *Annals of Mathematics Studies*, Vol. 99. Princeton Univ. Press, Princeton (1981)
5. Braverman, M.: Cohomology of Mumford Quotient, Quantization of Singular Symplectic Quotients. *Programs in Mathematics*, 198, pp. 47–59. Birkhäuser, Basel (2001)
6. Catlin, D.: The Bergman Kernel and a Theorem of Tian. *Analysis and Geometry in Several Complex Variables* (Katata, 1997), Trends in Mathematics, pp. 1–23. Birkhäuser, Boston (1999)
7. Charles, L.: Toeplitz operators and Hamiltonian torus actions. *J. Funct. Anal.* **236**(1), 299–350 (2006)
8. Engliš, M.: Weighted Bergman kernels and quantization. *Comm. Math. Phys.* **227**(2), 211–241; MR1903645. *Zbl* **1010**, 32002 (2002)
9. Geiges, H.: Constructions of contact manifolds. *Math. Proc. Camb. Philos. Soc.* **121**(3), 455–464 (1997)
10. Guillemin, V., Sternberg, S.: Geometric quantization and multiplicities of group representations. *Invent. Math.* **67**(3), 515–538 (1982)
11. Guillemin, V.: Star products on compact pre-quantizable symplectic manifolds. *Lett. Math. Phys.* **35**(1), 85–89 (1995)
12. Hörmander, L.: The Analysis of Linear Partial Differential Operators. I, *Classics in Mathematics*. Springer, Berlin (2003)
13. Hsiao, C.-Y.: Projections in Several Complex Variables, *Mémoires Society Mathematics. France, Nouv. Sér. Vol. 123*, p. 131 (2010)
14. Hsiao, C.-Y., Marinescu, G.: Asymptotics of spectral function of lower energy forms and Bergman kernel of semi-positive and big line bundles. *Comm. Anal. Geom.* **22**(1), 1–108 (2014)
15. Hsiao, C.-Y., Marinescu, G.: On the singularities of the Szegő projections on lower energy forms. *J. Differ. Geom.* **107**(1), 83–155 (2017)
16. Kohn, J.J.: The range of the tangential Cauchy-Riemann operator. *Duke Math. J.* **53**(2), 307–562 (1986)
17. Loose, F.: A remark on the reduction of Cauchy-Riemann manifolds. *Math. Nachr.* **214**, 39–51 (2000)
18. Ma, X., Marinescu, G.: The first coefficients of the asymptotic expansion of the Bergman kernel of the $spin^c$ Dirac operator. *Int. J. Math.* **17**(6), 737–759 (2006)
19. Ma, X., Marinescu, G.: *Holomorphic Morse Inequalities and Bergman Kernels*, Progress in Mathematics, 254. Birkhäuser, Basel (2007)
20. Ma, X., Zhang, W.: Bergman kernels and symplectic reduction. *Astérisque* **318**, 154 (2008)
21. Ma, X., Marinescu, G.: Generalized Bergman kernels on symplectic manifolds. *Adv. Math.* **217**(4), 1756–1815 (2008)
22. Ma, X.: Geometric Quantization on Kähler and Symplectic Manifolds, *International Congress of Mathematicians*, vol. II, Hyderabad, India, August 19–27, pp. 785–810 (2010)
23. Ma, X., Zhang, W.: Geometric quantization for proper moment maps: the Vergne conjecture. *Acta Math.* **212**(1), 11–57 (2014)

24. Meinrenken, E.: Symplectic surgery and the Spinc-Dirac operator. *Adv. Math.* **134**(2), 240–277 (1998)
25. Melin, A., Sjöstrand, J.: Fourier integral operators with complex-valued phase functions. *Springer Lect. Notes Math.* **459**, 120–223 (1975)
26. Paoletti, R.: Moment maps and equivariant Szegő kernels. *J. Symplectic Geom.* **2**(1), 133–175 (2003)
27. Puchol, M.: G-invariant holomorphic Morse inequalities. *J. Differ. Geom.* (**to appear**). [arXiv:1506.04526](https://arxiv.org/abs/1506.04526)
28. Rossi, H.: Attaching analytic spaces to an analytic space along a pseudoconcave boundary. In: *Proceedings of the Conference on Complex Manifolds (Minneapolis)*, pp. 242–256. Springer, New York (1965)
29. Shiffman, B., Zelditch, S.: Distribution of zeros of random and quantum chaotic sections of positive line bundles. *Comm. Math. Phys.* **200**, 661–683 (1999)
30. Teleman, C.: The quantization conjecture revisited. *Ann. Math. (2)* **152**(1), 1–43 (2000)
31. Tian, Y., Zhang, W.: An analytic proof of the geometric quantization conjecture of Guillemin-Sternberg. *Invent. Math.* **132**(2), 229–259 (1998)
32. Zelditch, S.: Szegő kernels and a theorem of Tian. *Int. Math. Res. Not.* **6**, 317–331 (1998)
33. Zhang, W.: Holomorphic quantization formula in singular reduction. *Commun. Contemp. Math.* **1**(3), 281–293 (1999)

Publisher's Note Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.